HOW PREVALENT ARE GALAXIES WITH SIMULTANEOUS STAR FORMATION AND BLACK HOLE GROWTH AT COSMIC NOON? PREDICTIONS FOR THE JAMES WEBB SPACE TELESCOPE

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1. INTRODUCTION

The galaxies most actively contributing to the buildup of stellar mass at cosmic noon $(z \sim 1-3)$ contain large amounts of dust. This dust obscures the majority of star formation, making it necessary to study these galaxies in infrared wavelengths (Madau & Dickinson 2014). Additionally, the majority of supermassive black hole growth at these redshifts is also heavily dust obscured (e.g., Hickox & Markevitch 2007). Many of the massive dusty galaxies contain a true mix of star formation and obscured black hole growth, the obscured signatures of which can be seen in their infrared spectral energy distribution (SED), making these galaxies ideal laboratories for understanding the physical link between star formation and active galactic nuclei (AGN). The AGN-star formation connection is an open question, particularly whether AGN feedback is a key component of star formation quenching, and whether all galaxies have a distinct star formation phase followed by an AGN phase before ultimately quenching (e.g. Sanders et al. 1988; Hopkins et al. 2006). The nature of AGN within strongly star forming galaxies (what we term "composites") is even more uncertain. Do these objects represent a unique phase between star forming galaxies and AGN? Unfortunately, due to limitations of previous space telescopes, detailed studies of the energetics of these objects were extremely limited, but the James Webb Space Telescope (JWST) will reveal their true nature.

Prior to JWST, the most reliable method for identifying composites and disentangling AGN emission from star formation was mid-IR spectroscopy from the Spitzer Space Telescope. The low resolution spectra can be modeled as a combination of star formation features (most notably the polycyclic aromatic hydrocarbons, or PAHs, that exist in photodissociation regions and in stellar/HII regions), and hot continuum emission primarily arising

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from a dusty torus surrounding the accreting black hole (Pope et al. 2008; Coppin et al. 2010; Kirkpatrick et al. 2012; Sajina et al. 2012; Hernán-Caballero et al. 2015; Kirkpatrick et al. 2015). In this way, a precise fraction of IR luminosity due to star formation and AGN can be quantified. The medium-resolution spectrometer (MRS Wells et al. 2015), which is part of the Mid-Infrared Instrument (MIRI) on JWST, will be enable separation of PAH emission from continuum in the same manner, but with higher resolution and on smaller spatial scales within host galaxies. It will also enable detection of high ionization gas lines excited by the AGN (Bonato et al. 2017), further improving our ability to detect and measure the physical properties (such as black hole accretion rate and Eddington ratios) of dust obscured black holes.

As there are only a few hundred Spitzer IRS spectra available for distant galaxies (Kirkpatrick et al. 2015), color techniques were also developed to identify large samples of luminous dust obscured AGN. The most popular color selection techniques are with Spitzer IRAC photometry (Lacy et al. 2004; Stern et al. 2005; Alonso Herrero et al. 2006; Donley et al. 2012), which separate AGN using different combinations of the 3.6, 4.5, 5.8, and $8.0 \,\mu m$ filters. The original techniques presented in Lacy et al. (2004) and Stern et al. (2005) were limited to the most luminous AGN and become increasingly contaminated with host galaxies when deeper IR data are used (Mendez et al. 2013). Moreover, with increasing redshift, the rest wavelengths of these bands decrease, causing contamination of the AGN regions by star forming galaxies to become significant so the original IRACbased criteria cannot be applied. Donley et al. (2012) propose more conservative IRAC criteria that, at cosmic noon, essentially separates galaxies which exhibit a so-called stellar bump (1.6 μ m emission from stars) from those which do not, where the torus radiation is strong enough to fill in the dip in the star forming spectrum around $3-5 \mu m$, producing power-law emission such as is typical of unobscured AGN (e.g. Elvis et al. 1994). The Donley et al. (2012) criteria increases the reliability of AGN color selection, although it is less complete due to excluding composites, where the IR emission of the AGN does not dominate over the star formation.

For the purposes of probing the AGN-star formation connection, the limitation of IRAC techniques is that AGN within strongly star forming galaxies can have different levels of host contamination. Then, many galaxies containing AGN signatures at longer wavelengths will also include a stellar bump and therefore be missed by IRAC techniques (Kirkpatrick et al. 2013, 2015). To alleviate host contamination, Messias et al. (2012) propose combining K-band with IRAC and $24 \,\mu \text{m}$ to separate AGN from host galaxies all the way out to $z \sim 7$.

Going further, including mid-IR and far-IR colors can greatly improve the selection of composite galaxies, since this will trace the contribution of warmer AGN-heated dust compared with cold dust from the diffuse interstellar medium in the host galaxy (Kirkpatrick et al. 2015). However, this requires observations from the Herschel Space Observatory, which have a large beam size and do not reach the same depths as Spitzer observations. MIRI will greatly improve color selection techniques due to the increased sensitivity and the number of transmission filters covering the mid-infrared (Bouchet et al. 2015; Glasse et al. 2015). Now, we will be able to separate AGN from SFGs by comparing PAH emission with the minimum emission from stars that occurs around $5 \mu m$: in AGN, the stellar minimum is not visible due to strong torus emission, and composites will lie in between strong AGN and pure SFGs in colorspace.

In this paper, we build on the Herschel and Spitzer color selection techniques initially presented in Kirkpatrick et al. (2013) to identify composite galaxies at $z \sim 1-2$ using the CANDELS and COSMOS survey. We present number counts of $24\,\mu\mathrm{m}$ sources classified as SFGs, AGN, or composites based on their IR colors, making this the first identified statistical sample of composites at cosmic noon. We use this sample to predict black hole and star formation properties of samples that JWST/MIRI will identify. We also present color diagnostics for identifying both AGN and composites using JWST/MIRI filters in three redshift bins. Throughout this paper, we assume a standard cosmology with $H_0 = 70\,\mathrm{km\,s^{-1}\,Mpc^{-1}},\,\Omega_\mathrm{M} = 0.3,\,\mathrm{and}\,\Omega_\Lambda = 0.7.$

2. CANDELS AND COSMOS CATALOGS

To calculate number counts, we use Spitzer and Herschel photometry from the COSMOS, EGS, GOODS-S, and UDS fields from the Cosmic Assembly NearIR Deep Extragalactic Survey (CANDELS, P.I. S. Faber and H. Ferguson; GOODS-Herschel, P.I. D. Elbaz; CANDELS-Herschel, P. I. M. Dickinson). We do not include GOODS-N as, at the time of the writing of this paper, the IR catalog does not have uniquely identified optical counterparts. We also use photometric redshifts (Dahlen et al. 2013) and M_* (Santini et al. 2015). We will be using the photometric redshifts to classify sources as AGN, SFGs, or composites. The CANDELS photo-zs are the median redshift determined through five separate codes that fit templates to the UV/optical/near-IR data. Taking the median of several methods improves the accuracy, and comparison of photo-zs with spectroscopic redshifts for a limited sample gives $\sigma = 0.03$ where σ is the rms of $(z_{phot} - z_{spec})/(1 + z_{spec})$. As we sort sources into redshift bins of $\Delta z = 0.5$, we do not expect the uncertainty on the photometric redshifts to be a dominant source of uncertainty in our results.

