

6.21 (a) From Problem 4.2, we know that the dipole's charges can be expressed formally as

$\rho(\vec{x}, t) = -(\vec{p} \cdot \nabla) \delta(\vec{x} - \vec{r}_0(t))$. Using the conservation of charge, $\frac{\partial \rho}{\partial t} + \nabla \cdot \vec{j} = 0$, we have

$$\frac{\partial \rho}{\partial t} = -(\vec{p} \cdot \nabla) [-\vec{v} \cdot \nabla \delta(\vec{x} - \vec{r}_0(t))] = (\vec{p} \cdot \nabla) [\nabla \cdot (\vec{v} \delta(\vec{x} - \vec{r}_0(t)))] = \nabla \cdot \{ \vec{v} (\vec{p} \cdot \nabla) \delta(\vec{x} - \vec{r}_0(t)) \}.$$

Therefore, $\vec{j}(\vec{x}, t) = -\vec{v} (\vec{p} \cdot \nabla) \delta(\vec{x} - \vec{r}_0(t))$.

(b) Since the dipole is moving non-relativistically, we can ignore the retardation effect and treat the potential as instantaneous. Therefore,

$$\begin{aligned} \Phi(\vec{x}, t) &= \frac{1}{4\pi\epsilon_0} \int d^3x' \frac{\rho(\vec{x}', t)}{|\vec{x} - \vec{x}'|} = -\frac{1}{4\pi\epsilon_0} \int d^3x' (\vec{p} \cdot \nabla') \frac{\delta(\vec{x} - \vec{r}_0)}{|\vec{x} - \vec{x}'|} \\ &= \frac{1}{4\pi\epsilon_0} \int d^3x' \delta(\vec{x}' - \vec{r}_0) \frac{\vec{p} \cdot (\vec{x} - \vec{x}')}{|\vec{x} - \vec{x}'|^3} = \frac{1}{4\pi\epsilon_0} \frac{\vec{p} \cdot (\vec{x} - \vec{r}_0)}{|\vec{x} - \vec{r}_0|^3}. \end{aligned}$$

Now, expanding the scalar potential in \vec{r}_0 , $|\vec{x} - \vec{r}_0|^{-3} = |\vec{x}|^{-3} + 3\vec{x} \cdot \vec{r}_0 / |\vec{x}|^5$, the potential becomes

$$\Phi(\vec{x}, t) = \frac{1}{4\pi\epsilon_0} (\vec{p} \cdot \vec{x} - \vec{p} \cdot \vec{r}_0) \left(\frac{1}{r^3} + \frac{3\vec{x} \cdot \vec{r}_0}{r^5} \right) = \frac{1}{4\pi\epsilon_0} \left(\frac{\vec{p} \cdot \vec{x}}{r^3} + \frac{3(\vec{p} \cdot \vec{r}_0)(\vec{r}_0 \cdot \vec{x}) - \vec{p} \cdot \vec{r}_0}{r^3} - \frac{3(\vec{p} \cdot \vec{r}_0)(\vec{r}_0 \cdot \vec{r}_0)}{r^4} \right)$$

Here, it is clear that, in addition to the dipole potential, $\frac{1}{4\pi\epsilon_0} \frac{\vec{p} \cdot \vec{x}}{r^3}$, we also have an contribution from quadrupole moment,

$$\frac{1}{4\pi\epsilon_0} \frac{1}{2} \sum_{ij} Q_{ij} \frac{n_i n_j}{r^3}, \quad \text{with } Q_{ij} = 3(x_{0i} r_{0j} + x_{0j} r_{0i}) - 2\vec{r}_0 \cdot \vec{r}_0 \delta_{ij}.$$

Similarly, for the vector potential,

$$\begin{aligned} \vec{A}(\vec{x}, t) &= \frac{\mu_0}{4\pi} \int d^3x' \frac{\vec{j}(\vec{x}', t)}{|\vec{x} - \vec{x}'|} = -\frac{\mu_0 \vec{v}}{4\pi} \int d^3x' (\vec{p} \cdot \nabla') \frac{\delta(\vec{x}' - \vec{r}_0)}{|\vec{x} - \vec{x}'|} = \frac{\mu_0}{4\pi} \vec{v} \int d^3x' \delta(\vec{x}' - \vec{r}_0) \frac{\vec{p} \cdot (\vec{x} - \vec{x}')}{|\vec{x} - \vec{x}'|^3} \\ &= \frac{\mu_0}{4\pi} \frac{\vec{v} [\vec{p} \cdot (\vec{x} - \vec{r}_0)]}{|\vec{x} - \vec{r}_0|^3}. \end{aligned}$$

