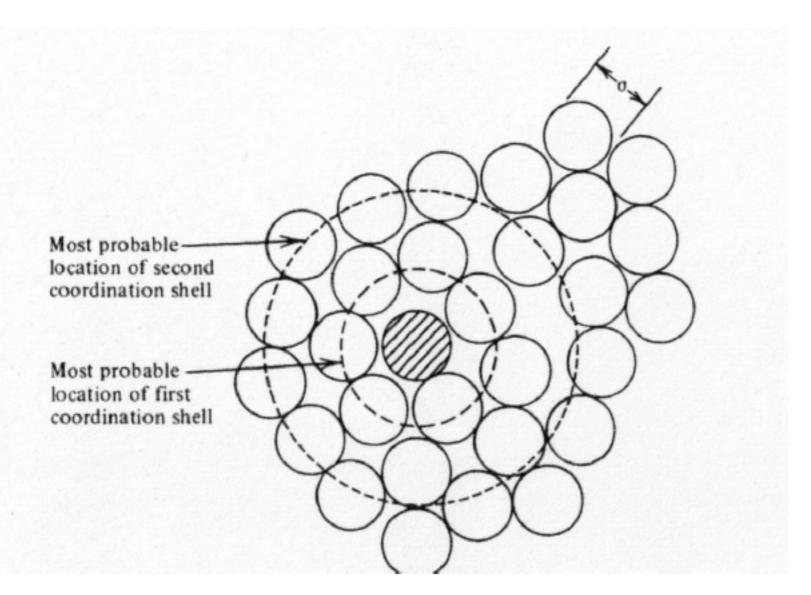
# Lecture 11

CHM695 Feb. 12

### Radial Distribution Function

$$g(r) = \frac{(N-1)}{4\pi\rho r^2} \langle \delta(r-r') \rangle_{r',\theta',\phi',\mathbf{R}}$$

(also called pair distribution function)



It measures the probability of finding two particles at a distance r in NVT ensemble

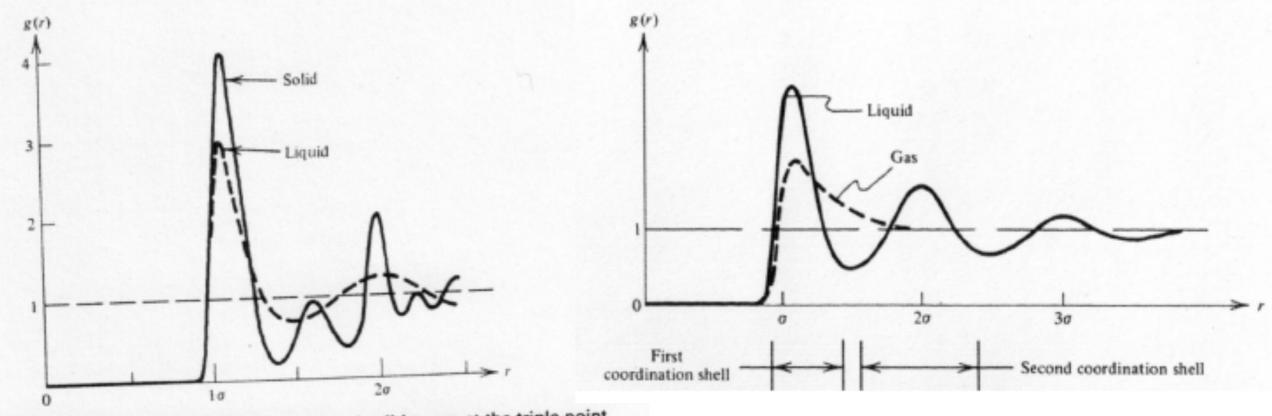


Fig. 7.5. Radial distribution functions for liquid and solid argon at the triple point  $(\sigma = 3.4 \text{ Å})$ .

