

# The double pendulum

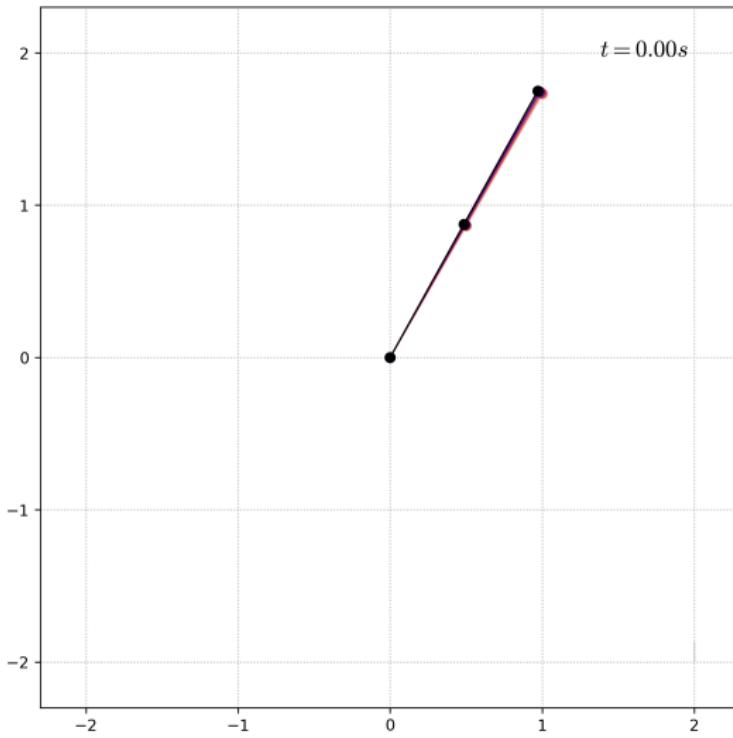
## Is it really that chaotic ?

Vincent Degrooff

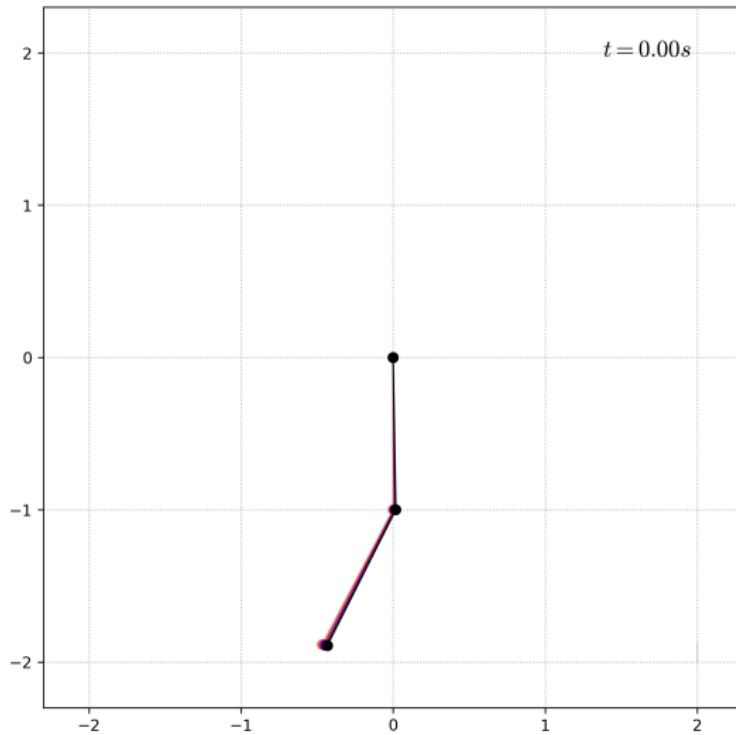
LINMA2361 - Nonlinear dynamical systems  
EPL UCLouvain

Project presentation, January 2022

# Introduction: A chaotic trajectory



# Introduction: A non-chaotic trajectory



# Table of Contents

- 1 Modeling of the mechanical system
- 2 Analysis of the equilibria
- 3 Quasiperiodicity
- 4 Poincaré sections

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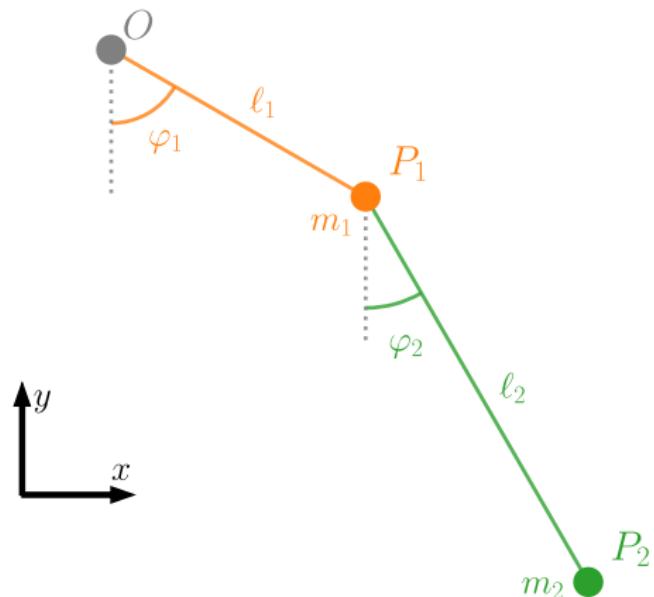
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# Description