MIPS $24\,\mu\mathrm{m}$ and Herschel PACS and SPIRE catalogs were built following the prior-based PSF fitting method described in Pérez-González et al. (2005, MIPS photometry) and Pérez-González et al. (2010, merged MIPS plus Herschel photometry). See also Rawle et al. (2016) for more details about the method followed for the Herschel catalogs. Briefly, the algorithm uses IRAC and MIPS data to extract photometry for sources in longer wavelength data using positional priors. We remark that the method considers a minimum separation between sources

introduced as position priors below which debelending is not possible. This separation is approximately 75% of the FWHM of the PSF in each band. The final product of the cataloging method is a list of IRAC sources with possible counterparts in all longer wavelength data. In this sense, several IRAC sources might be identified with the same MIPS or *Herschel* source. This is what we call multiplicity. The multiplicity for MIPS and PACS is in more than 95% of the cases equal to 1 (i.e., only one IRAC source is identified with a single MIPS and PACS source), but it is higher for SPIRE (on average, 6 IRAC sources are found within the FWHM of the SPIRE $250 \,\mu\mathrm{m}$ PSF). In order to identify the "right" IRAC counterpart for each far-IR sources, we follow the method described in Rodríguez-Muñóz et al. (2017, in prep). In practice, we choose the MIPS most probable counterpart as the brightest IRAC candidate. Then, we shift this methodology to longer wavelength bands. We identify the most likely PACS counterpart as the brightest source in MIPS $24 \,\mu\mathrm{m}$ among the different candidates. When MIPS is not available, we use the reddest IRAC band in which the source is detected. We note that using IRAC as a tracer of PACS emitters can lead to spurious identifications. For this reason, these cases are flagged in order to evaluate the possible impact in the results. Finally, we use the fluxes in PACS or MIPS (if PACS is not available) to find the counterpart of the SPIRE sources. The flux of each FIR source is assigned to a single IRAC counterpart.

We have also added sources from the COSMOS survey (Scoville 2007) which is necessary to boost the bright end of the number counts, due to the small survey area of CANDELS $(0.22 \,\mathrm{deg^2})$. We use the public COSMOS2015 catalog in Laigle et al. (2016), which presents multiwavelength data as well as stellar masses and photometric redshifts. The Spitzer IRAC data in this catalog originally comes from SPLASH COSMOS and S-COSMOS (Sanders et al. 2007) while the MIPS $24 \mu m$ observations are described in Le Floc'h et al. (2009). The catalog also contains *Herschel* observations from the PEP guaranteed time program (Lutz et al. 2011) and the HERMES consortium (Oliver et al. 2012). The counterpart identification and procedures for measuring stellar masses and photometric redshifts are fully described in Laigle et al. (2016).

The difficulty in matching MIPS, PACS, and SPIRE sources with their IRAC counterparts underscores the improvements that will be made by using MIRI color selection to identify AGN host galaxies, since the much smaller PSF (<1'' for all filters) and smaller spectral range used will obviate the need for counterpart identification for robust color diagnostics.

3. IR IDENTIFICATION OF AGN AND COMPOSITES

To identify SFGs, composites, and AGN, we build on the color techniques in Kirkpatrick et al. (2013, 2015) that sample the full IR spectral energy distribution (SED). At $z \sim 1-2$, the color $S_8/S_{3.6}$ separates sources with a strong stellar bump, present in SFGs, from those with hot torus emission, found in AGN. Composites span a range in this color, depending on the ratio of relative strengths of the AGN and host galaxy emission and the amount of obscuration of the AGN due to dust. In fact, heavily obscured AGN such as Mrk 231, NGC 1068, the

Circinus galaxy, and IRAS 08572+3915 have SEDs that drop rapidly from $10 \,\mu\mathrm{m}$ toward shorter wavelengths and will show the near IR stellar spectral peak characteristic of SFGs. Hereafter, we refer to 'AGN' with the understanding that the samples discussed may suffer from incompleteness of sources like these. This issue is discussed further in Section 4.1. S_{100} and S_{250} trace the peak of the IR SED, which is generally dominated by the cold dust in the diffuse ISM. S_{24} traces the PAH emission in SFGs or the warm dust emission heated by the AGN. Then, the color S_{250}/S_{24} or S_{100}/S_{24} will measure the relative amounts cold emission to warm dust or PAH emission, and this ratio is significantly higher in SFGs. However, significant scatter is introduced into color selection by redshift, since S_{24} will move over different PAH features and the 9.7 μ m silicate absorption, changing where SFGs lie in color space. We can more robustly identify SFGs, AGN, and composites if we introduce a redshift criterion.

The color diagnostics (S_{250}/S_{24}) vs. $S_8/S_{3.6}$ and S_{100}/S_{24} vs. $S_8/S_{3.6}$) were calibrated with a sample of 343 galaxies with Spitzer IRS spectroscopy spanning the range $z \sim 0.5 - 4$. We identified SFGs, composites, and AGN through spectral decomposition, where we fit the mid-IR spectrum (5–18 μ m restframe) with a model consisting of PAH features for star formation, a power-law continuum for the AGN, and extinction. We then quantified the AGN emission, $f(AGN)_{MIR}$, as the fraction of mid-IR luminosity $(5-15 \,\mu\text{m})$ due to the power-law continuum. We define three classes of galaxies based on $f(AGN)_{MIR}$, and we also report the fraction of MIR luminosity due solely to emission from the PAH features in the $5-15\,\mu\mathrm{m}$ range: (1) SFGs are dominated by PAH emission $(f(AGN)_{MIR} < 0.2, L_{PAH}/L_{MIR} > 0.6);$ (2) AGN have negligible PAH emission $(f(AGN)_{MIR} > 0.8,$ $L_{\rm PAH}/L_{\rm MIR}$ < 0.15); (3) Composites have a mix of PAH and continuum emission ($f(AGN)_{MIR} = 0.2 - 0.8$, $L_{PAH}/L_{MIR} = 0.15 - 0.6$). We note that below, we will redefine these thresholds for color selection. We relate the mid-IR classification to the full IR SED by creating empirical templates using data from Spitzer and Herschel. We sort sources into subsamples based on $f(AGN)_{MIR}$, and after normalization, determine the median L_{ν} in differential bin sizes of λ (Kirkpatrick et al. 2012, 2015). The Kirkpatrick et al. (2015) SEDs are the first comprehensive public library of IR templates specifically designed for high redshift galaxies that account for AGN emission.

We create a redshift dependent color diagnostic through use of the empirical MIR-based template Library from Kirkpatrick et al. (2015). We use a template library because our spectroscopic sample of 343 sources is not large enough to separate sources into multiple z bins. The MIR-based Library contains 11 templates created from our spectroscopic sources that demonstrate the change in IR spectral shape as the contribution of the AGN to the mid-IR luminosity increases, in steps of $\Delta f(\text{AGN})_{\text{MIR}} = 0.1$. We randomly redshift each tem-

plate 500 times, uniformly sampling a redshift distribution from z=0.75-2.25. For each redshifted template, we convolve it with the observed frame IRAC, MIPS, PACS, and SPIRE transmission filters to create photometry, and then we resample the photometry within the template uncertainties at that particular wavelength, following a Gaussian distribution. We now have a catalog of 5500 synthetic galaxies, where we know the intrinsic AGN contribution, that represent the scatter in colorspace of real galaxies.

Next, we create color diagrams in redshift bins of z = 0.75 - 1.25, z = 1.25 - 1.75, and z = 1.75 - 2.25. Beyond this redshift, it becomes too difficult to reliably separate composites from SFGs with these colors. Because only a fraction of CANDELS and COSMOS sources have a SPIRE or PACS detection, we also create a color diagnostic using the colors S_{24}/S_8 vs. $S_8/S_{3.6}$, although this is slightly less accurate. In each redshift bin, we divide the color space into regions of $0.2 \times 0.2 \,\mathrm{dex}$ and calculate the average $f(AGN)_{MIR}$ and standard deviation, σ_{AGN} of all the synthetic galaxies that lie in that region. In the Appendix, we show our three diagnostics: $S_{250}/S_{24}v.S_8/S_{3.6}$ (used when a galaxy has the appropriate photometry, as it is the most complete at selecting composite galaxies), $S_{100}/S_{24}vs.S_8/S_{3.6}$ (used when a galaxy doesn't have a $250 \,\mu\mathrm{m}$ detection), and $S_{24}/S_8vs.S_8/S_{3.6}$ (used for all galaxies without a longer wavelength detection).

