[0.13 (a) We can use Eq. (10.92) to express the diffracted electric field as $\widetilde{E}(\vec{x}) = \frac{ie^{ikr}}{4\pi r} \vec{k} \times \vec{p}_{s}, \left[\frac{c\vec{k} \times (\vec{n} \times \vec{B}(\vec{x}))}{k} - \vec{n} \times \vec{E}(\vec{x}) \right] e^{-i\vec{k} \cdot \vec{x}} da'$ Since $\widetilde{E} = E_{o}\left(\vec{\xi}_{s}\cos a - \vec{\xi}_{s}\sin a\right) e^{ik\left[\vec{\xi}\cos a + x\sin a\right]}, \quad \widetilde{B} = \frac{E_{o}}{c} \cdot \vec{\xi}_{s} e^{ik\left[\vec{\xi}\cos a + x\sin a\right]}, \quad then <math display="block">(\vec{n} \times \vec{E}_{s})_{\vec{\xi}=o} = E_{o} \cdot \vec{\xi}_{s}\cos a e^{ikx'\sin a}, \quad (\vec{n} \times \vec{B}_{s})_{\vec{\xi}=o} = -\frac{E_{o}}{c} \cdot \vec{\xi}_{s} e^{ikx'\sin a}.$ The diffracted electric field para becomes $\widetilde{E}(\vec{x}) = \frac{ie^{ikr}}{cxr} \cdot E_{o} \cdot \vec{k} \times \vec{p}_{s}, \quad e^{-i\vec{k} \cdot \vec{x}} \cdot \left[\frac{\vec{k} \times \vec{\xi}_{s}}{k} + \vec{\xi}_{s}\cos a \right] da'$

$$\tilde{E}(\tilde{x}) = \frac{ie^{ikr}}{4\pi r} E_0 \tilde{k} \times \int_{S_1}^{S_2} e^{-i\tilde{k}\cdot\tilde{x}'} \left[\frac{\tilde{k}\times\tilde{\epsilon}_1}{k} + \tilde{\xi}_2 \cos\theta \right] da'$$

$$= \frac{ie^{ikr}}{4\pi r} E_0 \tilde{k} \times \int_{S_1}^{S_2} e^{-i\tilde{k}\cdot\tilde{x}'} \left[\tilde{\xi}_2 \left(\cos\theta + \cos\theta \right) - \tilde{\xi}_3 \sin\theta \sin\phi \right] da'$$

Where we have used $\bar{k} \times \bar{\epsilon}_1 = k \left(\bar{\epsilon}_2 \cos \theta - \bar{\epsilon}_3 \sin \theta \sin \phi \right)$. Compare with Eq. (10.117), we can see that cosd is replaced by $(\cos \theta + \cos \theta)/z$, and there is an additional factor proportioner to $\bar{k} \times \bar{\epsilon}_3$.