

12.1 (a) From the Lagrangian $\mathcal{L} = -\frac{m}{2} u_\alpha u^\alpha - \frac{q}{c} u_\alpha A^\alpha$, we have

$$\frac{\partial \mathcal{L}}{\partial u_\alpha} = -m u^\alpha - \frac{q}{c} A^\alpha, \quad \frac{\partial \mathcal{L}}{\partial x^\alpha} = -\frac{q}{c} u_\beta \partial^\alpha A^\beta$$

and $\frac{d}{d\tau} \left(\frac{\partial \mathcal{L}}{\partial u_\alpha} \right) = -m \frac{du^\alpha}{d\tau} - \frac{q}{c} \frac{\partial A^\alpha}{\partial x^\beta} \frac{dx^\beta}{d\tau} = -m u^\alpha - \frac{q}{c} u_\beta \partial^\beta A^\alpha$.

The Euler-Lagrange equation $\frac{d}{d\tau} \left(\frac{\partial \mathcal{L}}{\partial u_\alpha} \right) - \frac{\partial \mathcal{L}}{\partial x^\alpha} = 0$ becomes

$$-m \frac{du^\alpha}{d\tau} - \frac{q}{c} u_\beta \partial^\beta A^\alpha + \frac{q}{c} u_\beta \partial^\alpha A^\beta = 0,$$

or $m \frac{du^\alpha}{d\tau} = \frac{q}{c} F^{\alpha\beta} u_\beta$, where $F^{\alpha\beta} = \partial^\alpha A^\beta - \partial^\beta A^\alpha$

(b) The conjugate 4-vector momentum is $p^\alpha = -\frac{\partial \mathcal{L}}{\partial u_\alpha} = m u^\alpha + \frac{q}{c} A^\alpha$, and the Hamiltonian is

$$H = p^\alpha u_\alpha + \mathcal{L} = m u_\alpha u^\alpha + \frac{q}{c} u_\alpha A^\alpha - \frac{m}{2} u_\alpha u^\alpha - \frac{q}{c} u_\alpha A^\alpha = \frac{m}{2} u_\alpha u^\alpha$$

Since by definition $u_\alpha u^\alpha = 1$, the effective Hamiltonian is $H = \frac{m}{2}$, which is a Lorentz invariant.