#### **Coordination Number:**

$$N_c = 4\pi\rho \int_0^{r_{\min}} r^2 g(r) dr$$

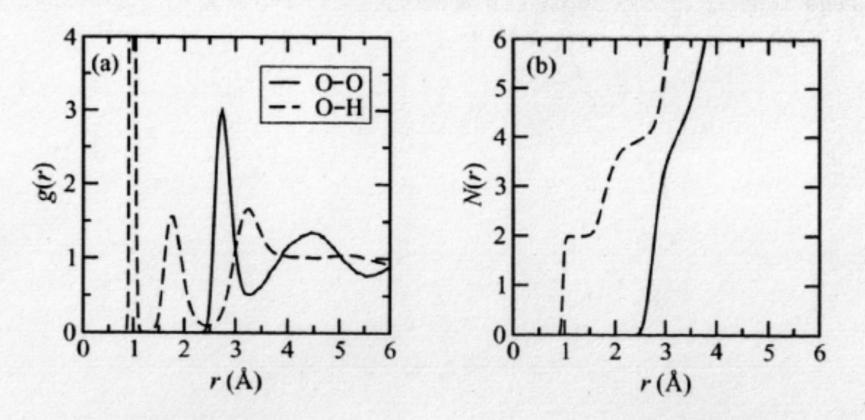
