7.16 (a) From  $\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$  and  $\nabla \times \vec{H} = \frac{\partial \vec{D}}{\partial t}$  where no current is present, we have  $\nabla \times (\nabla \times \vec{E}) = -\nabla \times \frac{\partial \vec{B}}{\partial t} = -\frac{\partial}{\partial t} (\nabla \times \vec{B}) = -M_0 \frac{\partial}{\partial t} (\nabla \times \vec{H}) = -M_0 \frac{\partial^2 \vec{D}}{\partial t}$ 

Assuming  $\vec{E} \propto e^{i\vec{k}\cdot\vec{v}-i\omega t}$  and consequently,  $\vec{D} \propto e^{i\vec{k}\cdot\vec{v}-i\omega t}$ , the above equation becomes  $i\vec{k} \times (i\vec{k} \times \vec{E}) = \mu_0 \omega^* \vec{D}$ , or  $\vec{k} \times (\vec{k} \times \vec{z}) + \mu_0 \omega^* \vec{D} = 0$ .

(b) Using the identity  $\vec{A} \times (\vec{b} \times \vec{c}) = (\vec{a} \cdot \vec{c})\vec{b} - (\vec{a} \cdot \vec{b})\vec{c}$ , the above equation can be transformed to  $\vec{k}(\vec{k} \cdot \vec{E}) - \vec{k} \cdot \vec{E} + \mu_0 \omega \cdot \vec{b} = 0$ . Let  $\vec{k} = k \vec{n}$ , the equation can be rewritten as  $|\vec{n}(\vec{n} \cdot \vec{E}) - \vec{E} + \mu_0 v \cdot \vec{b} = 0$ , where v = w/k.

Since  $D_i = E_i E_i$ , i = 1, 2, 3, the above equation has the following component-wise form  $N_i (\vec{n} \cdot \vec{E}) - E_i + \frac{v^*}{v_i} E_i = 0$ 

or  $\overline{E}_{i} = \frac{n_{i}(\vec{n} \cdot \vec{E})}{1 - v_{i}^{\prime}v_{i}^{2}}$ 

Multiply both sides by N: and take the sum, we are left with

$$\vec{N} \cdot \vec{E} = \sum_{i=1}^{3} \frac{\vec{N}_{i}}{|-\vec{V}/\vec{V}_{i}|^{2}} (\vec{n} \cdot \vec{E}), \quad \text{or} \quad \frac{3}{2} \frac{\vec{n}_{i}}{|-\vec{V}/\vec{V}_{i}|^{2}} = |-\frac{3}{2} \vec{n}_{i}|^{2}$$

Therefore,  $\frac{3}{2} \operatorname{Ni}\left(\frac{1}{1-V/v_i^2}-1\right)=0$ ,  $\Rightarrow \frac{3}{2} \operatorname{Ni}\left(\frac{v_i^2}{V_i^2-V_i^2}=0\right)$ 

which leads to  $\frac{3}{2} \frac{n_i^2}{v^2 - v_i^2} = 0$ 

The Fresnel condition is equivalent to a quadratic equation in "",

 $\mathcal{V}^4 - \left( n_1^2 \left( \nu_1^2 + \nu_3^2 \right) + n_2^2 \left( \nu_3^2 + \nu_1^2 \right) + n_3^2 \left( \nu_1^2 + \nu_2^2 \right) \right) \mathcal{V}^2 + n_1^2 \mathcal{V}_2^2 \mathcal{V}_1^2 + n_2^2 \mathcal{V}_2^2 \mathcal{V}_1^2 + n_3^2 \mathcal{V}_2^2 \mathcal{V}_2^2 = 0$ This has two distinct Solutions. as it clear from the coefficients.

(c) The component wise equation can also be written as

$$N_{i} \sum_{\hat{j}} N_{j} \frac{D_{j}}{\xi_{j}} - \frac{D_{\hat{k}}}{\xi_{i}} + \mu_{0} V^{\nu} D_{i} = 0 \implies N_{i} \sum_{\hat{j}} N_{j} V_{\hat{s}}^{\nu} D_{\hat{j}} - V_{\hat{k}}^{\nu} D_{\hat{k}} + V^{\nu} D_{i} = 0$$

Multiply by  $N_i$  and sum, noticing  $\tilde{Z}$   $N_i^*=1$ , we have  $\tilde{N}_i$ ,  $\tilde{D}=0$ , which means that the two distinct modes in the displacement are transverse. For the two modes  $\tilde{D}_n$  and  $\tilde{D}_b$ , the component wise equations are

$$N_{i} \left[ \sum_{j=1}^{n} N_{j}^{2} D_{aj} - N_{i}^{2} D_{ai} + V_{a}^{2} D_{ai} = 0 \right]$$
 (1)

$$N_i = \overline{l} N_j V_j^* D_{bj} - V_i^* D_{bi} + V_b^* D_{bi} = 0$$
 (3)

Multiply (1) by Do; and sum, and is by Dai and sum, we have

$$\left(\sum_{i} n_{i} D_{bi}\right) \cdot \left(\sum_{j} n_{j} V_{j}^{*} D_{aj}\right) - \sum_{i} N_{i}^{*} D_{ai} D_{bi} + V_{a}^{*} D_{ai} D_{bj} = 0$$

$$(3)$$

$$\left(\begin{array}{ccc} \overline{\lambda} & \Omega_{i} & D_{ai} \end{array}\right) & \left(\begin{array}{ccc} \overline{\lambda} & \eta_{j} & V_{j} & D_{bj} \end{array}\right) & - & \overline{\lambda} & V_{i}^{\dagger} & D_{ai} & D_{bi} & + & V_{b}^{\dagger} & D_{ai} & D_{bi} & = 0 \end{array} \tag{4}$$

Take the difference between 13) and (4) and use the transversality of  $\vec{D}$ , we are left with  $(\vec{Va} - \vec{V_b}) \vec{D_a} \cdot \vec{D_b} = 0$ 

Since  $V_a^{\dagger}$  and  $V_b^{\star}$  are distinct, we must have  $\tilde{D}_a \cdot \tilde{P}_b = 0$