

# How SHADOW manipulates electric vectors

M. Sanchez del Rio

March 15, 2013

## Abstract

This document explains how the electric fields and related parameters (intensity, polarization) are defined and modified in SHADOW. It also describes a bug fixed in March 2013.

## 1 Definitions

Most Monte-Carlo codes store a parameter “intensity” or weight for each ray that allow to track the energy or intensity of each traced event. As most of electromagnetic interactions are better described by the change in the electric field it is more appropriated to store directly the electric fields and not the intensity. The intensity is the squared-modulus of the electric fields, The electric fields permit not only to compute intensity or each ray, but also the polarization, phase changes, and coherence properties. In order to extract usable and accurate information on these important parameters, the electric fields must be correctly simulated and propagated. I try to summarize here some of the assumptions and manipulations of the electric fields in SHADOW.

A single ray  $\vec{R}$  in SHADOW is an array of 18 parameters (called “columns” in SHADOW’s jargon) stored in the following order:

$$\vec{R} = (\vec{x}, \vec{v}, \vec{E}_\sigma, f, k, n, opd, \phi_\sigma, \phi_\pi, \vec{E}_\pi) \quad (1)$$

where  $\vec{x}$  are the spatial coordinates,  $\vec{v}$  the director cosines (ray direction),  $f$  is a flag (good or lost ray),  $k$  is the wavevector modulus in  $cm^{-1}$  ( $k = 2\pi/\lambda$ , with  $\lambda$  the photon wavelength),  $n$  is a counter (from one to the total number of rays),  $opd$  is the optical path (the travelled distance times the refraction index). The complex electric vector is expressed as a function of the parallel ( $\sigma$ ) and perpendicular ( $\pi$ ) components:

$$\vec{E} = \vec{E}_\sigma e^{i\phi_\sigma} + \vec{E}_\pi e^{i\phi_\pi} \quad (2)$$

Obviously the storage in SHADOW is redundant, for example  $v_3 = \sqrt{v_1^2 + v_2^2}$  and regarding the electric fields:

- $\vec{E}_\sigma \perp \vec{E}_\pi$ ,  $\vec{E}_\sigma \perp \vec{v}$ ,  $\vec{E}_\pi \perp \vec{v}$
- $\vec{E}_\sigma$  and  $\vec{E}_\pi$  are real.
- $|\vec{E}|^2 = 1$  at the source, and  $|\vec{E}|$  is reduced as far as the ray interacts with the different optical elements.

The intensity of the ray is:

$$I = |\vec{E}|^2 = \vec{E} \cdot \vec{E}^* = (\vec{E}_\sigma e^{i\phi_\sigma} + \vec{E}_\pi e^{i\phi_\pi}) \cdot (\vec{E}_\sigma e^{-i\phi_\sigma} + \vec{E}_\pi e^{-i\phi_\pi}) = |\vec{E}_\sigma|^2 + |\vec{E}_\pi|^2 + 2\vec{E}_\sigma \cdot \vec{E}_\pi \cos(\phi_\sigma - \phi_\pi) = |\vec{E}_\sigma|^2 + |\vec{E}_\pi|^2 \quad (3)$$

## 2 A polarized source

When simulating a source, SHADOW first samples the ray energy, direction and positions and then define set electric vectors. A first step is to obtain orthogonal unitary vectors  $\vec{u}_\sigma$  and  $\vec{u}_\pi$  normal to  $\vec{v}$ . At the source level, there is not yet knowledge of any optical surface to define  $\sigma$  and  $\pi$  directions, so these directions are considered “mostly” horizontal and vertical, respectively. The unitary vectors are calculated as:

$$\begin{aligned} \vec{u}_\sigma &= \vec{v} \wedge (\vec{u}_1 \wedge \vec{v}) \\ \vec{u}_\pi &= \vec{u}_\sigma \wedge \vec{v} \end{aligned} \quad (4)$$

