



$$\mathbf{r}_1 = \mathbf{r} + \mathbf{u} / 2$$

$$\mathbf{r}_2 = \mathbf{r} - \mathbf{u} / 2$$

$$\begin{aligned} \boldsymbol{\tau} &= -\mathbf{r}_1 \times \left(q\mathbf{E} + q \frac{d\mathbf{r}_1}{dt} \times \mathbf{B} \right) + \mathbf{r}_2 \times \left(q\mathbf{E} + q \frac{d\mathbf{r}_2}{dt} \times \mathbf{B} \right) = \\ &= q \left\{ -(\mathbf{r} + \mathbf{u} / 2) \times \mathbf{E} + (\mathbf{r} - \mathbf{u} / 2) \times \mathbf{E} - (\mathbf{r} + \mathbf{u} / 2) \times \left(\frac{d\mathbf{r}}{dt} \times \mathbf{B} + \frac{1}{2} \frac{d\mathbf{u}}{dt} \times \mathbf{B} \right) + (\mathbf{r} - \mathbf{u} / 2) \times \left(\frac{d\mathbf{r}}{dt} \times \mathbf{B} - \frac{1}{2} \frac{d\mathbf{u}}{dt} \times \mathbf{B} \right) \right\} = \\ &= q \left\{ -(\mathbf{r} + \mathbf{u} / 2) \times \mathbf{E} + (\mathbf{r} - \mathbf{u} / 2) \times \mathbf{E} - (\mathbf{r} + \mathbf{u} / 2) \times \left(\frac{d\mathbf{r}}{dt} \times \mathbf{B} + \frac{1}{2} \frac{d\mathbf{u}}{dt} \times \mathbf{B} \right) + (\mathbf{r} - \mathbf{u} / 2) \times \left(\frac{d\mathbf{r}}{dt} \times \mathbf{B} - \frac{1}{2} \frac{d\mathbf{u}}{dt} \times \mathbf{B} \right) \right\} = \\ &= q \left\{ \underbrace{-\mathbf{u}}_{q(-\mathbf{u})=\mathbf{p}} \times \mathbf{E} - \mathbf{r} \times \left(\frac{d\mathbf{u}}{dt} \times \mathbf{B} \right) - \underbrace{\mathbf{u} \times \left(\frac{d\mathbf{r}}{dt} \times \mathbf{B} \right)}_{=0} \right\} = \mathbf{p} \times \mathbf{E} + \mathbf{r} \times \underbrace{\left(\frac{d\mathbf{p}}{dt} \times \mathbf{B} \right)}_{\text{2nd order small}} \end{aligned}$$

$$\begin{aligned} \mathbf{L} &= m\mathbf{r}_1 \times \frac{d\mathbf{r}_1}{dt} + m\mathbf{r}_2 \times \frac{d\mathbf{r}_2}{dt} = \\ &= m \left\{ (\mathbf{r} + \mathbf{u} / 2) \times \left(\frac{d\mathbf{r}}{dt} + \frac{1}{2} \frac{d\mathbf{u}}{dt} \right) + (\mathbf{r} - \mathbf{u} / 2) \times \left(\frac{d\mathbf{r}}{dt} - \frac{1}{2} \frac{d\mathbf{u}}{dt} \right) \right\} = \\ &= \frac{m}{q^2} \left\{ 2\mathbf{r} \times \underbrace{\frac{d\mathbf{r}}{dt}}_{=0} + \frac{1}{2} \mathbf{p} \times \frac{d\mathbf{p}}{dt} \right\} = \frac{m}{2q^2} \mathbf{p} \times \frac{d\mathbf{p}}{dt} \end{aligned}$$



$$\frac{d\mathbf{L}}{dt} = \frac{m}{2q^2} \frac{d}{dt} \left(\mathbf{p} \times \frac{d\mathbf{p}}{dt} \right)$$

