

12.3 (a) From the equation of motion, $\frac{d\vec{p}}{dt} = e\vec{u} \cdot \vec{E}$, we have

$$\frac{dp_{||}}{dt} = \frac{d}{dt} \left(\frac{m u_{||}}{\sqrt{1 - \frac{u^2}{c^2}}} \right) = e u_{||} E_0, \quad \frac{dp_{\perp}}{dt} = \frac{d}{dt} \left(\frac{m u_{\perp}}{\sqrt{1 - \frac{u^2}{c^2}}} \right) = 0,$$

Where $u^2 = u_{||}^2 + u_{\perp}^2$. From the normal component, we know that

$$\frac{m u_{\perp}}{\sqrt{1 - \frac{u^2}{c^2}}} = C_0 = \frac{m v_0}{\sqrt{1 - \frac{v_0^2}{c^2}}}. \quad \text{The longitudinal component can be directly integrated,}$$

$$\frac{m u_{||}}{\sqrt{1 - \frac{u^2}{c^2}}} = e E_0 t. \quad \text{Then, } m^2 u^2 / (1 - u^2/c^2) = C_0^2 + e^2 E_0^2 t^2, \text{ or } \frac{u^2}{c^2} = \frac{C_0^2 + e^2 E_0^2 t^2}{m^2 c^2 + C_0^2 + e^2 E_0^2 t^2}$$

$$\text{Therefore, } u_{\perp} = \frac{C_0 c}{m \sqrt{1 - \frac{u^2}{c^2}}} = \frac{C_0 c}{[m^2 c^2 + C_0^2 + e^2 E_0^2 t^2]^{1/2}}, \quad u_{||} = \frac{e E_0 c t}{[m^2 c^2 + C_0^2 + e^2 E_0^2 t^2]^{1/2}}.$$

The position of the particle, assume it starts from the origin, is

$$x_{\perp}(t) = \int_0^t u_{\perp}(s) ds = \frac{C_0 c}{e E_0} \operatorname{arcsinh} \left(\frac{e E_0 t}{\sqrt{m^2 c^2 + C_0^2}} \right), \quad x_{||}(t) = \int_0^t u_{||}(s) ds = \frac{c}{e E_0} \left(\sqrt{m^2 c^2 + C_0^2 + e^2 E_0^2 t^2} - \sqrt{m^2 c^2 + C_0^2} \right).$$

Notice that $m^2 c^2 + C_0^2 = m^2 c^2 \left(1 - \frac{v_0^2}{c^2} \right)^{-1} = \gamma_0^2 m^2 c^2$ where $\gamma_0 = \left(1 - \frac{v_0^2}{c^2} \right)^{-1/2}$, the position can be

$$\text{written as } x_{\perp}(t) = \frac{\gamma_0 m v_0 c}{e E_0} \operatorname{arcsinh} \left(\frac{e E_0}{\gamma_0 m c^2} t \right), \quad x_{||}(t) = \frac{\gamma_0 m c^2}{e E_0} \left(\sqrt{1 + \frac{e^2 E_0^2}{\gamma_0^2 m^2 c^2} t^2} - 1 \right).$$

(b) From the position of the particle, we know

$$t = \frac{\gamma_0 m c}{e E_0} \sinh \left(\frac{e E_0 x_{\perp}}{\gamma_0 m v_0 c} \right) = \frac{\gamma_0 m c}{e E_0} \left[\left(1 + \frac{e E_0 x_{||}}{\gamma_0 m c^2} \right)^2 - 1 \right]^{1/2},$$

$$\text{Which leads to } \left(1 + \frac{e E_0 x_{||}}{\gamma_0 m c^2} \right)^2 = 1 + \sinh^2 \left(\frac{e E_0 x_{\perp}}{\gamma_0 m v_0 c} \right) = \cosh^2 \left(\frac{e E_0 x_{\perp}}{\gamma_0 m v_0 c} \right), \text{ or}$$

$$x_{||} = \frac{\gamma_0 m c^2}{e E_0} \left(\cosh \left(\frac{e E_0 x_{\perp}}{\gamma_0 m v_0 c} \right) - 1 \right).$$

For small time, $x_{\perp} \ll 1$, expand cosh to second order, we have $x_{||} = \frac{e E_0}{2 \gamma_0 m v_0^2} x_{\perp}^2$, which is a parabola.

For large time, $x_{\perp} \gg 1$, $\cosh \left(\frac{e E_0 x_{\perp}}{\gamma_0 m v_0 c} \right) - 1 \approx \exp \left\{ \frac{e E_0 x_{\perp}}{\gamma_0 m v_0 c} \right\}$, and $x_{||} = \frac{\gamma_0 m c^2}{e E_0} \exp \left\{ \frac{e E_0 x_{\perp}}{\gamma_0 m v_0 c} \right\}$.

The short and long time regime can be identified by $\frac{e E_0 x_{\perp}}{\gamma_0 m v_0 c} \ll 1$.