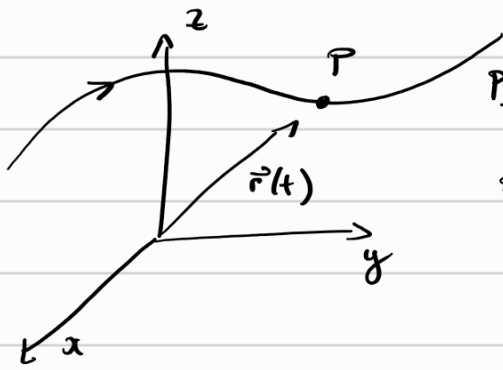


## Space curves



$$P = \vec{r}(t) = x(t)\vec{i} + y(t)\vec{j} + z(t)\vec{k}$$

For example

$$\vec{r}(t) = \cos(\omega t)\vec{i} + \sin(\omega t)\vec{j} + at\vec{k}$$

## Parameterisation of curve

Let  $u$  be a parameter

need to find expressions for  $x(u)$ ,  $y(u)$  &  $z(u)$

(4) Parametrise  $y = 3x + 4$

$$\vec{r}(u) = x(u)\vec{i} + y(u)\vec{j} + z(u)\vec{k}$$

(in general)



$$\vec{r}(u) = u\vec{i} + (3u + 4)\vec{j} + 0\vec{k}$$

Let  $v = 3x$  then  $y = 3x + 4$

$$\vec{r}(v) = \frac{v}{3}\vec{i} + (v + 4)\vec{j}$$

NO ONE SINGLE WAY TO PARAMETRISE

Express  $y = 3x^2$  as  $\vec{r}(t) = x(t)\vec{i} + y(t)\vec{j} + z(t)\vec{k}$

$$\vec{r}(t) = x\vec{i} + (3x^2)\vec{j} + z\vec{k}$$

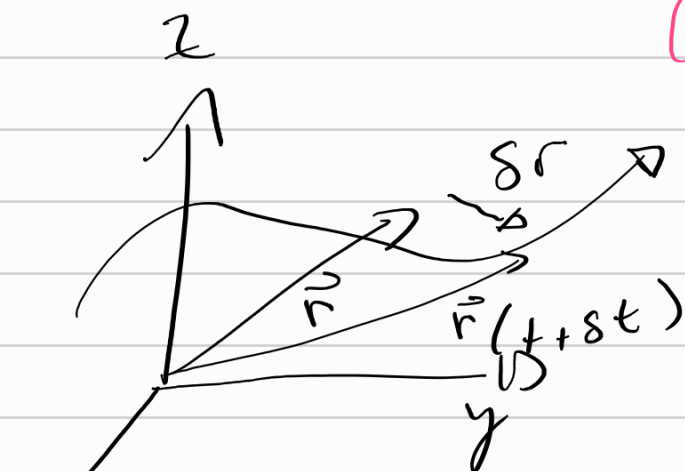
$$v = \sqrt{3}$$

$$u = 2x^2$$

$$\vec{r}(t) = \frac{v}{\sqrt{3}} \vec{i} + (v^2) \vec{j}$$

— fuch missed it, (correct)  
 ched video

pruned ~ 10 min  
 17:15



$$\begin{aligned} \delta \vec{r} &= \vec{r}(t + \delta t) - \vec{r}(t) \\ &\approx \frac{d\vec{r}}{dt} \delta t \end{aligned}$$

In the lim  $\delta t \rightarrow 0$ , the approximat<sup>o</sup> becomes exact

Then  $\delta \vec{r} \rightarrow //$  to the curve at P

(where P is the point with posit<sup>o</sup> vector  $\vec{r}(t)$ )

