

9.10 (a) From the definition of orbital magnetization, we know

$$\vec{M}(r, \theta, \phi, t) = \frac{1}{2} \vec{r} \times \vec{j} = -i \frac{v_0 a_0}{4Z} (\vec{r} \times \hat{z}) \rho(r, \theta, \phi, t),$$

Since $\vec{r} \times \vec{r} = 0$. Also,

$$\vec{r} \times \hat{z} = r \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \sin\theta\cos\phi & \sin\theta\sin\phi & \cos\theta \\ 0 & 0 & 1 \end{vmatrix} = r(\hat{x} \sin\theta\sin\phi - \hat{y} \sin\theta\cos\phi) = (\hat{x} \sin\phi - \hat{y} \cos\phi) r \sin\theta,$$

$$\text{Then, } \vec{M}(r, \theta, \phi, t) = -i \frac{v_0 a_0}{4} \frac{r \sin\theta}{r \cos\theta} (\hat{x} \sin\phi - \hat{y} \cos\phi) \rho(r, \theta, \phi, t) = -i \frac{v_0 a_0}{4} \tan\theta (\hat{x} \sin\phi - \hat{y} \cos\phi) \rho(r, \theta, \phi, t)$$

$$\begin{aligned} \text{Using the charge expression, } \rho(r, \theta, \phi, t) &= \frac{2e}{\sqrt{6} a_0^4} r e^{-3r/2a_0} Y_{00} Y_{10} e^{-i\omega_0 t} \\ &= \frac{\sqrt{2} e}{4\pi a_0^4} r \cos\theta e^{-3r/2a_0} e^{-i\omega_0 t}, \end{aligned}$$

the magnetization becomes

$$\vec{M}(r, \theta, \phi, t) = -i \frac{\sqrt{2} e a_0}{16\pi a_0^3} (y \hat{x} - x \hat{y}) e^{-3r/2a_0} e^{-i\omega_0 t}$$

Direct calculation shows that

$$\begin{aligned} \nabla \cdot \vec{M} &= \frac{\partial M_x}{\partial x} + \frac{\partial M_y}{\partial y} + \frac{\partial M_z}{\partial z} \\ &= -i \frac{\sqrt{2} e a_0}{16\pi a_0^3} \left(-\frac{3}{2a_0} e^{-3r/2a_0} \right) \frac{1}{r} (xy - yx) e^{-i\omega_0 t} \\ &= 0. \end{aligned}$$

$$\text{Also, } \vec{r} \times \vec{M} = -i \frac{\sqrt{2} e a_0}{16\pi a_0^3} \left(xz \hat{x} + yz \hat{y} - (x^2 + y^2) \hat{z} \right) e^{-3r/2a_0} e^{-i\omega_0 t}$$

and it is easy to show that $\nabla \cdot (\vec{r} \times \vec{M}) = 0$. Then, by Eqs. (9.170) and (9.172), only non-vanishing multipole moments will be

$$Q_{lm} = \int r^l Y_{lm}^* \rho d^3x,$$

which is non-zero only when $l=1$ and $m=0$, as $\rho(r, \theta, \phi, t) \propto P_1(\cos\theta)$. Performing the integral,

$$\begin{aligned} Q_{10} &= \int r \cdot \sqrt{\frac{3}{4\pi}} \cos\theta \cdot \frac{\sqrt{2} e}{4\pi a_0^4} r \cos\theta e^{-3r/2a_0} e^{-i\omega_0 t} d^3x \\ &= \sqrt{\frac{3}{4\pi}} \frac{\sqrt{2} e}{4\pi a_0^4} e^{-i\omega_0 t} \cdot 2\pi \cdot \int_{-1}^1 \cos^2\theta d(\cos\theta) \cdot \int_0^\infty r^4 e^{-3r/2a_0} dr \end{aligned}$$

$$= \sqrt{\frac{3}{4\pi}} \frac{\sqrt{e}}{4\pi a_0^2} e^{-i\omega_0 t} 2\pi \cdot \frac{2}{3} \cdot \frac{24}{(3/2a_0)^5} = \frac{2^8}{3^5} \sqrt{\frac{3}{2\pi}} e a_0 e^{-i\omega_0 t}$$

Therefore $a_E(1,0) = \frac{ck^2}{3i} \sqrt{2} D_{10} = -ick^2 e a_0 \frac{2^8}{3^6} \sqrt{\frac{3}{\pi}}$.