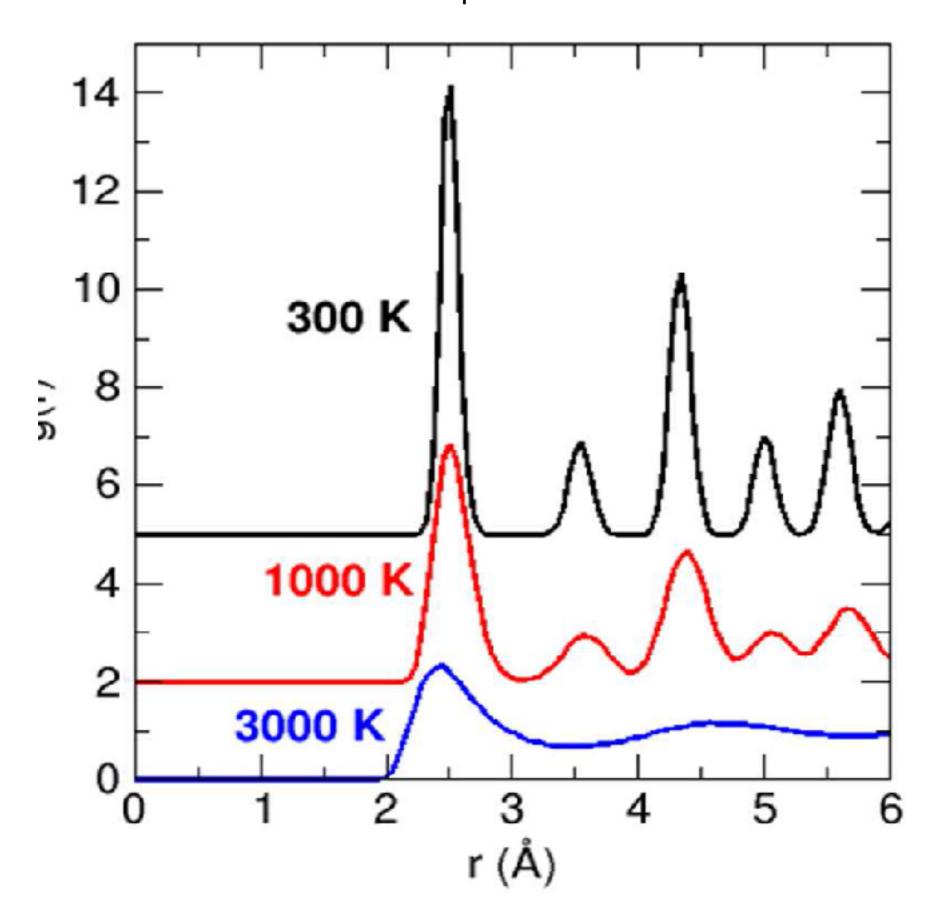
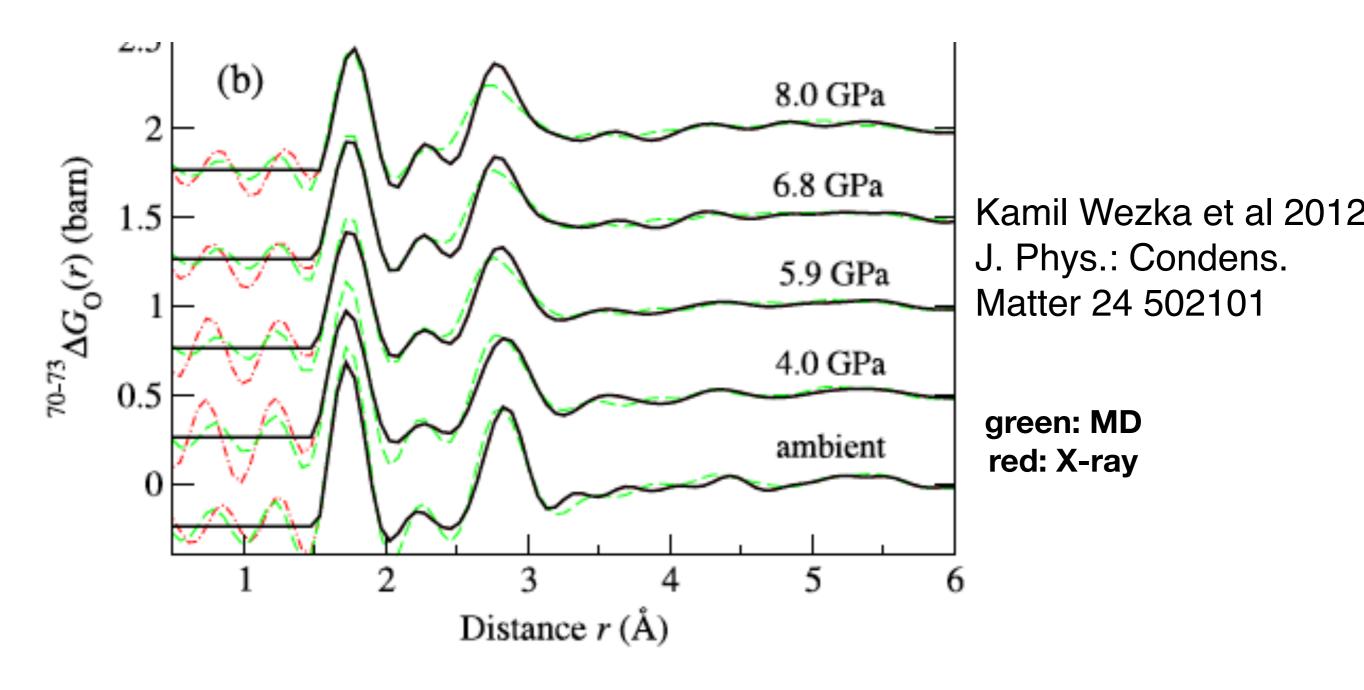


