

SIMULATION AND ADJOINT-BASED DESIGN FOR VARIABLE DENSITY INCOMPRESSIBLE FLOWS WITH HEAT TRANSFER

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Multiphysics Modeling and Simulation

Robert Bosch LLC

3rd Annual SU2 Developers Meeting

2018.09.17



Agenda

1. Background

2. Modeling & Implementation

3. Results (V&V)

4. Conclusions

BACKGROUND

Bosch – technology to enhance quality of life



- ▶ Some 59,000¹ researchers and developers work at Bosch: at 120² locations worldwide, in a single network.
- ▶ Bosch is one of the world's leading international providers of technology and services.
- ▶ Over the past six years, Bosch has invested more than 27 billion euros in research and development.
- ▶ Our objective: to develop innovative, useful, and exciting products and solutions to enhance quality of life – technology that is “Invented for life.”

¹ As of 12.16 ² R&D locations with >50 associates, as of 12.16

Bosch – Four business sectors

Key figures 2016*

Bosch Group

- ▶ 73.1 billion euros in sales
- ▶ 389,281 associates

Mobility Solutions

- ▶ One of the world's largest suppliers of mobility solutions

60% share of sales



Industrial Technology

- ▶ Leading in drive and control technology, packaging, and process technology



Energy and Building Technology

- ▶ One of the leading manufacturers of security and communication technology
- ▶ Leading manufacturer of energy-efficient heating products and hot-water solutions

40% share of sales



Consumer Goods

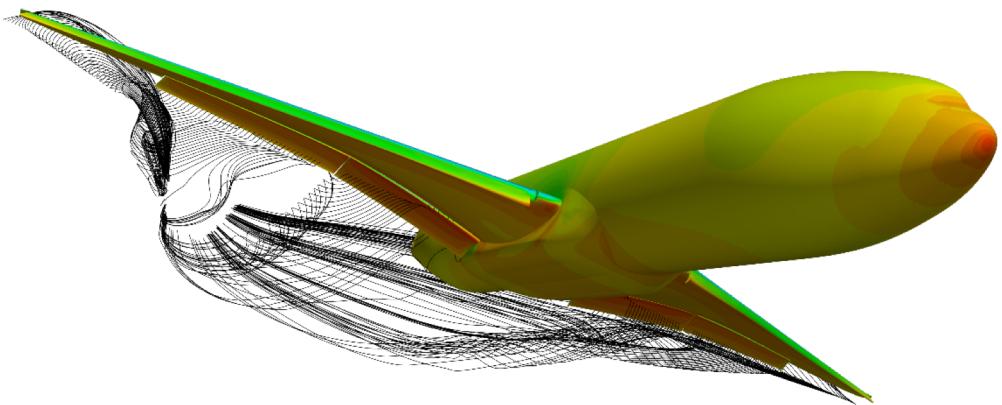
- ▶ Leading supplier of power tools and accessories
- ▶ Leading supplier of household appliances



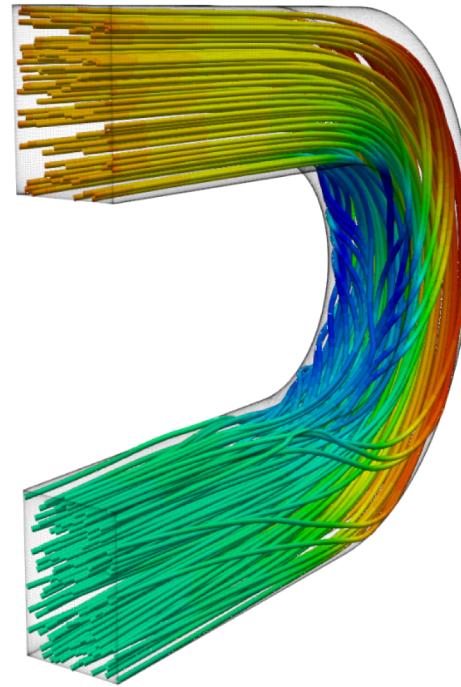
* As of 12.16

Background

Some SU2 History



External, Compressible Aerodynamics



Internal, Incompressible Flows with Heat Transfer



Background Coupled Approaches (Density-based) in SU2

(Reprint of 1967 paper)

JOURNAL OF COMPUTATIONAL PHYSICS 135, 118–125 (1997)
ARTICLE NO. CP975716

A Numerical Method for Solving Inviscid Flow Problems

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A numerical method for solving incompressible viscous flow problems is introduced. This method uses the velocities and the pressure as variables and is equally applicable to problems in two and three space dimensions. The principle of the method lies in the introduction of an artificial compressibility δ into the equations of motion, in such a way that the final results do not depend on δ . An application to thermal convection problems is presented. © 1997 Academic Press.

INTRODUCTION

The equations of motion of an incompressible viscous fluid are

$$\begin{aligned} \partial_t u_i + u_j \partial_j u_i &= -\frac{1}{\rho_0} \partial_j p + \nu \Delta u_i + F_i, \quad \Delta = \sum_j \partial_j^2, \\ \partial_j u_j &= 0, \end{aligned}$$

where u_i are the velocity components, p is the pressure, F_i are the components of the external force per unit mass, ρ_0 is the density, ν is the kinematic viscosity, t is the time, and the indices i, j refer to the space coordinates x_i, x_j , $j = 1, 2, 3$.

Let L be some reference length, and U some reference velocity; we write

$$\begin{aligned} u'_i &= \frac{u_i}{U}, \quad x'_i = \frac{x_i}{L}, \quad p' = \left(\frac{d}{\rho_0 U^2} \right) p, \\ F'_i &= \frac{vU}{L^2} F_i, \quad t' = \left(\frac{v}{L^2} \right) t \end{aligned}$$

and drop the primes, obtaining the dimensionless equations

$$\partial_t u_i + R u_j \partial_j u_i = -\partial_j p + \Delta u_i + F_i, \quad (1a)$$

$$\partial_j u_j = 0, \quad (1b)$$

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where Q is a quadratic function of u_i , the boundary conditions for (2) are the same as for (1). The fact that (1b) is satisfied with velocity U is always satisfied every step so that an ingenious formula

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PRECONDITIONING METHODS FOR LOW-SPEED FLOWS

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October 1996

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On entropy generation and dissipation in high-resolution shock-capturing methods

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Abstract

This paper addresses entropy generation and the corresponding dissipation in shock-capturing (Godunov) methods. Analytical formulae are derived for arbitrary jumps in primitive variables at a cell interface. It is demonstrated that the inherent numerical entropy change of Godunov methods is not proportional to the velocity jump squared. Furthermore, it is shown to be proportional to the velocity jump squared and the speed of sound associated with jumps in pressure, density and shear waves is detailed. Further dissipation due to the perpendicular velocity jumps which dominates. The Godunov methods at low Mach numbers. The analysis is also applied to high and all analytical results are validated with simple numerical experiments. © 2008 Elsevier Inc. All rights reserved.

Keywords: High-resolution methods; Godunov methods; Dissipation; Kinetic energy number

1. Introduction

The finite volume (FV) high-resolution, shock-capturing methods have proven extremely successful in the simulation of high-speed applications in the broader field of compressible fluid dynamics. The difficulties in compressible flows and in order to provide a stable and numerically accurate scheme for solving the Euler equations, one needs to

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A low-Mach number fix for Roe's approximate Riemann solver

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ABSTRACT

We present a low-Mach fix for Roe's Riemann solver at low Mach numbers. Methods tend to zero values for the incompressible equations. Yet, stability: the artificial viscosity grows like a discrete asymptotic analysis of the Riemann problem at low Mach numbers. The effect of decreasing time-step size on the numerical convergence of reducing this term by one order of magnitude is analyzed. The numerical convergence is shown to be independent of the time-step size. As a result, the artificial viscosity is reduced by a factor of 10. Finally, it is shown that all discontinuities disappear, while at the same time, checkerboard patterns appear. Numerical tests show, first, that the accuracy of the new scheme is comparable to the original Roe's scheme. Second, the new scheme does not exhibit oscillations near discontinuities. Finally, the new scheme demonstrates the fall back of Roe's scheme at intermediate and high Mach numbers.

Several new approaches to solving low-Mach-number flows were first addressed by pressure-based solution algorithms.¹ These methods that equate pressure to zero solve for the incompressible equations in a segregated manner. The first solution algorithm was developed to solve flows at high Reynolds numbers with a single code.² The second approach is based on the assumption that the flowfields of interest consist of more geometrically complex and contain flow features that span a broad spectrum of length scales. Because it is more practical to familiarize oneself with a single code, designers, experimentalists, and other end-users of CFD software look to a single flow solver that can handle this wide range of flow regimes.

Historically, incompressible low-Reynolds-number flows were first addressed by pressure-based solution algorithms.¹ These methods that equate pressure to zero solve for the incompressible equations in a segregated manner. The first solution algorithm was developed to solve flows at high Reynolds numbers with a single code.² The second approach is based on the assumption that the flowfields of interest consist of more geometrically complex and contain flow features that span a broad spectrum of length scales. Because it is more practical to familiarize oneself with a single code, designers, experimentalists, and other end-users of CFD software look to a single flow solver that can handle this wide range of flow regimes.

In their fundamental form, time-marching schemes are useless for solving incompressible flows because the incompressible system is not fully hyperbolic, and pressure cannot be updated from an equation of state.³ The first approach to solving incompressible density-based schemes were developed in the context of transonic, compressible flow applications.⁴ These methods employed time-stepping procedures, both implicit^{5,6} and explicit.⁷ Due to the hyperbolic system of governing equations, density-based methods have also been extended to solve low-Reynolds-number and incompressible flows.^{8,9} The similarities and differences between the pressure-based and density-based approaches are discussed in detail in Ref. 9.