where  $\vec{u}_1 = (1, 0, 0)$  the unitary vector along the horizontal direction at the source plane. This is done with the following code in the `sourceGeom` subroutine of the `shadow_kernel.F90` module:

```

A_VEC(1)      = 1.0D0
A_VEC(2)      = 0.0D0
A_VEC(3)      = 0.0D0
! C
! C   Rotate A_VEC so that it will be perpendicular to DIREC and
! C   with the
! C   right components on the plane.
! C
CALL CROSS    (A_VEC, DIREC, A_TEMP)
CALL CROSS    (DIREC, A_TEMP, A_VEC)
CALL NORM     (A_VEC, A_VEC)
CALL CROSS    (A_VEC, DIREC, AP_VEC)
CALL NORM     (AP_VEC, AP_VEC)

```

Then the electric vectors are these versors times the corresponding ( $\sigma$  or  $\pi$ ) modulus, in such a way that the intensity is one. The modulus of each component is a function of a “particular” polarization degree defined as:

$$P = \frac{\cos \chi}{\cos \chi + \sin \chi} \quad (5)$$

where  $\chi$  is the angle between the  $\vec{E}$  vector and  $\vec{u}_1$  when the beam is linearly polarized. The electric vector components are:

$$\begin{aligned} \vec{E}_\sigma &= \frac{P}{\sqrt{1 - 2P + 2P^2}} \vec{u}_\sigma \\ \vec{E}_\pi &= \frac{1 - P}{\sqrt{1 - 2P + 2P^2}} \vec{u}_\pi. \end{aligned} \quad (6)$$

The phase  $\phi_\sigma$  is set to zero for coherent beams or to a random angle in  $[0, 180^\circ)$  for an incoherent beam, and  $\phi_\pi = \phi_\sigma + \Phi$  with:

- $\Phi = 0$  for linearly polarized light

- $\Phi = 90^\circ$  for right (CW) elliptical polarization
- $\Phi = -90^\circ$  for left (CCW) elliptical polarization.

A few examples:

- $P = 1, \Phi = 0$  for horizontal linearly polarized light
- $P = 0, \Phi = 0$  for vertical linearly polarized light
- $P = 0.5, \Phi = 0$  for  $45^\circ$  linearly polarized light
- $P = 0.5, \Phi = 90^\circ$  for right circular polarization
- $P < 0.5, \Phi = -90^\circ$  for left elliptical polarization mostly horizontal

### 3 Change of reference frame

The beam created in the source reference frame must be converted into the optical element reference frame for performing the ray tracing. The rotation of the vectors from the source (or *previous* o.e. to the o.e. under study) is done in the subroutine **RESTART**. It applies sequential rotations of the vectors using the mirror orientation angle  $\alpha$ , incident angle  $\theta_i$ , etc. An example of code is the following, for changing the  $\vec{E}_\pi$ :

```
TEMP_1P(1) = AP(1,ITIME)*COSAL + AP(3,ITIME)*SINAL
TEMP_1P(2) = AP(2,ITIME)
TEMP_1P(3) = - AP(1,ITIME)*SINAL + AP(3,ITIME)*COSAL

TEMP_2P(1) = TEMP_1P(1)
TEMP_2P(2) = TEMP_1P(2)*SINTHR + TEMP_1P(3)*COSTHR
TEMP_2P(3) = - TEMP_1P(2)*COSTHR + TEMP_1P(3)*SINTHR

TEMP_1P(1) = TEMP_2P(1)*COSALS - TEMP_2P(2)*SINALS
TEMP_1P(2) = TEMP_2P(1)*SINALS + TEMP_2P(2)*COSALS
TEMP_1P(3) = TEMP_2P(3)

AP(1,ITIME) = TEMP_1P(1)
AP(2,ITIME) = TEMP_1P(2)
AP(3,ITIME) = TEMP_1P(3)
```