Fig. 4.3 (a) Oxygen-oxygen (O-O) and oxygen-hydrogen (O-H) radial distribution functions for a particular model of water (Lee and Tuckerman, 2006; Marx et al., 2010). (b) The corresponding running coordination numbers computed from eqn. (4.6.23).

For a system at different temperatures, the following g(r) is obtained: interpret these results.



Pair distribution function can be also extracted from XRD or neutron scattering (inverse Fourier transform)



## Structure Factor

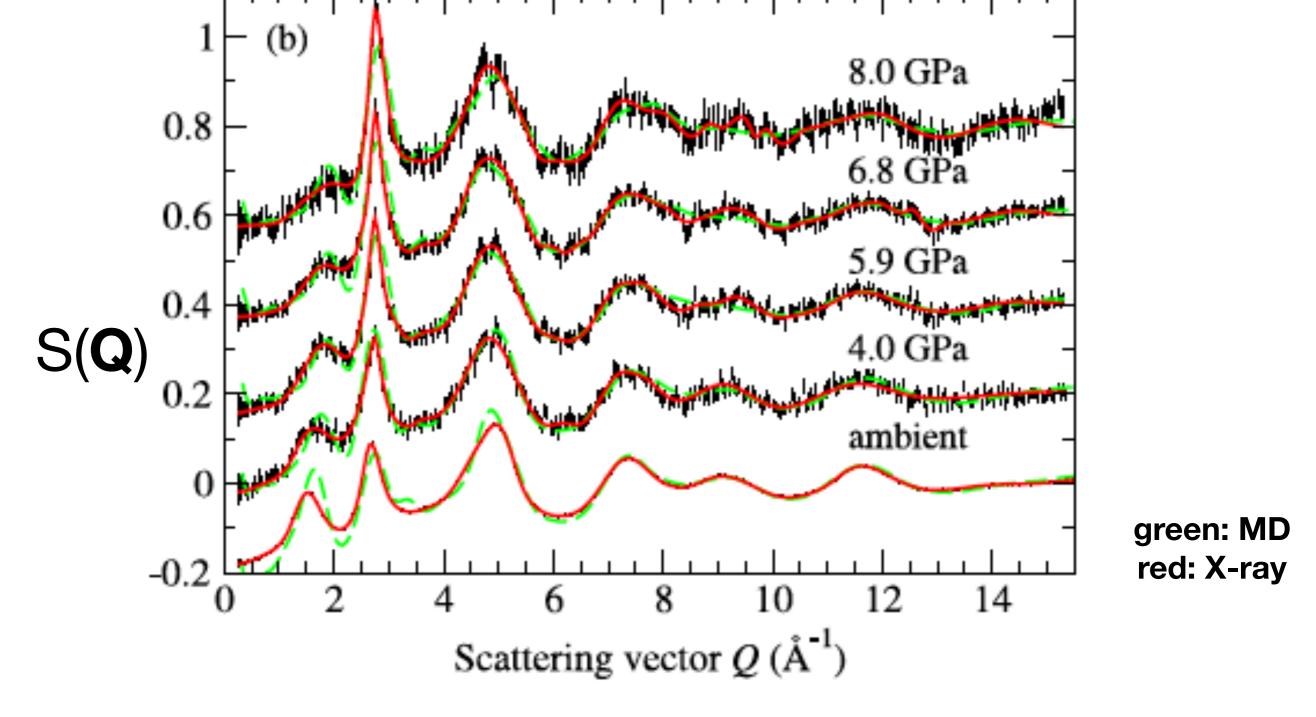
$$S(\mathbf{k}) = 1 + \frac{1}{N} \int d\mathbf{r} \int d\mathbf{r}' \exp(-i\mathbf{k} \cdot [\mathbf{r} - \mathbf{r}']) \rho(\mathbf{r}, \mathbf{r}')$$

$$\rho\left(\mathbf{r},\mathbf{r}'\right) = \left\langle \sum_{I}^{N} \sum_{J}^{N} \delta\left(\mathbf{r} - \mathbf{R}_{I}\right) \delta\left(\mathbf{r}' - \mathbf{R}_{J}\right) \right\rangle$$

For isotropic systems, this can be simplified to

$$S(k) = 1 + 4\pi\rho \int_0^\infty dr \, r^2 g(r) \frac{\sin(kr)}{kr}$$

sine transform of g(r)



oxygen structural factors in GeO2 computed from gO<sub>0</sub>(r)

Kamil Wezka et al 2012 J. Phys.: Condens. Matter 24 502101 For algorithm of g(r) see Frenkel & Smit

# Advanced Techniques

# Multiple Timestep Integrators

#### **Classical evolution operators:**

Consider a property  $a(\mathbf{X})$  that can be evaluated as a function of time.

