

2.17 Solution: (a) The Green function in three-dimensional space is

$$G(\mathbf{r}, \mathbf{r}') = \frac{1}{|\mathbf{r} - \mathbf{r}'|},$$

or equivalently, in Cartesian coordinates,

$$G(x, y, z; x', y', z') = \frac{1}{\sqrt{(x - x')^2 + (y - y')^2 + (z - z')^2}}.$$

Using the identity

$$\int \frac{dx}{\sqrt{x^2 + a^2}} = \log \left(x + \sqrt{x^2 + a^2} \right) + C,$$

integrating the three-dimensional Green function in $z' - z$ from $-Z$ to Z , we have

$$\begin{aligned} & \int_{-Z}^Z \frac{d(z' - z)}{\sqrt{(x - x')^2 + (y - y')^2 + (z - z')^2}} \\ &= \log \left(\frac{\sqrt{(x - x')^2 + (y - y')^2 + Z^2} + Z}{\sqrt{(x - x')^2 + (y - y')^2 + Z^2} - Z} \right) \\ &= 2 \log \left(\sqrt{(x - x')^2 + (y - y')^2 + Z^2} + Z \right) - \log \left((x - x')^2 + (y - y')^2 \right). \end{aligned} \quad (1)$$

When Z is large, the leading order correction is proportional to

$$\frac{(x - x')^2 + (y - y')^2}{Z^2}$$

and can be ignored compared to $\log \left((x - x')^2 + (y - y')^2 \right)$. Therefore, the two-dimensional free-space Green function becomes

$$\begin{aligned} G(x, y; x', y') &= -\log \left((x - x')^2 + (y - y')^2 \right) \\ &= -\log \left(\rho^2 + \rho'^2 - 2\rho\rho' \cos(\phi - \phi') \right). \end{aligned}$$

(b) For two-dimensional problem, the completeness relation in the angular direction is

$$\delta(\phi - \phi') = \frac{1}{2\pi} \sum_{m=-\infty}^{\infty} e^{im(\phi - \phi')},$$

and the two-dimensional Dirac delta function can therefore be written as

$$\delta(\boldsymbol{\rho} - \boldsymbol{\rho}') = \frac{1}{\rho} \delta(\rho - \rho') \cdot \frac{1}{2\pi} \sum_{m=-\infty}^{\infty} e^{im(\phi - \phi')}.$$

Thus, the Green function must have the following form,

$$G(\boldsymbol{\rho}, \boldsymbol{\rho}') = \frac{1}{2\pi} \sum_{m=-\infty}^{\infty} g_m(\rho, \rho') e^{im(\phi - \phi')}.$$

Applying the Laplacian operator on the Green function leads to the Dirac delta function,

$$\nabla_{\boldsymbol{\rho}'}^2 G(\boldsymbol{\rho}, \boldsymbol{\rho}') = -4\pi \delta(\boldsymbol{\rho} - \boldsymbol{\rho}'),$$

which is equivalent to

$$\frac{1}{2\pi} \sum_{m=-\infty}^{\infty} \left[\frac{1}{\rho'} \frac{\partial}{\partial \rho'} \left(\rho' \frac{\partial}{\partial \rho'} \right) - \frac{m^2}{\rho'^2} \right] g_m(\rho, \rho') e^{im(\phi-\phi')} = -4\pi \frac{\delta(\rho-\rho')}{\rho} \cdot \frac{1}{2\pi} \sum_{m=-\infty}^{\infty} e^{im(\phi-\phi')}.$$

The above equation must be valid term-wise, which gives the following differential equation for g_m ,

$$\frac{1}{\rho'} \frac{\partial}{\partial \rho'} \left(\rho' \frac{\partial g_m}{\partial \rho'} \right) - \frac{m^2}{\rho'^2} g_m = -4\pi \frac{\delta(\rho-\rho')}{\rho}. \quad (2)$$

(c) Solution for Eq. (2) must be linear combinations of ρ'^m and ρ'^{-m} , for $m \neq 0$. For $\rho' < \rho$, the regularity of the solution at $\rho' = 0$ dictates that

$$g_m \propto \begin{cases} \rho'^m, & m > 0, \\ \rho'^{-m}, & m < 0. \end{cases}$$

Similarly, for $\rho' > \rho$, the regularity at $\rho' \rightarrow \infty$ requires that

$$g_m \propto \begin{cases} \rho'^{-m}, & m > 0, \\ \rho'^m, & m < 0. \end{cases}$$

Since the solution must be symmetric in ρ and ρ' , the solution we are seeking should be in the form of

$$g_m(\rho, \rho') = C_m \left(\frac{\rho_{<}}{\rho_{>}} \right)^{|m|}, \quad (3)$$

where C_m is a yet to be determined constant, and $\rho_{<} (\rho_{>})$ is the smaller (larger) of ρ and ρ' .

With Eq. (3), we can now multiply both sides of Eq. (2) with ρ' and integrate on a small interval around ρ . This leads to an equation that connects the derivatives of g_m on both sides of ρ ,

$$\rho' \frac{\partial}{\partial \rho'} g_m(\rho, \rho') \Big|_{\rho+\varepsilon} - \rho' \frac{\partial}{\partial \rho'} g_m(\rho, \rho') \Big|_{\rho-\varepsilon} = -4\pi,$$

which gives

$$C_m = \frac{2\pi}{|m|}.$$

For $m = 0$, there is some ambiguity and arbitrariness in picking a form for g_0 , with $\log \rho_{<}$, $\log \rho_{>}$, and $\log(\rho_{<}/\rho_{>})$ are all valid choices. To recover the form given in the book, let us choose

$$g_0 = C_0 \log \rho_{>}.$$

By the connecting relation of the derivatives, it can be easily shown that

$$C_0 = -4\pi.$$

Putting all the coefficients together, the Green function can be written as

$$G(\rho, \phi; \rho', \phi') = -2 \log \rho_{>} + \sum_{m \neq 0} \frac{1}{m} \left(\frac{\rho_{<}}{\rho_{>}} \right)^{|m|} e^{im(\phi-\phi')}$$

$$= -\log \rho_{>}^2 + 2 \sum_{m=1}^{\infty} \frac{1}{m} \left(\frac{\rho_{<}}{\rho_{>}} \right)^m \cos [m(\phi - \phi')].$$

The second term can be summed. Using the identity

$$\log(1 - x) = - \sum_{n=1}^{\infty} \frac{x^n}{n},$$

the second term can be expressed as

$$\begin{aligned} & 2 \sum_{m=1}^{\infty} \frac{1}{m} \left(\frac{\rho_{<}}{\rho_{>}} \right)^m \cos [m(\phi - \phi')] \\ = & -2 \cdot \operatorname{Re} \left\{ \sum_{m=1}^{\infty} \frac{1}{m} \left(\frac{\rho_{<}}{\rho_{>}} \right)^m e^{im(\phi - \phi')} \right\} \\ = & -2 \cdot \operatorname{Re} \left\{ \log \left(1 - \frac{\rho_{<}}{\rho_{>}} e^{im(\phi - \phi')} \right) \right\} \\ = & -\log \left(1 + \frac{\rho_{<}^2}{\rho_{>}^2} - 2 \frac{\rho_{<}}{\rho_{>}} \cos(\phi - \phi') \right). \end{aligned}$$

Combine with the first term, we can recover the Green function of part (a).