$$\frac{d\mathbf{L}}{dt} = \boldsymbol{\tau} \Rightarrow \frac{d}{dt} \left(\mathbf{p} \times \frac{d\mathbf{p}}{dt} \right) = \frac{2q^2}{m} \mathbf{p} \times \mathbf{E}$$

$$\frac{d}{dt}\left(\mathbf{p}\times\frac{d\mathbf{p}}{dt}\right)=\frac{2q^2}{m}\mathbf{p}\times\mathbf{E}$$

$$\mathbf{p}(t)=\mathbf{p}_0+\mathbf{p}_1(t)$$

$$\frac{d\mathbf{p}_0}{dt}=0; \quad \mathbf{p}_1(t)=\mathbf{p}_0\times\delta\mathbf{a}(t) \quad (\delta\mathbf{a}\parallel\mathbf{z}, \quad \mathbf{p}_0\in[xoy])$$

$$\frac{d}{dt}\left(\mathbf{p}_0\times\frac{d\mathbf{p}_1}{dt}\right)=\frac{2q^2}{m}\mathbf{p}_0\times\mathbf{E}$$

$$\Downarrow$$

$$\frac{d^2}{dt^2}(\mathbf{p}_0\times\mathbf{p}_1)=\frac{2q^2}{m}\mathbf{p}_0\times\mathbf{E}$$

$$\Downarrow$$

$$\frac{d^2}{dt^2}(\mathbf{p}_0\times\mathbf{p}_1)=\frac{d^2}{dt^2}\left\{\mathbf{p}_0\times[\mathbf{p}_0\times\delta\mathbf{a}(t)]\right\}=\frac{d^2}{dt^2}\left\{\left(\underbrace{\mathbf{p}_0\cdot\delta\mathbf{a}}_{=0}\right)\cdot\mathbf{p}_0-(\mathbf{p}_0\cdot\mathbf{p}_0)\cdot\delta\mathbf{a}\right\}=$$

$$-\mathbf{z}|p_0|^2\frac{d^2(\delta\alpha)}{dt^2}=\mathbf{z}\omega^2(p_{x0}^2+p_{y0}^2)\cdot\delta\alpha$$

$$\text{-----}$$

$$\mathbf{p}_0\times\mathbf{E}=(\mathbf{x}p_{x0}+\mathbf{y}p_{y0})\times(\mathbf{x}E_x+\mathbf{y}E_y)=\mathbf{z}(p_{x0}E_y-p_{y0}E_x)$$

$$\text{-----}$$

$$\omega^2(p_{x0}^2+p_{y0}^2)\cdot\delta\alpha=\frac{2q^2}{m}(p_{x0}E_y-p_{y0}E_x)$$

$$\Downarrow$$

$$\delta\alpha=\frac{2q^2}{\omega^2m}\frac{p_{x0}E_y-p_{y0}E_x}{p_{x0}^2+p_{y0}^2}=A\cdot E_y-B\cdot E_x \quad \left\{A=\frac{2q^2}{\omega^2m}\frac{p_{x0}}{p_{x0}^2+p_{y0}^2}, \quad B=\frac{2q^2}{\omega^2m}\frac{p_{y0}}{p_{x0}^2+p_{y0}^2}\right\}$$

$$\text{-----}$$

$$\mathbf{p}_1(t) = (\mathbf{x}p_{x0} + \mathbf{y}p_{y0}) \times \mathbf{z} (A \cdot E_y - B \cdot E_x) = -\mathbf{y}p_{x0} (A \cdot E_y - B \cdot E_x) + \mathbf{x}p_{y0} (A \cdot E_y - B \cdot E_x)$$