$\delta \vec{r}$  becomes a tangent to the space curve in the limit  $\delta t \rightarrow 0$

$$\frac{\frac{d\vec{r}}{dt}}{\left| \frac{d\vec{r}}{dt} \right|} = \hat{T} \text{ unit tangent vectors}$$

$$\hat{G} = \frac{\vec{G}}{|\vec{G}|}$$

From last time we know that

$$\frac{d\vec{r}}{dt} = \vec{v}(t) \quad \& \quad \left| \frac{d\vec{r}}{dt} \right| = |\vec{v}| = \text{speed}$$

$$\hat{T} = \frac{\frac{d\vec{r}}{dt}}{\left| \frac{d\vec{r}}{dt} \right|} = \frac{\vec{v}}{|\vec{v}|} = \hat{v}$$

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Let  $S$  be distance,

consider  $\vec{r}(s)$  instead of  $\vec{r}(t)$

$$\hat{T} = \frac{d\vec{r}}{ds}$$

$$\text{But } \underbrace{\frac{d\vec{r}}{dt}} = \frac{d\vec{r}}{ds} \frac{ds}{dt} = \hat{T} \frac{ds}{dt}$$

$$\& \text{ speed} = |\vec{v}| = \left| \frac{d\vec{r}}{dt} \right| = \frac{ds}{dt}$$

ex

Find  $\hat{T}$  at the point  $(0, a)$  on the space curve  $\vec{r}(t) = a \cos(\omega t) \hat{i} + a \sin(\omega t) \hat{j}$

$$\hat{T} = \frac{\frac{d\vec{r}}{dt}}{\left| \frac{d\vec{r}}{dt} \right|}$$

$$\frac{d\vec{r}}{dt} = -aw \sin(\omega t) \vec{i} + aw \cos(\omega t) \vec{j}$$

$$\left| \frac{d\vec{r}}{dt} \right| = \left( (aw)^2 \underbrace{(-\sin^2(\omega t) + \cos^2(\omega t))}_{=1} \right)^{1/2}$$

$$= aw$$

so  $\hat{T} = -\sin(\omega t) \vec{i} + \cos(\omega t) \vec{j}$

at  $(0, a) = \vec{r} = a \cos(\omega t) \vec{i} + a \sin(\omega t) \vec{j}$

$$= 0 \vec{i} + a \vec{j}$$

$$\Rightarrow \left\{ \begin{array}{l} \cos(\omega t) = 0 \\ \sin(\omega t) = 1 \end{array} \right\} \Rightarrow \omega t = \pi/2$$

so  $\hat{T} \Big|_{(0, a)} = -\sin(\pi/2) \vec{i} + \cos(\pi/2) \vec{j}$

$$= -\vec{i}$$

Newton's 2<sup>nd</sup> law in vector form

$$F = ma \quad F = m \frac{d^2 r}{dt^2}$$

$$\vec{F} = m\vec{a} \quad \vec{F} = \frac{m d^2 \vec{r}}{dt^2}$$

$$\vec{r}(t) = x(t)\vec{i} + y(t)\vec{j} + z(t)\vec{k}$$

**Lorentz Force** (force particle in electromagnetic field)

$$\vec{F} = q (\vec{E} + \vec{v} \times \vec{B})$$

$$\vec{F} = \frac{m d^2 \vec{r}}{dt^2} = q (\vec{E} + \vec{v} \times \vec{B})$$

$$\vec{F} = q\vec{E} + q \left( \frac{d\vec{r}}{dt} \times \vec{B} \right)$$

$$\vec{F} = F_x \vec{i} + F_y \vec{j} + F_z \vec{k}$$

$$F_x = m \frac{d^2 x}{dt^2}$$

$$F_y = m \frac{d^2 y}{dt^2}$$

$$\vec{F} = m \frac{d^2 \vec{r}}{dt^2}$$

Problem The force on a particle of charge  $q$  moving in a uniform magnetic field  $\vec{B}$  is given by

$$\vec{F} = q(\vec{v} \times \vec{B})$$

Let  $m$  be the mass of the particle  
 so  $m\vec{a} = q(\vec{v} \times \vec{B})$  so  $\vec{a} = \frac{q}{m}(\vec{v} \times \vec{B})$

Explain why, the acceleration of the particle is always perpendicular to the  $\vec{B}$  field.