To meet the demand of CFD users cited earlier, a density-based flow solver is developed based on an unstructured, solution-adaptive mesh topology. The algorithm employed here are designed to compute steady-state and time-dependent flows of incompressible and variable density fluids at all speeds over a wide range of Reynolds numbers.

Unfortunately, schemes designed to capture shock waves face a number of regimes, i.e. if the Mach number tends to zero: stiffness of the equations, cancellation problem due to numerical diffusion of the order $O(1/\Delta t)$. The research effort and the approaches are diverse.

The cancellation problem was solved by Sesterhei et al. in [1]. They avoid

problems introduced in the wave propagation approach by LeVeque [2].

To overcome the cancellation problem, various flow solvers and different approaches have been developed and applied to compressible (and incompressible) flows. These schemes are generally used in conjunction with explicit methods, such as Turkel's approach, [3,4], or the characteristic time stepping approach by (almost) equalizing the propagation speeds of the different waves to steady state. At the same time, the artificial viscosity is tuned correctly for

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Preconditioning Applied to Variable and Constant Density Flows

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A time-derivative preconditioning of the Navier-Stokes equations, suitable for both variable- and constant-density fluids, is developed. The ideas of low-Mach-number preconditioning and artificial compressibility are combined into a unified approach to enhance convergence rates of density-based, time-marching schemes for solving flows of incompressible and variable density fluids at all speeds. The preconditioning is coupled with a dual time-stepping scheme implemented within an explicit, multistage algorithm for solving time-accurate flows. The resultant time integration scheme is used in conjunction with a finite volume discretization designed for unstructured, solution-adaptive mesh topologies. This method is shown to provide accurate steady-state solutions for transonic and low-speed flows of variable density fluids. The time-accurate solution of unsteady, incompressible flow is also demonstrated.

Introduction

The use of computational fluid dynamics (CFD) technologies has permeated throughout industry, academia, and the research community. As such, numerical flow solvers are relied upon to solve a wide class of problems ranging from incompressible, low-Reynolds-number flows to high-speed, variable density flows at high Reynolds numbers. The development of the field of interest has led to more geometrically complex and contain flow features that span a broad spectrum of length scales. Because it is more practical to familiarize oneself with a single code, designers, experimentalists, and other end-users of CFD software look to a single flow solver that can handle this wide range of flow regimes.

Historically, incompressible low-Reynolds-number flows were first addressed by pressure-based solution algorithms.¹ These methods that equate pressure to zero solve for the incompressible equations in a segregated manner. The first solution algorithm was developed to solve flows at high Reynolds numbers with a single code.² The second approach is based on the assumption that the flowfields of interest consist of more geometrically complex and contain flow features that span a broad spectrum of length scales. Because it is more practical to familiarize oneself with a single code, designers, experimentalists, and other end-users of CFD software look to a single flow solver that can handle this wide range of flow regimes.

In their fundamental form, time-marching schemes are useless for solving incompressible flows because the incompressible system is not fully hyperbolic, and pressure cannot be updated from an equation of state.³ The first approach to solving incompressible density-based schemes was developed in the context of transonic, compressible flow applications.⁴ These methods employed time-stepping procedures, both implicit^{5,6} and explicit.⁷ Due to the hyperbolic system of governing equations, the system becomes hyperbolic and a means to update pressure is required. Since solving the equations in a segregated manner is not possible, the pressure is added to the system in all of the equations.^{8,9} In either case, the pressure derivative is normalized by a pseudospectral speed (squared). The pseudospectral speed is typically set to twice the local velocity, such that a pseudospectral Mach number of one is achieved, thereby providing optimal convergence.

It is our objective to combine the ideas of low-Mach-number preconditioning and artificial compressibility into a unified approach and produce a preconditioned matrix that will provide for efficient solution of both steady-state and variable density flows at all speeds. The development of this new preconditioner follows closely with that given in Ref. 12 for compressible flows with three notable exceptions: 1) derivatives of density with respect to pressure and temperature are included within the preconditioner, 2) the density is a variable that is dependent on the pressure, 3) the density is normalized by a pseudospectral speed (squared). The pseudospectral speed is typically set to a local speed of sound and reference Mach number, and 3) derivatives of density with respect to temperature are retained. It has been our experience that this latter scheme is more robust for solving flows in which density is a function of temperature only.

Because time-derivative preconditioning destroys the time accuracy of the flow, it is necessary to take steps to ensure that the flow is not polluted by these means alone. To overcome this limitation, we employ a dual time-stepping^{10,11} procedure. This involves an inner iteration loop in pseudotime that is wrapped by an outer loop step through physical time. Thus, the flowfield at each physical time

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Background Coupled Approaches (Density-based) in SU2

(Reprint of 1967 paper)

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A Numerical Method for Solving Incompressible Viscous Flow Problems*

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A numerical method for solving incompressible viscous flow problems is introduced. This method uses the vorticity and the pressure as variables and is equally applicable to problems in two and three space dimensions. The principle of the method lies in the introduction of an artificial compressibility δ into the equations of motion, in such a way that the final results do not depend on δ . An application to thermal convection problems is presented. © 1997 Academic Press

INTRODUCTION

The equations of motion of an incompressible viscous fluid are

$$\partial_t u_i + u_j \partial_j u_i = -\frac{1}{\rho_0} \partial_j p + \nu \Delta u_i + F_i, \quad \Delta = \sum_j \partial_j^2,$$

$$\partial_i u_j = 0,$$

where u_i are the velocity components, p is the pressure, F_i are the components of the external force per unit mass, ρ_0 is the density, ν is the kinematic viscosity, Δ is the time, and the indices i, j refer to the space coordinates $x_i, x_j, i, j = 1, 2, 3$.

Let d be some reference length, and U some reference velocity; we write

$$u'_i = \frac{u_i}{U}, \quad x'_i = \frac{x_i}{d}, \quad p' = \left(\frac{d}{\rho_0 \nu U} \right) p,$$

$$F'_i = \frac{vU}{d^2} F_i, \quad t' = \left(\frac{v}{d^2} \right) t$$

and drop the primes, obtaining the dimensionless equations

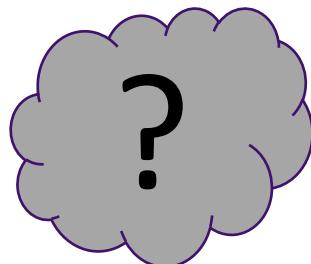
$$\partial_t u_i + R u_j \partial_j u_i = -\partial_i p + \Delta u_i + F_i, \quad (1a)$$

$$\partial_i u_j = 0, \quad (1b)$$

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where R is a quadratic function of the velocities and, eventually, a function also of the external forces. Boundary conditions for (1a) can be obtained from (1a) applied at the boundary. There remains, however, the task of ensuring that (1b) is satisfied. This is done by starting the calculation with velocity fields satisfying (1b), making sure that (1b) is always satisfied at the boundary, and solving (2) at every step so that (1b) remains satisfied as time is advanced. An ingenious formulation of the finite difference form of



Let's develop a robust coupled method more general than Artificial Compressibility w/out complexity of fully compressible N-S.

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Preconditioning Applied to Variable and Constant Density Flows

Jonathan M. Weiss* and Wayne A. Smith*
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A time-derivative preconditioning of the Navier–Stokes equations, suitable for both variable and constant density flows, is developed. The idea of low-Mach-number preconditioning and artificial compressibility are combined into a unified approach designed to enhance convergence rates of density-based, time-marching schemes for solving flows of incompressible and variable density fluids at all speeds. The preconditioning scheme implemented within an explicit, multistage algorithm for solving time-accurate flows, the resultant time-integration scheme is used in conjunction with a finite volume discretization designed for unstructured, solution-adaptive mesh topologies. This method is shown to provide accurate steady-state solutions for transonic and low-speed flow of variable density fluids. The time-accurate solution of unsteady, incompressible flow is also demonstrated.

Introduction

The use of computers in fluid dynamics (CFD) techniques is expanding throughout industry, academia, and the research community. As such, numerical flow solvers are relied upon to solve a wide class of flow problems ranging from incompressible, low-Reynolds-number flows to high-speed, compressible flows at high Reynolds numbers. In addition, the flow fields of interest are often numerically complex and contain flow features that span a range of length scales. Because it is more practical to work with a single code, designers, experimentalists, and other end-users of CFD software look to a single flow solver that can handle all of these requirements.

Historically, incompressible low-Reynolds-number flows were first addressed by pressure-based solution algorithms.^{1–3} In these methods the equations of motion are solved in a segregated uncoupled manner relying on diagonal dominance to converge. Pressure-based algorithms have since been extended to solve flows at high Reynolds numbers and compressible flows^{4–6} as well. Alternatively, density-based schemes were developed in the context of transonic flow calculations. These density-based schemes are time-marching procedures, both implicit^{7–9} and explicit,^{10–12} that solve the hyperbolic system of governing equations. Density-based methods have also been extended to solve low-Reynolds-number and incompressible flows.^{7,8} The similarities and differences between the pressure-based and density-based approaches are discussed in further detail in Ref. 9.

