

```
In [41]: import timeit
import numpy as np
import matplotlib.pyplot as plt #and the useful plotting library
from numba import jit
```

Diffusion part 1: the fundamental solution

The one-dimensional diffusion equation is:

$$\frac{\partial u}{\partial t} = \nu \frac{\partial^2 u}{\partial x^2}$$

The first thing you should notice is that —unlike the previous two simple equations we have studied— this equation has a second-order derivative. We first need to learn what to do with it!

Discretizing $\frac{\partial^2 u}{\partial x^2}$

The second-order derivative can be represented geometrically as the line tangent to the curve given by the first derivative. We will discretize the second-order derivative with a Central Difference scheme: a combination of Forward Difference and Backward Difference of the first derivative. Consider the Taylor expansion of u_{i+1} and u_{i-1} around u_i :

$$u_{i+1} = u_i + \Delta x \frac{\partial u}{\partial x} \Big|_i + \frac{\Delta x^2}{2} \frac{\partial^2 u}{\partial x^2} \Big|_i + \frac{\Delta x^3}{3!} \frac{\partial^3 u}{\partial x^3} \Big|_i + O(\Delta x^4)$$

$$u_{i-1} = u_i - \Delta x \frac{\partial u}{\partial x} \Big|_i + \frac{\Delta x^2}{2} \frac{\partial^2 u}{\partial x^2} \Big|_i - \frac{\Delta x^3}{3!} \frac{\partial^3 u}{\partial x^3} \Big|_i + O(\Delta x^4)$$

If we add these two expansions, you can see that the odd-numbered derivative terms will cancel each other out. If we neglect any terms of $O(\Delta x^4)$ or higher (and really, those are very small), then we can rearrange the sum of these two expansions to solve for our second-derivative.

$$u_{i+1} + u_{i-1} = 2u_i + \Delta x^2 \frac{\partial^2 u}{\partial x^2} \Big|_i + O(\Delta x^4)$$

Then rearrange to solve for $\frac{\partial^2 u}{\partial x^2} \Big|_i$ and the result is:

$$\frac{\partial^2 u}{\partial x^2} = \frac{u_{i+1} - 2u_i + u_{i-1}}{\Delta x^2} + O(\Delta x^4)$$

Discretizing both $\frac{\partial u}{\partial t}$ and $\frac{\partial^2 u}{\partial x^2}$

We can now write the discretized version of the diffusion equation in 1D:

$$\frac{u_i^{n+1} - u_i^n}{\Delta t} = \nu \frac{u_{i+1}^n - 2u_i^n + u_{i-1}^n}{\Delta x^2}$$

As before, we notice that once we have an initial condition, the only unknown is u_i^{n+1} , so we re-arrange the equation solving for our unknown:

$$u_i^{n+1} = u_i^n + \underbrace{\frac{\nu \Delta t}{\Delta x^2}}_{\beta} (u_{i+1}^n - 2u_i^n + u_{i-1}^n)$$

$$u_i^{n+1} = \beta u_{i-1}^n + u_i^n (1 - 2\beta) + \beta u_{i+1}^n$$

The above discrete equation allows us to write a program to advance a solution in time. But we need an initial condition. Let's continue using our favorite: the hat function. So, at $t = 0$, $u = 1$ in the interval $0.25 \leq x \leq 0.5$ and $u = 0$ everywhere else. We are ready to number-crunch!

```
In [42]: nx = 128
domain_length = 64
dx = domain_length / (nx-1)
xspace = np.linspace(0, domain_length, nx)

nt = 200 # the number of timesteps we want to calculate
nu = 5 # the value of viscosity
sigma = .2 # sigma is a parameter, we'll learn more about it later
dt = sigma * dx**2 / nu # dt is defined using sigma ... more later!

u_IC = 0*np.ones(nx) # numpy function ones()
u_IC[int(nx/4):int(nx/2)] = 1 # setting u = 1 between 0.25 and 0.5 as per our I.C.s

plt.plot(xspace, u_IC)
```

```
Out[42]: [matplotlib.lines.Line2D object at 0x7f870cdabc70>]
```

```
In [43]: def calc_diffusion_FD_btcs_naive(IC,nx,nt,nu,dt):
# backward time, central space
u = IC.copy()
un = IC.copy() #our placeholder array, un, to advance the solution in time
beta = nu * dt / dx**2

for n in range(nt): #iterate through time
un = u.copy() #copy the existing values of u into un
for i in range(0, nx):
# this is slow (index operations & branching)
if i == nx-1:
u[i] = beta*un[i-1]+ (1-2*beta)*un[i] + beta*un[0] # periodic BC
else:
u[i] = beta*un[i-1]+ (1-2*beta)*un[i] + beta*un[i+1]
return u

# @jit(nopython=True)
# @jit(nopython=True, parallel=True)
def calc_diffusion_FD_btcs(IC,nx,nt,nu,dt):
# backward time, central space
u = IC.copy()
un = IC.copy() #our placeholder array, un, to advance the solution in time
beta = nu * dt / dx**2

c_ind = np.arange(0, nx)
l_ind = np.roll(c_ind, -1)
r_ind = np.roll(c_ind, 1)

for n in range(nt): #iterate through time
un = u.copy() # copy the existing values of u into un

lap_u = un[l_ind] - 2 * un[c_ind] + un[r_ind] # periodic BC
u = un + beta* lap_u
return u

starttime = timeit.default_timer()
u_FD = calc_diffusion_FD_btcs(u_IC,nx,nt,nu,dt)
# u_FD = calc_diffusion_FD_btcs_naive(u_IC,nx,nt,nu,dt)
print("The time difference is :", timeit.default_timer() - starttime)
plt.plot(xspace, u_FD)
```

```
Out[43]: The time difference is : 0.0007448280084645376
[matplotlib.lines.Line2D object at 0x7f870da23d0>]
```

