

# Gravitational Lensing Expressions

September 6, 2016

## 1 —

In the following I use O, S and L to denote the Observer, Source and Lens positions, respectively. P denotes any point in the ray of light we deal with. The lensing equation is

$$\frac{d^2u}{d\theta^2} + u = \frac{3}{2}r_g u^2 \quad (1)$$

where  $u \equiv 1/r$ , the radius  $r \equiv LP$ , and  $\theta \equiv \angle SLP$ . The gravitational radius  $r_g$  is the Schwarzschild radius

$$r_g = \frac{2MG}{c^2} \quad (2)$$

where  $c$  is the speed of light and  $M$  the lens mass.

For thin lens approximation I use a couple of definitions: the "breaking point" (B) is a point on the incoming ray at which  $\angle SBL = \pi/2$ , the incoming ray angle  $\phi \equiv \angle BSL$  and the impact parameter  $b \equiv BL$ . Given some  $\phi$ , the breaking point B and the impact parameter  $b$  are calculated, the incoming ray is deflected at B (stays the same before B) such that  $\phi_{afterB} = \phi - 2r_g/b$ . This assumes that the impact parameter  $b$  is much larger than the gravitational radius  $r_g$  (this is the thin lens approximation assumption)

Calculating the time of travel is done as follows:

$$\Delta t_{ray} = \int_{Ray} \frac{dp}{c} \left(1 - \frac{r_g}{r(p)}\right)^{-\frac{1}{2}} \quad (3)$$

the square root term is the regular gravitational time dilation.

For extended masses I think changing  $r_g$  to depend on the inner mass at P, the numerics should work just fine for both thin and thick lenses. The metric outside a static spherical symmetric mass doesn't care of the exact inner structure.

Thoughts, comments?