Our color diagnostics assign sources an $f(AGN)_{MIR}$ in bins of $\Delta f(AGN)_{MIR} = 0.1$, but the σ_{AGN} of each region is often larger than this (see the Appendix for a visual representation). Therefore, it is more accurate to broadly group sources as SFGs, Composites, and AGN. We determine how to group sources by comparing the $f(AGN)_{MIR}$ assigned to each synthetic galaxy by the three different color diagnostics. There is a good correlation, and the residuals have a standard deviation of 0.15. Accordingly, we classify as SFGs sources with $f(AGN)_{MIR} < 0.30$, while the AGN have $f(AGN)_{MIR} > 0.70$, and Composites are everything in between. We assess the completeness and reliability of our color technique by determining how many of our synthetic galaxies are correctly identified as SFGs, Composites, and AGN in each diagnostic, and we list the completeness and reliability in Table 1. Completeness is defined as the fraction of AGN (for example) selected: n_{AGN}/N_{AGN} , and reliability is the fraction of all the sources selected by the diagnostic as AGN (for example) that actually are, intrinsically, AGN: $n_{\rm AGN}/n_{\rm all}$). The lower completeness and reliability of the Composites and SFGs is due to these sources being more easily confused with each other with the limited coverage of the mid-IR provided by Spitzer. By adding more bands, MIRI will allow for a more nuanced measurement of the strength of the PAH emission compared with continuum and stellar bump emission. It is also important to note that we are missing AGN with extreme obscuration, whose IR colors could mimic those of SFGs. We discuss this issue more fully in Section 4.1.

We assign each CANDELS or COSMOS source with z=0.75-2.25 an $f(AGN)_{MIR}$ and associated uncertainty (σ_{AGN}) and then broadly group sources into SFGs, Composites, and AGN. Overall, from CANDELS (COSMOS), 534 (6426) sources have been classified with $S_{250}/S_{24} v. S_8/S_{3.6}$, 864 (175) with $S_{100}/S_{24} v. S_8/S_{3.6}$,

 $^{^{10}}$ There are many AGN templates in the literature. In the $1-20\,\mu\mathrm{m}$ (rest wavelength) range critical for most of our color sorting the AGN templates agree well (Lyu & Rieke 2017). At wavelengths longer than $20\,\mu\mathrm{m}$, there is considerable divergence; fortunately for our goals, the star forming output is so dominant by 100 and $250\,\mu\mathrm{m}$ that the range of possibilities for AGN output has little effect on our results.

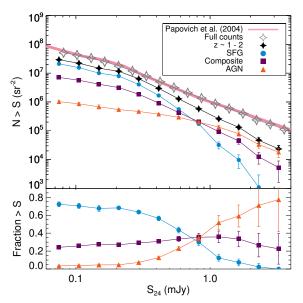


Fig. 1.— Top panel—Cumulative 24 $\mu \rm m$ number counts for the CANDELS and COSMOS fields. The open grey stars show the number counts of all 24 $\mu \rm m$ detected sources, and these agree with the published number counts (Papovich et al. 2004, pink solid line). The filled stars show the 24 $\mu \rm m$ counts from z=0.75-2.25. These are incomplete, particularly at the bright end, which is due to the relatively small field sizes. We then show the contribution of SFGs $(f(\rm AGN)_{\rm MIR} < 0.3$; blue circles), Composites $(0.3 \ge f(\rm AGN)_{\rm MIR} < 0.7$; purple squares), and AGN $(f(\rm AGN)_{\rm MIR} \ge 0.7$; orange triangles) to the $z\sim 1-2$ number counts. Bottom panel—The percentage of each subsample above a given flux threshold. At $S_{24} > 1\,\rm mJy$, Composites and AGN dominate samples. Even at fainter fluxes, JWST/MIRI samples will contain at least 20% identifiable Composites.

and 871 (5360) with $S_{24}/S_8 v. S_8/S_{3.6}$. From CANDELS, we also fit an additional 111 sources, which lie slightly beyond the regions (within 0.2 dex) in our color classification scheme, with the Kirkpatrick et al. (2015) template library to determine the classification.

3.1. Galaxy and AGN Counts

Now, we determine how traditional $24 \,\mu\mathrm{m}$ number counts break down into the SFG, Composite, and AGN categories. We only consider sources with $S_{24} > 80 \,\mu\text{Jy}$, which is the 80% completeness limit (Pérez-González et al. 2005). We measure directly the EGS, COSMOS, GOODS-S, and UDS field sizes covered by our sources. We show the total CANDELS+COSMOS $24 \,\mu\mathrm{m}$ number counts as the open grev stars in Figure 1, and these counts are in agreement with the counts from (Papovich et al. 2004). We plot the $24 \,\mu\mathrm{m}$ counts at cosmic noon (z = 0.75 - 2.25) as the filled black stars. The incompleteness from applying the redshift cut chiefly affects the bright end $(\mathring{S}_{24} > 1 \, \text{mJy})$, which is where the AGN will dominate the counts (Kirkpatrick et al. 2013). This incompleteness is a result of the small field sizes of CANDELS $(0.22 \,\mathrm{deg^2})$ and COSMOS $(1.5 \,\mathrm{deg^2})$. We show how the cosmic noon counts break down into SFGs (blue), Composites (purple), and AGN (orange). We have calculated uncertainties on the counts using a Monte Carlo technique, where we vary the $f(AGN)_{MIR}$ for each source within its associated uncertainty and recount sources. We follow this procedure 1000 times. The counts in Figure 1 represent the mean from the Monte Carlo simulations, and the error bars encompass

the standard deviation from the Monte Carlo trials and the standard Poisson errors.

Below 0.8 mJy, SFGs dominate the counts, but AGN become more prevalent with increasing brightness. In the bottom panel of Figure 1, we show the percentage of sources above a given flux threshold. Although AGN are frequently assumed not to be abundant in fainter IR samples, the presence of AGN hosts at $S_{24} < 100 \,\mu\text{Jy}$ was also seen in a small Spitzer/IRS spectroscopic sample of lensed galaxies at $z \sim 2$, where the authors found that 30% of the sample had IR AGN signatures and 40% had X-ray AGN signatures (Rigby et al. 2008). The Composites comprise >25\% of a sample down to the faintest flux threshold at 63% completeness, which we determined by applying the completeness estimates listed in Table 1 to the number of sources classified with each method. Then, at least 25% of a JWST/MIRI sample will be Composite galaxies, providing a rich data set for probing the AGN/star formation connection at cosmic noon.

4. JWST COLOR SELECTION

Color selection is a powerful technique for identifying likely AGN, Composites, and SFGs. We have done an exhaustive search to identify the best MIRI filter combinations for separating galaxies into these three classes at cosmic noon by creating synthetic photometry in the JWST/MIRI filters from the Kirkpatrick et al. (2015) MIR based Library following the Monte Carlo technique outlined in Section 3. As many JWST/MIRI observations will be carried out in fields with available photometric redshifts, or in parallel with NIRcam and NIRspec observations, we include redshift information in our color diagnostics to improve reliability and completeness. We identify three diagnostics covering the ranges $z \sim 1 \ (z = 0.75 - 1.25), \ z \sim 1.5 \ (z = 1.25 - 1.75),$ and $z \sim 2$ (z = 1.75 - 2.25). These three diagnostics, shown in Figure 2, are different combinations of the $S_{21}, S_{18}, S_{15}, S_{12.8}, S_{10}$, and $S_{7.7}$ filters, which cover the 6.2 and 7.7 μ m PAH complexes and the 3 – 5 μ m stellar minimum at these redshifts.

We present two methods for separating SFGs, Composites, and AGN. First, we have determined the optimal AGN, Composite, and SFG regions, labeled in Figure 2. The boundaries of each region are circles, with AGN lying inside the inner circle, SFGs lying outside the outer circle, and Composites lying in between.