Using chain rule:

$$\frac{da}{dt} = \sum_{\alpha=1}^{3N} \left[ \frac{\partial a}{\partial q_\alpha} \dot{q}_\alpha + \frac{\partial a}{\partial p_\alpha} \dot{p}_\alpha \right] \qquad \text{poisson bracket}$$

$$\dot{q}_{\alpha} = \frac{\partial H}{\partial p_{\alpha}}$$

$$\dot{p}_{\alpha} = -\frac{\partial H}{\partial q_{\alpha}}$$

We know that 
$$\dot{q}_{\alpha}=\dfrac{\partial H}{\partial p_{\alpha}}$$
 and  $\dot{p}_{\alpha}=-\dfrac{\partial H}{\partial q_{\alpha}}$  Liouville operator 
$$\dfrac{da}{dt}=\sum_{\alpha=1}^{3N}\left[\dfrac{\partial a}{\partial q_{\alpha}}\dfrac{\partial H}{\partial p_{\alpha}}-\dfrac{\partial a}{\partial p_{\alpha}}\dfrac{\partial H}{\partial q_{\alpha}}\right]=\{a,H\} =iLa$$

$$iL = \sum_{\alpha=1}^{3N} \left[ \frac{\partial H}{\partial p_{\alpha}} \frac{\partial}{\partial q_{\alpha}} - \frac{\partial H}{\partial q_{\alpha}} \frac{\partial}{\partial p_{\alpha}} \right]$$

$$\frac{da}{dt} = iLa$$

has the formal solution:

$$a(\mathbf{X}_t) = \exp(iLt)a(\mathbf{X}_0)$$

$$iL_1 = \sum_{\alpha=1}^{3N} \frac{\partial H}{\partial p_{\alpha}} \frac{\partial}{\partial q_{\alpha}} \qquad iL_2 = -\sum_{\alpha=1}^{3N} \frac{\partial H}{\partial q_{\alpha}} \frac{\partial}{\partial p_{\alpha}}$$

$$iL = iL_1 + iL_2$$

$$iL_1 iL_2 f(\mathbf{X}) \neq iL_2 iL_1 f(\mathbf{X})$$

#### **Trotter's Theorem**

$$\exp(A+B) = \lim_{P \to \infty} \left[ \exp(B/2P) \exp(A/P) \exp(B/2P) \right]^{P}$$

$$\exp(iLt) = \exp(iL_1t + iL_2t)$$

$$= \lim_{P \to \infty} \left[ \exp(iL_2t/2P) \exp(iL_1t/P) \exp(iL_2t/2P) \right]^P$$

Let us define:  $\Delta t = t/P$  Then,  $P \to \infty \Rightarrow \Delta t \to 0$ 

For small  $\Delta t$ 

 $\exp(iLt/P) \equiv \exp(iL\Delta t) \approx [\exp(iL_2\Delta t/2) \exp(iL_1\Delta t) \exp(iL_2\Delta t/2)]$ 

$$e^{iL\Delta t}x(0) = x(\Delta t)$$

$$e^{iL\Delta t}p(0) = p(\Delta t)$$

$$e^{iL\Delta t} \approx \exp\left(\frac{\Delta t}{2}F(x)\frac{\partial}{\partial p}\right) \exp\left(\Delta t \frac{p}{m}\frac{\partial}{\partial x}\right) \exp\left(\frac{\Delta t}{2}F(x)\frac{\partial}{\partial p}\right)$$

$$\exp\left(c\frac{d}{dx}\right)g(x) = g(x+c)$$

#### First operator:

$$\exp\left(\frac{\Delta t}{2}F(x(0))\frac{\partial}{\partial p}\right) \left(\begin{array}{c} x(0) \\ p(0) \end{array}\right) = \left(\begin{array}{c} x(0) \\ p(0) + \frac{\Delta t}{2}F(x(0)) \end{array}\right)$$

#### Second operator:

$$\exp\left(\Delta t \frac{p_{(0)}\partial}{m} \frac{\partial}{\partial x}\right) \begin{pmatrix} x(0) \\ p(0) + \frac{\Delta t}{2} F(x(0)) \end{pmatrix} = \begin{pmatrix} x(0) + \Delta t \frac{p_{(0)}}{m} \\ p(0) + \frac{\Delta t}{2} F(x(0)) + \Delta t \frac{p_{(0)}}{m} \end{pmatrix}$$