# Equations of motion

The Lagrangian of the system:

$$\mathcal{L} = \mathcal{K} - \mathcal{U}$$

$$\mathcal{K} = \frac{m_1}{2} \ell_1^2 \dot{\varphi}_1^2 + \frac{m_2}{2} \left[ \ell_1^2 \dot{\varphi}_1^2 + \ell_2^2 \dot{\varphi}_2^2 + 2\ell_1\ell_2 \dot{\varphi}_1 \dot{\varphi}_2 \cos(\overbrace{\varphi_1 - \varphi_2}^{:=\theta}) \right]$$

$$\mathcal{U} = [(m_1 + m_2) g \ell_1 (1 - \cos \varphi_1) + m_2 g \ell_2 (1 - \cos \varphi_2)]$$

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The Euler-Lagrange equations:

$$\frac{d}{dt} \left( \frac{\partial \mathcal{L}}{\partial \dot{\varphi}_i} \right) = \frac{\partial \mathcal{L}}{\partial \varphi_i}$$

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$$(m_1 + m_2) \ell_1 \ddot{\varphi}_1 + m_2 \ell_2 \ddot{\varphi}_2 \cos \theta = -m_2 \ell_2 \dot{\varphi}_2^2 \sin \theta - (m_1 + m_2) g \sin \varphi_1$$

$$\ell_1 \ddot{\varphi}_1 \cos \theta + \ell_2 \ddot{\varphi}_2 = \ell_1 \dot{\varphi}_1^2 \sin \theta - g \sin \varphi_2$$

# State-space formulation

A choice of nondimensionalization ?

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$$\lambda = \frac{\ell_2}{\ell_1} \quad \mu = \frac{m_2}{m_1 + m_2} \quad \tau = \sqrt{\frac{g}{\ell_1}} t$$

# State-space formulation

A choice of nondimensionalization ?

$$\lambda = \frac{\ell_2}{\ell_1} \quad \mu = \frac{m_2}{m_1 + m_2} \quad \tau = \sqrt{\frac{g}{\ell_1}} t$$

We obtain the following system

$$f(x) = \frac{d}{dt} \begin{bmatrix} \varphi_1 \\ \varphi_2 \\ \omega_1 \\ \omega_2 \end{bmatrix} = \begin{bmatrix} \omega_1 \\ \omega_2 \\ [\mu \cos \theta (\sin \varphi_2 - \omega_1^2 \sin \theta) - \lambda \mu \omega_2^2 \sin \theta - \sin \varphi_1] \frac{1}{1 - \mu \cos^2 \theta} \\ \frac{1}{\lambda} \sin \theta [\cos \varphi_1 + \omega_1^2 + \lambda \mu \omega_2^2 \cos \theta] \frac{1}{1 - \mu \cos^2 \theta} \end{bmatrix}$$

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# What about the equilibria

- ▶ How many ?
- ▶ What coordinates ?
- ▶ Which one(s) is/are stable ?

