

8.4 (a) For TM wave,  $E_z(\rho, \phi) = E_0 J_m\left(\frac{\chi_{mn}}{R} \rho\right) e^{\pm im\phi}$ , where  $\chi_{mn}$  is the  $n$ -th solution of  $J_m(x)$ . The cutoff frequency is  $\omega_{mn} = \chi_{mn}/R\sqrt{\mu\epsilon}$ .

Using Eq. (8.51), the transmitted power is

$$\begin{aligned} P &= \frac{1}{2} \int_A \frac{\epsilon}{\mu} \left(\frac{\omega}{\omega_{mn}}\right)^2 \left(1 - \frac{\omega_{mn}^2}{\omega^2}\right)^{1/2} |E_z|^2 dA \\ &= \pi E_0^2 \int_0^R \frac{\epsilon}{\mu} \left(\frac{\omega}{\omega_{mn}}\right)^2 \left(1 - \frac{\omega_{mn}^2}{\omega^2}\right)^{1/2} \rho J_m\left(\frac{\chi_{mn}}{R} \rho\right)^2 d\rho \\ &= \frac{\pi R^2 E_0^2}{2} \int_0^1 \frac{\epsilon}{\mu} \left(\frac{\omega}{\omega_{mn}}\right)^2 \left(1 - \frac{\omega_{mn}^2}{\omega^2}\right)^{1/2} J_{m+1}(\chi_{mn})^2 \chi_{mn}^2 d\chi_{mn} \end{aligned}$$

From Eq. (8.59), the power loss is

$$\begin{aligned} -\frac{dP}{dz} &= \frac{1}{2\sigma\delta} \left(\frac{\omega}{\omega_{mn}}\right)^2 \oint_C \frac{1}{\mu^2 \omega_{mn}^2} \left|\frac{\partial E_z}{\partial \rho}\right|^2 d\rho \\ &= \frac{1}{2\sigma\delta} \left(\frac{\omega}{\omega_{mn}}\right)^2 \cdot 2\pi R \cdot \frac{\chi_{mn}^2}{\mu^2 \omega_{mn}^2 R^2} E_0^2 J_m'(\chi_{mn})^2 \\ &= \frac{\epsilon}{\mu} \pi R E_0^2 \frac{1}{\sigma\delta} \left(\frac{\omega}{\omega_{mn}}\right)^2 J_m'(\chi_{mn})^2 \end{aligned}$$

Therefore, the attenuation constant is

$$\begin{aligned} \beta_{mn} &= -\frac{1}{2P} \frac{dP}{dz} = \frac{\pi R E_0^2 \frac{\epsilon}{\mu} \frac{1}{\sigma\delta} \left(\frac{\omega}{\omega_{mn}}\right)^2 J_m'(\chi_{mn})^2}{\pi R^2 E_0^2 \int_0^1 \frac{\epsilon}{\mu} \left(\frac{\omega}{\omega_{mn}}\right)^2 \left(1 - \frac{\omega_{mn}^2}{\omega^2}\right)^{1/2} J_{m+1}(\chi_{mn})^2 \chi_{mn}^2 d\chi_{mn}} \\ &= \int_0^1 \frac{\epsilon}{\mu} \frac{1}{R\sigma\delta} \left(1 - \frac{\omega_{mn}^2}{\omega^2}\right)^{-1/2} \left[ \frac{J_m'(\chi_{mn})}{J_{m+1}(\chi_{mn})} \right]^2 d\chi_{mn} = \int_0^1 \frac{\epsilon}{\mu} \frac{1}{R\sigma\delta} \left(1 - \frac{\omega_{mn}^2}{\omega^2}\right)^{-1/2} d\chi_{mn} \quad \text{Since } J_m'(\chi_{mn}) = -J_{m+1}(\chi_{mn}) \end{aligned}$$

(b) For TE wave,  $H_z(\rho, \phi) = H_0 J_m\left(\frac{\chi'_{mn}}{R} \rho\right) e^{\pm im\phi}$ , where  $\chi'_{mn}$  is the  $n$ -th root of  $J_m'(x)$ . The cutoff frequency is  $\omega_{mn} = \chi'_{mn}/R\sqrt{\mu\epsilon}$ .

