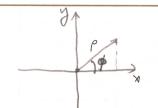
5.15 (a) For a sixyle wire, the magnetic field is given by

$$\vec{l} = \frac{\vec{l}}{\partial n} \hat{\phi}$$



Therefore, one scalar potential is

Then, in the two works setting, assuming the current at n= o/2 along the +2-direction

$$\oint_{M} : \frac{I}{J\pi} \left(\arctan \left(\frac{y}{y + d/2} \right) - \arctan \left(\frac{y}{y - d/2} \right) \right)$$

For dec x, y,

$$\hat{\Psi}_{m} = \frac{\text{Id}}{2\pi} \frac{\partial}{\partial x} \left(\arctan \left(\frac{y}{x} \right) \right) = -\frac{\text{Id}}{2\pi} \frac{y}{x^{2} + y^{2}} = -\frac{\text{Id} \sin \varphi}{2\pi \rho}$$

(b) Following the same procedure as 5.14, the potential on the three regions are

$$\hat{\Phi}_{m} = -\frac{\text{Jd sin}\phi}{2\pi\rho} + \sum_{m=1}^{100} \text{Am } (^{m} \text{sin}(m\phi))_{m} \quad \rho < \alpha$$

$$\underline{\hat{\psi}}_{m} = \sum_{m=1}^{\infty} \left(b_{m} \rho^{m} + C_{m} \rho^{-m} \right) \operatorname{Sim}(m\phi). \quad \alpha \in \rho \in \mathcal{B}$$

$$\frac{1}{2} \sum_{m=1}^{\infty} d_m e^{-m} \sin(m\phi)$$

The continuity conditions

$$\frac{\partial \hat{\Phi}_{m}}{\partial \rho}\Big|_{\rho=\alpha_{-}} = M_{n} \frac{\partial \hat{\Phi}_{m}}{\partial \rho}\Big|_{\rho=\alpha_{+}}, \qquad \frac{\partial \hat{\Phi}_{m}}{\partial \rho}\Big|_{\rho=L_{-}} = \frac{\partial \hat{\Phi}_{m}}{\partial \rho}\Big|_{\rho=L_{+}}$$

$$\frac{\partial \hat{Q}_{m}}{\partial \phi}\Big|_{\rho=0} = \frac{\partial \hat{Q}_{m}}{$$

 $\begin{cases}
-1 + a^2 a_1 = \mu r a^2 b_1 - \mu r C_1 \\
-1 + a^2 a_1 = a^2 b_1 + C_1 \\
\mu r b^2 b_1 - \mu r C_1 = -d_1
\end{cases}$

These equations can be solved, which leads to

$$b_1 = \frac{Mr-1}{2Mr} \frac{d_1}{b^2}$$

$$A_{r} = \frac{\mu_{r}^{2}-1}{4\mu_{r}}\left(\frac{1}{b^{2}}-\frac{1}{b^{2}}\right)d_{r}$$

Thus, the scalar potential outside of the cylinder is reduced in strength

$$F = \frac{4 \mu r b}{(\mu r + 1)^2 b^2 - (\mu r - 1)^2 (b - t)^2} \rightarrow \frac{4 \mu r b}{\mu r} = \frac{2b}{\mu r t}$$