# Eigenvalues of the 4 equilibria

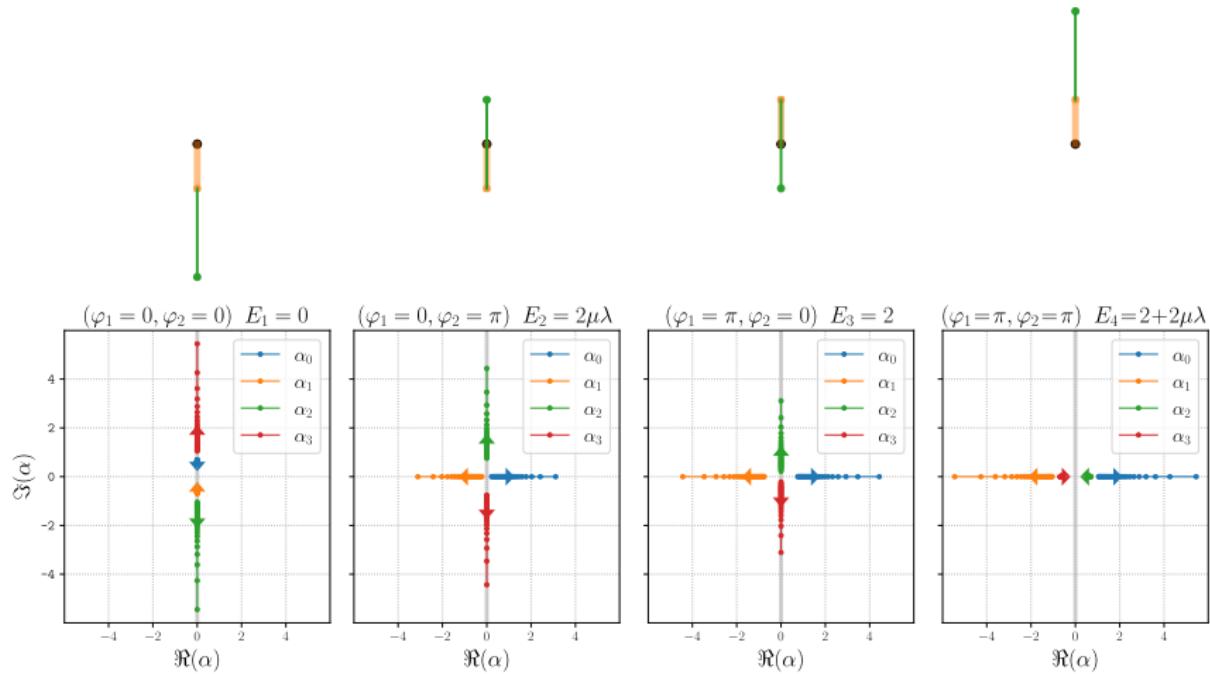


Figure: Eigenvalues for  $\lambda = 2$  and  $\mu \in [0.05, 0.95]$ .