\Downarrow

$$p_{1x} = p_{y0} (A \cdot E_y - B \cdot E_x) = -\frac{2q^2}{\omega^2 m} \frac{(p_{y0})^2}{p_{x0}^2 + p_{y0}^2} \cdot E_x + \frac{2q^2}{\omega^2 m} \frac{p_{y0} \cdot p_{x0}}{p_{x0}^2 + p_{y0}^2} \cdot E_y$$

$$p_{1y} = -p_{x0} (A \cdot E_y - B \cdot E_x) = \frac{2q^2}{\omega^2 m} \frac{p_{x0} \cdot p_{y0}}{p_{x0}^2 + p_{y0}^2} \cdot E_x - \frac{2q^2}{\omega^2 m} \frac{(p_{x0})^2}{p_{x0}^2 + p_{y0}^2} \cdot E_y$$

\Downarrow

$$\begin{pmatrix} p_{1x} \\ p_{1y} \end{pmatrix} = \begin{pmatrix} -\frac{2q^2}{\omega^2 m \varepsilon_0} \frac{(p_{y0})^2}{p_{x0}^2 + p_{y0}^2} & \frac{2q^2}{\omega^2 m \varepsilon_0} \frac{p_{y0} \cdot p_{x0}}{p_{x0}^2 + p_{y0}^2} \\ \frac{2q^2}{\omega^2 m \varepsilon_0} \frac{p_{x0} \cdot p_{y0}}{p_{x0}^2 + p_{y0}^2} & -\frac{2q^2}{\omega^2 m \varepsilon_0} \frac{(p_{x0})^2}{p_{x0}^2 + p_{y0}^2} \end{pmatrix} \cdot \begin{pmatrix} E_x \\ E_y \end{pmatrix} = \varepsilon_0 \bar{\bar{\varepsilon}} \begin{pmatrix} E_x \\ E_y \end{pmatrix}$$

\Downarrow

$$\begin{pmatrix} p_{1x} \\ p_{1y} \end{pmatrix} = \begin{pmatrix} -\varepsilon_1 & \varepsilon_2 \\ \varepsilon_2 & -\varepsilon_3 \end{pmatrix} \cdot \begin{pmatrix} E_x \\ E_y \end{pmatrix} = \varepsilon_0 \bar{\bar{\varepsilon}} \begin{pmatrix} E_x \\ E_y \end{pmatrix}$$

Maxwell's curl equations for the grating layer:

$$\nabla \times \mathbf{H}(x, y) = -i\omega \varepsilon_0 \tilde{\varepsilon} \mathbf{E}(x, y) \quad \parallel \mathbf{D}(x, y) = \varepsilon_0 \mathbf{E} + \mathbf{P} = \varepsilon_0 (\mathbf{E} + \bar{\bar{\varepsilon}} \mathbf{P})$$

$$\nabla \times \mathbf{E}(x, y) = i\omega \mu_0 \mathbf{H}(x, y)$$

where

$$\tilde{\varepsilon} = \begin{pmatrix} 1 - \varepsilon_1 & \varepsilon_2 \\ \varepsilon_2 & 1 - \varepsilon_3 \end{pmatrix} = \begin{pmatrix} 1 - \frac{2q^2}{\omega^2 m \varepsilon_0} \sin^2 \varphi & \frac{2q^2}{\omega^2 m \varepsilon_0} \sin \varphi \cdot \cos \varphi \\ \frac{2q^2}{\omega^2 m \varepsilon_0} \sin \varphi \cdot \cos \varphi & 1 - \frac{2q^2}{\omega^2 m \varepsilon_0} \cos^2 \varphi \end{pmatrix} \equiv \begin{pmatrix} \tilde{\varepsilon}_1 & \tilde{\varepsilon}_2 \\ \tilde{\varepsilon}_2 & \tilde{\varepsilon}_3 \end{pmatrix}$$

$$\{\mathbf{E}(x, y), \mathbf{H}(x, y)\} \sim e^{-i\omega t + ik_x x \pm ik_y y}$$

$$\frac{\partial}{\partial z} \{\mathbf{E}(x, y), \mathbf{H}(x, y)\} = 0$$

(No field variations along the z-direction). We get the following expressions for Maxwell equation:

$$\frac{\partial H_z}{\partial y} = -i\omega \varepsilon_0 (\tilde{\varepsilon}_1 E_x + \tilde{\varepsilon}_2 E_y)$$

$$-\frac{\partial H_z}{\partial x} = -i\omega \varepsilon_0 (\tilde{\varepsilon}_2 E_x + \tilde{\varepsilon}_3 E_y)$$

$$\frac{\partial E_y}{\partial x} - \frac{\partial E_x}{\partial y} = i\omega \mu_0 H_z$$

