# Analytic computations of an effective lattice theory for heavy quarks

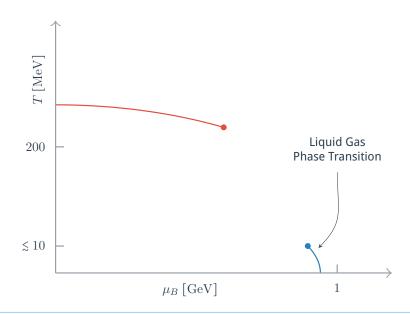


Jonas R. Glesaaen Mathias Neuman, Owe Philipsen Lattice Conference 2015 - July 16th 1 The Effective Theory

2 Results

3 Conclusion

## Heavy QCD Phase Diagram



## Advantages of the Effective Theory

- Dimensionally reduced theory
  - $4D \rightarrow 3D$
  - $U_{\mu}(x) \rightarrow L(x)$
- Very mild sign problem, most gauge fields integrated analytically
- Want to study the very dense limit, liquid gas transition

## The Effective Theory

## The Effective Lattice Theory

#### **Effective Theory**

• Integrate out all spatial gauge links

$$egin{aligned} \mathcal{Z} &= \int extit{D} U_{\mu} \, ext{exp} \left\{ - extit{S}_{ ext{action}} 
ight\} \ &= \int extit{D} U_{0} \, ext{exp} \left\{ - extit{S}_{ ext{effective action}} 
ight\} \end{aligned}$$

#### Using:

- The strong coupling expansion
- The hopping parameter expansion

#### **Effective Theory**

$$\mathcal{Z} = \int \prod_{x} dL(x) \exp \{-S_{\text{eff action}}\} \qquad (\dagger)$$

- Previous Talk: Monte Carlo simulations of (†)
- ullet Current Talk: Analytic calculation of  ${\mathcal Z}$

## The Effective Theory Action

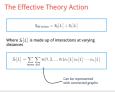
$$S_{\rm eff\ action} = S_0[L] + S_I[L]$$

Where  $S_I[L]$  is made up of interactions at varying distances

$$S_{I}[L] = \sum_{\text{terms dof}} \sum_{\text{dof}} v_{i}(1, 2, ..., n_{i}) \phi_{1}[L] \phi_{2}[L] \cdots \phi_{n_{i}}[L]$$

Can be represented

with connected graphs



Remember to say that the effective couplings  $V_i$  are themselves functions of  $\kappa$  and  $\beta$ , and we can therefore carry out a consistent expansion.

Talk about the expansion point. What is  $\emph{S}_{0}$  physically? Free hadron gas.

## The Effective Theory Action

$$S_{I}[L] = \sum_{\text{terms dof}} \sum_{\text{dof}} v_{i}(1, 2, ..., n_{i}) \phi_{1}[L] \phi_{2}[L] \cdots \phi_{n_{i}}[L]$$

#### In our theory:

- $v_i(1,2,...n_i) \rightarrow \{\lambda_i,h_i\} \times \text{geometry}$
- $\bullet \ \phi_i \to \{L_i, L_i^*, W_i\}$

## **Analytic Calculations**

N-point Linked Cluster Expansion

#### **Classical Linked Cluster Expansion**

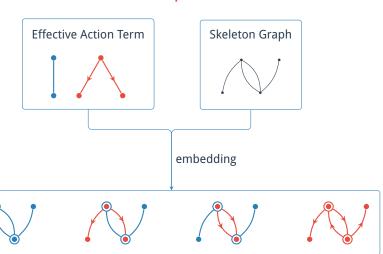
The action consists of two-point interactions which can be expanded in a set of connected graphs.

