5.36 (a) From 
$$\vec{E} = -\frac{3\vec{A}}{3t}$$
, we have
$$\vec{E}_{b} = \frac{3\vec{B}_{0}\vec{a}}{7} \sin \theta \int_{0}^{t} v k^{2} e^{-vt k^{2}} j_{i}(k) j_{i}(kr/a) dk.$$

In the limit of t >0+, the electric field becomes

$$E_{\phi} = \frac{3B_0 aV}{\pi} Sino \int_{0}^{4\pi o} k^2 j_1(k) j_1(k) k k dk = \frac{3B_0 aV}{\pi} Sino \frac{\pi}{2} f(1-\frac{r}{a}) = \frac{3B_0}{2} Sino \frac{a^2}{\mu \sigma a^2} f(r-a)$$

$$= \frac{3B_0}{2} Sino \delta(r-a)$$

and the current density is  $\vec{J} = \sigma \vec{E} = \frac{3B_0}{2\mu} \sin \delta(r \cdot a) \hat{q}$ . This agrees with prob. 5.35 (a)

For vt >>1, applying the same approximation as in Prob. 5.35 (c), we have

$$E_{\theta} = \frac{3B_{\theta}a}{\pi} \sin \frac{\nu}{(\nu t)^{3h}} \int_{0}^{t_{10}} e^{-u^{2}} u^{2} \int_{0}^{t_{1}} \left(\frac{u}{|w|}\right) \int_{0}^{t_{$$

(b) The power dissipated is given by

$$P = \int \vec{J} \cdot \vec{E} \, d^3x = \frac{90 \, B_0^2 \, a^2}{\pi^2} \int d^3x \, \sin^3\theta \int_0^{+\infty} \nu \, k_i^2 \, e^{-\nu t \, k_i^2} \, j_1(k_i) \, j_1(k_i) \, dk_i$$

$$= \frac{90 \, B_0^2 \, a^2}{\pi^2} \, \nu^2 \int_0^{+\infty} dk_i \int_0^{+\infty} dk_i \, k_i^2 \, e^{-\nu t \, k_i^2} \, j_1(k_i) \, j_1(k_i) \, dk_i$$

$$= \frac{90 \, B_0^2 \, a^2}{\pi^2} \, \nu^2 \int_0^{+\infty} dk_i \int_0^{+\infty} dk_i \, k_i^2 \, e^{-\nu t \, k_i^2} \, j_1(k_i) \, k_i^2 \, e^{-\nu t \, k_i^2} \, k_i^2 \, k_i$$

Obviously, P= - dWm/olt.

For VESSI, We have

Which decays asymptotically as (vt)-5h and is the same order as the induced current.