

**University of Bath**  
**Department of Physics**  
**PH30030**  
**Quantum Mechanics**

**19<sup>th</sup> January 2023**  
**13:30 – 16:30**  
**180 minutes**

**Answer all questions**

Only calculators provided by the University may be used

University Formula Book with fundamental Constant tables provided by the University may be used.

**During this exam you are not permitted to communicate with any person(s) except an invigilator or an assigned support worker.**

**You must not have any unauthorised devices or materials with you.**

**You must keep your Library card on your desk at all times.**

**Please fill in the details on the front of your answer book/cover and sign in the section on the right of your answer book/cover, peel away adhesive strip and seal.**

**Take care to enter the correct candidate number as detailed on your desk label.**

**Do not turn over your question paper until instructed to by the chief invigilator.**

1. (a) Give the definition of a Hermitian operator. (1)  
 (b) Prove that the eigenvalues of Hermitian operators are real. (4)
  
2. A particle moving along the  $x$  axis in a region of space where the potential energy is zero everywhere is described by the wavefunction  $\psi(x) = A \exp(ikx)$  where  $A$  is a normalisation constant.  
 (a) Find the momentum and total energy eigenvalues. Show that these values are consistent with the de Broglie relation.  $\hbar k$ ;  $\hbar^2 k^2 / 2m$  (3)  
 (b) By confining the particle to a box of finite size, find the normalisation constant  $A$ . (1)
  
3. Consider the observable associated with the operator  $\hat{Q}$ . Show that the observables associated with the operators  $\hat{Q}$ ,  $\hat{Q}^2$ ,  $\hat{Q}^3$  etc., are all compatible. Hence show that the linear momentum of a particle in one-dimension can always be measured compatibly with the non-relativistic kinetic energy. (4)
  
4. For a spin-half system, the spin operator  $\hat{S} = (\hat{S}_x, \hat{S}_y, \hat{S}_z)$ , where  $\hat{S}_i = \frac{\hbar}{2} \sigma_i$  and the Pauli spin matrices are given by  $\sigma_x = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$ ,  $\sigma_y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$ ,  $\sigma_z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$ .  
 The eigenvectors of  $\hat{S}_z$  are  $\begin{pmatrix} 1 \\ 0 \end{pmatrix}$  and  $\begin{pmatrix} 0 \\ 1 \end{pmatrix}$  with eigenvalues of  $\frac{\hbar}{2}$  and  $-\frac{\hbar}{2}$ , respectively.  
 (a) For the orbital angular momentum, the ladder operators are defined by  $\hat{L}_+ = \hat{L}_x + i\hat{L}_y$  and  $\hat{L}_- = \hat{L}_x - i\hat{L}_y$ .  
 (i) Using this formalism, obtain matrices representing the raising and lowering operators for a spin-half system.  $\hat{S}_+ = \hbar \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix}$ ;  $\hat{S}_- = \hbar \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix}$  (2)  
 (ii) By applying these operators to the eigenstates of  $\hat{S}_z$ , verify that they act as ladder operators. (4)  
 (iii) Verify that both eigenstates of  $\hat{S}_z$  are also eigenstates of  $\hat{S}^2$  and correspond to the same eigenvalue of this operator. (3)

- (b) Calculate the expectation values  $\langle S_x^2 \rangle$  and  $\langle S_y^2 \rangle$  for a spin-half particle known to be in an eigenstate of  $\hat{S}_z$ . Show that their product is consistent with the uncertainty principle, which states that

$$\langle S_x^2 \rangle \langle S_y^2 \rangle \geq \frac{1}{4} \hbar^2 \langle S_z^2 \rangle. \quad (4)$$

5. Consider two non-interacting spin-half fermions in a one-dimensional infinite square well potential.

- (a) Describe the symmetry condition for the two-particle wavefunction and how this wavefunction can be built from spatial- and spin-dependent parts. (2)

- (b) The orthogonal and normalised one-particle states are given by  $u_n(x_i)$  and the spin states are given by  $\alpha_i$  and  $\beta_i$ , with  $i = 1, 2$  numbering the particles. Find the ground state wavefunction if the spins of the fermions are

(i) parallel  $\phi(1, 2) = u_-(1, 2) \chi_{1, 1/2}(1, 2)$  (3)

(ii) anti-parallel  $\phi(1, 2) = u_+(1, 2) \chi_{0, 0}(1, 2)$  (3)

Note that: -

$$\chi_{00}(1, 2) = \frac{1}{\sqrt{2}}(\alpha_1\beta_2 - \alpha_2\beta_1),$$

$$\chi_{10}(1, 2) = \frac{1}{\sqrt{2}}(\alpha_1\beta_2 + \alpha_2\beta_1), \quad \chi_{11}(1, 2) = \alpha_1\alpha_2, \quad \chi_{1-1}(1, 2) = \beta_1\beta_2$$

- (c) For both the parallel and anti-parallel spin states, describe what happens to the ground state wavefunction if both particles try to occupy the same region of space.

(4)

6. The potential energy of a one-dimensional harmonic oscillator at position  $x$  about the mean  $\langle x \rangle$  is given by

$$V(x) = \frac{1}{2}k(x - \langle x \rangle)^2$$

where  $k = m\omega^2$  is the spring constant.

- (a) Explain why energy quantisation is not observed for a macroscopic system. (2)

- (b) For a quantum system: -

- (i) Briefly explain why  $\langle \hat{p}_x \rangle = 0$ . (1)

- (ii) Use the uncertainty relation to place a lower bound on the mean kinetic energy of the oscillator in terms of  $\Delta x$ .  $\langle \hat{T} \rangle \geq \hbar^2 / 8m (\Delta x)^2$  (5)

- (iii) Hence, derive the lower bound on the total energy of the oscillator in terms of  $\Delta x$ .  $\langle \hat{H} \rangle_{\min} = \hbar^2 / 8m (\Delta x)^2 + m\omega^2 (\Delta x)^2 / 2$  (2)

- (iv) Use this expression to estimate the zero-point energy of the harmonic oscillator.  $\hbar\omega/2$  (4)

- (c) Discuss briefly how time-independent perturbation theory is used in the approximate evaluation of the ground-state energies of non-degenerate systems, for which the exact quantum mechanical solutions cannot be found. (4)

- (d) Consider an *anharmonic* oscillator with mean position  $\langle x \rangle = 0$  described by the Hamiltonian

$$\hat{H} = \frac{\hat{p}_x^2}{2m} + \frac{1}{2}kx^2(1 + \alpha x^2)$$

where  $\alpha$  is a constant. Derive an approximate (first-order) expression for the ground state energy of this oscillator, given that the wave function for the harmonic oscillator is  $\phi_0(x) = (m\omega/\pi\hbar)^{1/4} \exp(-m\omega x^2/2\hbar)$ . (4)

$$\text{Note: } \int_{-\infty}^{\infty} x^4 e^{-bx^2} dx = \frac{3}{4} \left[ \frac{\pi}{b^5} \right]^{1/2} \quad E_0 = \frac{\hbar\omega}{2} + \frac{3}{8} \alpha \frac{\hbar^2}{m}$$

**THE END**

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