The components of the electromagnetic fields are periodic functions of x and can be approximately expanded in truncated generalized Fourier series:

$$\{E_x(x, y), H_z(x, y)\} = \sum_{n=-\infty}^{\infty} \{E_{x,n}(y), H_{z,n}(y)\} e^{i(k_x + \frac{2\pi}{h}n)x}$$

where h is a period of the grating. The relation among the Fourier coefficients can be written as:

$$\begin{aligned} \frac{\partial \mathbf{h}_z(y)}{\partial y} &= -i\omega\epsilon_0 \{ [\tilde{\epsilon}_1] \mathbf{e}_x(y) + [\tilde{\epsilon}_2] \mathbf{e}_y(y) \} \\ -i\mathbf{K} \mathbf{h}_z(y) &= -i\omega\epsilon_0 \{ [\tilde{\epsilon}_2] \mathbf{e}_x(y) + [\tilde{\epsilon}_3] \mathbf{e}_y(y) \} \\ i\mathbf{K} \mathbf{e}_y(y) - \frac{\partial \mathbf{e}_x(y)}{\partial y} &= i\omega\mu_0 \mathbf{h}_z(y) \end{aligned}$$

with

$\mathbf{K} = [k_{xn} \delta_{nm}]$ - diagonal matrix

$$[\mathbf{e}_x(y)]_n = E_{x,n}(y)$$

$$[\mathbf{e}_y(y)]_n = E_{y,n}(y)$$

$$[\tilde{\epsilon}_p(x)]_{n,m} = \frac{1}{h} \int_0^h dx \tilde{\epsilon}_p(x) e^{-i(n-m)\frac{2\pi}{h}x} \text{ - (Toeplitz matrix), } p=1, 2, 3$$

From the 2nd expression we get:

$$\begin{aligned} [\tilde{\epsilon}_3] \mathbf{e}_y &= \frac{\mathbf{K}}{\omega\epsilon_0} \mathbf{h}_z - [\tilde{\epsilon}_2] \mathbf{e}_x \\ \Downarrow \\ \mathbf{e}_y &= [\tilde{\epsilon}_3]^{-1} \frac{\mathbf{K}}{\omega\epsilon_0} \mathbf{h}_z - [\tilde{\epsilon}_3]^{-1} [\tilde{\epsilon}_2] \mathbf{e}_x \end{aligned}$$

$$\frac{\partial \mathbf{h}_z(y)}{\partial y} = -i\omega\epsilon_0 \llbracket \tilde{\epsilon}_1 \rrbracket \mathbf{e}_x - i\omega\epsilon_0 \llbracket \tilde{\epsilon}_2 \rrbracket \left\{ \llbracket \tilde{\epsilon}_3 \rrbracket^{-1} \frac{\mathbf{K}}{\omega\epsilon_0} \mathbf{h}_z - \llbracket \tilde{\epsilon}_3 \rrbracket^{-1} \llbracket \tilde{\epsilon}_2 \rrbracket \mathbf{e}_x \right\} =$$

$$-i\omega\epsilon_0 \left(\llbracket \tilde{\epsilon}_1 \rrbracket - \llbracket \tilde{\epsilon}_2 \rrbracket \llbracket \tilde{\epsilon}_3 \rrbracket^{-1} \llbracket \tilde{\epsilon}_2 \rrbracket \right) \mathbf{e}_x - i \llbracket \tilde{\epsilon}_2 \rrbracket \llbracket \tilde{\epsilon}_3 \rrbracket^{-1} \mathbf{K} \mathbf{h}_z$$