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# Simple system

## Example

$$\dot{\theta}_1 = \omega_1$$

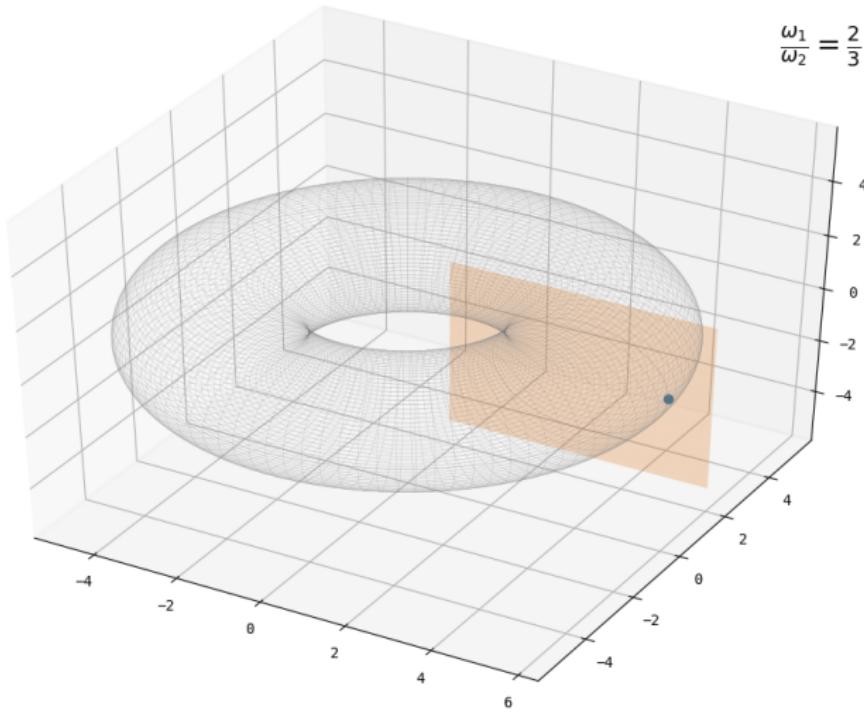
$$\dot{\theta}_2 = \omega_2$$

Two possibilities

- ▶  $\omega_1/\omega_2 = p/q$  rational
- ▶  $\omega_1/\omega_2$  irrational

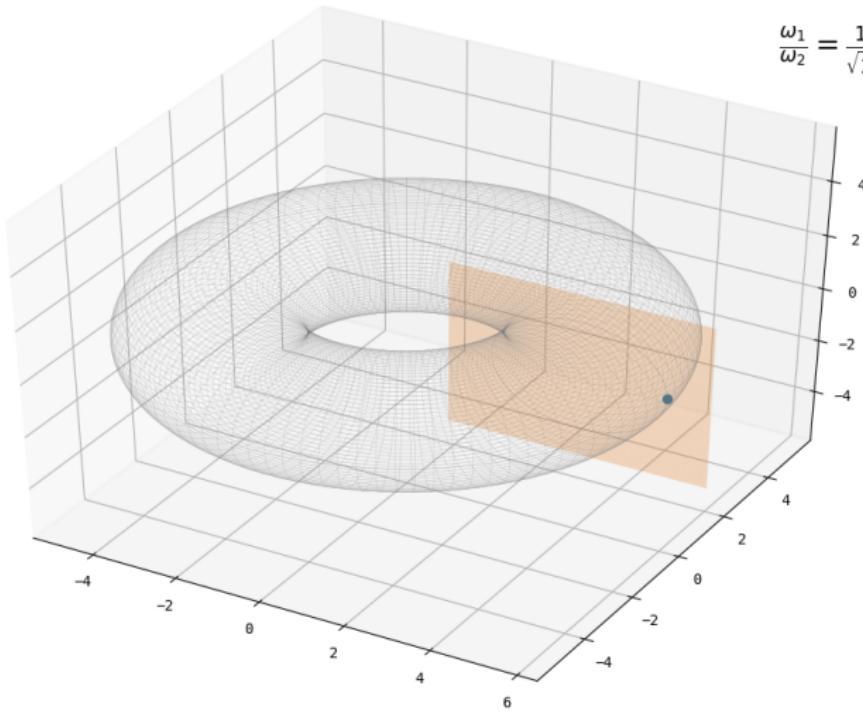
# Commensurable frequencies

$$\frac{\omega_1}{\omega_2} = \frac{2}{3}$$



# Incommensurable frequencies

$$\frac{\omega_1}{\omega_2} = \frac{1}{\sqrt{2}}$$



# Recap

- ▶ Periodic orbits: finite number of points
- ▶ Quasiperiodic orbits: curve-filling points
- ▶ Chaotic orbits: area-filling points

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# What is a Poincaré section

- ▶ Cannot visualize 4-dimensional state-space
- ▶ Pick a lower-dimensional slice of the state-space and look for intersection points with the trajectory

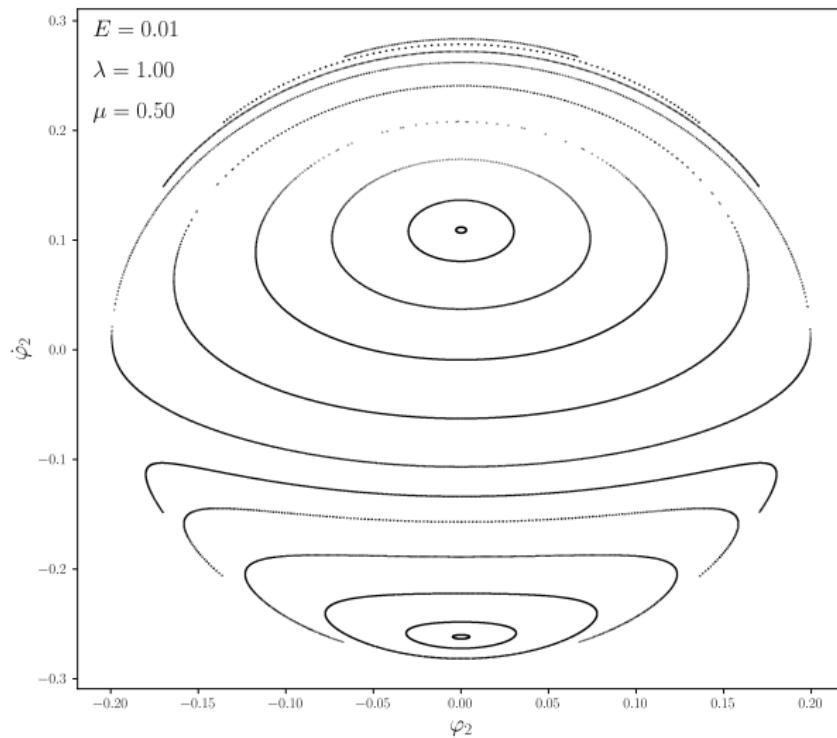
# What is a Poincaré section

- ▶ Cannot visualize 4-dimensional state-space
- ▶ Pick a lower-dimensional slice of the state-space and look for intersection points with the trajectory

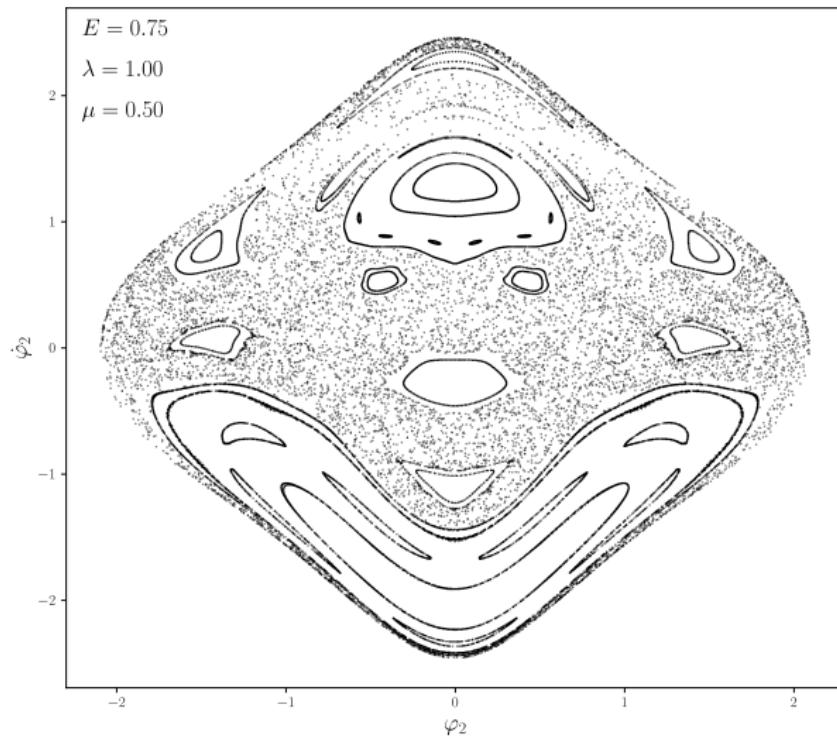
In the case of the double pendulum, section can be defined:

