

8.12 (a) Here, we choose the normal direction to be inwardly pointing. Given the unperturbed solution, the perturbed mode is given by $\psi = \psi_0 + \frac{\partial \psi_0}{\partial n} \delta(x, y)$, where $\delta(x, y)$ is positive if the deformation is away from the surface. For the new mode, the eigenequation is $(\nabla^2 + \gamma^2) \psi = 0$.

Applying the Green's theorem on ψ and ψ_0 , we have

$$\int_{S_0} (\psi \nabla_t^2 \psi_0^* - \psi_0^* \nabla_t^2 \psi) da = \oint_{C_0} \left(\psi_0^* \frac{\partial \psi}{\partial n} - \psi \frac{\partial \psi_0^*}{\partial n} \right) dl.$$

With inward normal direction. The L.H.S. of the equation is

$$(\gamma^2 - \gamma_0^2) \int_{S_0} \psi \psi_0^* da = (\gamma^2 - \gamma_0^2) \int_{S_0} \left(\psi_0 + \frac{\partial \psi_0}{\partial n} \delta(x, y) \right) \psi_0^* da \approx (\gamma^2 - \gamma_0^2) \int_{S_0} |\psi_0|^2 da,$$

The R.H.S. becomes

$$\begin{aligned} \oint_{C_0} \left[\psi_0^* \frac{\partial \psi}{\partial n} + \psi_0^* \frac{\partial^2 \psi_0}{\partial n^2} \delta(x, y) - \psi_0 \frac{\partial \psi_0^*}{\partial n} - \left| \frac{\partial \psi_0}{\partial n} \right|^2 \delta(x, y) \right] dl \\ = - \oint_{C_0} \delta(x, y) \left[\left| \frac{\partial \psi_0}{\partial n} \right|^2 - \psi_0^* \frac{\partial^2 \psi_0}{\partial n^2} \right] dl, \end{aligned}$$

where terms in the form of $\psi_0^* \frac{\partial \psi}{\partial n}$ and $\psi_0 \frac{\partial \psi_0^*}{\partial n}$ vanish due to the boundary condition.

Then,

$$\gamma^2 - \gamma_0^2 = - \frac{\oint_{C_0} \delta(x, y) \left[\left| \frac{\partial \psi_0}{\partial n} \right|^2 - \psi_0^* \frac{\partial^2 \psi_0}{\partial n^2} \right] dl}{\int_{S_0} |\psi_0|^2 da}$$

(b) It is straightforward to show that, for $TM_{m,n}$ and $TE_{m,n}$ modes,

$$\int_{S_0} |\psi_0|^2 da = \begin{cases} ab/4, & m \neq 0, n \neq 0 \\ ab/2, & m = 0 \text{ or } n = 0 \end{cases}$$

Now, for $TM_{m,n}$ mode, $m \neq 0, n \neq 0$, and $\psi_0(x, y) = \sin\left(\frac{m\pi x}{a}\right) \sin\left(\frac{n\pi y}{b}\right)$. Only the first term in the numerator will contribute. For this configuration,

$$\delta(x, y) = \frac{y}{b} \delta$$

and

$$\frac{\partial \psi_0}{\partial n} = \begin{cases} \frac{\partial \psi_0}{\partial x} \Big|_{x=0} = \frac{m\pi}{a} \sin\left(\frac{n\pi y}{b}\right) \\ - \frac{\partial \psi_0}{\partial x} \Big|_{x=a} = - \frac{m\pi}{a} \sin\left(\frac{n\pi y}{b}\right) \end{cases}$$

The line integral is given by

$$\oint_C \delta(x,y) \left| \frac{\partial \psi_0}{\partial n} \right|^2 d\ell = 2 \int_0^b \frac{y}{b} \delta \cdot \frac{m^2 \pi^2}{a^2} \cdot \sin^2\left(\frac{\pi y}{b}\right) dy = \frac{2\delta}{b} \cdot \frac{m^2 \pi^2}{a^2} \cdot \frac{b^2}{4} = \frac{b \pi^2 \delta}{2a^2} m^2.$$

The correction is

$$\gamma^2 - \gamma_0^2 = - \frac{b \pi^2 \delta}{2a^2} m^2 / \frac{ab}{4} = - \frac{2 \pi^2 \delta}{a^3} m^2,$$

and $\gamma^2 = \gamma_0^2 - \frac{2 \pi^2 \delta}{a^3} = \left(\frac{1}{a^2} + \frac{1}{b^2} \right) \pi^2 - \frac{2 \delta}{a^3} \pi^2$, for $TM_{1,1}$ mode.

For $TE_{1,0}$ mode, $\psi_0(x,y) = \cos\left(\frac{\pi x}{a}\right)$, and

$$\frac{\partial^2 \psi}{\partial n^2} = \frac{\partial^2 \psi}{\partial x^2} = -\frac{\pi^2}{a^2} \cos\left(\frac{\pi x}{a}\right) = \begin{cases} -\pi^2/a^2, & x=0 \\ \pi^2/a^2, & x=a \end{cases}$$

Also, $\psi_0|_{x=0} = 1$, $\psi_0|_{x=a} = -1$, the line integral is

$$\oint_C \delta(x,y) \psi_0^2 \frac{\partial^2 \psi_0}{\partial n^2} d\ell = -2 \int_0^b \frac{y}{b} \delta \cdot \frac{\pi^2}{a^2} dy = -\frac{b \pi^2 \delta}{a^2}.$$

And the correction is

$$\gamma^2 - \gamma_0^2 = - \frac{b \pi^2 \delta}{a^2} / \frac{ab}{2} = - \frac{2 \pi^2 \delta}{a^3}.$$

and $\gamma^2 = \gamma_0^2 - \frac{2 \pi^2 \delta}{a^3} = \left(\frac{1}{a^2} - \frac{2 \delta}{a^3} \right) \pi^2.$