Then the subroutine **MIRROR1** performs the reflection/refraction/diffraction by the optical element (o.e.), changing the beam direction and electric vectors. This will be discussed in detail in the next paragraphs. Last, after the reflection, the rays are projected into the image plane. For that it projects the vectors onto the versors of the new reference. For example, for  $\vec{E}_\pi$

```
AP_VEC(1) = AP(1,J)
AP_VEC(2) = AP(2,J)
AP_VEC(3) = AP(3,J)
CALL DOT (AP_VEC,UXIM,A_1)
CALL DOT (AP_VEC,VNIMAG,A_2)
CALL DOT (AP_VEC,VZIM,A_3)
AP(1,J) = A_1 * rr_reflectivity
AP(2,J) = A_2 * rr_reflectivity
AP(3,J) = A_3 * rr_reflectivity
```

It is important that these changes of reference system (before and after interaction with the o.e.) are not affecting the phase  $\phi_\sigma$  and  $\phi_\pi$  and conserve

the modulus of each electric field component  $\vec{E}_\sigma$  and  $\vec{E}_\pi$ . Also, they do not alter the orthogonality relationships:

$$\begin{aligned}\vec{E}_\sigma &\perp \vec{E}_\pi \\ \vec{E}_\sigma &\perp \vec{v} \\ \vec{E}_\pi &\perp \vec{v}\end{aligned}\tag{7}$$

## 4 Modifications of electric fields in “local” o.e.

The physics of the X-ray reflection and diffraction in an optical surface affects in a different way the *local* parallel ( $\sigma$ ) and perpendicular ( $\pi$ ) components. The “local” parallel component is a vector that sits on the o.e. surface and it is not coincident with the incident  $\vec{E}_\sigma$ . Therefore, one must calculate the “local”  $\vec{u}'_\sigma$  and  $\vec{u}'_\pi$  vectors that verify 1) they are orthogonal to  $\vec{v}$  and ii)  $\vec{u}'_\sigma$  is contained in the o.e. surface. This is done here ( MIRROR1 subroutine):

```
! * Check for reflectivity. If this mode is "on", we have to compute
! * some angles, namely the sine of the incidence angle and the sine
! * of the A vector with the normal. Also, the polarized light is
! * treated as a superposition of two orthogonal A vectors with the
! * appropriate
! * phase relation. These two incoming vectors have to be resolved into
! * the
! * local S- and P- component with a new phase relation.
! * A.VEC will be rotated later, once the amplitude will have been
! * determined.
      CALL CROSS.M.FLAG (VVIN,VNOR,AS.TEMP,M.FLAG)      ! vector
      pp. to inc.pl.
! print *, 'M.FLAG: ', M.FLAG

      IF (M.FLAG.EQ.1) THEN
      CALL DOT (AS.VEC,AS.VEC,AS2)
      CALL DOT (AP.VEC,AP.VEC,AP2)
      IF (AS2.NE.0) THEN
499        AS.TEMP(I) = AS.VEC(I)
      ELSE
      DO 599 I=1,3
599        AS.TEMP(I) = AP.VEC(I)
      END IF
      END IF

      CALL NORM (AS.TEMP,AS.TEMP)      ! Local unit As vector
      CALL CROSS (AS.TEMP,VVIN,AP.TEMP)
      CALL NORM (AP.TEMP,AP.TEMP)      ! Local unit Ap vector
```

The electric vector can thus be expressed in two orthonormal vectors  $\vec{E}'_\sigma$  and  $\vec{E}'_\pi$  along these new directions  $\vec{u}'_\sigma$  and  $\vec{u}'_\pi$ . Physically, it is a rotation of the old electric vectors around the  $\vec{v}$  direction to put the  $\sigma$  component on top of the surface, but the rotation affect not only the moduli, but also the phases. The new (complex) electric vectors are build from the projection of the old ones onto the new axes, thus originating a transformation in both moduli and phases:

$$\vec{E}'_\sigma e^{i\phi'_\sigma} = [(\vec{E}_\sigma e^{i\phi_\sigma}).\vec{u}'_\sigma + (\vec{E}_\pi e^{i\phi_\pi}).\vec{u}'_\sigma] \vec{u}'_\sigma \tag{8}$$

$$\vec{E}'_\pi e^{i\phi'_\pi} = [(\vec{E}_\sigma e^{i\phi_\sigma}).\vec{u}'_\pi + (\vec{E}_\pi e^{i\phi_\pi}).\vec{u}'_\pi] \vec{u}'_\pi \tag{9}$$

