Lab 2

Lukas Finkbeiner

March 6, 2020

Abstract

We need more narrative drive. Consider using powerful transition phrases

We investigated the hyperfine hydrogen transition by calibrating our measurements to thermal noise and by averaging over many blocks of data taken in optimal time windows of the day. We investigated the speed of light using waveguides (How do they work?). Our Chi-Squared analysis is completely deficient because our data were taken too broadly to show noticeable inconsistencies.

This is a great place to have a to-do section, because I prefer to write abstracts at the end of the writing process.

* Ask Professor about the permissions thing (page 3), to make sure there are no technical difficulties.

Figure out how to make the section headers smaller. They are taking up far too much space.

1   Introduction and Background

Block diagram of the telescope electronics. Fortunately, it seems like the professor took care of most of this for you in lecture, February 11.

The data that we sample from the Big Horn via the PicoScope will feature arbitrary units which depend on our setup. To calibrate the intensity of the spectrum, we first calculate the gain:

$$G = \frac{T_{\text{sys, cal}} - T_{\text{sys, cold}}}{\sum(s_{\text{cal}} - s_{\text{cold}})} \sum s_{\text{cold}} \approx \frac{300 \text{ K}}{\sum(s_{\text{cal}} - s_{\text{cold}})} \sum s_{\text{cold}} \tag{1}$$

Where $T_{\text{sys, cal}} \approx 98.6° \text{ F} \approx 310 \text{ K}$ represents the temperature of the human flesh that we use to cover most of the telescope. $T_{\text{sys, cold}}$ represents the temperature of the cold sky. We expect $T_{\text{sys, cold}} \ll T_{\text{sys, cal}}$, whereby we get the approximated form on the right. s_{cal} represents the spectrum for which we attempt to maximize thermal noise in the collector, and s_{cold} represents the spectrum for which we attempt to minimize.

To remove constant sources of interference in our data, we take an 'on-spectrum' where Line shape:

$$s_{\text{line}} = \frac{s_{\text{on}}}{s_{\text{off}}} \tag{2}$$

Putting it all together:

$$T_{\text{line}} = s_{line} \times G \tag{3}$$

For our first experiment with waveguides, we seek to calculate to calculate the wavelength λ_{sl} of a signal inside the waveguide based on the positions of the minima in the squared voltage, which we call nulls. We use the following equation for this first experiment; we expect a linear dependence of null spacings on λ_{sl}.

To plot our calibrated spectrum against Doppler velocity, we use the following relation:

$$v = -c\frac{\Delta\nu}{\nu_0} \tag{4}$$

$$x_m = A + m\frac{\lambda_{\text{sl}}}{2} \tag{5}$$

m is the index of the null, x_m is the position of the null m, and A represents a constant offset which can represent whether the far end of the waveguide is open or shorted.

X-band waveguide (this one requires more motivation):

$$\lambda_g = \frac{\lambda_{\text{fs}}}{[1 - (\frac{\lambda_{\text{fs}}}{2a})^2]^{1/2}} \tag{6}$$

λ_g is the guide wavelength, which we directly measure. λ_{fs} is the free-space wavelength of the signal, which we will calculate using the known input frequency and the relation $c = f\lambda_{\text{fs}}$. a is the width of the waveguide.

Here we define reduced chi-squared as follows:

$$\chi_r^2 = \frac{1}{N - M} \sum_i \frac{|y_i - \hat{y}_i|^2}{\sigma_i^2} \tag{7}$$

2 Methods

To place our test signal in the upper and lower sidebands, we use these two frequencies. To put the hydrogren in the upper and lower sidebands, we set the first local oscillator to these other two frequencies.

We measure λ_g, the guide wavelength, as the distance between the nulls in the guide output. We measure a with a set of calipers and compare this with our results from a least-squares fit to the data based on equation 6.

3 Observations

As a preface, I may want to include commentary on how I had to manually alter the data, based on readings from the oscilloscope, in order to account for the imbalance in picoscope inputs?? I think not!
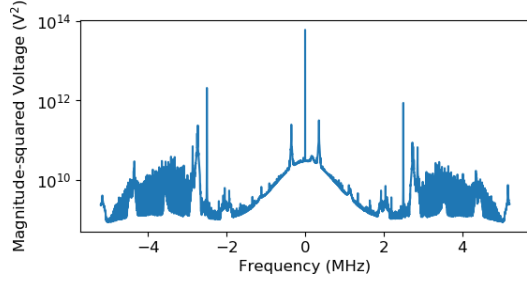
Figure 1: A semi-log plot over the range of frequencies sampled. We assume that our 2 MHz low-pass filter works, so we will dismiss the signals farther than 2 MHz from the center. Additionally, the large central spike and smaller 'bunny ear' spikes ($\pm.5$ MHz) appear on all data sets as persistent interference; we partially ignore these by limiting the y-axis. Maybe you should use an initial test sample, to discuss this.
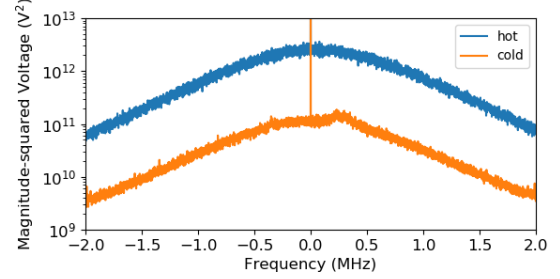

Figure 2: The 'hot' data correspond to three humans standing in front of the collector. The 'cold' data correspond to the collector pointing up at the cold sky. The noise is not an issue because of its regularity: observe the even discrepancy between the curves over the domain of interest.
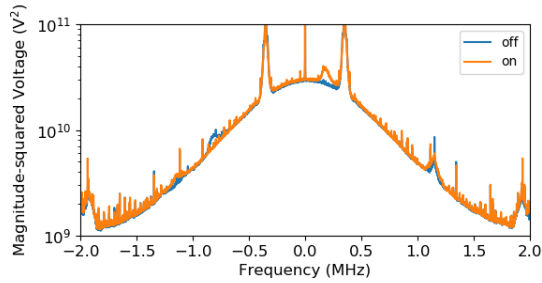

Figure 3: Combined plot of the 'on' and 'off' (LO1 at 1230 and 1231 MHz, respectively) power spectra. As we expect, the HI signal shifts by about 1 MHz between the two plots. This also supports our interpretation of the other patterns as interference: these patterns do not move between spectra.