Third operator:

$$\exp\left(\frac{\Delta t}{2}F(x(0))\frac{\partial}{\partial p}\right) \left(\begin{array}{c} x(0) + \frac{\Delta t}{m}p(0) \\ p(0) + \frac{\Delta t}{2}F\left(x(0) + \Delta t\frac{p(0)}{m}\right) \end{array}\right)$$

$$= \left( \frac{x(0) + \frac{\Delta t}{m} \left[ p(0) + \frac{\Delta t}{2} F(x(0)) \right]}{p(0) + \frac{\Delta t}{2} F(x(0)) + \frac{\Delta t}{2} F(x(0)) + \frac{\Delta t}{2} F(x(0)) \right]} \right)$$

$$x(\Delta t) = x(0) + \frac{\Delta t}{m} \left[ p(0) + \frac{\Delta t}{2} F(x(0)) \right]$$
$$= x(0) + v(0) \Delta t + \frac{\Delta t^2}{2m} F(x(0))$$

$$p(\Delta t) = p(0) + \frac{\Delta t}{2} F(x(0)) + \frac{\Delta t}{2} F\left(x(0) + \frac{\Delta t}{m} \left[p(0) + \frac{\Delta t}{2} F(x(0))\right]\right)$$

$$= p(0) + \frac{\Delta t}{2} F(x(0)) + \frac{\Delta t}{2} F(x(\Delta t))$$

$$v(\Delta t) = v(0) + \frac{\Delta t}{2m} \left[F(x(0)) + F(x(\Delta t))\right]$$

These are the velocity Verlet update of positions and velocities!

$$F(x) = F_{\text{fast}}(x) + F_{\text{slow}}(x)$$

$$iL = \frac{p}{m} \frac{\partial}{\partial x} + [F_{\text{fast}}(x) + F_{\text{slow}}(x)] \frac{\partial}{\partial p}$$

$$iL_{\text{fast}} = \frac{p}{m} \frac{\partial}{\partial x} + F_{\text{fast}}(x) \frac{\partial}{\partial p}$$

$$iL_{\text{slow}} = F_{\text{slow}}(x) \frac{\partial}{\partial p}$$

$$\exp(iL\Delta t) = \exp(iL_{\text{slow}}\Delta t/2) \exp(iL_{\text{fast}}\Delta t) \exp(iL_{\text{slow}}\Delta t/2)$$

factorise this time step to smaller time steps

$$\exp(iL_{\text{fast}}\Delta t) = \left[\exp(\frac{\delta t}{2}F_{\text{fast}}\frac{\partial}{\partial p})\exp(\delta t \frac{p}{m}\frac{\partial}{\partial x})\exp(\frac{\delta t}{2}F_{\text{fast}}\frac{\partial}{\partial p})\right]^n$$

#### Reference System Propagator Algorithm (RESPA)

# Molecular dynamics algorithm for multiple time scales: Systems with long range forces

Mark E. Tuckerman<sup>a)</sup> and Bruce J. Berne Department of Chemistry, Columbia University, New York, New York 10027

Glenn J. Martyna

Department of Chemistry, University of Pennsylvania, Philadelphia, Pennsylvania 19104

The Journal of Chemical Physics **94**, 6811 (1991); <a href="https://doi.org/10.1063/1.460259">https://doi.org/10.1063/1.460259</a>

```
do j=1,n_outer
 call calculate slow force
 p=p+0.5*dt large*F slow
  do i=1,n
    p=p+0.5*dt_small*F_fast
    x=x+dt_small*p/m
     call calculate fast force
    p=p+0.5*dt small*F fast
   end do
 p=p+0.5*dt_large*F_slow
```

end do

### Home work

#### 1. For a harmonic oscillator

$$U(x,y) = \frac{1}{2}k_1x^2 + \frac{1}{2}k_2y^2$$

$$k_1 = 10 k_2$$

- a) integrate EOMs using one timestep (which is of course decided by k1)
- b) integrate EOMs using multiple timesteps note that k1 > k2, thus the force on x is fast and that on y is slow