$$\begin{aligned} \varphi_1 = 0 \quad &\text{and} \quad \dot{\varphi}_1 + \mu\lambda\dot{\varphi}_2 \cos \varphi_2 > 0 \\ \text{or} \quad \varphi_2 = 0 \quad &\text{and} \quad \dot{\varphi}_2 + \frac{1}{\lambda}\dot{\varphi}_1 \cos \varphi_1 > 0 \end{aligned}$$

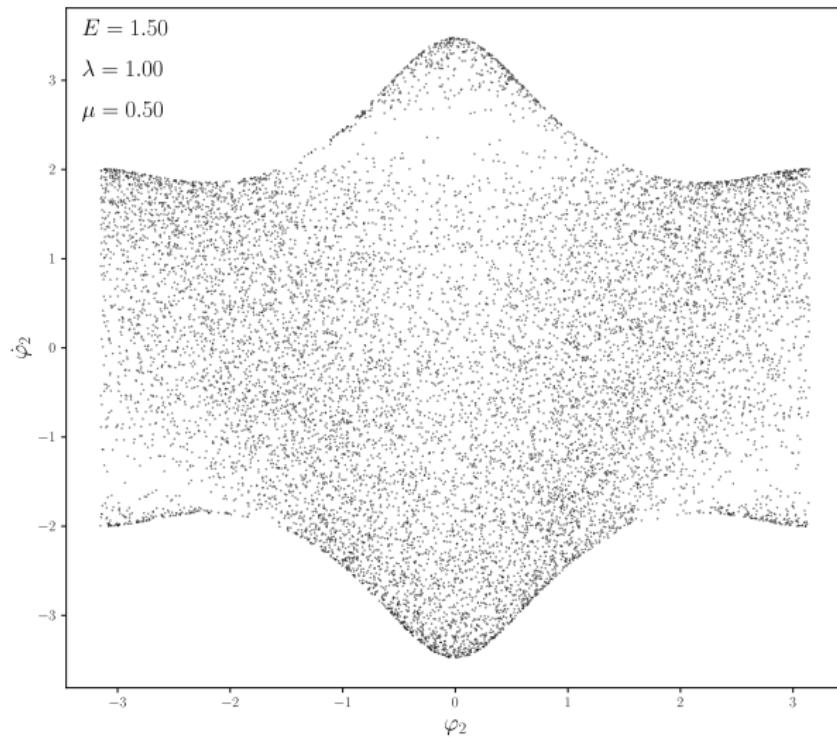
# Numerical experiment - low energy



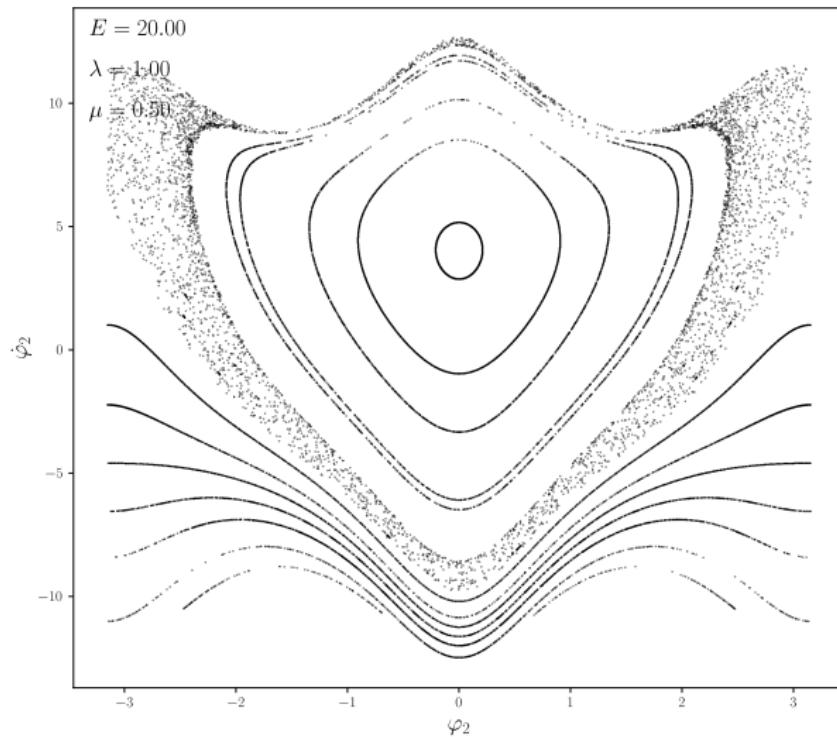
# Numerical experiment - mid-low energy



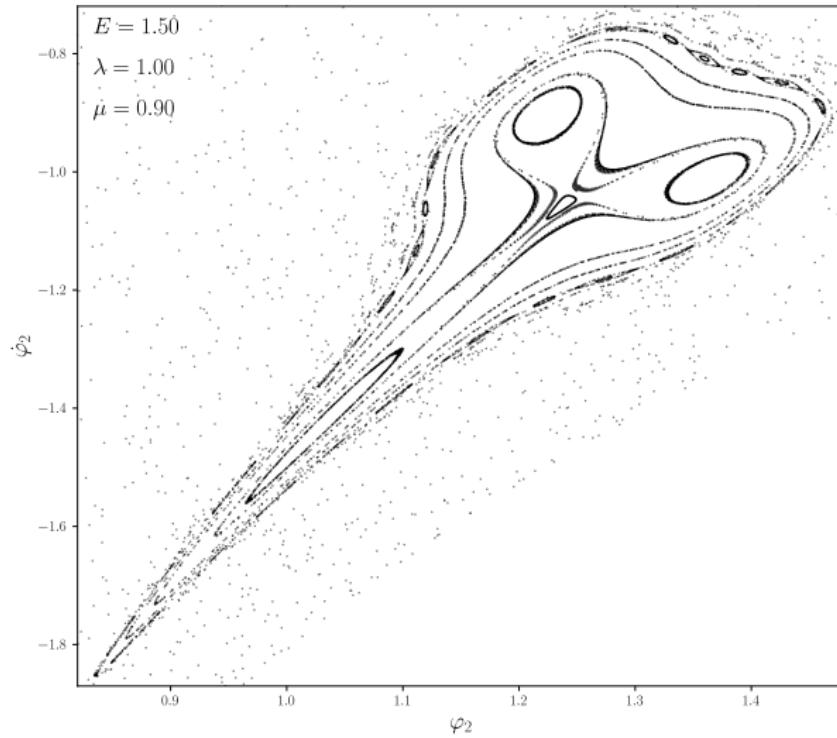