Defining the projection of the (real) electric field components onto these versors as:

$$\begin{aligned} a_{11} &= \vec{E}_\sigma \cdot \vec{u}'_\sigma, & a_{12} &= \vec{E}_\sigma \cdot \vec{u}'_\pi \\ a_{21} &= \vec{E}_\pi \cdot \vec{u}'_\sigma, & a_{22} &= \vec{E}_\pi \cdot \vec{u}'_\pi \end{aligned} \quad (10)$$

we get:

$$\begin{aligned} \vec{E}'_\sigma e^{i\phi'_\sigma} &= (a_{11}e^{i\phi_\sigma} + a_{12}e^{i\phi_\pi})\vec{u}'_\sigma \\ \vec{E}'_\pi e^{i\phi'_\pi} &= (a_{21}e^{i\phi_\sigma} + a_{22}e^{i\phi_\pi})\vec{u}'_\pi \end{aligned} \quad (11)$$

And the moduli are:

$$\begin{aligned} |\vec{E}'_\sigma|^2 &= \vec{E}'_\sigma \cdot \vec{E}'_\sigma = a_{11}^2 + a_{12}^2 + 2a_{11}a_{12}\cos(\phi_\sigma - \phi_\pi) \equiv M_\sigma^2 \\ |\vec{E}'_\pi|^2 &= \vec{E}'_\pi \cdot \vec{E}'_\pi = a_{21}^2 + a_{22}^2 + 2a_{21}a_{22}\cos(\phi_\sigma - \phi_\pi) \equiv M_\pi^2 \end{aligned} \quad (12)$$

Therefore we can construct the new local (real) electric fields as:

$$\begin{aligned} \vec{E}'_\sigma &= M_\sigma \vec{u}'_\sigma \\ \vec{E}'_\pi &= M_\pi \vec{u}'_\pi \end{aligned} \quad (13)$$

To compute the new phases, we insert Eq. 13 in Eq. 11 and get:

$$\begin{aligned} e^{i\phi'_\sigma} &= M_\sigma^{-1}(a_{11}e^{i\phi_\sigma} + a_{12}e^{i\phi_\pi}) \\ e^{i\phi'_\pi} &= M_\pi^{-1}(a_{21}e^{i\phi_\sigma} + a_{22}e^{i\phi_\pi}) \end{aligned} \quad (14)$$

Therefore:

$$\begin{aligned} \tan \phi'_\sigma &= \frac{\text{Im}(a_{11}e^{i\phi_\sigma} + a_{12}e^{i\phi_\pi})}{\text{Re}(a_{11}e^{i\phi_\sigma} + a_{12}e^{i\phi_\pi})} = \frac{a_{11}\sin \phi_\sigma + a_{12}\sin \phi_\pi}{a_{11}\cos \phi_\sigma + a_{12}\cos \phi_\pi} \\ \tan \phi'_\pi &= \frac{\text{Im}(a_{21}e^{i\phi_\sigma} + a_{22}e^{i\phi_\pi})}{\text{Re}(a_{21}e^{i\phi_\sigma} + a_{22}e^{i\phi_\pi})} = \frac{a_{21}\sin \phi_\sigma + a_{22}\sin \phi_\pi}{a_{21}\cos \phi_\sigma + a_{22}\cos \phi_\pi} \end{aligned} \quad (15)$$