To meet the demand of CFD users cited earlier, a density-based flow solver is developed based on an unstructured, solution-adaptive mesh topology. The algorithms herein are designed to compute steady-state and time-dependent flows of incompressible and compressible fluids at all speeds over a wide range of Reynolds numbers.¹³

We have chosen to implement the density-based flow solver for incompressible flows because the incompressible problem is not fully hyperbolic, and pressure cannot be updated from an equation of state. This deficiency is overcome by employing an artificial compressibility scheme. The density term in the time-derivative is introduced into the continuity equation. With the artificial pressure term, the system becomes hyperbolic and a means to update pressure is provided. When solving the equations in conservation form, it has been recommended that the time step be included in the form of a pseudotime.¹⁴ In this case, the present time-stepping is normalized by a pseudosonic speed (squared). The pseudosonic speed is typically set to about twice local velocity, such that a pseudo-Mach number of one-half is achieved, thereby providing optimal convergence.

It is our objective to combine the ideas of low-Mach-number preconditioning and artificial compressibility into a unified approach and produce a preconditioning matrix that will provide for efficient solution of both steady-state and time-dependent flows at all speeds. The derivation of this new preconditioning follows the approach that gives in Ref. 12 for compressible flows with three notable exceptions: 1) derivatives of density with respect to pressure and temperature are carried through without the assumption of an ideal gas law; 2) the pseudosonic speed is used to calculate the eigenvalues of the system; 3) the eigenvalues of the system is written in terms of a characteristic reference velocity as opposed to a local speed of sound and reference Mach number; and 3) derivatives of density with respect to temperature are retained. It has been our experience that this latter feature is important for solving flows in which density is a function of temperature only.

Because time-derivative preconditioning destroys the time accuracy of the governing equations, the solution of a steady flow is not possible by these means. However, after a few iterations, we enter a dual time-stepping^{14,15} procedure. This involves an inner iteration loop in pseudotime that is wrapped by an outer loop stepping through physical time. Thus, the flowfield at each physical time

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MODELING & IMPLEMENTATION

Modeling & Implementation

Governing Equations: low-Mach N-S (variable density)

- Eqns. in conservative form in a domain Ω with solid wall boundary S and inlet/outlet boundaries:

$$\left\{ \begin{array}{ll} R(U) = \frac{\partial U}{\partial t} + \nabla \cdot \bar{F}^c(U) - \nabla \cdot \bar{F}^v(U, \nabla U) = 0, & \text{in } \Omega, \\ \bar{v} = 0, & \text{on } S, \\ T = T_S, & \text{on } S, \\ (W)_{\bar{v}} = W_{in}, & \text{on } \Gamma_{in}, \\ (W)_{\bar{P}} = W_{out}, & \text{on } \Gamma_{out}, \end{array} \right.$$

$$U = \{\rho, \rho\bar{v}, \rho c_p T\}^\top \quad \bar{F}^c(U) = \begin{Bmatrix} \rho\bar{v} \\ \rho\bar{v} \otimes \bar{v} + \bar{I}p \\ \rho c_p T\bar{v} \end{Bmatrix}, \quad \bar{F}^v(U, \nabla U) = \begin{Bmatrix} \dot{\tau} \\ \bar{\tau} \\ \kappa \nabla T \end{Bmatrix},$$

$$\rho = \frac{p_o}{R T}$$

$$\bar{\tau} = \mu (\nabla \bar{v} + \nabla \bar{v}^T) - \mu \frac{2}{3} \bar{I} (\nabla \cdot \bar{v}).$$

Modeling & Implementation

Coupled Approach I/II

- We follow the approach of Weiss & Smith [1995] with some notable differences... in the end we have:

$$R(V) = \Gamma \frac{\partial V}{\partial t} + \nabla \cdot \bar{F}^c(V) - \nabla \cdot \bar{F}^v(V, \nabla V) = 0$$

$$\Gamma = \begin{bmatrix} \frac{1}{\beta^2} & 0 & 0 & 0 & \rho_T \\ \frac{u}{\beta^2} & \rho & 0 & 0 & \rho_T u \\ \frac{v}{\beta^2} & 0 & \rho & 0 & \rho_T v \\ \frac{w}{\beta^2} & 0 & 0 & \rho & \rho_T w \\ \frac{c_p T}{\beta^2} & 0 & 0 & 0 & \rho_T c_p T + \rho c_p \end{bmatrix}$$

Modeling & Implementation

Coupled Approach II/II

- ▶ What exactly is Beta?
 - ▶ It is an artificial sound speed
 - ▶ We can see this clearly by comparing the eigenvalues of convective flux Jacobian to the compressible case.
 - ▶ **Draws clear link between the Artificial Compressibility and preconditioning approaches.**

$$\lambda \left(\Gamma^{-1} \frac{\partial \bar{F}^c}{\partial V} \right) = \lambda \left(\Gamma^{-1} \bar{A}^c \right) = \begin{bmatrix} \bar{v} \cdot \bar{n} & 0 & 0 & 0 & 0 \\ 0 & \bar{v} \cdot \bar{n} & 0 & 0 & 0 \\ 0 & 0 & \bar{v} \cdot \bar{n} & 0 & 0 \\ 0 & 0 & 0 & \bar{v} \cdot \bar{n} - \beta |\bar{n}| & 0 \\ 0 & 0 & 0 & 0 & \bar{v} \cdot \bar{n} + \beta |\bar{n}| \end{bmatrix}$$

$$\beta^2 = \epsilon^2 (\bar{v} \cdot \bar{v})_{max}$$

Modeling & Implementation

Present Developments

- ▶ A **custom preconditioning method** for low speed flows with heat transfer:
 - ▶ Simplicity: fully compressible N-S not required
 - ▶ Euler, N-S, and RANS (low-Mach, decoupled energy, or isothermal)
 - ▶ Conservative formulation
 - ▶ Primitive variable-based, $V = \{p, u, v, w, T\}$
 - ▶ Custom Flux Difference Splitting (upwind) and centered schemes
 - ▶ Implicit & explicit time integration for steady relaxation, time-accurate flows with dual time-stepping
- ▶ Enables **variable density incompressible flows**:
 - ▶ Introduces 2 new fluid models: constant density fluid, incompressible ideal gas
- ▶ Includes **energy equation**:
 - ▶ New energy eqn. options: disable, solve decoupled, apply Boussinesq approximation, or couple for variable density

Simulation and Adjoint-based Design for Variable Density Incompressible Flows with Heat Transfer

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This article details the development and implementation of an incompressible solver for simulation and design in variable density incompressible flows with heat transfer. In the low-Mach approximation of the Navier-Stokes equations, density can vary as a function of transported scalars, and in this case, density varies with temperature from a coupled energy equation. These governing equations are spatially discretized using a finite volume method on unstructured grids and solved in a primal manner with a custom preconditioning approach. The implementation is within the SU2 suite for multiphysics simulation and design, and it has been algorithmically differentiated to construct a discrete adjoint for efficient sensitivity analysis. Results demonstrating the primal solver on a set of standard verification and validation cases and adjoint-based shape optimization are presented.

I. Introduction

In this article, we pursue the development and implementation of a solver for variable density incompressible flows with heat transfer. In many applications, such as natural (buoyancy-driven) or forced convection problems, environmental flows, fire simulations, or for reacting flows, such as combustion simulations. In these situations, the Mach number can be very small, but the effects of heat transfer and the accompanying variations in density remain important.

One approach for this regime is to apply the fully compressible form of the Navier-Stokes equations for conservation of mass, momentum, and energy. Unfortunately, it is well known that the equations become very stiff at low Mach numbers, resulting in poor convergence behavior for density-based, compressible codes, and the numerical methods applied typically suffer from accuracy issues due to artificial dissipation that is poorly scaled at small Mach numbers (related to disparate scales of convection/acoustics). Preconditioning approaches for the compressible Navier-Stokes equations at low speeds can be a remedy, and they have been successfully demonstrated in literature by many authors.^{1–4} These approaches can be an ideal choice for flows with mixed high and low Mach numbers. However, they carry more complexity than necessary for purely low-speed flows, which could lead to convergence or performance issues or more restrictions on the numerics.

On the other hand, the flow can be treated as incompressible. In order to include heat transfer effects in incompressible flows, the energy equation, or a temperature evolution equation, must be solved in addition to the continuity and momentum equations for the fluid. The specific coupling of the energy equation will depend on the situation. For constant density fluids, the energy equation can be solved with a one-way coupling, essentially as a passive scalar, or be two-way coupled through the Boussinesq approximation for problems with suitably small temperature variations. However, for some problems, large density variations are critical even at very small Mach numbers, such as in reacting flows, and a more elaborate model is necessary. Here, the low-Mach number formulation of the equations is an attractive choice.^{9–11}

The appeal of the low-Mach Navier-Stokes equations is the ability to treat incompressible fluids that feature large density variations while avoiding the complexity of the fully compressible form of the Navier-Stokes equations. Density is decoupled from pressure and determined from an equation of state that is a function of transported scalars, such as temperature. This decoupling of the thermodynamic pressure removes the acoustics from the equations. Practically speaking, the low-Mach approximation, arrived at