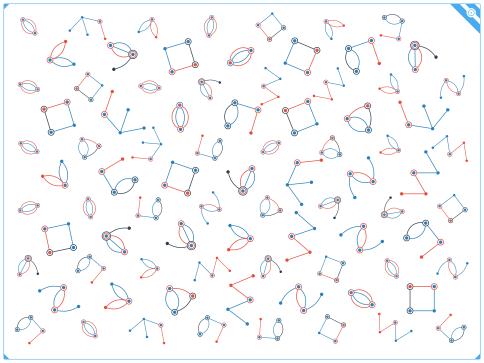
#### **Our Problem**

The action contains *n*-point interactions that we can embed on a set of connected graphs.

## **Analytic Calculations**

#### N-point Linked Cluster Expansion





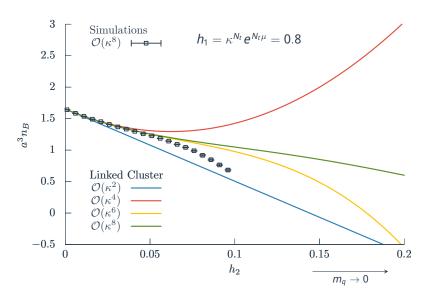
#### The power of resummations

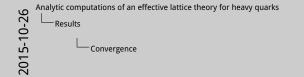
Using the resummed Linked Cluster Expansion as motivation

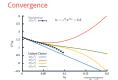
We can do the same resummation for the effective action itself, incorporating long-range effects



### Convergence

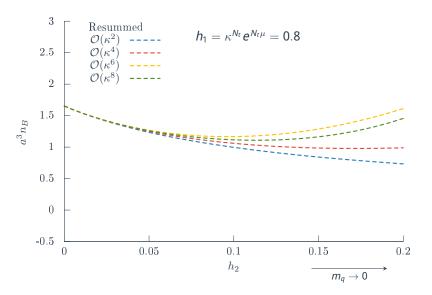




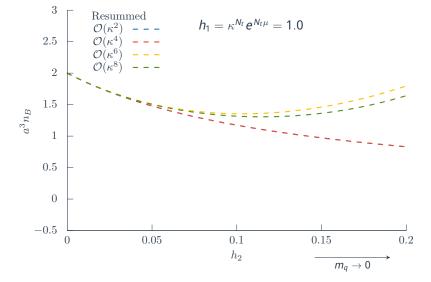


Say something about the fact that higher  $h_2$  means smaller quark masses.

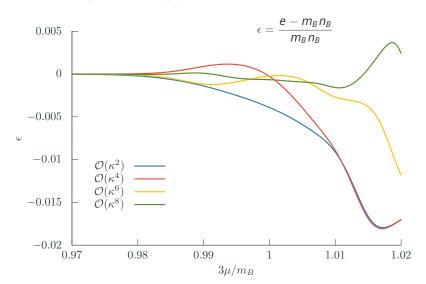
#### Effect of the resummations

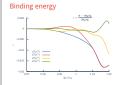


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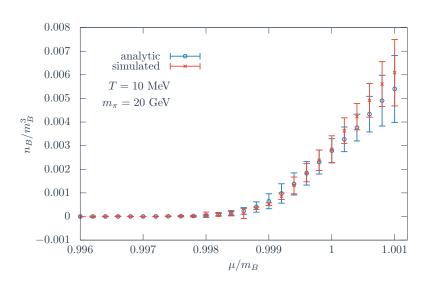
## Binding energy



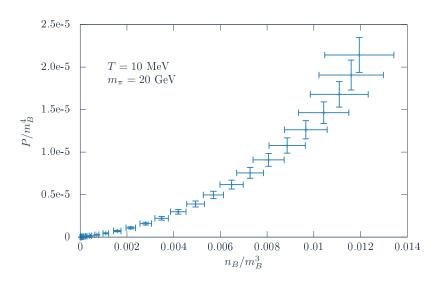


Plot parameters:  $\kappa =$  0.08,  $N_t =$  50,  $\beta =$  0

#### Continuum comparison



### Continuum Equation of State





## **Summary & Outlook**

#### **Summary**

- Introduced the effective dimensionally reduced lattice theory
- Looked at how a consistent analytic calculation could be carried out
- Demonstrated convergence and comparisons with numerics

