

9.5 (a) From Eq. (9.16), $\vec{A}(\vec{x}) = -i \frac{\mu_0 \omega}{4\pi} \frac{e^{ikr}}{r} \vec{p}$. For the scalar potential,

$$\Phi(\vec{x}) = \frac{1}{4\pi\epsilon_0} \int \rho(\vec{x}') \frac{e^{ik|\vec{x}-\vec{x}'|}}{|\vec{x}-\vec{x}'|} d^3x'$$

Since $|\vec{x}-\vec{x}'| = r - \vec{n} \cdot \vec{x}'$, and $|\vec{x}-\vec{x}'|^{-1} = \frac{1}{r} + \frac{\vec{n} \cdot \vec{x}'}{r^2}$, we have

$$\begin{aligned} \Phi(\vec{x}) &= \frac{e^{ikr}}{4\pi\epsilon_0} \int \rho(\vec{x}') (1 - ik \vec{n} \cdot \vec{x}') \left(\frac{1}{r} + \frac{\vec{n} \cdot \vec{x}'}{r^2} \right) d^3x' \\ &= \frac{e^{ikr}}{4\pi\epsilon_0} \int \rho(\vec{x}') \left(\frac{1}{r} - ik \frac{\vec{n} \cdot \vec{x}'}{r} + \frac{\vec{n} \cdot \vec{x}'}{r^2} + \dots \right) d^3x' \\ &= \frac{e^{ikr}}{4\pi\epsilon_0} \left[\frac{q}{r} + \frac{\vec{n} \cdot \vec{p}}{r^2} (1 - ikr) + \dots \right] \end{aligned}$$

where the leading term is the electric monopole contribution and can be dropped for radiation purposes.

Therefore,
$$\Phi(\vec{x}) = \frac{e^{ikr}}{4\pi\epsilon_0 r^2} (\vec{n} \cdot \vec{p}) (1 - ikr)$$

(b) Using the relation $\nabla \times (\psi \vec{a}) = \nabla \psi \times \vec{a} + \psi \nabla \times \vec{a}$, the magnetic field is

$$\begin{aligned} \vec{H} &= \frac{1}{\mu_0} \nabla \times \vec{A} = -i \frac{\omega}{4\pi} \nabla \left(\frac{e^{ikr}}{r} \right) \times \vec{p} = -i \frac{\omega}{4\pi} e^{ikr} \left(\frac{ik\vec{n}}{r} - \frac{\vec{n}}{r^2} \right) \times \vec{p} \\ &= \frac{k\omega}{4\pi} (\vec{n} \times \vec{p}) \frac{e^{ikr}}{r} \left(1 - \frac{1}{ikr} \right) = \frac{ck^2}{4\pi} (\vec{n} \times \vec{p}) \frac{e^{ikr}}{r} \left(1 - \frac{1}{ikr} \right) \end{aligned}$$

For the electric field, $\vec{E} = -\nabla\Phi - \frac{\partial\vec{A}}{\partial t}$. The vector potential has $e^{-i\omega t}$ time dependence, and

$$\frac{\partial\vec{A}}{\partial t} = -\frac{\mu_0 \omega^2}{4\pi} \frac{e^{ikr}}{r} \vec{p} = -\frac{k^2}{4\pi\epsilon_0} \frac{e^{ikr}}{r} \vec{p}. \text{ For the scalar potential part,}$$

$$\begin{aligned} \nabla\Phi &= \frac{-ik e^{ikr}}{4\pi\epsilon_0 r^2} \vec{n}(\vec{n} \cdot \vec{p}) + \frac{ik e^{ikr}}{4\pi\epsilon_0 r^2} \vec{n}(\vec{n} \cdot \vec{p})(1 - ikr) - \frac{e^{ikr}}{2\pi\epsilon_0 r^3} \vec{n}(\vec{n} \cdot \vec{p})(1 - ikr) \\ &\quad + \frac{e^{ikr}}{4\pi\epsilon_0 r^2} (1 - ikr) \nabla(\vec{n} \cdot \vec{p}) \quad \left(\nabla(\vec{n} \cdot \vec{p}) = (\vec{p} \cdot \nabla) \vec{n} = \frac{1}{r} [\vec{p} - \vec{n}(\vec{n} \cdot \vec{p})] \right) \\ &= \frac{k^2}{4\pi\epsilon_0} \frac{e^{ikr}}{r} \vec{n}(\vec{n} \cdot \vec{p}) - \frac{e^{ikr}}{2\pi\epsilon_0 r^3} \vec{n}(\vec{n} \cdot \vec{p})(1 - ikr) + \frac{e^{ikr}}{4\pi\epsilon_0 r^3} (1 - ikr) [\vec{p} - \vec{n}(\vec{n} \cdot \vec{p})] \\ &= \frac{k^2}{4\pi\epsilon_0} \frac{e^{ikr}}{r} \vec{n}(\vec{n} \cdot \vec{p}) + \frac{e^{ikr}}{4\pi\epsilon_0} \left(\frac{1}{r^3} - \frac{ik}{r^2} \right) [\vec{p} - 3\vec{n}(\vec{n} \cdot \vec{p})] \end{aligned}$$

Then,
$$\begin{aligned} \vec{E} &= -\nabla\Phi - \frac{\partial\vec{A}}{\partial t} = \frac{k^2}{4\pi\epsilon_0} \frac{e^{ikr}}{r} [\vec{p} - \vec{n}(\vec{n} \cdot \vec{p})] + \frac{e^{ikr}}{4\pi\epsilon_0} \left(\frac{1}{r^3} - \frac{ik}{r^2} \right) [3\vec{n}(\vec{n} \cdot \vec{p}) - \vec{p}] \\ &= \frac{e^{ikr}}{4\pi\epsilon_0 r} \left[k^2 (\vec{n} \times \vec{p}) \times \vec{n} + \left(3\vec{n}(\vec{n} \cdot \vec{p}) - \vec{p} \right) \left(\frac{1}{r^2} - \frac{ik}{r} \right) \right] \end{aligned}$$