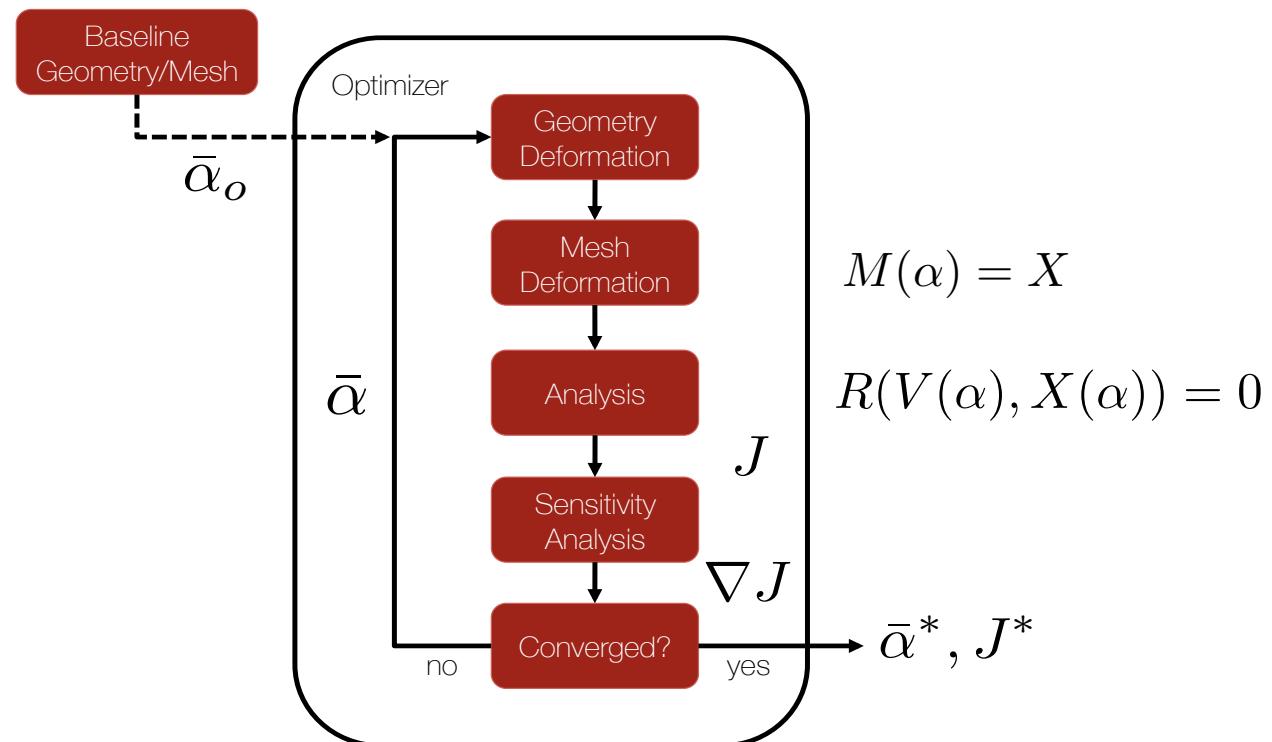
*Senior Research Scientist, Multiphysics Modeling and Simulation, AIAA Senior Member.

Modeling & Implementation

Gradient-based Optimal Shape Design

- ▶ Input: a baseline geometry/mesh and a chosen parameterization (α) controlling shape, and J .
- ▶ Primal gives us J , adjoint gives us the gradient efficiently.
- ▶ Meshes are deformed with pseudo-structural approach (operator M).
- ▶ Numerical optimizer drives the problem to a local optimum J^* with final geometry α^* .
- ▶ See Albring et al. 2015, Albring et al. 2016 for full details of discrete adjoint in SU2 with CoDiPack.

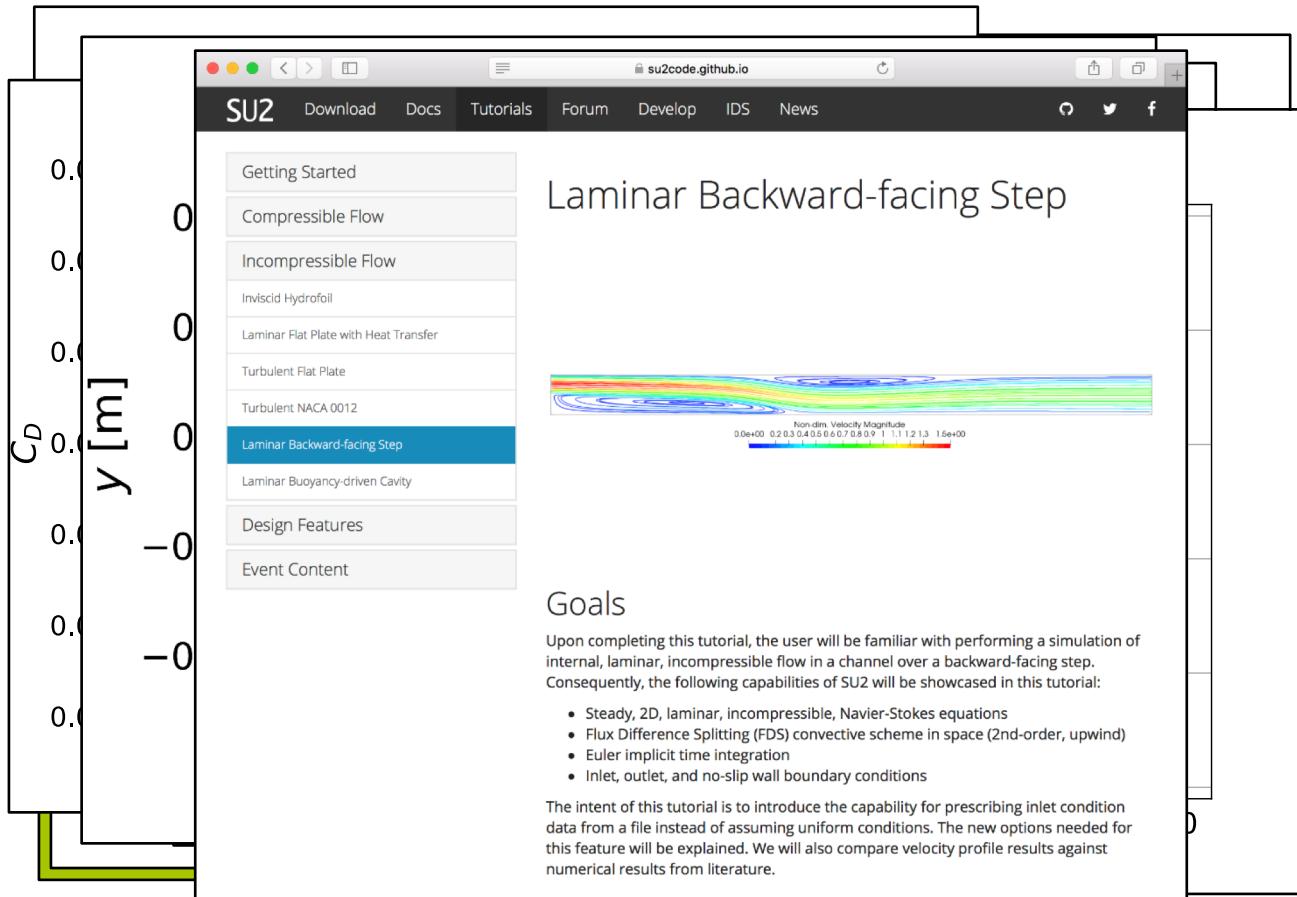
$$\begin{aligned} \min_{\alpha} \quad & J(V(\alpha), X(\alpha)) \\ \text{subject to} \quad & R(V(\alpha), X(\alpha)) = 0 \end{aligned}$$



RESULTS (V&V)

Results

Verification & Validation (V&V)



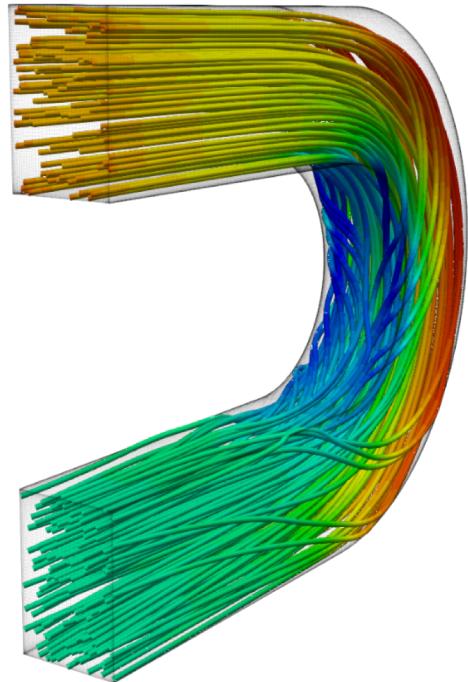
- ▶ Inviscid Hydrofoil
- ▶ Buoyancy-driven Cavity
- ▶ Laminar Flat Plate
- ▶ Turbulent Flat Plate
- ▶ Turbulent NACA 0012
- ▶ Turbulent 3D Bump-in-Channel
- ▶ Axisymmetric Pipe
- ▶ Laminar Backward-facing Step

- ▶ Excellent agreement for all comparisons against theory, well established codes, and experiment.

