3.11 (a) Take the Bessel equation

$$\frac{1}{\rho} \frac{d}{d\rho} \left[\rho \frac{dJ_{\nu}(k\rho)}{d\rho} \right] + \left(k^2 - \frac{\nu^2}{\rho^2} \right) J_{\nu}(k\rho) = 0 ,$$

multiply both sides by $\beta J_{\nu}(k'\rho)$, and integrate from o to a, we will get

$$\int_{0}^{\alpha} J_{\nu}(k'p) \frac{d}{dp} \left[p \frac{dJ_{\nu}(kp)}{dp} \right] dp + \int_{0}^{\alpha} \left(k^{2} - \frac{\nu^{2}}{p^{2}} \right) p J_{\nu}(k'p) J_{\nu}(kp) dp = 0.$$

Perform integration by parts on the first term, the equation becomes

$$|| J_{\nu}(k')| = \int_{0}^{a} \left| \int_{0}^{a} \left|$$

boundary term becomes - A Julk'a) Tulka). Therefore, we will get

Exchange k and k',

and take the difference with the previous equation, we can see

This means, for k + k', we must have

Which proves the orthogonality of Bessel functions with different eigenvalues.

(b) The Bessel equation can also be written as

Where primes means derivative with respect to P. Muttiply by Jilke) and integrate from v to a,

we have

The firs term is simply
$$\frac{1}{2} \left(\rho J_{\nu}(kp) \right)^{2} \left|_{0}^{A} = \frac{1}{2} a^{2} \left(\frac{dJ_{\nu}(kp)}{d\rho} \right)^{2} \right|_{\rho=0} = \frac{1}{2} k^{2} a^{2} \left[J_{\nu}'(ka) \right]^{2}$$
,

The second term can be manipulated by integration by parts,

$$\int_{0}^{\Omega} (k^{2} \rho^{2} - v^{2}) \int_{0}^{1} (k\rho) \int_{0}^{1} (k\rho) d\rho = \frac{1}{2} \int_{0}^{\Omega} (k^{2} \rho^{2} - v^{2}) d(J_{0}(k\rho)^{2})$$

$$= \frac{1}{2} (k^{2} \alpha^{2} - v^{2}) \int_{0}^{1} (k\alpha)^{2} - k^{2} \int_{0}^{\Omega} \rho J_{0}(k\rho)^{2} d\rho$$

Therefore,
$$\int_{0}^{a} \rho \operatorname{J}_{V}(k\rho)^{\nu} d\rho = \frac{Q^{\nu}}{2} \left(\operatorname{J}_{V}'(k\alpha) \right)^{2} + \frac{Q^{\nu}}{2} \left(1 - \frac{\nu^{\nu}}{k\alpha^{\nu}} \right) \operatorname{J}_{V}(k\alpha)^{2}$$

$$= \frac{Q^{\nu}}{2} \left[\left(1 - \frac{\nu^{\nu}}{y \tilde{v}_{n}} \right) \operatorname{J}_{V}(y_{n})^{2} + \left(\frac{d \operatorname{J}_{V}(y_{n})}{d y_{n}} \right)^{\nu} \right]$$

Thin, for the Bessel-Fourier series

$$f(\rho) = \sum_{n=1}^{10} A_n J_{\nu}(\frac{y_{\nu n}}{a}\rho)$$

we can determine An by multiplying both sides by pJv (young) and integrating from i no a.

By the second boundary condition, we can relate the derivate of the Besser function to the function.

Specifically, we know
$$\frac{dJ_{\nu}(y_{n})}{dy_{\nu n}} = -\frac{\lambda}{y_{\nu n}} J_{\nu}(y_{\nu n})$$
, and $J_{\nu}(y_{\nu n}) = -\frac{y_{\nu n}}{\lambda} \frac{dJ_{\nu}(y_{\nu n})}{dy_{\nu n}}$. Then,
$$\left(1 - \frac{\nu^{\nu}}{y_{\nu n}^{\nu}}\right) J_{\nu}(y_{\nu n})^{\nu} + \left(\frac{dJ_{\nu}(y_{\nu n})}{dy_{\nu n}}\right)^{\nu} = \left(1 - \frac{\nu^{\nu}}{y_{\nu n}^{\nu}}\right) J_{\nu}(y_{\nu n})^{\nu} + \frac{\lambda^{\nu}}{y_{\nu n}^{\nu}} J_{\nu}(y_{\nu n})^{\nu} = \left(1 + \frac{\lambda^{\nu} - \nu^{\nu}}{y_{\nu n}^{\nu}}\right) J_{\nu}(y_{\nu n})^{\nu}$$

$$= \left(1 - \frac{\nu}{3\nu_n}\right) \frac{3\nu_n}{\lambda^{\nu}} \left(\frac{dJ_{\nu}(3\nu_n)}{d3\nu_n}\right)^2 + \left(\frac{dJ_{\nu}(3\nu_n)}{d3\nu_n}\right)^2$$

$$= \left(1 + \frac{y_{vn} - v^2}{\lambda^2}\right) \left(\frac{\partial J_v(y_m)}{\partial y_{vn}}\right)^2$$

From the identity $\frac{dJ_{V}(N)}{dN} = \frac{1}{2} \left(J_{V+1}(N) - J_{V-1}(N) \right)$, we have

$$\left(\frac{dJ_{\nu}(y_{\nu n})}{dy_{\nu n}}\right)^{2} = \frac{1}{4}\left(J_{\nu+1}(y_{\nu n}) - J_{\nu-1}(y_{\nu n})\right)^{2} = \frac{1}{4}\left(J_{\nu+1}(y_{\nu n}) + J_{\nu-1}(y_{\nu n})\right)^{2} - J_{\nu+1}(y_{\nu n})J_{\nu-1}(y_{\nu n})$$

Where we have use $J_{v+1}(k) + J_{v+1}(k) = \frac{2\nu}{\varkappa} J_{\nu}(\varkappa)$. Then,

$$\left(1-\frac{v^2}{y_{vn}^2}\right)J_v(y_{vn})^2+\left(\frac{dJ_v(y_{vn})}{dy_{vn}}\right)^2=J_v(y_{vn})^2-J_{v+1}(y_{vn})J_{v-1}(y_{vn}).$$

For $\lambda \to \infty$, the Seward boundary condition becomes $J_{\nu}(x) = 0$. Its normalization becomes $\left(\frac{dJ_{\nu}(J_{\nu n})}{dJ_{\nu n}}\right)^2 = \left(J_{\nu+1}(J_{\nu n})\right)^2 = \left(J_{\nu-1}(J_{\nu n})\right)^2$, as discussed in the text.

In the above proof, we have assumed $\rho J_{\nu}^{\dagger} l_{\nu} \rho_{\rho} = 0$, and $J_{\nu} l_{\nu} \rho_{\rho} \rho_{\rho} = 0$. This can be seen from the first boundary condition. Since the Bessel function is regular at $\rho = 0$, we must have $J_{\nu}(k\rho) = O(\rho^{\alpha})$, with d>0, and $\frac{dJ_{\nu}(k\rho)}{d\rho} = O(\rho^{\alpha-1})$. Then, first boundary condition requires that d>0. From this asymptotic behavior, we will obtain the necessary vanishing lower boundary values.