The code that implements Eq. 13 and Eq. 15 is in MIRROR1:

```
CALL DOT (AS_VEC, AS_TEMP, A11)
CALL DOT (AP_VEC, AS_TEMP, A12)
CALL DOT (AS_VEC, AP_TEMP, A21)
CALL DOT (AP_VEC, AP_TEMP, A22)
! ** Now recompute the amplitude and phase of the local S- and P-
! component.
AS_NEW = SQRT(ABS(A11**2 + A12**2 + 2.0D0*A11*A12*COS(PHS-PHP)))
AP_NEW = SQRT(ABS(A21**2 + A22**2 + 2.0D0*A21*A22*COS(PHS-PHP)))
CALL DOT(AP_VEC, AS_VEC, TMP11)

CALL SCALAR (AS_TEMP, AS_NEW, AS_VEC) ! Local As vector
CALL SCALAR (AP_TEMP, AP_NEW, AP_VEC) ! Local Ap vector

PHTS = A11*SIN(PHS) + A12*SIN(PHP)
PHBS = A11*COS(PHS) + A12*COS(PHP)
PHTP = A21*SIN(PHS) + A22*SIN(PHP)
PHBP = A21*COS(PHS) + A22*COS(PHP)
CALL ATAN_2 (PHTS, PHBS, PHS) ! Phase of local As vector
CALL ATAN_2 (PHTP, PHBP, PHP) ! Phase of local Ap vector
```

Note that these transformations imply a polarization mixing, or not conservation of the intensity of each component during the projection ( $|\vec{E}'_\sigma| \neq |\vec{E}_\sigma|$ ), but the total intensity is conserved ( $I = |\vec{E}_\sigma|^2 + |\vec{E}_\pi|^2 = |\vec{E}'_\sigma|^2 + |\vec{E}'_\pi|^2$ ). The phases also change, but the difference of phase is conserved  $\Phi = \phi_\sigma - \phi_\pi = \phi'_\sigma - \phi'_\pi$  (I think, but I have not demonstrated it). The new electric vectors are orthogonal to  $\vec{v}$  by constructions.

## 5 Reflection/refraction in the o.e. and subsequent changes in electric vectors

Once the electric vectors are expressed in the o.e. local coordinates ( $\vec{E}'_\sigma$  and  $\vec{E}'_\pi$ ) they are multiplied by the mirror/crystal reflectivities and the phases also affected: :

$$\begin{aligned}\vec{E}'_\sigma^{new} &= \vec{E}'_\sigma R_\sigma \\ \vec{E}'_\pi^{new} &= \vec{E}'_\pi R_\pi \\ \phi_\sigma^{new} &= \phi'_\sigma \Sigma_\sigma \\ \phi_\pi^{new} &= \phi'_\pi \Sigma_\pi\end{aligned}\tag{16}$$

where  $R_\sigma$  and  $R_\pi$  are the o.e. reflectivities and  $\Sigma_\sigma$  and  $\Sigma_\pi$  are the phases added in the reflection/refraction/diffraction. This process is accompanied by a change of direction of the ray from the incident direction  $\vec{v}$  to a new output direction  $\vec{v}'$ . All these operations are done in MIRROR1 subroutine.