## Summary & Outlook

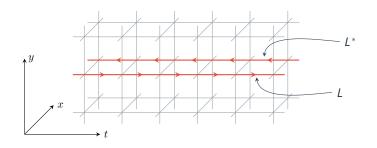
#### Outlook

- Use the analytic results as a tool to study the characteristics of the effective theory
- Find analytic resummation schemes to incorporate long-range effects



# Backup slides

## The Effective Lattice Theory Pure gluon contributions



What remains is an interaction between Polyakov Loops

As mentioned earlier, we are only interested in quantities that contribute to the thermodynamic of the system. For our system, that means quantities which span the full temporal direction, and thus picks up a temperature dependent component in the infinite volume limit.

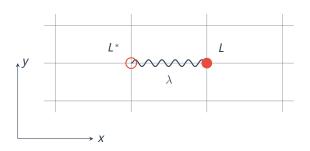
For pure gauge action, in the strong coupling expansion, this means a plane of plaquettes spanning the full temporal direction, as shown in the figure. We then integrate out the spatial links present in this strip of plaquettes, and are left with the Polyakov loops. The final result will thus be a nearest neighbour interaction between two Polyakov loops, or a continuos spin-system on a three dimensional lattice.

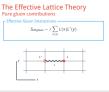
#### The Effective Lattice Theory

Pure gluon contributions

**Effective Gluon Interactions** 

$$S_{\mathrm{eff gluon}} \sim \lambda \sum_{(x,y)} L(x) L^*(y)$$



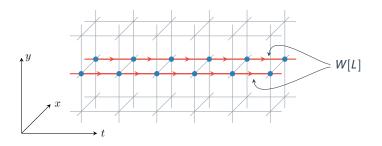


As mentioned earlier, we are only interested in quantities that contribute to the thermodynamic of the system. For our system, that means quantities which span the full temporal direction, and thus picks up a temperature dependent component in the infinite volume limit.

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#### The Effective Lattice Theory

Pure quark contributions



Producing an interaction between the W objects

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Analytic computations of an effective lattice theory for heavy quarks

Backup slides

The Effective Lattice Theory
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For the fermions we are in very much the same situation. We need a quantity that spans the temporal direction of the lattice. The simplest such object is of course a single quark line, exclusively jumping in the temporal direction, going around the lattice. This is the contribution from static quarks, and this is of course included.

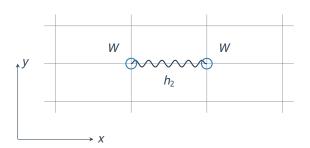
The next order term would be a quark that loops, jumps to a neighbouring site, loops, then jumps back, which is the term depicted in the figure. Here again we integrate out the spatial links, and are left with the interaction of two *W*-terms, the mathematical structure of which is not important for now.

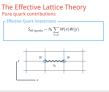
## The Effective Lattice Theory

Pure quark contributions

**Effective Quark Interactions** 

$$S_{
m eff\,quarks} \sim h_2 \sum_{\langle x,y \rangle} W(x) W(y)$$



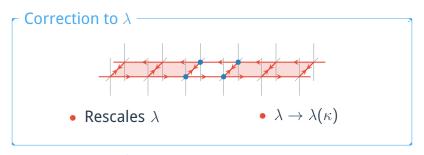


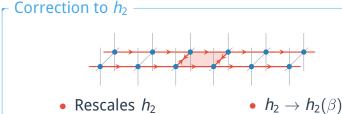
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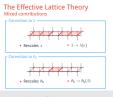
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## The Effective Lattice Theory

#### Mixed contributions







One can of course mix the terms from the two different expansions we are carrying out. Two examples of the possible terms can be seen in the figures. For these two simple cases, the mixing will only contribute as a shift in the nearest neighbour couplings, and can thus be absorbed by those.

There are higher order mixed terms that create entierly new interactions, but those are of much higher order of what we have shown here.

I will probably skip this slide.

#### EoS in lattice units