# Numerical experiment - mid-high energy



# Numerical experiment - high energy



# Numerical experiment - Fractal



# Recap

- ▶ Low energy: 2 normal modes of the linear model which are periodic and their associated quasiperiodic orbits
- ▶ Mid-range energy: new types of periodic orbits being progressively replaced by chaotic regions
- ▶ High energies: chaotic regions replaced by (quasi)periodic orbits similar to a simple pendulum

# Extra - Energy influence - 1

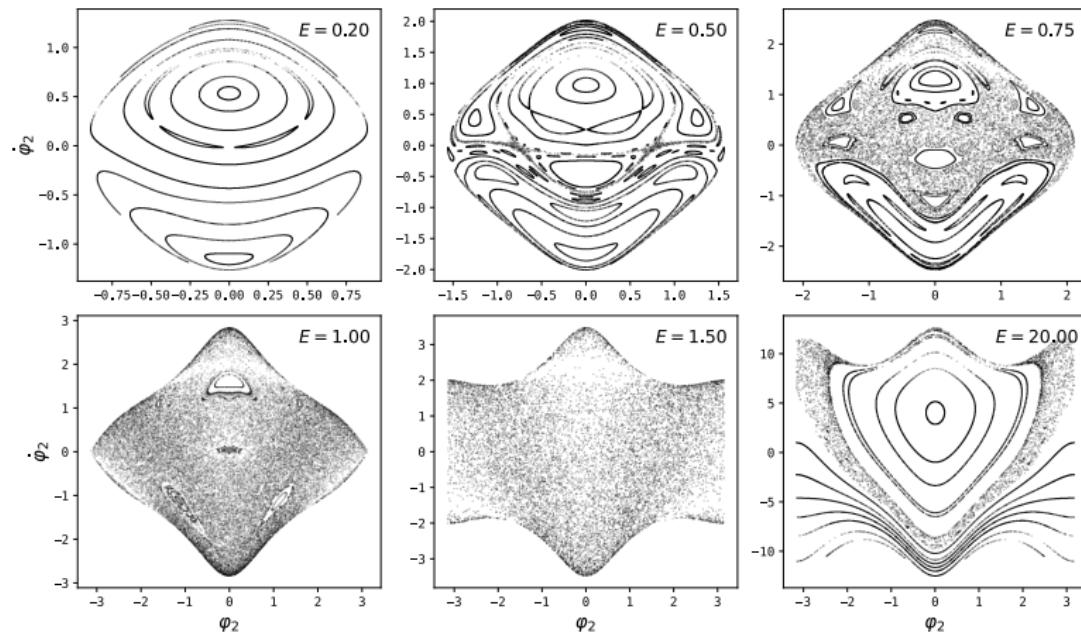


Figure: Poincaré sections for  $\lambda = 1$  and  $\mu = 0.5$ , with condition  $\varphi_1 = 0$ .