Results

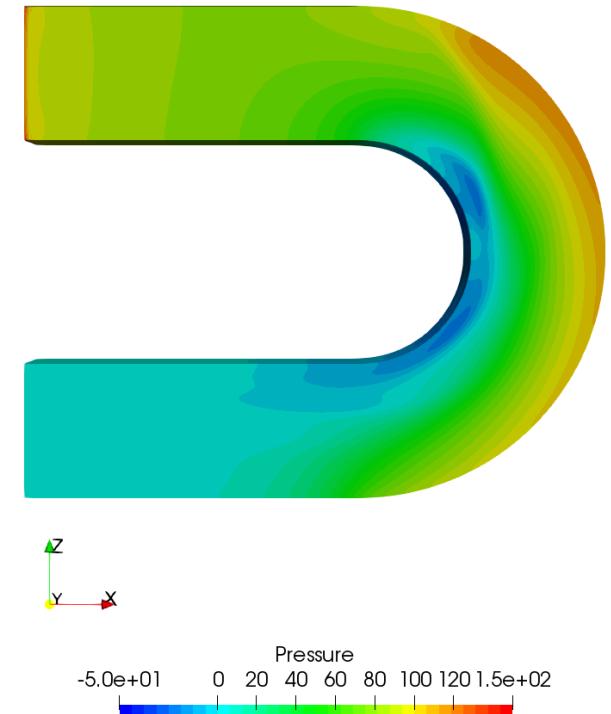
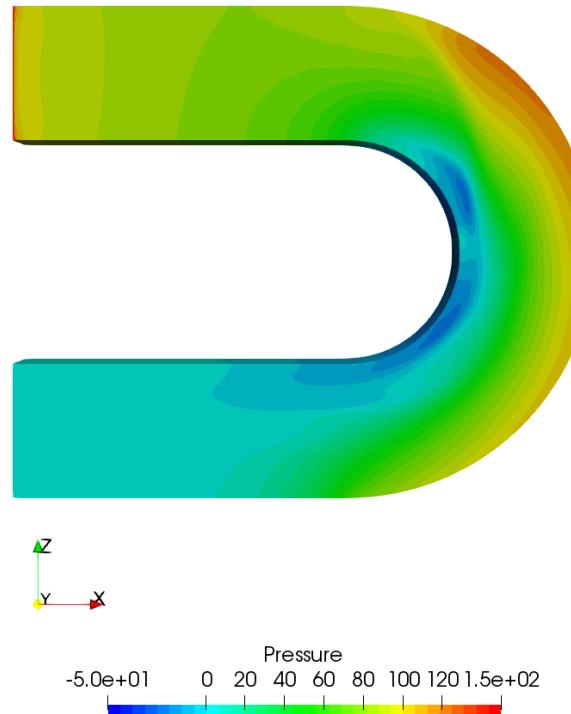
Code Comparison

SU2 | Fluent



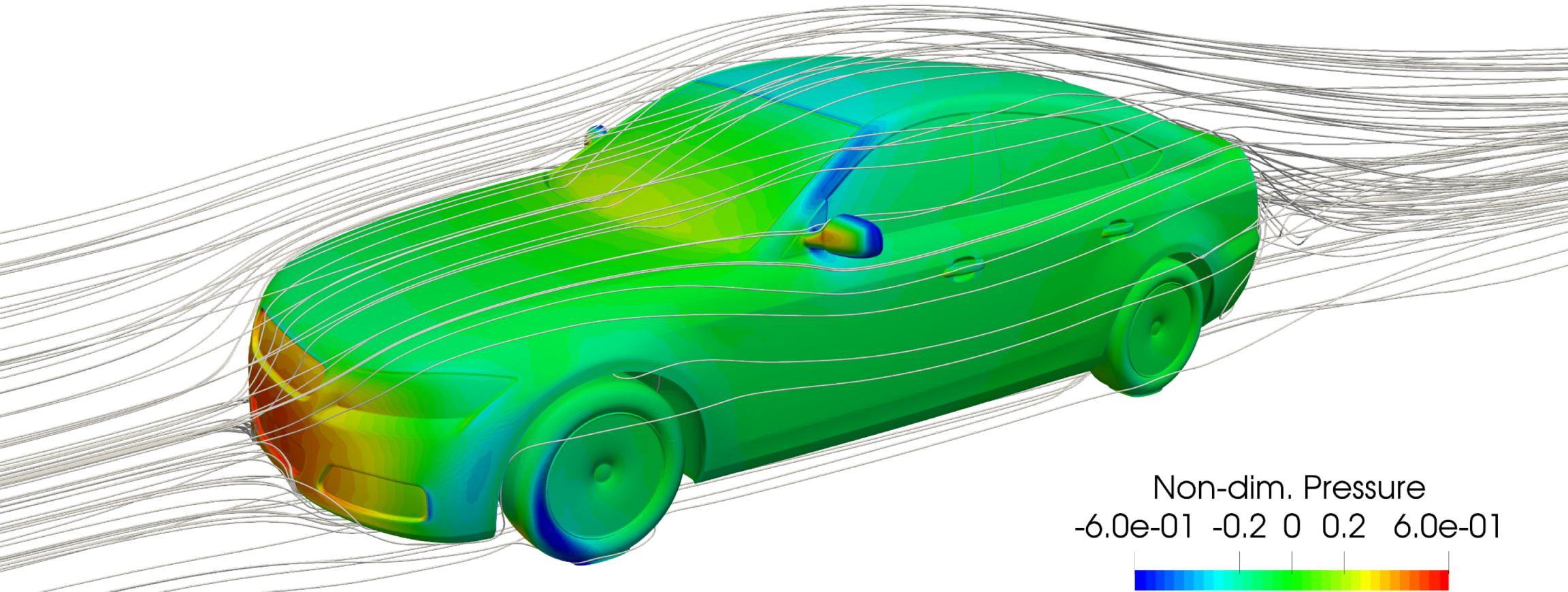
Turbulent flow through a rectangular u-bend.

Re ~ 4300.