The convolution - part I

<https://numpy.org/doc/stable/reference/generated/numpy.convolve.html>

<https://en.wikipedia.org/wiki/Convolution>

The discrete convolution operation is known as:

$$(a \star v)[n] = \sum_{m=-\infty}^{\infty} a[m]v[n-m]$$

Notice, that single step of the explicit algorithm implemented before can be expressed as convolution with a $[\beta, 1 - 2\beta, \beta]$ filter. To compute more time steps, one have to convolve many times.

Task

Solve the diffusion equation by convolving the initial condition with the filter in each iteration.

```
In [44]: def calc_diffusion_iterate_convolutions(IC,nx,nt,nu,dt):
u = IC.copy()
un = IC.copy() # our placeholder array, un, to advance the solution in time
beta = nu * dt / dx**2
filter = np.array((beta,1-2*beta,beta))
for n in range(nt): # iterate through time
un = u.copy() # copy the existing values of u into un
u = np.convolve(filter,un, 'same')
return u

u_iter_conv = calc_diffusion_iterate_convolutions(u_IC,nx,nt,nu,dt)
plt.plot(xspace, u_iter_conv)
```

```
Out[44]: [matplotlib.lines.Line2D object at 0x7f8707c8c040>]
```

The convolution - part II

The fundamental solution of the heat equation is the Gaussian function (impulse response).

Consider "diffusion" of a single particle. The probability of finding a particle after T time steps follows the normal (a.k.a Gaussian) distribution. To compute it, one can convolve the initial position with a Gaussian. The result will be equivalent with repeated convolutions with small filter.

This means that convolving with a Gaussian tells us the solution to the diffusion equation after a fixed amount of time. This is the same as low pass filtering an image. So smoothing, low pass filtering, diffusion, all mean the same thing.

Task

Solve the diffusion equation by convolving the initial condition with the Gaussian.

```
In [45]: def calc_diffusion_single_convolution(IC,x,nt,nu,dt):
u = IC.copy()
def get_gaussian(x, alfa, t):
g = -(x-domain_length/2.)***2
g /= (4*alfa*t)
g = np.exp(g)
g /= np.sqrt(4*np.pi*alfa*t)
g *= domain_length/(nx-1) # normalize --> sum(g)=1
return g

time_spot = dt*nt
fundamental_solution = get_gaussian(x, nu, time_spot)

u = np.convolve(fundamental_solution, u, 'same')
# plt.plot(x, fundamental_solution, marker='v', linestyle="", markevery=5)
return u, fundamental_solution

u_single_conv, fs = calc_diffusion_single_convolution(u_IC,xspace,nt,nu,dt)
plt.plot(xspace, u_single_conv)
```

```
Out[45]: [matplotlib.lines.Line2D object at 0x7f870766c820>]
```

```
In [46]: # Now plot the solutions obtained using 3 different approaches on the same plot

plt.rcParams.update({'font.size': 16})
figure, axis = plt.subplots(1, 1, figsize=(10, 8))
plt.subplots_adjust(hspace=1)
axis.set_title('Diffusion')
axis.plot(xspace, u_FD, label=r'$u_{FD}$', linewidth="3")
axis.plot(xspace, u_iter_conv, label=r'$u_{multiple \ ; \ convolutions}$', marker='o', linestyle="", markevery=1)
axis.plot(xspace, u_single_conv, label=r'$u_{convolution}$', marker='x', linestyle="", markevery=1)
axis.set_xlabel('x')
axis.set_ylabel('Concentration')
axis.legend(loc="upper right")
```

```
Out[46]: <matplotlib.legend.Legend object at 0x7f87071f7100>
```

Analytical solution: Advection - Diffusion of a Gaussian Hill

In case of an isotropic diffusion, the analytical solution describing evolution of a Gaussian Hill can be expressed as

$$C(\boldsymbol{x}, t) = \frac{(2\pi\sigma_0^2)^{D/2}}{(2\pi(\sigma_0^2 + 2kt))^{D/2}} C_0 \exp\left(-\frac{(\boldsymbol{x} - \boldsymbol{x}_0 - \boldsymbol{u}t)^2}{2(\sigma_0^2 + 2kt)}\right)$$

where:

- C_0 - initial concentration,
- D - number of dimensions,
- t - time,
- k - conductivity,
- \boldsymbol{u} - velocity of advection
- σ_0 the initial variance of the distribution.