# Extra - Energy influence - 2

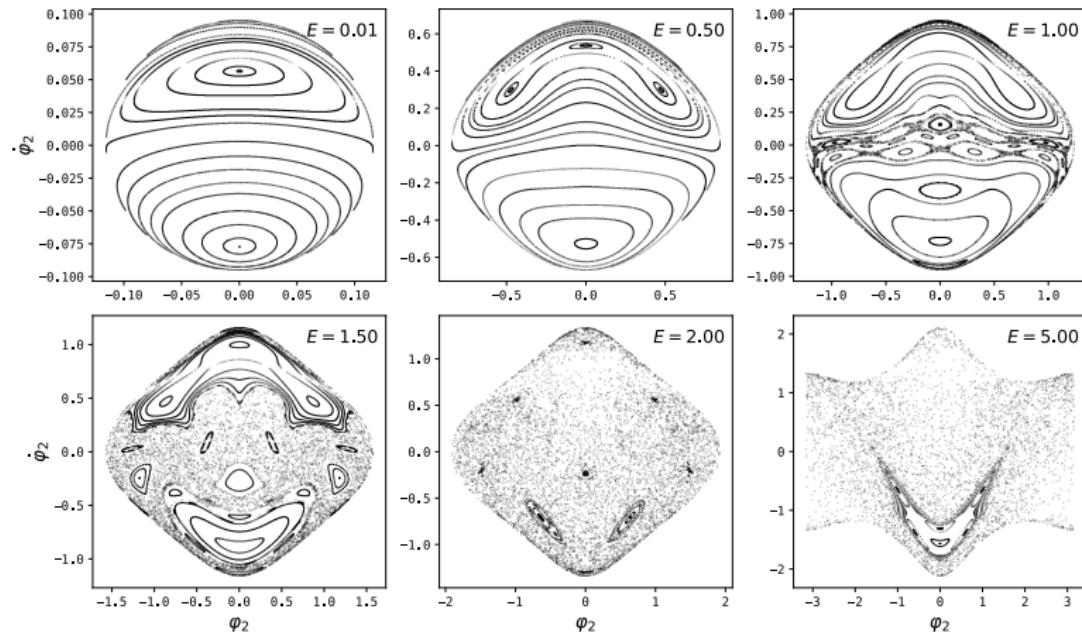


Figure: Poincaré sections for  $\lambda = 3$  and  $\mu = 0.5$ , with condition  $\varphi_1 = 0$ .

# Extra - Energy influence - 3

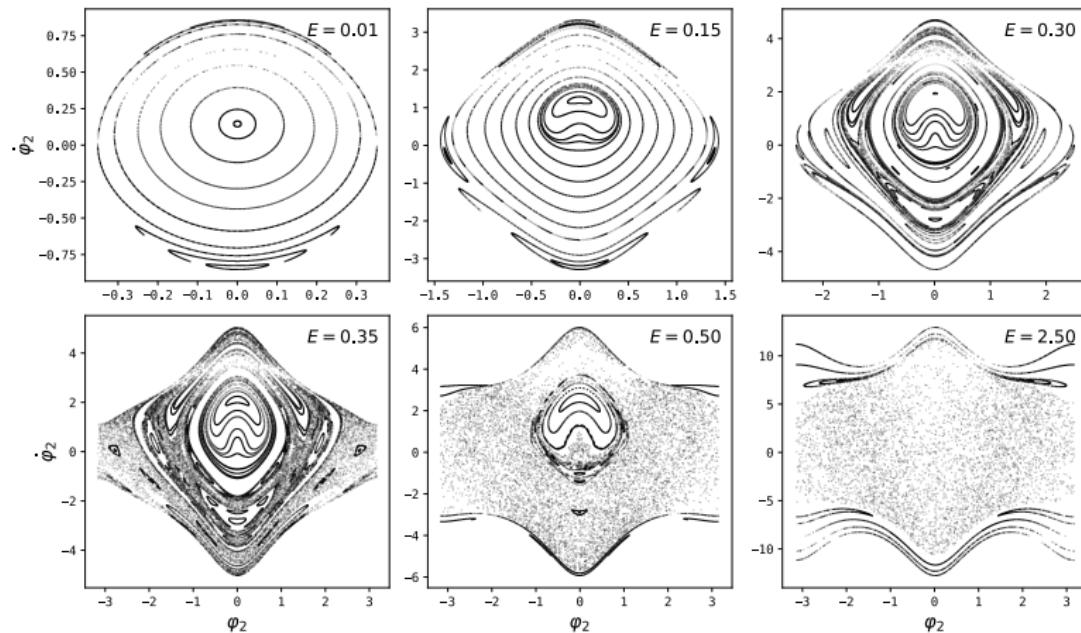


Figure: Poincaré sections for  $\lambda = 1/3$  and  $\mu = 0.5$ , with condition  $\varphi_1 = 0$ .

