

8.6 (a) For TM mode, $E_z(\rho, \phi) = E_0 J_m(\chi_{mn} \rho/R) e^{\pm i m \phi}$, and the resonant frequency is

$$\omega_{mnp} = \frac{1}{\sqrt{\mu \epsilon}} \left(\frac{\chi_{mn}^2}{R^2} + \frac{p^2 \pi^2}{L^2} \right)^{1/2}, \text{ with } p \geq 0 \text{ and } \chi_{mn} \text{ being the } n\text{-th root of } J_m(x).$$

Similarly, for TE mode, $H_z(\rho, \phi) = H_0 J_m(\chi'_{mn} \rho/R) e^{\pm i m \phi}$, and

$$\omega_{mnp} = \frac{1}{\sqrt{\mu \epsilon}} \left(\frac{\chi'^2_{mn}}{R^2} + \frac{p^2 \pi^2}{L^2} \right)^{1/2}, \text{ with } p \geq 0 \text{ and } \chi'_{mn} \text{ being the } n\text{-th root of } J'_m(x).$$

(b) The lowest resonant frequency is achieved with TM₀₁₀ mode. For this mode, the stored energy is

$$U = \frac{L}{2} \epsilon \int_A |E_z|^2 da = \pi \epsilon L E_0^2 \int_0^R \rho J_0^2 \left(\chi_{01} \frac{\rho}{R} \right)^2 d\rho = \frac{\epsilon L E_0^2}{2} \pi R^2 J_1^2(\chi_{01}),$$

and the power loss is

$$P_{\text{loss}} = \frac{1}{2\sigma\delta} \left[\oint_C dl \int_0^L dz |\vec{n} \times \vec{H}|^2_{\text{sides}} + 2 \int_A da |\vec{n} \times \vec{H}|^2_{\text{ends}} \right]$$

$$\text{Since } \nabla_t E_z = \left(\hat{\rho} \frac{\partial}{\partial \rho} + \hat{\phi} \frac{1}{\rho} \frac{\partial}{\partial \phi} \right) E_z = E_0 e^{\pm i m \phi} \left[\hat{\rho} \frac{\chi_{mn}}{R} J'_m \left(\chi_{mn} \frac{\rho}{R} \right) + \hat{\phi} \frac{i m}{\rho} J_m \left(\chi_{mn} \frac{\rho}{R} \right) \right]$$

$$= E_0 \frac{\chi_{01}}{R} J'_0 \left(\chi_{01} \frac{\rho}{R} \right) \hat{\rho}, \text{ for } m=0, n=1,$$

the magnetic field is

$$\vec{H}_t = \frac{i}{\mu \omega} E_0 \frac{\chi_{01}}{R} J'_0 \left(\chi_{01} \frac{\rho}{R} \right) \hat{\phi}, \text{ for } p=0.$$

The power loss now becomes

$$P_{\text{loss}} = \frac{1}{2\sigma\delta} \frac{E_0^2}{\mu^2 \omega_{010}^2} \frac{\chi_{01}^2}{R^2} \left[2\pi R \cdot L \cdot J_0'^2(\chi_{01}) + 4\pi \int_0^R \rho J_0' \left(\chi_{01} \frac{\rho}{R} \right)^2 d\rho \right]$$

$$= \frac{1}{2\sigma\delta} \frac{E_0^2}{\mu^2 \omega_{010}^2} \frac{\chi_{01}^2}{R^2} \left[2\pi R \cdot L \cdot J_1^2(\chi_{01}) + 4\pi R^2 \int_0^1 t J_1^2(\chi_{01} t)^2 dt \right]$$

$$= \frac{1}{2\sigma\delta} \frac{E_0^2}{\mu^2 \omega_{010}^2} \frac{\chi_{01}^2}{R^2} \cdot 2\pi R J_1^2(\chi_{01}) (R+L) = \frac{1}{2\sigma\delta} \frac{\epsilon}{\mu} E_0^2 \cdot 2\pi R J_1^2(\chi_{01}) (R+L)$$

Where we have used the identity

$$\int_0^1 x J_1(ax)^2 dx = \frac{1}{2} \left(J_1(a)^2 - \frac{2J_1(a)J_0(a)}{a} + J_0(a)^2 \right)$$

and $J_0(a) = 0$ for $a = \chi_{01}$. Then,

$$Q = \omega_{010} \frac{U}{P_{\text{loss}}} = \omega_{010} \cdot \mu \sigma \delta \frac{L}{2 \left(1 + \frac{L}{R} \right)} = \mu \sqrt{2\sigma \omega_{010} / \mu \epsilon} \frac{L}{2 \left(1 + \frac{L}{R} \right)} = \frac{\mu}{\mu \epsilon} \frac{L}{\sqrt{\frac{2}{\mu \epsilon \omega_{010}}}} \frac{1}{1 + \frac{L}{R}} = \frac{\mu}{\mu \epsilon} \frac{L}{\delta} \frac{1}{1 + \frac{L}{R}}$$