(b) Using Eq. (9.151), we have

$$\begin{aligned} \frac{dP(1,0)}{d\Omega} &= \frac{Z_0}{2k^2} |a_E(1,0)|^2 |\hat{X}_{10}|^2 = \frac{Z_0}{2k^2} c^2 k^6 e^2 a_0^2 \frac{2^{16}}{3^{12}} \frac{3}{\pi} \cdot \frac{3}{8\pi} \sin^2\theta \\ &= \frac{Z_0 c^2 k^4 e^2 a_0^2}{\pi^2} \frac{2^{12}}{3^{10}} \sin^2\theta, \end{aligned}$$

and total power radiated is

$$P = \frac{Z_0 c^2 k^4 e^2 a_0^2}{\pi^2} \frac{2^{12}}{3^{10}} \frac{8\pi}{3} = \frac{Z_0 c^2 k^4 e^2 a_0^2}{\pi} \frac{2^{15}}{3^{11}} = \frac{2^{15}}{3^{11}} \cdot \frac{\omega_0^4 e^2 a_0^2}{\pi \epsilon_0 c^3},$$

where we have used that $k = \omega_0/c$, and $Z_0 c = \sqrt{\frac{\mu_0}{\epsilon_0}} \cdot \frac{1}{\sqrt{\mu_0 \epsilon_0}} = \frac{1}{\epsilon_0}$. After some tedious algebra, it

can be shown that $P = \left(\frac{2}{3}\right)^8 (\hbar\omega_0) \left(\frac{d^4c}{a_0}\right)$

(c) Skipped

(d) Under the assumption of circular orbits, the charge density can be expressed as

$$\rho(\vec{r}) = \frac{e}{r^2} \delta(r - 2a_0) \delta(\cos\theta) \delta(\phi - \omega_0 t).$$

Then the dipole moment is

$$\begin{aligned} \vec{p} &= \int \rho \vec{r} d^3x = \int \rho \begin{Bmatrix} x \\ y \\ z \end{Bmatrix} d^3x = \int_0^{2\pi} d\phi \int_{-1}^1 d(\cos\theta) \int_0^{2\pi} \delta(r - 2a_0) \delta(\cos\theta) \delta(\phi - \omega_0 t) e r \begin{Bmatrix} \sin\theta \cos\phi \\ \sin\theta \sin\phi \\ \cos\theta \end{Bmatrix} \\ &= 2ea_0 \int_0^{2\pi} \begin{Bmatrix} \cos(\omega_0 t) \\ \sin(\omega_0 t) \\ 0 \end{Bmatrix} d\phi = 2ea_0 \operatorname{Re} \left[\begin{Bmatrix} 1 \\ i \\ 0 \end{Bmatrix} e^{-i\omega_0 t} \right]. \end{aligned}$$

Therefore, by Eq. (9.24), the total power radiated is

$$\begin{aligned} P &= \frac{Z_0 c^2 k^4}{12\pi} |\dot{\vec{p}}|^2 = \frac{\omega_0^4}{12\epsilon_0 \pi c^3} 4e^2 a_0^2 \cdot 2 = \frac{2}{3} \frac{\omega_0^4 e^2 a_0^2}{\pi \epsilon_0 c^3} = \frac{2}{3} \cdot \frac{3^3}{2^7} (\hbar\omega_0) \left(\frac{d^4c}{a_0}\right) \\ &= \frac{9}{64} (\hbar\omega_0) \left(\frac{d^4c}{a_0}\right) \end{aligned}$$

Compared with part (b), the ratio is $\frac{3^2/2^6}{2^8/3^8} = \frac{3^{10}}{2^{14}} = 3.604065$.