# Extra - Energy influence - 4

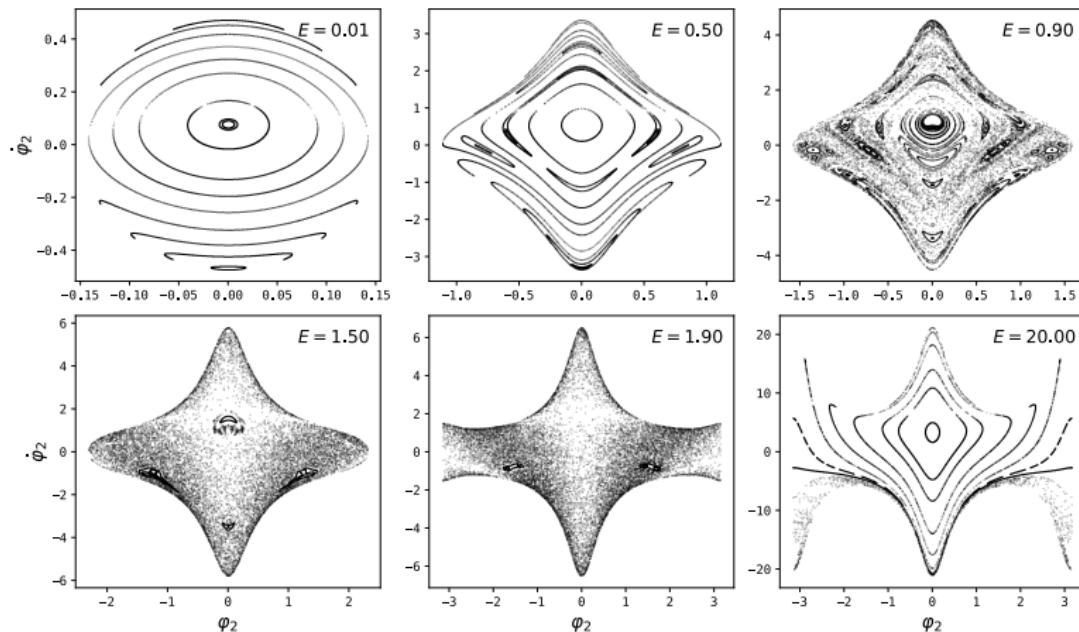


Figure: Poincaré sections for  $\lambda = 1$  and  $\mu = 0.9$ , with condition  $\varphi_1 = 0$ .

# Extra - Energy influence - 5

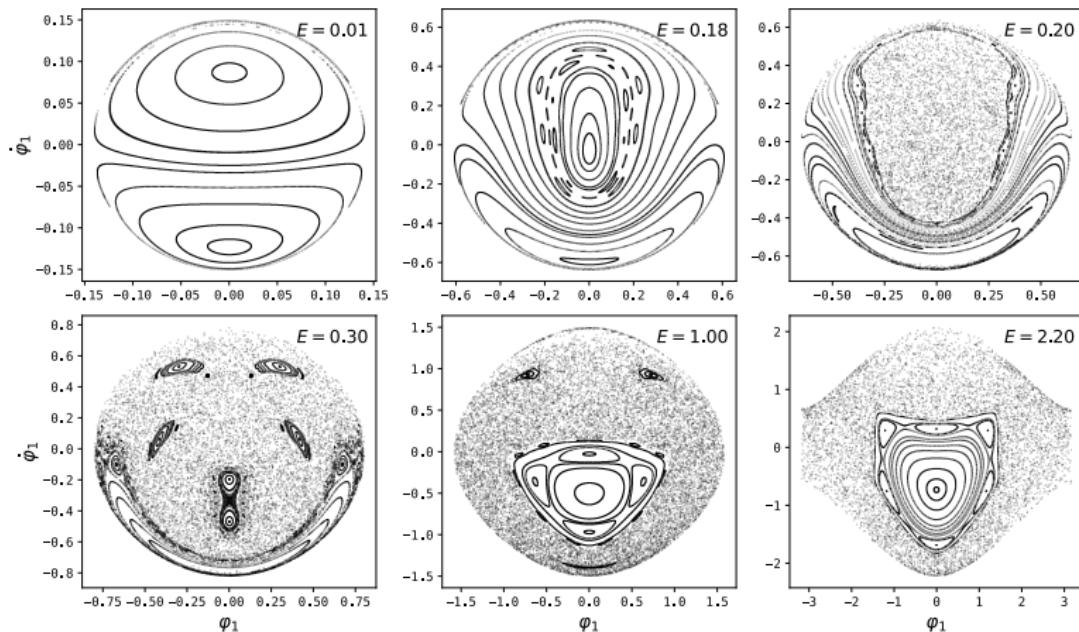


Figure: Poincaré sections for  $\lambda = 1$  and  $\mu = 0.1$ , with condition  $\varphi_2 = 0$ .