Here, we show that the subdiagrams (a) to (c) on p.6 vanish (at least (a) and (b) are properly investigated here):

- subdiagram (a):

$$
\begin{aligned}
h_v(0)\bar{h}_v(x) &= \Theta(-v\cdot x)\delta^{(d-1)}(x_\perp)P_+ ig_s \int_{v\cdot x} \mathrm{d}s v^\mu A_\mu(sv) \\
&= \Theta(-v\cdot x)\delta^{(d-1)}(x_\perp)P_+ ig_s \int_{v\cdot x} \mathrm{d}s\ v^\mu \int_0^1 \mathrm{d}u\ usv^\nu G_{\nu\mu}(usv) \\
&= 0 \qquad (1)
\end{aligned}
$$

  since

$$
v^\mu v^\nu G_{\nu\mu}(...) = -v^\mu v^\nu G_{\mu\nu}(...) = -v^\nu v^\mu G_{\mu\nu}(...) = -v^\mu v^\nu G_{\nu\mu}(...) \qquad (2)
$$

  Generally, a product of an antisymmetric and a symmetric tensor leads 0. We immediately started this investigation at $\mathcal{O}(g_s)$, because an emission of a background gluon field is an NLO effect.

- subdiagram (b):

  Here, the heavy quark propagator is again expanded to $\mathcal{O}(g_s)$ and the appearing backgound field gluon is contracted with the field strength tensor from the three-body current.

$$
\begin{aligned}
&\Theta(-v\cdot x)\delta^{(d-1)}(x_\perp)P_+ \langle 0|\, T\{\bar{q}\Gamma_1 g_s G_{\mu\nu} ig_s \int_{v\cdot x} \mathrm{d}s\ v^\lambda A_\lambda(sv)\Gamma_2 q\}\,|0\rangle \\
&= \Theta(-v\cdot x)\delta^{(d-1)}(x_\perp)P_+ \langle 0|\, T\{\bar{q}\Gamma_1\Gamma_2 q\}\,|0\rangle\, ig_s^2 \int_{v\cdot x} \mathrm{d}s\ \frac{\Gamma\left(\frac{d}{2}\right)\delta^{ab}}{2\pi^{\frac{d}{2}}(-s^2v^2+i0^+)^{\frac{d}{2}}} s(g_{\nu\lambda}v_\mu - g_{\mu\lambda}v_\nu)v^\lambda \\
&= 0 \qquad (3)
\end{aligned}
$$

  since

$$
(g_{\nu\lambda}v_\mu - g_{\mu\lambda}v_\nu)v^\lambda = v_\mu v_\nu - v_\nu v_\mu = 0 \qquad (4)
$$

  The only missing piece seems to be subdiagram (c). The proof might be as simple as to take equation (30) and set z=0, because the gluon line is connected to the same on-shell quark line.