Using Eq. (8.51), the transmitted power is

$$\begin{aligned} P &= \frac{1}{2} \int_A \frac{\mu}{\epsilon} \left(\frac{\omega}{\omega_{mn}}\right)^2 \left(1 - \frac{\omega_{mn}^2}{\omega^2}\right)^{1/2} |H_z|^2 dA \\ &= \pi H_0^2 \int_0^R \frac{\mu}{\epsilon} \left(\frac{\omega}{\omega_{mn}}\right)^2 \left(1 - \frac{\omega_{mn}^2}{\omega^2}\right)^{1/2} \rho J_m\left(\frac{\chi'_{mn}}{R} \rho\right)^2 d\rho \end{aligned}$$

$$= \frac{\pi R^2 H_0^2}{2} \left[ \frac{\mu}{\epsilon} \left( \frac{\omega}{\omega_{mn}} \right)^2 \left( 1 - \frac{\omega_{mn}^2}{\omega^2} \right)^{1/2} \left( 1 - \frac{m^2}{\chi_{mn}^{\prime 2}} \right) J_m(\chi_{mn}')^2 \right]$$

From Eq (8.59), the power loss is

$$- \frac{dP}{dz} = \frac{1}{2\sigma\delta} \left( \frac{\omega}{\omega_{mn}} \right)^2 \oint \left[ \frac{1}{\mu\epsilon\omega_{mn}^2} \left( 1 - \frac{\omega_{mn}^2}{\omega^2} \right) |\vec{n} \times \nabla_t H_z|^2 + \frac{\omega_{mn}^2}{\omega^2} |H_z|^2 \right] dl$$

Since  $\nabla_t H_z = \left( \hat{\rho} \frac{\partial}{\partial \rho} + \hat{\phi} \frac{1}{\rho} \frac{\partial}{\partial \phi} \right) H_z = H_0 e^{\pm im\phi} \left[ \hat{\rho} \frac{\chi_{mn}'}{R} J_m' \left( \frac{\chi_{mn}'}{R} \rho \right) \pm \hat{\phi} \frac{im}{\rho} J_m \left( \frac{\chi_{mn}'}{R} \rho \right) \right]$ ,

at the boundary  $\rho = R$ , the first term is simply zero, and

$$|\vec{n} \times \nabla_t H_z|^2 \Big|_{\rho=R} = \frac{m^2 H_0^2}{R^2} J_m(\chi_{mn}')^2$$

Then, 
$$- \frac{dP}{dz} = \frac{1}{2\sigma\delta} \left( \frac{\omega}{\omega_{mn}} \right)^2 \left[ \frac{1}{\mu\epsilon\omega_{mn}^2} \left( 1 - \frac{\omega_{mn}^2}{\omega^2} \right) \cdot \frac{m^2 H_0^2}{R^2} J_m(\chi_{mn}')^2 + \frac{\omega_{mn}^2}{\omega^2} H_0^2 J_m(\chi_{mn}')^2 \right] \times 2\pi R$$

$$= \pi R H_0^2 \frac{1}{\sigma\delta} \left( \frac{\omega}{\omega_{mn}} \right)^2 \left[ \frac{1}{\mu\epsilon\omega_{mn}^2} \left( 1 - \frac{\omega_{mn}^2}{\omega^2} \right) \frac{m^2}{R^2} + \frac{\omega_{mn}^2}{\omega^2} \right] J_m(\chi_{mn}')^2$$

and the attenuation constant is

$$\beta_{mn} = - \frac{1}{2P} \frac{dP}{dz} = \sqrt{\frac{\epsilon}{\mu}} \frac{1}{R\sigma\delta} \frac{\frac{1}{\mu\epsilon\omega_{mn}^2} \left( 1 - \frac{\omega_{mn}^2}{\omega^2} \right) \frac{m^2}{R^2} + \frac{\omega_{mn}^2}{\omega^2}}{\left( 1 - \frac{\omega_{mn}^2}{\omega^2} \right)^{1/2} \left( 1 - \frac{m^2}{\chi_{mn}^{\prime 2}} \right)}$$

$$= \sqrt{\frac{\epsilon}{\mu}} \frac{1}{R\sigma\delta} \left( 1 - \frac{\omega_{mn}^2}{\omega^2} \right)^{-1/2} \left[ \frac{m^2}{\chi_{mn}^{\prime 2} - m^2} + \frac{\omega_{mn}^2}{\omega^2} \right]$$