The $z \sim 1$ boundaries are

inner:
$$(\log \frac{S_{15}}{S_{7.7}} - 0.40)^2 + (\log \frac{S_{18}}{S_{10}} - 0.38)^2 = 0.25^2$$
 (1)

outer:
$$(\log \frac{S_{15}}{S_{7.7}} - 0.35)^2 + (\log \frac{S_{18}}{S_{10}} - 0.45)^2 = 0.65^2$$
(2)

(3)

TABLE 1
COMPLETENESS (RELIABILITY) OF REDSHIFT DEPENDENT
COLOR SELECTION

Region	$z \sim 1$	$z \sim 1.5$	$z \sim 2$
	S_{250}/S_2	$v. S_8/S_{3.6}$	
AGN	93 (85)%	92 (89)%	97 (86)%
Composite	67 (67)%	81 (63)%	56 (66)%
SFG	66 (77)%	41 (75)%	64 (64)%
	S_{100}/S_{2}	$v. S_8/S_{3.6}$	
AGN	97 (89)%	94 (89)%	98 (81)%
Composite	69 (71)%	76(64)%	42 (59)%
SFG	67 (75)%	46 (69)%	62 (57)%
	S_{24}/S_{8}	$v. S_8/S_{3.6}$	
AGN	93 (85)%	90 (88)%	97 (83)%
Composite	69 (65)%	67 (66)%	48 (82)%
SFG	60 (77)%	61 (65)%	87 (68)%

Completeness is defined as the percentage of sources of a given intrinsic classification that are also selected by the color diagnostic. Reliability (shown in parenthesis) is the percentage of all sources classified in a given category where the intrinsic classification agrees.

The $z \sim 1.5$ boundaries are

inner:
$$(\log \frac{S_{21}}{S_{10}} - 0.49)^2 + (\log \frac{S_{18}}{S_{12.8}} - 0.18)^2 = 0.21^2$$
(4)

outer:
$$(\log \frac{S_{21}}{S_{10}} - 0.60)^2 + (\log \frac{S_{18}}{S_{12.8}} - 0.03)^2 = 0.65^2$$
(5)

(6)

The $z \sim 2$ boundaries are

inner:
$$(\log \frac{S_{18}}{S_{10}} - 0.43)^2 + (\log \frac{S_{21}}{S_{15}} - 0.18)^2 = 0.18^2$$

outer:
$$(\log \frac{S_{18}}{S_{10}} - 0.50)^2 + (\log \frac{S_{21}}{S_{15}} - 0.12)^2 = 0.52^2$$
(8)

(9)

These regions are useful for broadly classifying large numbers of sources or identifying targets for followup observations. We use these regions to asses the reliability and completeness of our color diagnostic, where again, we classify all synthetic sources as SFGs when $f(AGN)_{MIR} < 0.3$, Composites where $0.3 \ge f(AGN)_{MIR} < 0.7$, and AGN when $0.7 \ge f(AGN)_{MIR}$. Table 2 list these values for all three redshift regimes. Comparison with Table 1 shows an improvement over what we were able to reliably classify with the *Herschel* and Spitzer diagnostics, particularly for separating Composites from SFGs. The spacing of the MIRI filters allows us to sensitively trace the strength of the PAH features relative to the stellar minimum, where the proportionate amount of PAH emission will be lower for Composite galaxies as the power-law emission from the AGN begins to outshine the stellar minimum (see the insets in Figure 2 for a visual guide).

Perhaps, instead of broad classifications, the reader would rather have an estimate of $f(AGN)_{MIR}$. With-

TABLE 2 Completeness (Reliability) of MIRI color selection

Region	$z \sim 1$	$z \sim 1.5$	$z \sim 2$
AGN	87 (90)%	87 (89)%	87 (80)%
Composite	77 (72)%	79 (74)%	71 (71)%
SFG	76 (81)%	81 (86)%	79 (89)%

out mid-IR spectroscopy, robust decomposition into an AGN and star forming component still is not feasible, even with 6 photometry filters. However, we have determined how to linearly combine the colors in each redshift regime in order to estimate $f(AGN)_{MIR}$, and we also measure the standard deviation (σ_{AGN}) of the residuals when each equation is applied to our synthetic sources so that the reader has a measure of the uncertainty. At $z \sim 1$

$$f(AGN)_{MIR} = -0.97 \times (\log \frac{S_{15}}{S_{7.7}}) - 0.10 \times (\log \frac{S_{18}}{S_{10}}) + 1.29$$
(10)

and $\sigma_{AGN} = 0.15$.

$$f(\text{AGN})_{\text{MIR}} = -0.56 \times (\log \frac{S_{21}}{S_{10}}) - 0.85 \times (\log \frac{S_{18}}{S_{12.8}}) + 1.29 \tag{11}$$

and $\sigma_{AGN} = 0.13$.

$$f(\text{AGN})_{\text{MIR}} = -0.55 \times (\log \frac{S_{18}}{S_{10}}) - 1.01 \times (\log \frac{S_{21}}{S_{15}}) + 1.25$$
(12)

and $\sigma_{AGN} = 0.16$.

4.1. Mid-IR concerns: Metallicity and Obscuration

At cosmic noon, the bulk of the star formation is occurring in massive, dusty galaxies with $M_* > 10^{10}~M_{\odot}$ (e.g. Murphy et al. 2011; Madau & Dickinson 2014; Pannella et al. 2015), which is the type of galaxies that our MIRI

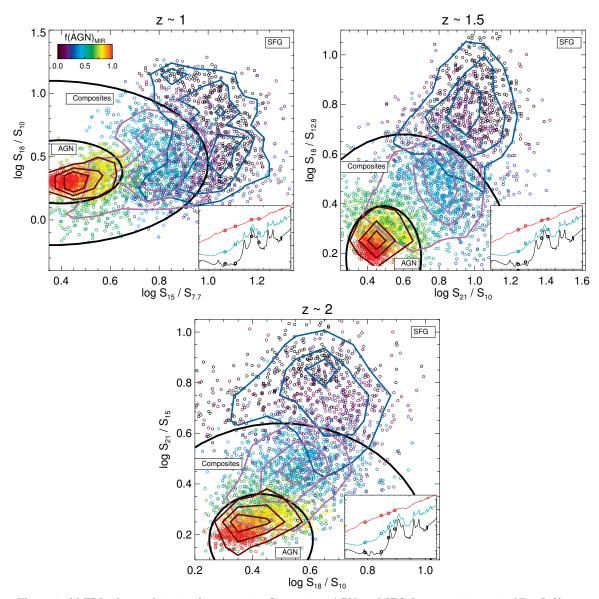


Fig. 2.— The optical MIRI color combination for separating Composites, AGN, and SFG from z=0.75-1.25 ($Top\ Left$), z=1.25-1.75 ($Top\ Right$), and z=1.75-2.25 (Bottom). We show the synthetic galaxies (shaded according to $f(AGN)_{MIR}$) from the Kirkpatrick et al. (2015) used to determine the best AGN and Composite selection regions (black lines), based on completeness and reliability. We also overplot the contours of all the synthetic galaxies classified as SFG (blue lines), Composites (purple lines), and AGN (red lines) to allow easier viewing of where each category predominantly lies. In the bottom right corner of each diagram, we demonstrate where the photometry filters fall on an SFG (black), Composite (blue), and AGN (red) template at z=1,1.5,2.

diagnostics were created from (Kirkpatrick et al. 2012; Sajina et al. 2012; Kirkpatrick et al. 2015). For studying the AGN-star formation connection, we expect these types of galaxies to form the most appealing targets. Nevertheless, the sensitivity of JWST/MIRI will enable studies of lower mass galaxies, which tend to have lower metallicties (Ma et al. 2016, and references therein). Decreasing gas phase metallicities have been linked with decreasing PAH strengths (e.g. Engelbracht et al. 2008; Sandstrom et al. 2012; Shivaei et al. 2017), which is a source of concern since we are effectively detecting AGN hosts based on the strength of PAH features compared with the stellar minimum at $3-5\,\mu\mathrm{m}$. Shipley et al. (2016) find that below $Z < 0.7 Z_{\odot}$, PAH emission no longer scales linearly with $L_{\rm IR}$, which based on the massmetallicity relation, could be a source of concern for contamination of our Composite regions at $M_* < 3 \times 10^9 M_{\odot}$,

up to $z\sim 2.3$ (Erb et al. 2006; Zahid et al. 2013; Sanders et al. 2015). Recently, using the MOSDEF optical spectroscopic survey, Shivaei et al. (2017) found that at $z \sim 2$, $L_{7.7}/L_{\rm IR}$ is lower for galaxies with $M_* < 10^{10} M_{\odot}$ with a behavior similar to that seen for local galaxies (?Shipley et al. 2016). At $z \sim 2$, an hour of integration time at $21 \,\mu\text{m}$ will produce 10σ detections of galaxies at roughly $8 \mu \text{Jy}$. At $z \sim 2$, $21 \mu \text{m}$ is tracing the $7.7 \mu \text{m}$ PAH feature, so applying Equation 11 from Shipley et al. (2016) gives a star formation rate of $12 M_{\odot}/\text{yr}$, and assuming this is a main sequence galaxy, $M_* \sim 3 \times 10^9 M_{\odot}$, which is below the threshold where we expect low metallicity galaxies might contaminate the Composite regime. As such, our color diagnostics may require recalibration for low metallicity galaxies when using observations below $S_{21} \lesssim 30 \,\mu\text{Jy}.$