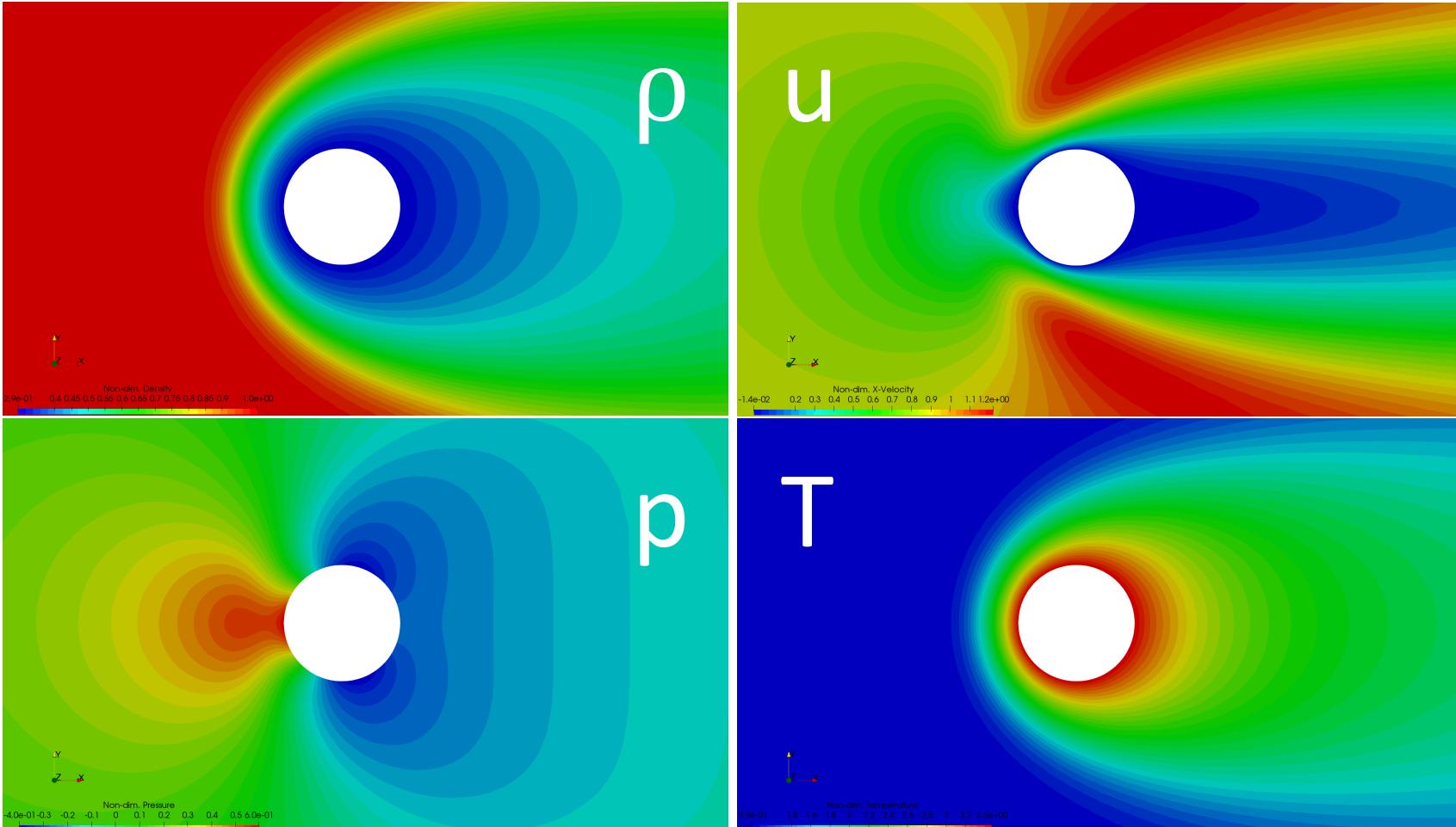
Results

Complex Geometry



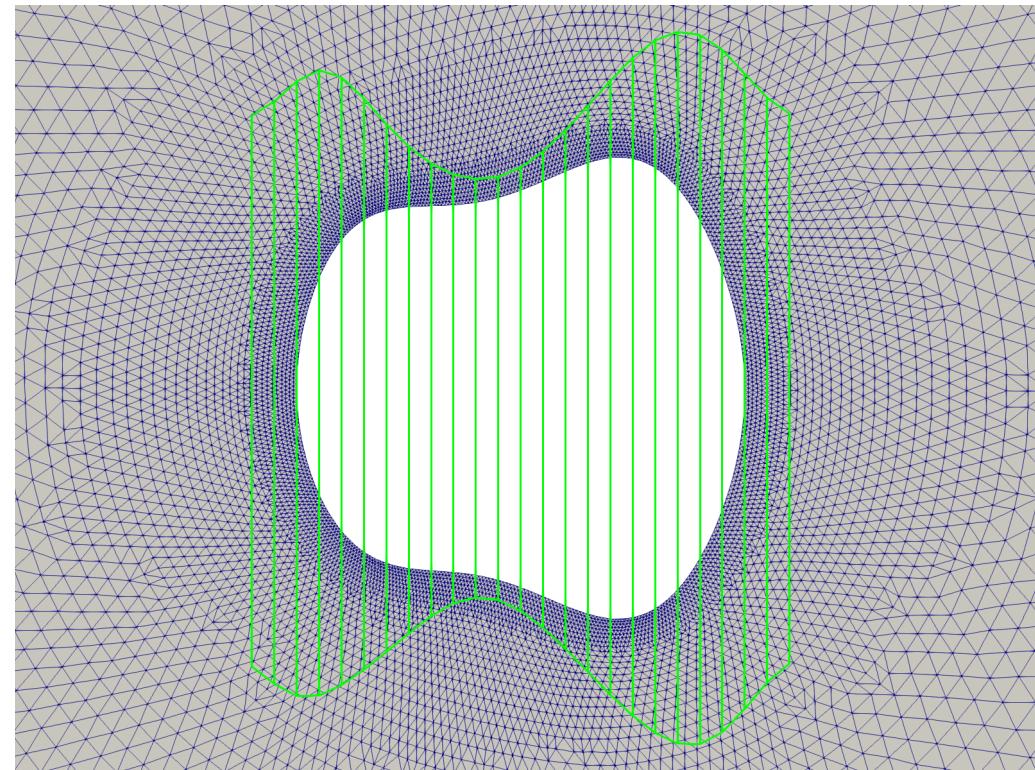
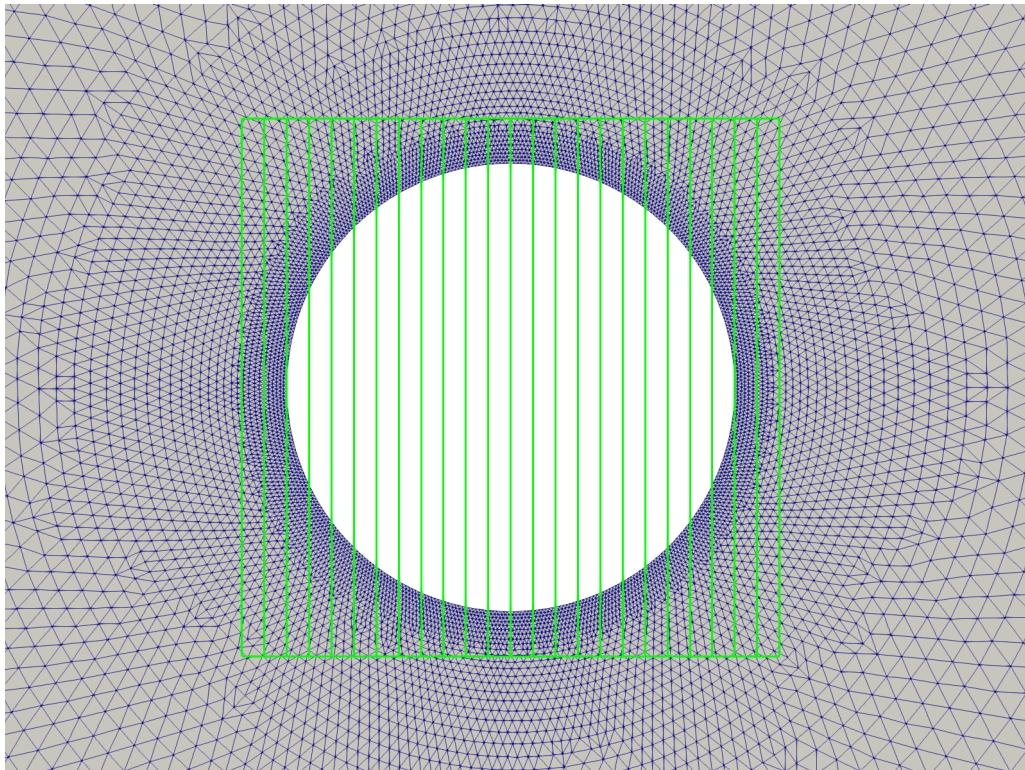
Results