Ignoring all the \vec{r}_0 terms, we have

$$\vec{A}(\vec{x}, t) = \frac{\mu_0}{4\pi} \frac{\vec{v} (\vec{p} \cdot \vec{x})}{r^3} = \frac{\mu_0}{4\pi} \left[\frac{1}{2} \frac{(\vec{p} \times \vec{v}) \times \vec{x}}{r^3} + \frac{1}{2} \frac{\vec{p}(\vec{x} \cdot \vec{v}) + \vec{v}(\vec{x} \cdot \vec{p})}{r^3} \right]$$

and therefore contains a magnetic dipole field

$$\vec{A} = \frac{\mu_0}{4\pi} \frac{\vec{m} \times \vec{x}}{r^3}, \quad \text{with } \vec{m} = \frac{1}{2} \vec{p} \times \vec{v}.$$

(c) To calculate the quadrupole field, use the definition,

$$\begin{aligned}\vec{E} &= -\nabla\Phi = -\hat{e}_k \partial_k \left(\frac{1}{8\pi\epsilon_0} \sum_{ij} Q_{ij} \frac{x_i x_j}{r^3} \right) = -\frac{1}{8\pi\epsilon_0} \hat{e}_k \sum_{ij} \partial_k \left(Q_{ij} \frac{x_i x_j}{r^5} \right) \\ &= -\frac{1}{8\pi\epsilon_0} \hat{e}_k \sum_{ij} Q_{ij} \partial_k \left(\frac{x_i x_j}{r^5} \right).\end{aligned}$$

Since $\partial_k \left(\frac{x_i x_j}{r^5} \right) = \frac{\delta_{ik} x_j + \delta_{jk} x_i}{r^5} - \frac{5 x_i x_j x_k}{r^7}$, we have

$$\begin{aligned}\vec{E} &= -\frac{1}{8\pi\epsilon_0} \hat{e}_k \sum_{ij} \left(\frac{\delta_{ik} x_j + \delta_{jk} x_i}{r^5} - \frac{5 x_i x_j x_k}{r^7} \right) \left[3(x_{0i} p_{0j} + x_{0j} p_{0i}) - 2 \vec{p} \cdot \vec{r}_0 \delta_{ij} \right] \\ &= -\frac{1}{8\pi\epsilon_0} \hat{e}_k \left[\frac{3(x_{0k} p_{0j} x_j + x_{0i} p_{0k} x_i)}{r^5} - \frac{2(\vec{p} \cdot \vec{r}_0) x_k}{r^5} + \frac{3(x_{0i} p_{0k} x_i + x_{0k} p_{0i} x_k)}{r^5} - \frac{2(\vec{p} \cdot \vec{r}_0) x_k}{r^5} \right. \\ &\quad \left. - \frac{15(x_{0i} p_{0j} + x_{0j} p_{0i}) x_i x_j x_k}{r^7} + \frac{10(\vec{p} \cdot \vec{r}_0) x_i x_j x_k}{r^7} \right] \\ &= -\frac{1}{8\pi\epsilon_0} \left(\frac{3\vec{r}_0(\vec{p} \cdot \vec{r}) + 3\vec{p}(\vec{r} \cdot \vec{r}_0)}{r^5} - \frac{2\vec{r}(\vec{p} \cdot \vec{r}_0)}{r^5} + \frac{3\vec{p}(\vec{r} \cdot \vec{r}_0) + 3\vec{r}(\vec{p} \cdot \vec{r})}{r^5} - \frac{2\vec{r}(\vec{p} \cdot \vec{r}_0)}{r^5} \right. \\ &\quad \left. - \frac{15\vec{r}(\vec{r} \cdot \vec{r}_0)(\vec{p} \cdot \vec{r}) + 15\vec{p}(\vec{r} \cdot \vec{r}_0)(\vec{p} \cdot \vec{r})}{r^7} + \frac{10\vec{r}(\vec{p} \cdot \vec{r}_0)|\vec{r}|^2}{r^7} \right) \\ &= \frac{1}{8\pi\epsilon_0} \left[\frac{30\vec{n}(\vec{n} \cdot \vec{r}_0)(\vec{n} \cdot \vec{p})}{r^4} - \frac{6\vec{r}_0(\vec{n} \cdot \vec{p}) + 6\vec{p}(\vec{n} \cdot \vec{r}_0) + 6\vec{n}(\vec{p} \cdot \vec{r}_0)}{r^4} \right] \\ &= \frac{1}{4\pi\epsilon_0 r^4} \left[15\vec{n}(\vec{n} \cdot \vec{r}_0)(\vec{n} \cdot \vec{p}) - 3\vec{r}_0(\vec{n} \cdot \vec{p}) - 3\vec{p}(\vec{n} \cdot \vec{r}_0) - 3\vec{n}(\vec{p} \cdot \vec{r}_0) \right]\end{aligned}$$