As a visual check, we demonstrate in Figure 3 where

SFGs with different PAH strengths will lie in our $z \sim 1.5$ diagnostic. To accomplish this, we use the Small Magellanic Cloud (SMC) dust model (PAH fraction $q_{PAH} =$ 0.10%) and a Milky Way dust model with $q_{PAH} = 0.47\%$ from Draine & Li (2007), which is included to show where a galaxy with a low SFR will lie. The Draine & Li (2007) models are also parameterized in terms of the strength of the radiation field, U_{\min} and U_{\max} . We set these values to $U_{\min} = 1$ and $U_{\max} = 1e5$, although these parameters have little effect on the final colors. Also, we note that we add in a stellar blackbody with $T = 5000 \,\mathrm{K}$ to complete the near-IR portion of the spectrum. Even with a low PAH fraction, the Milky Way template still lies in our SFG region, while the SMC template lies directly on the Composite/SFG border. We also plot the template from Chary & Elbaz (2001) which corresponds to an $L_{\rm IR}=10^{10}\,L_{\rm IR}$, as this is an order of magnitude less luminous than the Kirkpatrick et al. (2015) library. A galaxy of this luminosity also lies in the Composite region.

If stellar masses of MIRI samples are known (possibly through NIRcam observations), low mass galaxies that lie in our Composite regions provide excellent targets for follow-up spectroscopy observations, to distinguish between AGN or metallicity as the underlying cause of the diminished PAH emission.

The other prominent concern in a mid-IR diagnostic is how obscuration can affect the detection of AGN. Our template library was built assuming the AGN can be represented as a power law, and we empirically measure the power law component to have an average slope of $F_{\nu} \propto \lambda^{1.5}$, but individual sources will show a range of slopes, and a range of dust obscurations. The AGN templates in the Kirkpatrick et al. (2015) library is derived from AGN where 75% of the sample is also detected in the X-ray, implying that they are largely unobscured. Of the composite sources in Kirkpatrick et al. (2015), 35% are X-ray detected, indicating that they contain more heavily obscured AGN. We now explore the effects of dust obscuration by examining where different galaxies will lie in the $z \sim 1.5$ colorspace (Figure 3).

Arp 220 (orange bowtie) is a local Ultra Luminous Infrared Galaxy (ULIRG) that is heavily dust obscured and may host an AGN (Veilleux et al. 2009; Teng et al. 2015). Its position near the SMC and at the edge of the Composite region indicates another possible ambiguity, that the aromatic bands tend to be suppressed in the most luminous and compact infrared galaxies. How many such objects exist at cosmic noon is not well quantified, as most galaxies of the same luminosity $(L_{\rm IR}>10^{12}\,L_{\odot})$ have extended ISMs (Younger et al. 2009; Ivison et al. 2012; Rujopakarn et al. 2011, 2016). NGC 1068 (red cross) is an archetypal local Compton thick Seyfert II AGN. Despite its extreme obscuration, it lies securely in our Composite region, close to the AGN boundary. We also use the AGN library of Siebenmorgen et al. (2015) to examine what extinction conditions would push an AGN into our SFG region. These AGN templates are calculated assuming the AGN IR emission arises from 2-phase dust region consisting of a torus and disk, a torus radius R, viewing angle, and cloud filling factor. The optical depth of the clouds in the torus is primarily what causes the AGN to move into the Com-

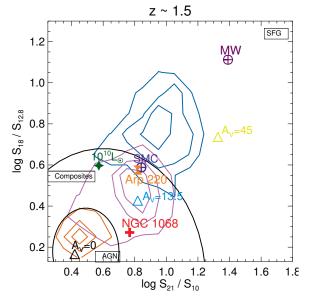


Fig. 3.— We use our $z\sim 1.5$ MIRI diagnostic to explore where sources with different luminosities, metallicities, and obscurations than the Kirkpatrick et al. (2015) library will lie in colorspace; we show the distribution of our synthetic SFGs (blue), Composites (purple), and AGN (orange) as well as the black lines marking the AGN, Composite, and SFG regions. We use the Draine & Li (2007) library to calculate where an SMC galaxy (dark purple circle with cross) and MW galaxy with a lower PAH fraction (purple circle with cross) will lie. The SMC galaxy lies in the Composite region, as well as a galaxy with $L_{\rm IR}=10^{10}\,L_{\odot}$ (dark green star) from the Chary & Elbaz (2001) library. Our diagnostics were calibrated using galaxies with $L_{\rm IR} > 10^{11} L_{\odot}$ and $M_* > 10^{10} M_{\odot}$, so they should be applied with caution to lower mass, lower luminosity objects. We also look at what affect obscuration may have on our ability to detect AGN using local obscured ULIRG Arp 220 (orange bowtie) and Compton Thick AGN NGC 1068 (red cross), both of which lie in the Composite Region. We also use the AGN template library (triangles) of Siebenmorgen et al. (2015) to show how obscuration in the torus (measured in A_V) will push AGN into the Composite and SFG region.

posite and SFG regions, so we hold all other parameters fixed (for reference, we use the model with viewing angle= 67° , $R=1545\times 10^{15}$ cm, $A_d=300, V_c=77.7\%$). This model is a pure AGN, with no star formation, but when the optical depth of the torus is $A_V = 13.5$ (blue triangle), the AGN model lies in our Composite region, and when $A_V = 45$ (yellow triangle), the AGN lies in the SFG region. Detecting such an obscured AGN at other wavelengths would also be extremely challenging. Currently, identifying true Type II obscured AGN requires measuring X-ray column densities or classification with optical spectra. By comparison with deep *Chandra* 6 Ms data in the GOODS-S field Mendez et al. (2013) estimate that only 10% of IR identified AGN will be heavily obscured. In Kirkpatrick et al. (2015), we found that 65% of objects identified as Composites were not detected in the same Chandra 6 Ms data, indicating heavy obscuration. But beyond these estimates, it is difficult to say how many heavily obscured AGN there are that would not be selected as such in the X-ray or the mid-IR. Identifying these very obscured AGN will require detailed SED modeling using a full suite of NIRcam+MIRI observations, which is beyond the scope of this paper.

If we have a good understanding of the typical full IR SED of high redshift galaxies, as well as the scatter in the population, then a single photometric point can be used in conjunction with representative templates to estimate $L_{\rm IR}$ and star formation rates (SFRs). Since PAH molecules are illuminated by the UV/optical photons from young stars, they are a natural SFR indicator and have been extensively used in the literature to probe SFR and $L_{\rm IR}$ (Peeters et al. 2004; Brandl et al. 2006; Pope et al. 2008; Battisti et al. 2015; Shipley et al. 2016).

Given the coverage of the MIRI filters, we will now examine how an AGN can affect the 7.7 μ m PAH feature for the Kirkpatrick et al. (2015) templates used in this work, as any AGN contribution will need to be corrected for before converting a PAH luminosity to a SFR. We remind the reader that for these templates, the AGN component is represented as a power-law with a slope of $F_{\nu} \propto \lambda^{1.5}$. We measure the intrinsic $L_{7.7}$ of each template using PAHFIT (Smith et al. 2007). Then, we measure $L_{\rm MIRI}$, which is the photometry of the template through the following MIRI filters at the given redshifts:

$$z = 0, 7.7 \,\mu\text{m}$$

$$z = 0.95, 15.0 \,\mu\text{m}$$

$$z = 1.34, 18.0 \,\mu\text{m}$$

$$z = 1.73, 21.0 \,\mu\text{m}$$

$$z = 2.31, 25.5 \,\mu\text{m}$$
(13)

The redshifts mark where the rest frame central wavelength of each filter is $7.7\,\mu\mathrm{m}$.

