

8.10 (a) Using the identity $\nabla \cdot (\vec{A} \times \vec{B}) = \vec{B} \cdot (\nabla \times \vec{A}) - \vec{A} \cdot (\nabla \times \vec{B})$, we have

$$\vec{E}^* \cdot [\nabla \times (\nabla \times \vec{E})] = -\nabla \cdot [\vec{E}^* \times (\nabla \times \vec{E})] + (\nabla \times \vec{E}^*) \cdot (\nabla \times \vec{E}),$$

then the numerator becomes

$$\int_V \vec{E}^* \cdot [\nabla \times (\nabla \times \vec{E})] d^3x = \int_V (\nabla \times \vec{E}^*) \cdot (\nabla \times \vec{E}) d^3x - \oint_S [\vec{E}^* \times (\nabla \times \vec{E})] \cdot \vec{n} da$$

Since $[\vec{E}^* \times (\nabla \times \vec{E})] \cdot \vec{n} = (\vec{n} \times \vec{E}^*) \cdot (\nabla \times \vec{E})$, where $\vec{n} \times \vec{E}^* = 0$ on the boundary, the surface integral term is identically 0. Therefore, the variational principle becomes

$$k^2 = \frac{\int_V (\nabla \times \vec{E}^*) \cdot (\nabla \times \vec{E}) d^3x}{\int_V \vec{E}^* \cdot \vec{E} d^3x}$$

(b) For the TE wave, $\vec{E}_t = \frac{-i}{\mu\epsilon\omega - k^2} \omega \hat{z} \times \nabla_t B_z$, where

$$\nabla_t B_z = \hat{\rho} \frac{\partial}{\partial \rho} B_z + \hat{\phi} \frac{1}{\rho} \frac{\partial}{\partial \phi} B_z$$

$$= B_0 \sin\left(\frac{\pi z}{a}\right) \left[\hat{\rho} \frac{1}{R} \left(1 - \frac{\rho}{R}\right) \cos\phi - \hat{\phi} \frac{1}{R} \left(1 - \frac{\rho}{2R}\right) \sin\phi \right]$$

$$\text{Then, } \vec{E}_t = -\frac{i\omega B_0}{(\mu\epsilon\omega - k^2)R} \left[\hat{\rho} \left(1 - \frac{\rho}{2R}\right) \sin\phi \sin\left(\frac{\pi z}{a}\right) + \hat{\phi} \left(1 - \frac{\rho}{R}\right) \cos\phi \sin\left(\frac{\pi z}{a}\right) \right]$$

where we have used $\hat{z} \times \hat{\rho} = \hat{\phi}$, and $\hat{z} \times \hat{\phi} = -\hat{\rho}$. Identifying E_0 as $-i\omega B_0 / (\mu\epsilon\omega - k^2)R$, we can write the transverse components of \vec{E} as

$$E_\rho = E_0 \left(1 - \frac{\rho}{2R}\right) \sin\phi \sin\left(\frac{\pi z}{a}\right), \quad E_\phi = E_0 \left(1 - \frac{\rho}{R}\right) \cos\phi \sin\left(\frac{\pi z}{a}\right).$$

(c) In cylindrical coordinates,

$$\begin{aligned} \nabla \times \vec{E} &= -\frac{\partial E_\phi}{\partial z} \hat{\rho} + \frac{\partial E_\rho}{\partial z} \hat{\phi} + \frac{1}{\rho} \left[\frac{\partial}{\partial \rho} (\rho E_\phi) - \frac{\partial}{\partial \phi} E_\rho \right] \hat{z} \\ &= E_0 \left[-\frac{\pi}{a} \left(1 - \frac{\rho}{R}\right) \cos\phi \cos\left(\frac{\pi z}{a}\right) \hat{\rho} + \frac{\pi}{a} \left(1 - \frac{\rho}{2R}\right) \sin\phi \cos\left(\frac{\pi z}{a}\right) \hat{\phi} \right. \\ &\quad \left. + \frac{1}{\rho} \left\{ \left(1 - \frac{2\rho}{R}\right) \cos\phi \sin\left(\frac{\pi z}{a}\right) - \left(1 - \frac{\rho}{2R}\right) \cos\phi \sin\left(\frac{\pi z}{a}\right) \right\} \hat{z} \right] \end{aligned}$$

$$\begin{aligned} \text{then } (\nabla \times \vec{E}^*) \cdot (\nabla \times \vec{E}) &= |E_0|^2 \left[\frac{\pi^2}{a^2} \left(1 - \frac{\rho}{R}\right)^2 \cos^2\phi \cos^2\left(\frac{\pi z}{a}\right) + \frac{\pi^2}{a^2} \left(1 - \frac{\rho}{2R}\right)^2 \sin^2\phi \cos^2\left(\frac{\pi z}{a}\right) \right. \\ &\quad \left. + \frac{9}{4R^2} \cos^2\phi \sin^2\left(\frac{\pi z}{a}\right) \right] \end{aligned}$$

$$\text{Also, } \vec{E}^* \cdot \vec{E} = |\vec{E}_0|^2 \left\{ \left(1 - \frac{\rho}{2R}\right)^2 \sin^2 \phi \sin^2\left(\frac{\pi \vec{z}}{d}\right) + \left(1 - \frac{\rho}{R}\right)^2 \cos^2 \phi \sin^2\left(\frac{\pi \vec{z}}{d}\right) \right\}$$

The integration in the numerator is

$$\begin{aligned} \frac{1}{|\vec{E}_0|^2} \int_V (\nabla \times \vec{E}^*) (\nabla \times \vec{E}) d^3x &= \frac{d}{2} \cdot \pi \cdot \int_0^R \left[\frac{\pi^2}{d^2} \left(1 - \frac{\rho}{R}\right)^2 + \frac{\pi^2}{d^2} \left(1 - \frac{\rho}{2R}\right)^2 + \frac{9}{4R^2} \right] \rho d\rho \\ &= \frac{d}{2} \cdot \pi \left[\frac{\pi^2}{d^2} \cdot \frac{R^3}{12} + \frac{\pi^2}{d^2} \cdot \frac{11}{48} R^3 + \frac{9}{8} \right] = \frac{\pi d}{2} \left[\frac{\pi^2}{d^2} \cdot \frac{5}{16} R^3 + \frac{9}{8} \right] \end{aligned}$$

and the denominator is

$$\frac{1}{|\vec{E}_0|^2} \int_V \vec{E}^* \cdot \vec{E} d^3x = \frac{d}{2} \cdot \pi \int_0^R \left[\left(1 - \frac{\rho}{2R}\right)^2 + \left(1 - \frac{\rho}{R}\right)^2 \right] \rho d\rho = \frac{d}{2} \cdot \pi \left(\frac{1}{12} R^3 + \frac{11}{48} R^3 \right) = \frac{\pi d}{2} \cdot \frac{5}{16} R^3$$

where we have replaced the $\sin^2(\dots)$ and $\cos^2(\dots)$ integration with the average. Then,

$$R^3 = \frac{\frac{\pi^2}{d^2} \cdot \frac{5}{16} R^3 + \frac{9}{8}}{\frac{5}{16} R^3} = \frac{18}{5R^2} + \frac{\pi^2}{d^2}$$

(d) The form $\vec{E}^* \cdot [\nabla \times (\nabla \times \vec{E})] = \vec{E}^* \cdot [\nabla(\nabla \cdot \vec{E}) - \nabla^2 \vec{E}]$ is convenient in the TM case, as

we only need to calculate the Laplacian of one component. When \vec{E} has transverse components, $\nabla \times \vec{E}$ is more convenient.