

7.26. A strange convention is chosen in performing the Fourier transform, but the final result after integration in momentum space does not depend on the convention chosen.

(a) The charge density in the real space is $\rho(\vec{x}, t) = Ze \delta(\vec{x} - \vec{v}t)$, assuming the charged particle starts at $\vec{x} = 0$. Then,

$$\begin{aligned} \tilde{\rho}(\vec{q}, \omega) &= \int \frac{d^3x dt}{(2\pi)^4} \rho(\vec{x}, t) e^{-i(\vec{q} \cdot \vec{x} - \omega t)} = Ze \int \frac{dt}{(2\pi)^4} e^{-i(\vec{q} \cdot \vec{v} - \omega)t} \\ &= \frac{Ze}{(2\pi)^3} \delta(\omega - \vec{q} \cdot \vec{v}) \end{aligned}$$

(b) The Coulomb's law states that $\nabla \cdot \vec{D}(\vec{x}, t) = \rho(\vec{x}, t)$. Also,

$$\vec{D}(\vec{x}, t) = \int \frac{d^3x' dt'}{(2\pi)^4} \epsilon(\vec{x} - \vec{x}', t - t') \vec{E}(\vec{x}', t') = - \int \frac{d^3x' dt'}{(2\pi)^4} \epsilon(\vec{x} - \vec{x}', t - t') \nabla' \phi(\vec{x}', t')$$

Introduce the Fourier transform,

$$\begin{aligned} \vec{D}(\vec{x}, t) &= - \int \frac{d^3x' dt'}{(2\pi)^4} \int d^3q_1 d\omega_1 \epsilon(\vec{q}_1, \omega_1) e^{i\vec{q}_1 \cdot (\vec{x} - \vec{x}') - i\omega_1(t - t')} \\ &\quad \times \left(\nabla' \int d^3q_2 d\omega_2 \phi(\vec{q}_2, \omega_2) e^{i\vec{q}_2 \cdot \vec{x}' - i\omega_2 t'} \right) \\ &= - \int d^3q_1 d\omega_1 \int d^3q_2 d\omega_2 (i\vec{q}_2) \epsilon(\vec{q}_1, \omega_1) \phi(\vec{q}_2, \omega_2) e^{i\vec{q}_1 \cdot \vec{x} - i\omega_1 t} \\ &\quad \times \int \frac{d^3x' dt'}{(2\pi)^4} e^{i(\vec{q}_2 - \vec{q}_1) \cdot \vec{x}' - i(\omega_2 - \omega_1)t'} \\ &= - \int d^3q_1 d\omega_1 \int d^3q_2 d\omega_2 (i\vec{q}_2) \epsilon(\vec{q}_1, \omega_1) \phi(\vec{q}_2, \omega_2) \delta(\vec{q}_1 - \vec{q}_2) \delta(\omega_1 - \omega_2) e^{i\vec{q}_1 \cdot \vec{x} - i\omega_1 t} \\ &= - \int d^3q d\omega e^{i\vec{q} \cdot \vec{x} - i\omega t} i\vec{q} \epsilon(\vec{q}, \omega) \phi(\vec{q}, \omega) \end{aligned}$$

$$\text{then, } \nabla \cdot \vec{D}(\vec{x}, t) = \int d^3q d\omega e^{i\vec{q} \cdot \vec{x} - i\omega t} \vec{q} \cdot \epsilon(\vec{q}, \omega) \phi(\vec{q}, \omega).$$

$$\text{Also, } \rho(\vec{x}, t) = \int d^3q d\omega e^{i\vec{q} \cdot \vec{x} - i\omega t} \rho(\vec{q}, \omega). \text{ Therefore, for each component,}$$

$$q^i \epsilon(\vec{q}, \omega) \phi(\vec{q}, \omega) = \rho(\vec{q}, \omega), \quad \text{or} \quad \phi(\vec{q}, \omega) = \frac{\rho(\vec{q}, \omega)}{q^i \epsilon(\vec{q}, \omega)}$$

$$\begin{aligned}
(c) \quad -\frac{dW}{dt} &= -\int \ddot{\mathbf{J}} \cdot \vec{\mathbf{E}} \, d^3x = \int \ddot{\mathbf{J}} \cdot \nabla \phi \, d^3x = \int [\nabla(\phi \ddot{\mathbf{J}}) - \phi(\nabla \cdot \ddot{\mathbf{J}})] \, d^3x \\
&= \int \phi \frac{\partial}{\partial t} d^3x = \int d^3x \int d^3q \, d\omega_1 \, \phi(\vec{q}, \omega_1) e^{i(\vec{q}_1 \cdot \vec{x} - \omega_1 t)} \\
&\quad \times \frac{\partial}{\partial t} \int d^3q_2 \, d\omega_2 \, \rho(\vec{q}_2, \omega_2) e^{i(\vec{q}_2 \cdot \vec{x} - \omega_2 t)} \\
&= \int d^3q_1 \, d\omega_1 \int d^3q_2 \, d\omega_2 (-i\omega_2) \phi(\vec{q}_1, \omega_1) \rho(\vec{q}_2, \omega_2) \int d^3x \, e^{i(\vec{q}_1 + \vec{q}_2) \cdot \vec{x} - i(\omega_1 + \omega_2)t} \\
&= \int d^3q \, d\omega_1 \, d\omega_2 (-i\omega_2) \phi(\vec{q}, \omega_1) \rho(-\vec{q}, \omega_2) \cdot (2\pi)^3 e^{-i(\omega_1 + \omega_2)t} \\
&= (2\pi)^3 \int d^3q \, d\omega_1 \, d\omega_2 (-i\omega_2) \frac{\rho(\vec{q}, \omega_1)}{q^2 \epsilon(\vec{q}, \omega_1)} \rho(-\vec{q}, \omega_2) e^{-i(\omega_1 + \omega_2)t} \\
&= \frac{Z^2 e^2}{(2\pi)^3} \int d^3q \, d\omega_1 \, d\omega_2 (-i\omega_2) \frac{1}{q^2 \epsilon(\vec{q}, \omega_1)} \delta(\vec{q} \cdot \vec{v} - \omega_1) \delta(-\vec{q} \cdot \vec{v} - \omega_2) e^{-i(\omega_1 + \omega_2)t} \\
&= \frac{Z^2 e^2}{(2\pi)^3} \int d^3q \, d\omega \, (i\omega) \frac{1}{q^2 \epsilon(\vec{q}, \omega)} \delta(\omega - \vec{q} \cdot \vec{v})
\end{aligned}$$

Only the real part will contribute to the energy loss. Therefore,

$$-\frac{dW}{dt} = \frac{Z^2 e^2}{(2\pi)^3} \int d^3q \, d\omega \, \frac{\omega}{q^2} \text{Im} \left[\frac{1}{\epsilon(\vec{q}, \omega)} \right] \delta(\omega - \vec{q} \cdot \vec{v})$$

Since $\text{Im}[\epsilon(\vec{q}, \omega)]$ is odd in ω , we can write the final result as

$$-\frac{dW}{dt} = \frac{Z^2 e^2}{4\pi^2} \int \frac{d^3q}{q^2} \int_0^{+\infty} d\omega \, \omega \, \text{Im} \left[\frac{1}{\epsilon(\vec{q}, \omega)} \right] \delta(\omega - \vec{q} \cdot \vec{v})$$