In the top panel of Figure 4, we demonstrate how much of the 7.7 μ m feature each filter covers at the above listed redshifts. In the bottom panel, we show the relationship $L_{7.7}/L_{\rm MIRI}$ as a function of $f({\rm AGN})_{\rm MIR}$ for each filter at the listed redshifts. The decreasing fractions with increasing $f({\rm AGN})_{\rm MIR}$ are due to the increased contribution of the warm dust continuum to the measured photometry. We fit a quadratic relationship to all the points and measure

$$\frac{L_{7.7}}{L_{\text{MIRI}}} = (-1.09 \pm 0.20) \times f(\text{AGN})_{\text{MIR}}^{2}
- (0.50 \pm 0.21) \times f(\text{AGN})_{\text{MIR}}
+ (1.86 \pm 0.04)$$
(14)

This equation, in conjunction with estimating $f(AGN)_{MIR}$ from MIRI colors, can be used for first order corrections to $L_{7.7}$ before converting to a SFR. Similarly, in Kirkpatrick et al. (2015), we demonstrated that there is a quadratic relationship between $f(AGN)_{MIR}$ and the total contribution of an AGN to L_{IR} that can be used to correct L_{IR} for AGN emission:

$$f(AGN)_{IR} = 0.66 \times f(AGN)_{MIR}^2 - 0.035 \times f(AGN)_{MIR}$$
(15)

where $f(AGN)_{IR}$ is the fraction of $L_{IR}(8-1000\,\mu\text{m})$ due to AGN heating. Then, the portion of L_{IR} due to star formation is $L_{IR}^{SF} = L_{IR} \times (1 - f(AGN)_{IR})$. Once the AGN contribution is accounted for, L_{IR} can be converted to a SFR using standard equations (e.g., Murphy et al. 2011). For a strong AGN $(f(AGN)_{MIR} \ge 0.9)$, at least 50% of L_{IR} needs to be removed before converting to a SFR, and the same is true if using 7.7 μ m to calculate SFR. Then, the strongest AGN will have SFRs that are

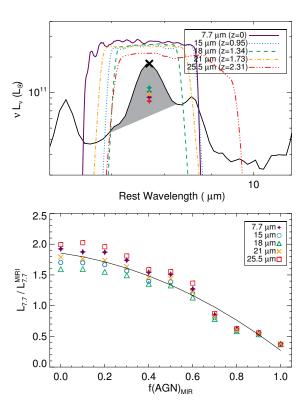


Fig. 4.— Top panel-We demonstrate, using the MIR0.0 template from Kirkpatrick et al. (2015), how much of the 7.7 μ m feature each MIRI filter covers at the redshifts where the central wavelength of each filter is $7.7 \,\mu\mathrm{m}$ (listed in the legend). The grey shaded region indicates which part of the spectrum is integrated to calculated $L_{7.7}$ (black cross), while the filled stars show the photometry (in units of νL_{ν}) measured through each of the MIRI filters. The photometry is lower because it includes more of the spectrum at lower luminosity. Bottom panel—We show the ratio of the MIRI photometry $(L_{
m MIRI})$ measured through different MIRI filters depending on redshift to the intrinsic $L_{7.7}$ of the PAH feature. The lower fractions with increasing $f(AGN)_{MIR}$ are due to an increased warm dust continuum due to heating by an AGN. If $L_{7,7}$ is going to be used to calculate SFRs, corrections need to be made for an AGN contribution. The black line is the empirical relationship between $L_{\text{MIRI}}/L_{7.7}$ and $f(\text{AGN})_{\text{MIR}}$.

overestimated by at least a factor of 2. Of more concern is Composites, which are routinely misidentified as SFGs. For a Composite with $f(\text{AGN})_{\text{MIR}} = 0.5$, an L_{IR} based SFR will be overestimated by $\sim 15\%$. But, if one uses L_{MIRI} to approximate $L_{7.7}$, then the resulting SFR will be overestimated by $\sim 35\%$.

5. DISCUSSION: PHYSICAL PROPERTIES OF A MIRI SAMPLE

We now return to our CANDELS + COSMOS sample to investigate the physical properties of galaxies that MIRI color selection will identify as being AGN hosts. First, we illustrate the predicted number counts at cosmic noon with the MIRI $10\,\mu\mathrm{m}$ filter, which is chosen for its sensitivity ($\sim 0.6\,\mu\mathrm{Jy}$ at 10σ in <3 min; Glasse et al. 2015) and because we use it in all three color diagnostics. We calculate the $10\,\mu\mathrm{m}$ flux for all CANDELS+COSMOS galaxies at z=0.75-2.25 and with $M_*>10^8\,M_\odot$ by scaling the appropriate Kirkpatrick et al. (2015) template (based on the source's $f(\mathrm{AGN})_{\mathrm{MIR}}$ determined through color classification) to the available IR photometry and convolving with the $10\,\mu\mathrm{m}$ transmission filter. By template fitting, we are also able to cal-

culate $L_{\rm IR}$ and $f({\rm AGN})_{\rm IR}$. The total 10 $\mu{\rm m}$ counts are plotted as the black stars in Figure 5. By including lower mass galaxies, we push below the 80% completeness in Figure 1 and down to the 20% completeness limit (measured at $24\,\mu{\rm m}$).

 $f({\rm AGN})_{\rm MIR}$ is an observational qualifier of black hole emission, and now we examine a physically motivated quantity, the Eddington ratio. The Eddington ratio is defined as $\lambda_{\rm Edd} = L_{\rm bol}/L_{\rm Edd}$, where $L_{\rm bol}$ is the bolometric luminosity of the AGN and $L_{\rm Edd}$ is the Eddington luminosity. In this way, $\lambda_{\rm Edd}$ is a measure of how efficiently a black hole is accreting material. $L_{\rm bol}$ is commonly estimated from $L_X(2-10{\rm keV})$. Due to obscuration and varying depths of the the *Chandra* catalogs in the CANDELS fields, we do not have L_X for all of our IR identified AGN and Composites. As a first step towards calculating $L_{\rm bol}$, we estimate L_X from $L_{\rm IR}^{\rm AGN}$ for all sources. We empirically determine the scaling between these luminosities to be

$$\log L_X [\text{erg s}^{-1}] = 38.8 + 0.45 \times \log L_{\text{IR}}^{\text{AGN}} [L_{\odot}]$$
 (16)

which we measured directly using Chandra observations of the GOODS-S field, which is the only field where the Chandra data is complete down to $L_X=10^{42}\,\mathrm{erg/s}$ out to z=2 [REFERENCE]. We then apply Equation 16 to all sources in the CANDELS and COSMOS fields. Next, we convert L_X to L_{bol} using the relation in Hopkins et al. (2007). Finally, we calculate $L_{\mathrm{Edd}}[\mathrm{erg\,s^{-1}}]=1.3\times10^{38}\times M_{\mathrm{BH}}[M_{\odot}],$ where $M_{\mathrm{BH}}=0.002\,M_{*}$ following the convention in Marconi & Hunt (2003); Aird et al. (2012).

We show the $10 \,\mu \text{m}$ number counts in bins of λ_{Edd} , which is useful since with the techniques outlined in Section 4, we will be able to calculate this parameter for samples with M_* or $M_{\rm BH}$ measurements. The $\lambda_{\rm Edd}$ < 0.01 curve (pink circles) are essentially SFGs, while the $\lambda_{\rm Edd} > 0.1$ curve (yellow) have accretion rates typical of sources identified as AGN at IR and X-ray wavelengths. The purple curve are those objects that we would most likely expect to be identified as Composites. The MIRI field of view is $1.2' \times 1.9'$, so we also illustrate the number counts in a MIRI FOV on the right axis of Figure 5. We expect nearly 100 objects per MIRI FOV at $10 \,\mu\text{m}$, achievable in a few minutes of integration time. Of these objects, >50% may be AGN hosts where we can detect and measure the black hole accretion. Below $S_{10}=0.01\,\mathrm{mJy}$, the counts become dominated by sources with $M_*<10^9\,M_\odot$. Of the galaxies with $\lambda_\mathrm{Edd}>0.01$, 30% have $M_*<10^9\,M_\odot$ and comprise an prime population for followup studies to more concretely pin down the AGN fraction in low mass galaxies at $z \sim 1-2$.

