

12-19

(a) For β and γ as space indices, the tensor becomes $M^{\alpha ij} = \Theta^{\alpha i} x^j - \Theta^{\alpha j} x^i$.

The conservation condition $\partial_\alpha M^{\alpha ij} = 0$ is $\frac{1}{c} \frac{\partial}{\partial t} M^{\alpha ij} + \partial_k M^{kij} = 0$, where k is also a space index. Integrate this condition, we will have $\frac{1}{c} \frac{d}{dt} \int (\Theta^{\alpha i} x^j - \Theta^{\alpha j} x^i) d^3x + \int \partial_k M^{kij} d^3x = 0$.

The second integral over the entire space is zero, as it is a boundary term while the EM field is localized. Then, we have $\frac{1}{c} \frac{d}{dt} \int (\Theta^{\alpha i} x^j - \Theta^{\alpha j} x^i) d^3x = 0$. However, $\Theta^{\alpha i} = c g^i$ is just the momentum of the EM field. The condition becomes $\frac{d}{dt} \int \vec{x} \times \vec{p}_{em} d^3x = 0$, which is the statement of angular momentum conservation.

(b) For $\beta = 0$, $\partial_\alpha M^{\alpha \beta \gamma} = \partial_\alpha (\Theta^{\alpha 0} x^\gamma - \Theta^{\alpha \gamma} x^0) = \partial_\alpha \Theta^{\alpha 0} x^\gamma$, where we have used the conservation law to drop the second term. Then, take γ as a space index j .

$$\partial_\alpha \Theta^{\alpha 0} x^j = \frac{1}{c} \frac{\partial}{\partial t} \Theta^{\alpha 0} x^j + \partial_k \Theta^{k0} x^j = \frac{1}{c} \frac{\partial}{\partial t} \Theta^{\alpha 0} x^j + \partial_k (\Theta^{k0} x^j) - \Theta^{j0} = 0$$

Integrate over the entire space, the second term will drop out, we have

$$\frac{1}{c} \frac{d}{dt} \int \Theta^{\alpha 0} x^j d^3x = \int \Theta^{j0} d^3x$$

Since $\Theta^{00} = u$, $\Theta^{j0} = c p^j$, we can write the above equation as

$$\frac{1}{c} \frac{d}{dt} \int \vec{x} u d^3x = c \vec{p}_{em}$$

Define $\int \vec{x} u d^3x = \vec{x} \int u d^3x = \vec{x} E_{em}$, then we have

$$\frac{d\vec{x}}{dt} = \frac{c^2 \vec{p}_{em}}{E_{em}}$$