Task

1) Implement the `GaussianHillAnal` class. It shall have a method `get_concentration_ND(self, X, t)`, which will return the concentration at given time and space.

2) Benchmark the FD code against analytical solution.

```
In [47]: from sympy.matrices import Matrix
import sympy as sp

class GaussianHillAnal:
def __init__(self, C0, X0, Sigma2_0, k, U, D):
"""
:param C0: initial concentration
:param X0: initial position of the hill's centre = Matrix([x0, y0])
:param U: velocity = Matrix([ux, uy])
:param Sigma2_0: initial width of the Gaussian Hill
:param k: conductivity
:param dimenions: number of dimensions
"""
self.C0 = C0
self.X0 = X0
self.U = U
self.Sigma2_0 = Sigma2_0
self.k = k
self.dim = D

def get_concentration_ND(self, X, t):
decay = 2.*self.k*t
L = X - self.X0 - self.U*t
C = self.C0
C *= pow(2. * np.pi * self.Sigma2_0, self.dim / 2.)
C /= pow(2. * np.pi * (self.Sigma2_0 + decay), self.dim / 2.)
C *= sp.exp(-(L.dot(L)) / (2.*(self.Sigma2_0 + decay)))
return C
```

```
In [48]: time_0 = dt*nt/2 # initial contidion for FD
time_spot = dt*nt # time to be simulated (by FD and analytically)

X0 = Matrix([domain_length/2.]) # center of the hill
C0 = 1. # concentration
variance = 30 # initial variance
reference_level = 0

T_0 = np.zeros(nx)
T_anal = np.zeros(nx)

gha = GaussianHillAnal(C0, X0, variance, nu, Matrix([0]), D=1)

for i in range(nx):
T_0[i] = reference_level + gha.get_concentration_ND(Matrix([xspace[i]]), time_0)
T_anal[i] = reference_level + gha.get_concentration_ND(Matrix([xspace[i]]), time_spot)

T_FD = calc_diffusion_FD_btcs(T_0,nx,nt,nu,dt)
T_single_conv, fs = calc_diffusion_single_convolution(T_0,xspace,nt,nu,dt)

plt.rcParams.update({'font.size': 16})
figure, axis = plt.subplots(1, 1, figsize=(8, 6))
plt.subplots_adjust(hspace=1)
axis.set_title('Diffusion of a Gaussian Hill')
axis.plot(xspace, T_0, label=r'$T_{0}$')
axis.plot(xspace, T_anal, label=r'$T_{anal}$')
axis.plot(xspace, T_FD, label=r'$T_{FD}$', marker='x', linestyle="", markevery=1)
axis.plot(xspace, T_single_conv, label=r'$T_{conv}$', marker='v', linestyle="", markevery=1)
axis.set_xlabel('x')
axis.set_ylabel('Concentration')
axis.legend(loc="upper right")
```

```
Out[48]: <matplotlib.legend.Legend object at 0x7f8706acec70>
```

Questions:

- How do you find the FD solution compared to analytical one? Experiment with different dx, dt.
- How would you asses that your mesh is fine enough in a real CFD simulation (without analytical solution)?

Answers

- Mesh convergence study

Learn More

Inspiration

<http://www.cs.umd.edu/~djacobs/CMSC828seg/Diffusion.pdf>

<https://web.math.ucsb.edu/~helena/teaching/math124b/heat.pdf>

Some of the text has been inspired by the **12 steps to Navier–Stokes**, the practical module taught in the interactive CFD class of [Prof. Lorena Barba](#), (provided under a Creative Commons Attribution license, CC-BY. All code is made available under the FSF-approved BSD-3 license. (c) Lorena A. Barba, Gilbert F. Forsyth 2017. Thanks to NSF for support via CAREER award #1149784. [@LorenaABarba](#))

For a careful walk-through of the discretization of the diffusion equation with finite differences (and all steps from 1 to 4), watch **Video Lesson 4** by Prof. Barba on YouTube.

```
In [49]: from IPython.display import YouTubeVideo
YouTubeVideo('y2WaK7_iMRI')
```

Out[49]: 