Finally, we demonstrate the host galaxy properties of CANDELS AGN and Composites selected with different techniques at $z \sim 1-2$ in Figure 6. In red, we plot the distribution of sSFR = SFR/ M_* for those galaxies with an X-ray detection ($L_X(2-10\,\mathrm{kev} \geq 10^{42}\,\mathrm{erg\,s^{-1}};$ note that only GOODS-S is complete down to this limit). Then, we plot the distribution of sSFR in blue for those galaxies that will be selected as either AGN or Composites by our MIRI color diagnostics, based on our estimation of their JWST colors through template fitting. To mitigate completeness effects of the X-ray catalogs, we normalize both distributions to have a peak at one. Our cosmic noon CANDELS AGN hosts have higher sSFR

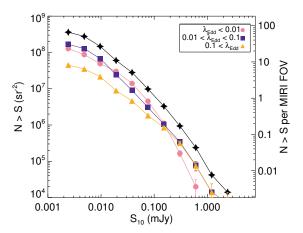


FIG. 5.— We predict the MIRI 10 $\mu \rm m$ number counts at z=0.75-2.25 as a function of $\lambda_{\rm Edd}$. We have scaled the 24 $\mu \rm m$ emission of the CANDELS and COSMOS galaxies shown in Figure 1 using the appropriate template for SFGs, Composites, and AGN. We then calculate $\lambda_{\rm Edd}$ by scaling $L_{\rm IR}^{\rm AGN}$ to $L_{\rm bol}$ and M_* to $M_{\rm BH}$ and then $L_{\rm Edd}$ for each galaxy. In one MIRI FOV, we will detect nearly 100 galaxies, and many of these will be AGN hosts.

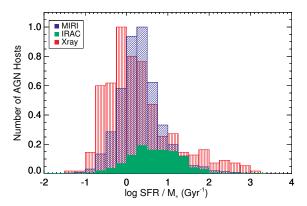


FIG. 6.— We show the host galaxy property sSFR=SFR/ M_{\ast} of CANDELS galaxies identified as AGN by X-ray detection (red histogram) and MIRI color selection (blue histogram). We have normalized each distribution to mitigate completeness effects of the Chandra X-ray observations in each field. The MIRI selection is sensitive to host galaxies with higher sSFRs. We also illustrate which of the MIRI selected AGN hosts would also be selected with the Donley et al. (2012) IRAC selection criteria. Due to the sensitivity and coverage of the MIRI filters, we will be able to select larger samples of AGN hosts than was possible with IRAC selection.

than the X-ray selected CANDELS galaxies, indicating that we are able to select galaxies prior to quenching of star formation.

Prior to JWST, the most popular way of identifying larges samples of IR AGN is with IRAC color techniques (Lacy et al. 2004; Stern et al. 2005; Donley et al. 2012). The Donley et al. (2012) IRAC diagnostic is the most reliable, since it eliminates host galaxy contamination, but it is only sensitive to the most actively accreting AGN as it is based on a power-law selection criterion. Of the CANDELS sources selected by our MIRI diagnostics, we show in green the sources that are also selected as AGN by the Donley et al. (2012) criteria. Clearly, due to the sensitivity and spacing of the MIRI filters, we will be able to detect >4 times as many AGN hosts as would be identified with IRAC alone. MIRI color selection will enable identification of statistical samples of AGN hosts in their star forming prime (as measured by sSFR), allowing as-

tronomers to trace the star formation-AGN connection at the peak period of stellar and black hole growth in the Universe.

6. CONCLUSIONS

We identify SFGs, AGN, and composites in four CANDELS fields and in the full COSMOS field using three different redshift dependent color identification techniques. We present the first $24\,\mu\mathrm{m}$ number counts of star forming+AGN Composite galaxies at $z\sim1-2$. We find that IR AGN and composites dominate $24\,\mu\mathrm{m}$ samples at $S_{24}>1\,\mathrm{mJy}$. Any $24\,\mu\mathrm{m}$ selected sample contains at least 20% identifiable composites.

We use a library of SFG, AGN, and Composite templates to create synthetic galaxies, and we use these synthetic galaxies to create JWST/MIRI color selection techniques for three redshift bins, $z\sim 1$, $z\sim 1.5$, and $z\sim 2$. At these redshift, our color selection techniques cover the $6.2\,\mu\mathrm{m}$ and $7.7\,\mu\mathrm{m}$ PAH features.

Finally, we predict the Eddington ratios ($\lambda_{\rm Edd}$), a measure of black hole accretion efficiencies, that we will observe with MIRI imaging. Our MIRI color selection diagnostic can identify samples of AGN and Composite galaxies with $\lambda_{\rm Edd} > 0.01$ that are four times larger than samples of AGN selected by Spitzer/IRAC techniques. We also use our new $24\,\mu{\rm m}$ number counts to predict the number counts at $10\,\mu{\rm m}$ in different bins of $\lambda_{\rm Edd}$. With MIRI color identification, we will be able to probe the star formation - AGN connection in dusty galaxies at cosmic noon.

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APPENDIX

In this appendix, we show our redshift dependent color diagnostic to find SFGs, Composites, and AGN using *Spitzer* and *Herschel* photometry. We create a catalog of 5500 synthetic galaxies from 11 templates where we know the intrinsic AGN contribution. We resample each photometric point within the uncertainties of the template from which it was created, so that we can represent the scatter in colorspace of real galaxies, which is an improvement upon using so-called redshift tracks alone to explore where SFGs, Composites, and AGN lie in colorspace.

We create color diagrams in redshift bins of z=0.75-1.25, z=1.25-1.75, and z=1.75-2.25. In each redshift bin, we divide the color space into regions of 0.2×0.2 dex and calculate the average $f(\text{AGN})_{\text{MIR}}$ and standard deviation, σ_{AGN} of all the synthetic galaxies that lie in that region. In the plots below, we show our three diagnostics: $S_{250}/S_{24}v.S_8/S_{3.6}$, $S_{100}/S_{24}vs.S_8/S_{3.6}$, and $S_{24}/S_8vs.S_8/S_{3.6}$.

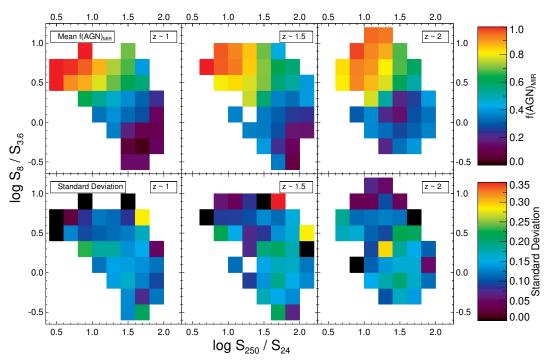


Fig. 7.— $S_{250}/S_{24}v.S_8/S_{3.6}$ in three redshift bins. The top panels are shaded according to the average $f(AGN)_{MIR}$ measured in each bin, and the bottom panels show the standard deviation of $f(AGN)_{MIR}$ for the sources in each bin.