Figure 4: This one is not helpful. Replace it with the final, calibrated spectrum, thermal power units.

Figure 1 shows a host of interfering signals, a consequence of the imperfections in our setup (such as low signal-to-noise ratio from low power on the local oscillators). How can I justify ignoring the 2 MHz spikes? Aren't those bad?

Data from the XBand waveguide section, frequency and resulting null positions.:

f (GHz)	Null Positions (cm)
7.0	8.8, 15.15
7.2	11.95, 17.15
7.5	10.05, 14.4
8.0	10.15, 13.35, 16.7
8.5	9.85, 12.7, 15.25
9.0	9.3, 12.3, 14.25, 16.7

f (GHz)	Null Positions (cm)
9.5	8.9, 11.15, 13.35, 15.45, 17.7
10.0	8.45, 10.45, 12.3, 14.35, 16.5
10.5	9.95, 11.75, 13.5, 15.45, 17.3
11.0	9.35, 11.4, 12.65, 14.45, 16.2, 17.85
11.5	8.95, 10.45, 12.0, 13.8, 15.15, 16.8
12.0	8.4, 9.85, 11.35, 12.85, 14.4, 15.9, 17.45

4   Analysis

Since the Horn was pointing straight up, we may say that the declination is equal to the latitude of the telescope. In our case, that is $37.873199° \approx 0.661012$ radians. The right ascension is equal to the local sidereal time (LST) of data-capture.

We labeled all of our samples with unix times. Included below is a table including examples of date conversions to Doppler corrections for the events of interest. For the sake of brevity, we omit the counterparts for the Cassiopeia measurements.

Label	Unix Time	Julian Date	LST
On-line start	1582937209.9447305	2458908.5325225084	0.8334479646673989
On-line stop	1582937704.4486437	2458908.5382459336	0.8695077621535888
Off-line start	1582937837.5090616	2458908.5397859844	0.8792106794915676
Off-line stop	1582938387.1761112	2458908.5461478718	0.9192930359644231

Label	PST	Doppler Correction (m / s)
On-line start	2/28/20 @ 16:46:49	-27648.19266498681
On-line stop	2/28/20 @ 16:55:04	-27897.548197380296
Off-line start	2/28/20 @ 16:57:17	-27960.009947759394
Off-line stop	2/28/20 @ 17:06:27	-28197.027228226034

The Doppler correction exhibits, at first glance, a critical time sensitivity. Observe that, just for one source (directly overhead), the final Doppler correction differed from the initial by about 550 m / s. However, as we shall soon see, the velocity axis will span many megameters, so taking the average of the Doppler corrections should suffice for a global correction.

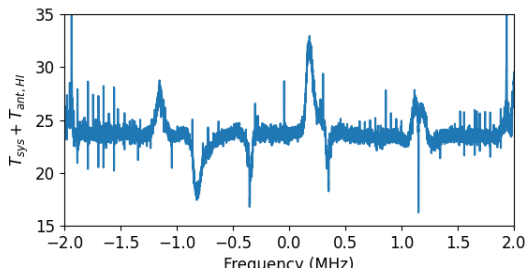If I have extra space, put the results from the Cass. Doppler correction calculation

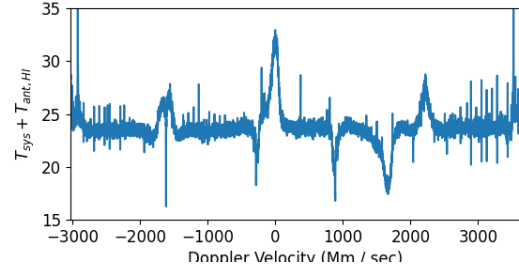Figure 5: Here is a fully calibrated line. We see HI for sure!

Figure 6: Here is that same line, now with a velocity axis. What is the point I am trying to make with this plot?
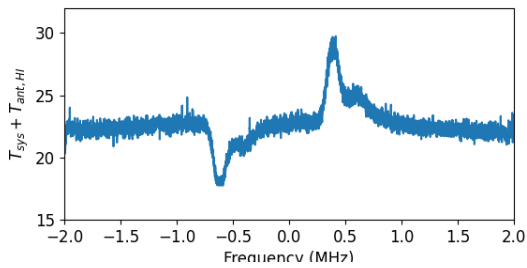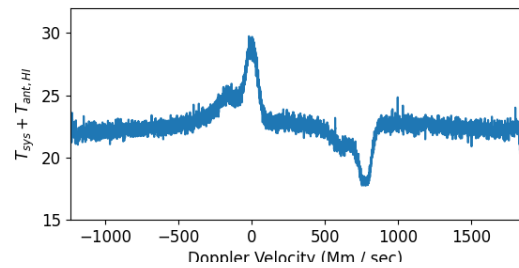

Figure 7: ...

Figure 8: ...

5   Conclusions

On the bright side, if you cannot reach ten pages, that means that you can expand all graphs until the font size is no longer a problem!

6   Acknowledgments

You have to remedy your complete ignorance of BibTex