\Downarrow

$$\frac{\partial \tilde{\mathbf{h}}_z(y)}{\partial y} = -ik_0 \left(\llbracket \tilde{\epsilon}_1 \rrbracket - \llbracket \tilde{\epsilon}_2 \rrbracket \llbracket \tilde{\epsilon}_3 \rrbracket^{-1} \llbracket \tilde{\epsilon}_2 \rrbracket \right) \mathbf{e}_x - ik_0 \llbracket \tilde{\epsilon}_2 \rrbracket \llbracket \tilde{\epsilon}_3 \rrbracket^{-1} \frac{\mathbf{K}}{k_0} \tilde{\mathbf{h}}_z$$

$$\frac{\partial \mathbf{e}_x(y)}{\partial y} = i\mathbf{K} \mathbf{e}_y(y) - i\omega\mu_0 \mathbf{h}_z(y) =$$

$$i\mathbf{K} \left\{ \llbracket \tilde{\epsilon}_3 \rrbracket^{-1} \frac{\mathbf{K}}{\omega\epsilon_0} \mathbf{h}_z - \llbracket \tilde{\epsilon}_3 \rrbracket^{-1} \llbracket \tilde{\epsilon}_2 \rrbracket \mathbf{e}_x \right\} - i\omega\mu_0 \mathbf{h}_z =$$

$$i \left\{ \llbracket \tilde{\epsilon}_3 \rrbracket^{-1} \frac{\mathbf{K}^2}{\omega\epsilon_0} - \omega\mu_0 \right\} \mathbf{h}_z - i\mathbf{K} \llbracket \tilde{\epsilon}_3 \rrbracket^{-1} \llbracket \tilde{\epsilon}_2 \rrbracket \mathbf{e}_x$$

\Downarrow

$$\frac{\partial \mathbf{e}_x(y)}{\partial y} = ik_0 \left\{ \llbracket \tilde{\epsilon}_3 \rrbracket^{-1} \frac{\mathbf{K}^2}{k_0^2} - \mathbf{I} \right\} \tilde{\mathbf{h}}_z - ik_0 \frac{\mathbf{K}}{k_0} \llbracket \tilde{\epsilon}_3 \rrbracket^{-1} \llbracket \tilde{\epsilon}_2 \rrbracket \mathbf{e}_x$$

$$\frac{\partial}{\partial y} \begin{pmatrix} \mathbf{e}_x(y) \\ \tilde{\mathbf{h}}_z(y) \end{pmatrix} = ik_0 \begin{pmatrix} M_{11} & M_{12} \\ M_{21} & M_{22} \end{pmatrix} \begin{pmatrix} \mathbf{e}_x(y) \\ \tilde{\mathbf{h}}_z(y) \end{pmatrix}$$

$$\mathbf{M}_{11} = -\frac{\mathbf{K}}{k_0} \llbracket \tilde{\epsilon}_3 \rrbracket^{-1} \llbracket \tilde{\epsilon}_2 \rrbracket, \quad \mathbf{M}_{12} = \llbracket \tilde{\epsilon}_3 \rrbracket^{-1} \frac{\mathbf{K}^2}{k_0^2} - \mathbf{I}$$

$$\mathbf{M}_{21} = -\left(\llbracket \tilde{\epsilon}_1 \rrbracket - \llbracket \tilde{\epsilon}_2 \rrbracket \llbracket \tilde{\epsilon}_3 \rrbracket^{-1} \llbracket \tilde{\epsilon}_2 \rrbracket \right), \quad \mathbf{M}_{22} = -\llbracket \tilde{\epsilon}_2 \rrbracket \llbracket \tilde{\epsilon}_3 \rrbracket^{-1} \frac{\mathbf{K}}{k_0}$$

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