REFERENCES

Aird, J., Coil, A., Moustakas, J., et al. 2012, ApJ, 746, 90 Alonso-Herrero, A., Pérez-González, P. G., Alexander, D. M., et al. 2006, ApJ, 640, 167

Barro, G., Pérez-González, P. G., Gallego, J., et al. 2011, ApJS, 193, 13

Battisti, A. J., Calzetti, D., Johnson, B. D., & Elbaz, D. 2015, ApJ, 800, 143

Béthermin, M., Dole, H., Beelen, A., Aussel, H. 2010, A&A, 512, 78

Bonato, M., Sajina, A., De Zotti, G., et al. 2017 ApJ, 836, 171 Bouchet et al. 2015, PASP, 127, 612

Brandl, B. R., et al. 2006, ApJ, 653, 1129

Chary, R. & Elbaz, D. 2001, ApJ, 556, 562

Coppin, K., Pope, A., Menéndez-Delemstre, K. et al. 2010, ApJ, 713, 503

Dahlen, T., Mobasher, B., Faber, S. M., et al. 2013, ApJ, 775, 93 Draine, B. T. & Li, A. 2007, ApJ, 657, 810

Donley, J. L., Koekemoer, A. M., Brusa, M., et al. 2012, ApJ, 748, 142

Elvis, M., Wilkes, B. J., McDowell, J. C., et al. 1994, ApJS, 95, 1
 Engelbracht, C. W., Rieke, G. H., Gordon, K. D., et al. 2008,
 ApJ, 685, 678

Erb, D. K., Shapley, A. E., Pettini, M., et al. 2006, ApJ, 644, 813 Glasse et al. 2015, PASP, 127, 686

Hernán-Caballero, A., Alonso-Herrero, A., Hatziminaoglou, E., et al. ApJ, 803, 109

Hickox, R. C. & Markevitch, M. 2007, ApJ, 743, 159

Hopkins, P. F., Hernquist, L., Cox, T. J., Di Matteo, T., Robertson, B., & Springel, V. 2006, ApJS, 163, 1

Hopkins, P. F., Richards, G. T., & Hernquist, L. 2007, ApJ, 654, 731

Ivison, R. J., et al. 2012, MNRAS, 425, 1320

Kirkpatrick, A., Pope, A., Alexander, D. M., et al. 2012, ApJ, 759, 139

Kirkpatrick, A., Pope, A., Charmandaris, V., et al. 2013a, ApJ, 763, 123

Kirkpatrick, A., Pope, A., Sajina, A., et al. 2013, ApJ, 814, 9
 Lacy, M., Storrie-Lombardi, L. J., Sajina, A., et al. 2004, ApJS, 154, 166

Laigle, C., McCracken, H. J., Ilbert, O., et al. 2016, ApJS, 224, 24
Le Floc'h, E., Aussel, H., Ilbert, O., et al. 2009, ApJ, 703, 222
Lutz, D., Poglitsch, A., Altieri, B., et al. 2011, A&A, 532, A90
Lyu & Rieke 2017, ApJ, 841, 76

Ma, X., Hopkins, P. F., Faucher-Giguère, C.-A., et al. 2016, MNRAS, 456, 2140

Madau, P. & Dickinson, M. 2014 ARA&A, 52, 415

Marconi, A. & Hunt, L. K. 2003, ApJ, 589, L21

Mendez, A. J., Coil, A. L., Aird, J., et al. 2013, ApJ, 770, 40

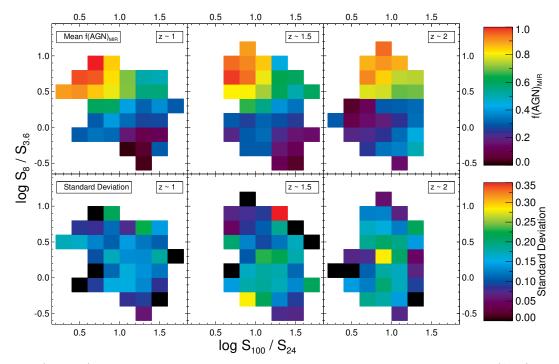


Fig. 8.— $S_{100}/S_{24}v.S_8/S_{3.6}$ in three redshift bins. The top panels are shaded according to the average $f(AGN)_{MIR}$ measured in each bin, and the bottom panels show the standard deviation of $f(AGN)_{MIR}$ for the sources in each bin.

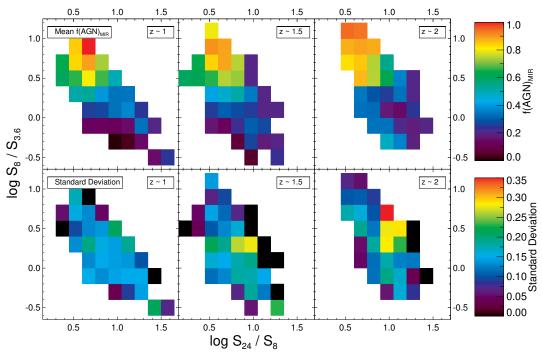


Fig. 9.— $S_{24}/S_8v.S_8/S_{3.6}$ in three redshift bins. The top panels are shaded according to the average $f(AGN)_{MIR}$ measured in each bin, and the bottom panels show the standard deviation of $f(AGN)_{MIR}$ for the sources in each bin.

Messias, H., Afonso, J., Salvato, M., Mobasher, B., Hopkins, A. M. 2012, ApJ, 754, 120

Mullaney, J. R., Alexander, D. M., Goulding, A. D., Hickox, R. C. 2011, MNRAS, 414, 1082

Murphy, E. J., Chary, R.-R., Dickinson, M., et al. 2011, ApJ, 732, 126

Oliver, S. J., Bock, J., Altieri, B. et al. 2012, MNRAS, 424, 1614 Pannella, M., Elbaz, D., Daddi, E., et al. 2015, ApJ, 807, 141 Papovich, C., Dole, H., Egami, E., et al. 2004, ApJS, 154, 70 Pérez-González, P. G., Rieke, G. H., Egami, E., et al. 2005, ApJ, $630,\,82$

Pérez-González, P. G., Rieke, G. H., Villar, V., et al. 2008, ApJ, 675, 234

Pérez-González, P. G., Egami, E., Rex, M., et al. 2010, A&A, 518, L15

Peeters, E., Spoon, H. W. W., Tielens, A. G. G. M. 2004, ApJ, 613, 986

Pope, A., Chary, R.-R., Alexander, D. M., et al. 2008, ApJ, 675, 1171

- Rawle, T. D., Altieri, B., Egami, E., et al. 2016, MNRAS, 459, 1626
- Rigby, J. R., Marcillac, D., Egami, E., et al. 2008, ApJ, 675, 262 Rujopakarn. W., Rieke, G. H., Eisenstein, D. J., Juneau, S. 2011, ApJ, 726, 93
- Rujopakarn, W., Dunlop, J. S., Rieke, G. H., et al. 2016, ApJ, 833, 12
- Sajina, A., Yan, L., et al. 2012, ApJ, 757, 13
- Sanders, D. B., Soifer, B. T., Elias, J. H., et al. 1998, ApJ, 325, 74 Sanders, D. B., Salvato, M., Aussel, H., et al. 2007, ApJS, 172, 86 Sanders, B. L., Shaplay, A. F., Kriek, M., et al. 2015, ApJ, 799
- Sanders, R. L., Shapley, A. E., Kriek, M., et al. 2015, ApJ, 799, 138
 Sandstrom, K. M., Bolatto, A. D., Bot, C., et al. 2012, ApJ, 744,
- Santini, P., Ferguson, H. C., Fontana, A., et al. 2015, ApJ, 801, 97
 Scoville, N. 2007, AIPC, 943, 221

- Shipley, H. V., Papovich, C., Rieke, G. H., et al. 2016, ApJ, 818, 60
- Shivaei, I., Reddy, N. A., Shapley, A. E., et al. 2017, ApJ, 837, 157 Siebenmorgen, R., Heymann, F., Efstathiou, A. 2015, A&A, 583, 120
- Smith, J. D. T., Draine, B. T., Dale, D. A., et al. 2007, ApJ, 656, 770
- Stern, D., Eisenhardt, P., Gorjian, V., et al. 2005, ApJ, 631, 163 Stern, D., Assef, R. J., Benford, D. J., et al. 2012, ApJ, 753, 30
- Teng, S. H., Rigby, J. R., Stern, D. et al. 2015, ApJ, 814, 56Veilleux, S., Rupke, D. S. N., Kim, D.-C., et al. 2009, ApJS, 182, 628
- Wells et al. 2015, PASP, 127, 646
- Younger, J. D., et al. 2009, ApJ, 704, 803
- Zahid, H. J., Geller, M. J., Kewley, L. J., et al. 2013, ApJ, 771, L19