The resulting electric vectors and phases after applying Eq. 16 must now hold the orthogonality relations (Eq. 7) with respect to  $\vec{v}'$  so they must be changed. This part was incorrectly done in the following code, which assumes conservation of the  $\vec{E}'_\sigma$  component (except for gratings) and the  $\pi$  component is “mirrored”. These operations do not guarantee that the resulting vectors are normal to  $\vec{v}'$ . In some cases, in particular for  $\sigma$  polarized light onto o.e.’s with mirror orientation angle  $90^\circ$  can result in a loss of beam intensity in subsequent changes of frames. The **wrong** code was (in MIRROR1):

```
! ** So far we have the new amplitude of the two components. We have now
! ** to 'reflect' A_VEC onto the mirror. For this, notice that the s-
! comp
! ** is geometrically unchanged, while the p-comp is changed. The angles
! ** are exchanged with respect to VVIN. Things are more complicated in
! ** the case of a grating, due to the vectorial nature of the
! diffraction,
! ** not treated here. We make the simplifying assumption that the
! ** diffraction will not change the degree of polarization. This means
! that
! ** A_VEC will have the same components referred to the ray as before
! the
! ** diffraction.

VVOUT(1)      =   RAY(4,ITIK)
VVOUT(2)      =   RAY(5,ITIK)
VVOUT(3)      =   RAY(6,ITIK)
! C
! C The following IF block applies only to the GRATING case.
! C The binormal is redefined in terms of the diffraction
! C plane.
! C
      IF (F_GRATING.NE.0.OR.F_BRAGG.A.EQ.1) THEN
        CALL PROJ (VVOUT,VNOR,VTEMP)
        CALL SCALAR (VTEMP,-2.0D0,VTEMP)
        CALL SUM (VTEMP,VVOUT,VTEMP)
        !CALL CROSS (VTEMP,VNOR,AS_TEMP)
        CALL CROSS_M_FLAG (VTEMP,VNOR,AS_TEMP,M_FLAG)
      IF (M_FLAG.EQ.1) THEN
        CALL DOT (AS_VEC,AS_VEC,AS2)
        CALL DOT (AP_VEC,AP_VEC,AP2)
        IF (AS2.NE.0) THEN
          DO 899 I=1,3
899          AS_TEMP(I) = AS_VEC(I)
        ELSE
          DO 999 I=1,3
```

```

999      AS_TEMP(I) = AP_TEMP(I)
      END IF
    END IF
    CALL NORM (AS_TEMP,AS_TEMP) ! Local unit As vector
    CALL CROSS (AS_TEMP,VTEMP,AP_TEMP)
    CALL NORM (AP_TEMP,AP_TEMP) ! Local unit Ap vector
    CALL DOT (AS_VEC,AS_VEC,RES)
    RES = SQRT (RES)

    CALL SCALAR (AS_TEMP,RES,AS_VEC)
    CALL DOT (AP_VEC,AP_VEC,RES)
    RES = SQRT (RES)
    CALL SCALAR (AP_TEMP,RES,AP_VEC)
  END IF

CALL PROJ (AP_VEC,VNOR,VTEMP)
CALL VECTOR (VTEMP,AP_VEC,VTEMP)
CALL SCALAR (VTEMP,-2.0D0,VTEMP)

CALL SUM (AP_VEC,VTEMP,AP_VEC)

```

This code has been replaced by a new one that simply calculates the new  $\sigma$  and  $\pi$  versors orthogonal to  $\vec{v}$  and affects them by the electric field moduli that do not change. Also, the phases are not changed. This guarantees that the intensity is conserved, as well as orthogonality. The new code is:

```

! Electric vectors are changed to assure orthogonality with the new
! direction VVOUT
! To conserve intensity, the moduli of Es and Ep must not change
! AS_VEC and VVOUT are not orthogonal so a projection of S and P
! coordinates into the
! new ones do not work as it may be a component of the electric field
! along VVOUT

CALL CROSS_MFLAG (VVOUT,VNOR,AS_TEMP,MFLAG) ! vector pp. to inc.pl.
CALL DOT (AS_VEC,AS_VEC,AS2)
CALL DOT (AP_VEC,AP_VEC,AP2)

IF (MFLAG.EQ.1) THEN
  IF (AS2.NE.0) THEN
    DO I=1,3
      AS_TEMP(I) = AS_VEC(I)
    END DO
  ELSE
    DO I=1,3
      AS_TEMP(I) = AP_VEC(I)
    END DO
  END IF
END IF

CALL NORM (AS_TEMP,AS_TEMP) ! Local unit As vector perp to vvout
CALL CROSS (AS_TEMP,VVOUT,AP_TEMP)
CALL NORM (AP_TEMP,AP_TEMP) ! Local unit Ap vector perp to vvout

do i=1,3
  as_vec(i) = as_temp(i) * sqrt(as2)
  ap_vec(i) = ap_temp(i) * sqrt(ap2)
end do

```