Shape Optimization of a Heated Cylinder: Primal



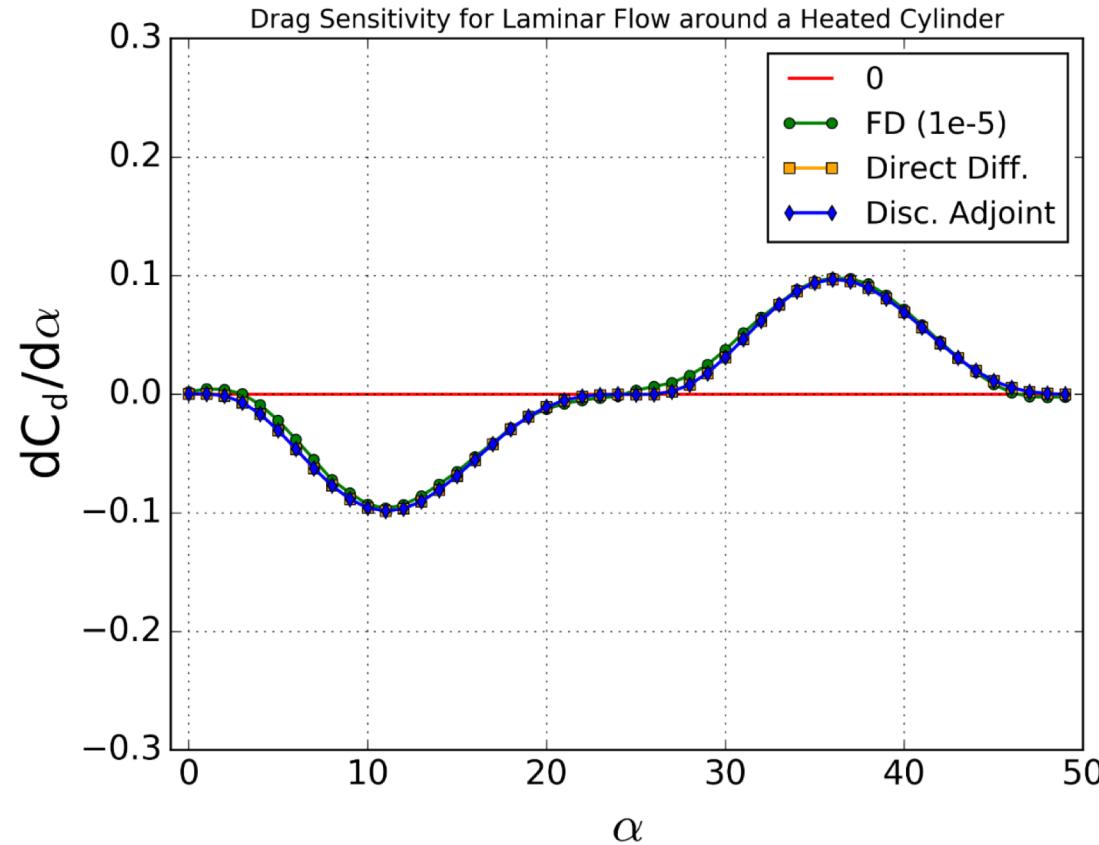
Results

Shape Optimization of a Heated Cylinder: Parameterization



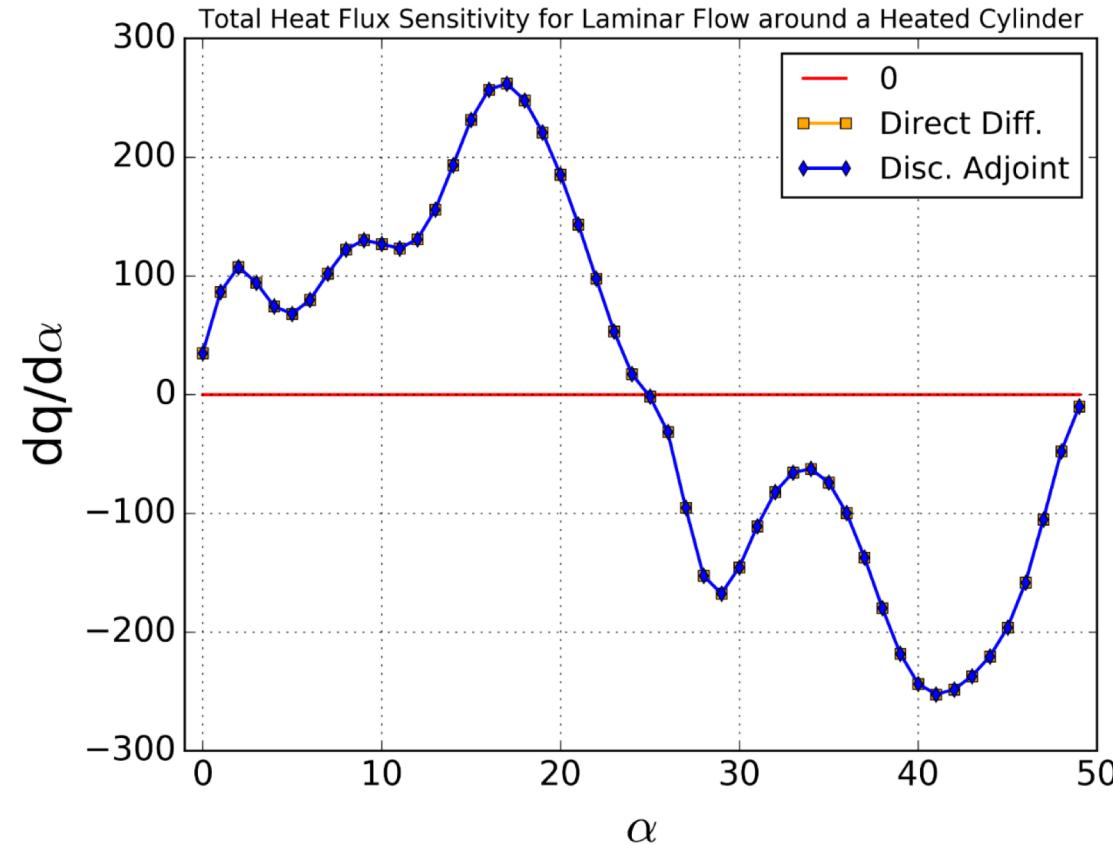
Results

Shape Optimization of a Heated Cylinder: Drag Sens. Verification



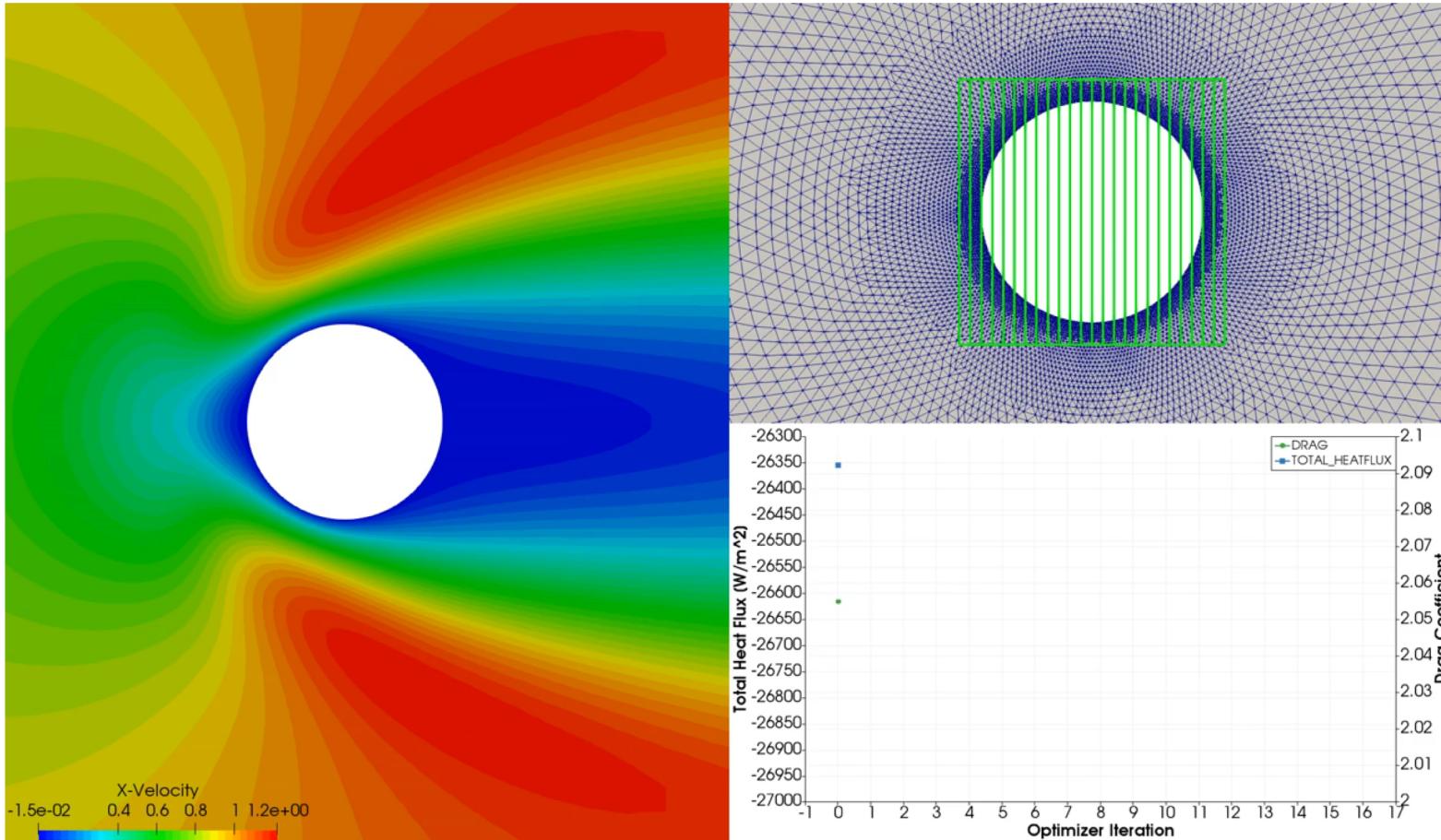
Results

Shape Optimization of a Heated Cylinder: Heat Flux Sens. Verification



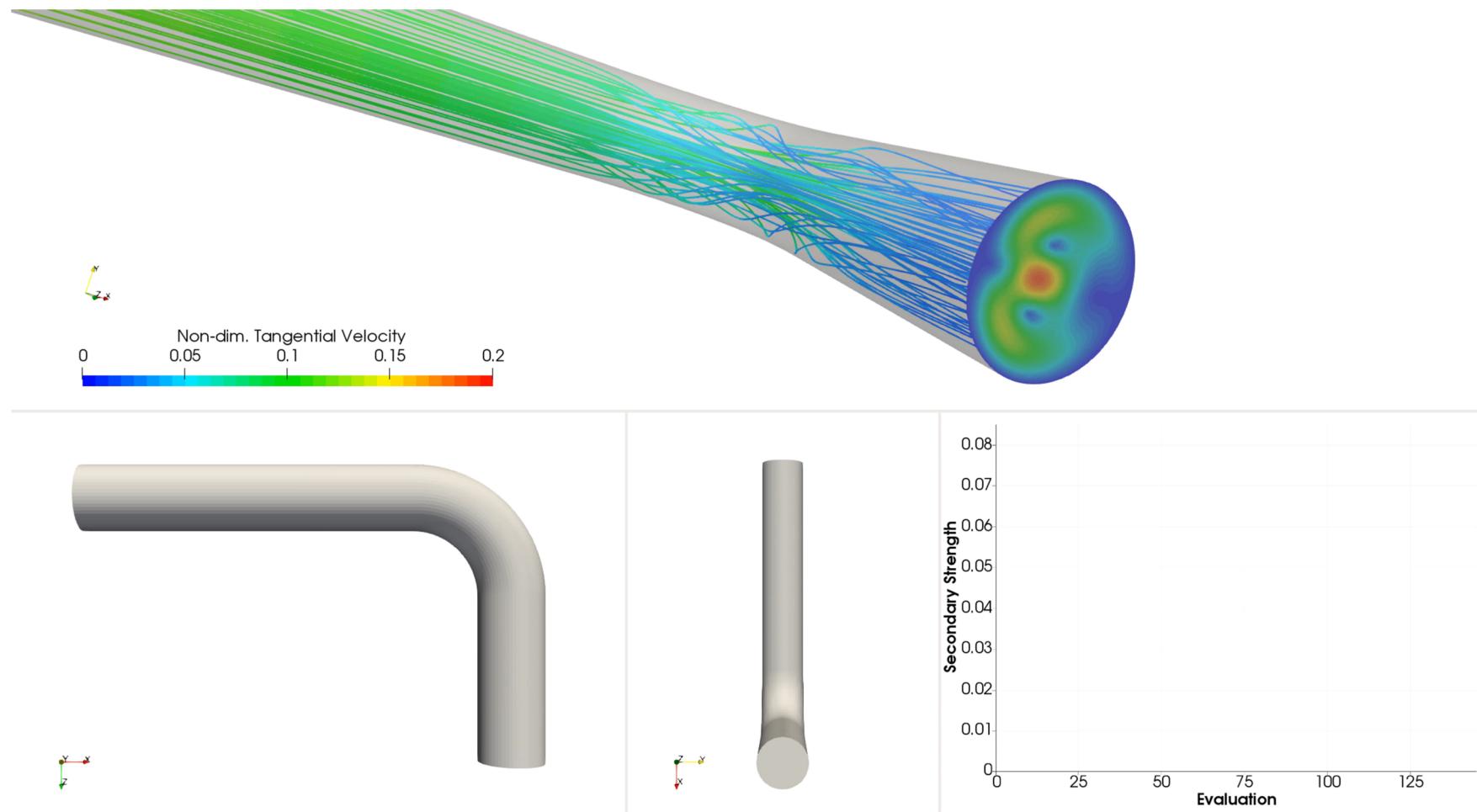
Results

Shape Optimization of a Heated Cylinder: Heat Flux w/ Cd Constraint



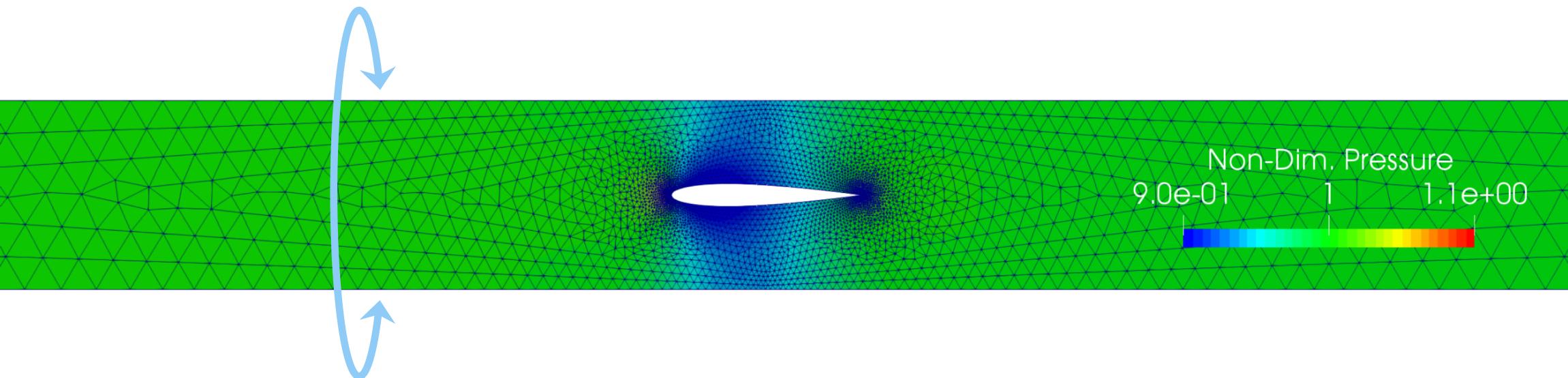
V&V

90 Degree Bend Optimization Example



One more thing...

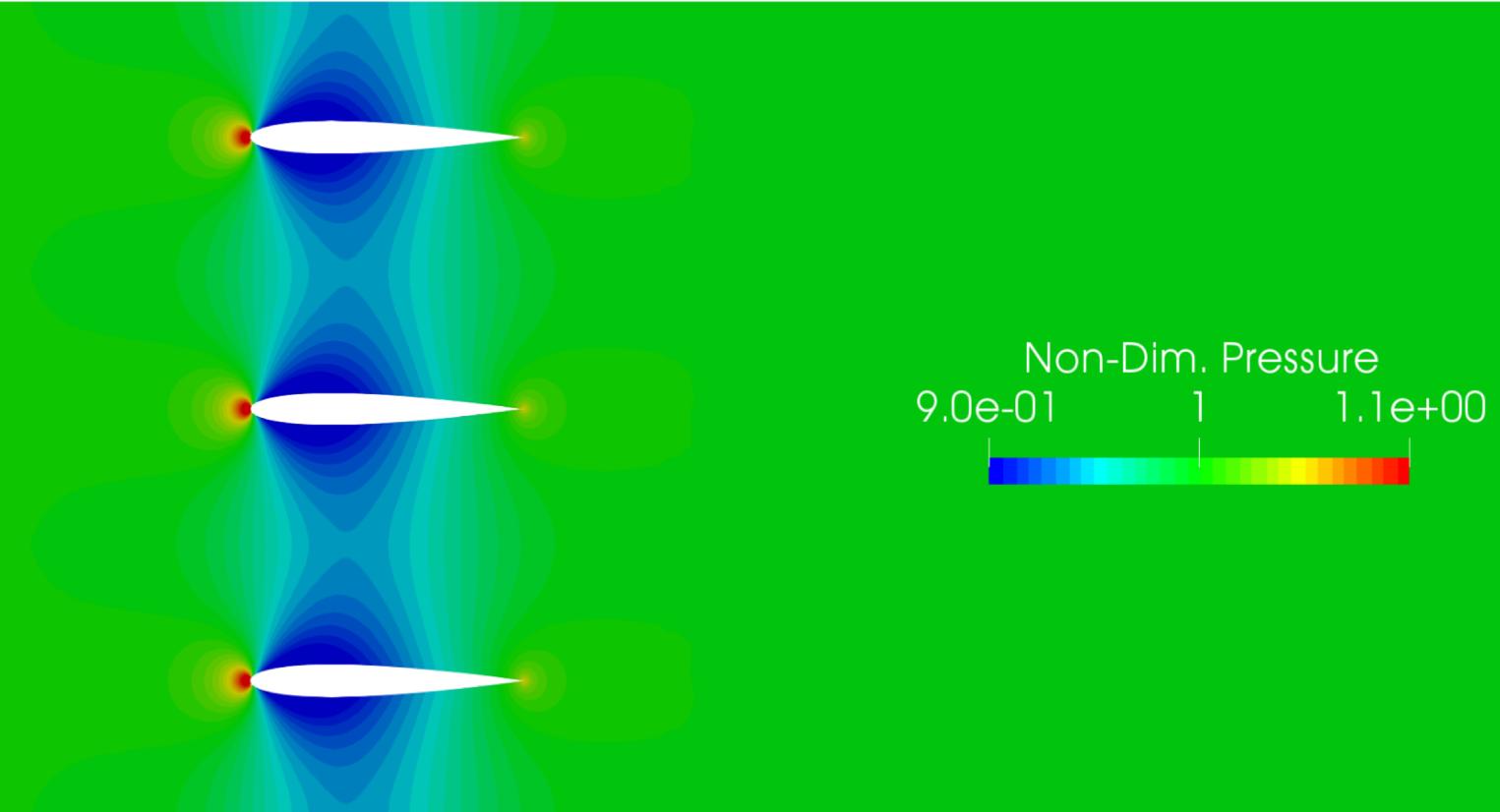
Periodic Boundary Condition



periodic boundary condition now based on concept of completing the residual rather than halo cells

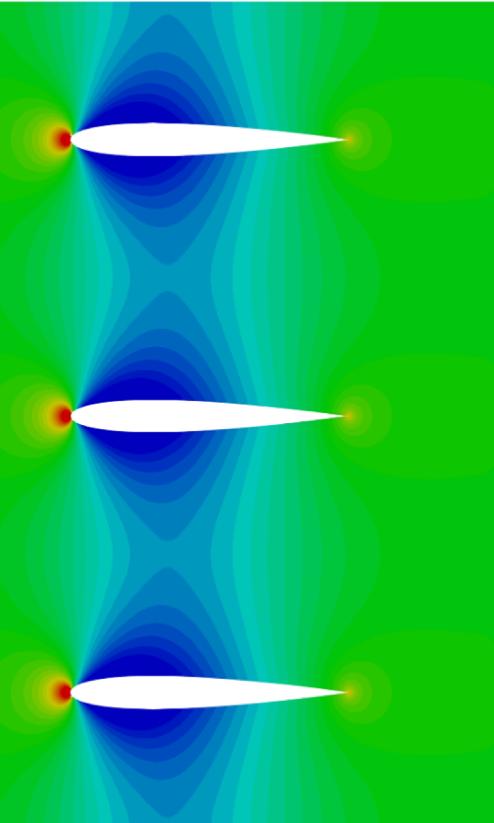
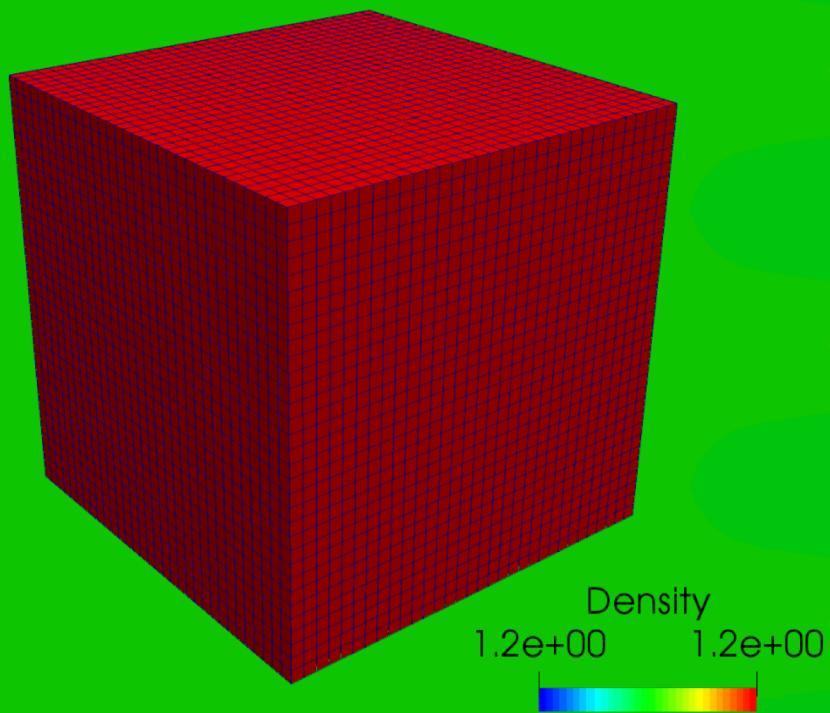
One more thing...

Periodic Boundary Condition



One more thing...

Periodic Boundary Condition



*ParaView binary output now available
(parallel implementation with MPI-IO)*

Conclusions

Key Messages from Today

- ▶ Showed a density-based preconditioning approach for a range of incompressible flows.
 - ▶ Seen as either a generalization of Artificial Compressibility or simplification of Weiss & Smith [1995].
 - ▶ In author's experience, preconditioning all eqns. is critical for robustness when adding energy eqn. or turbulence model.
- ▶ Showed V&V of method with classic and NASA turbulence modeling cases. Demonstrated shape design.
- ▶ V&V results are reproducible with open data, source code available to public, tutorials covering usage online.
- ▶ Meant as a reference for the SU2 community to build on for incompressible flows.

