Here, we would like to prove eq. (25). After contracting the heavy-quark propagator, we start at the following point:

$$
\begin{aligned}
&\Theta(-v\cdot x)\delta^{(3)}(x_\perp)\langle 0|\,\bar q(0)\Gamma_1 P_+\Gamma_2 q(x)\,|0\rangle \\
=\,&\Theta(-v\cdot x)\delta^{(3)}(x_\perp)\Big[\langle 0|\,\bar q(0)\Gamma_1 P_+\Gamma_2 q(0)\,|0\rangle + x^\mu\langle 0|\,\bar q(0)\Gamma_1 P_+\Gamma_2 D_\mu q(0)\,|0\rangle + \\
&\frac{1}{2}x^\mu x^\nu\langle 0|\,\bar q(0)\Gamma_1 P_+\Gamma_2 D_\mu D_\nu q(0)\,|0\rangle + ...\Big]
\end{aligned}
\tag{1}
$$

Now we start by investigating the first term in the sum:

$$\langle 0|\,\bar q^i_\alpha(0)\Gamma_{1,\alpha\beta}P_{+,\beta\gamma}\Gamma_{2,\gamma\delta}q^j_\delta(0)\,|0\rangle = A\delta^{ij}\delta_{\alpha\delta} \tag{2}$$

$$\Leftrightarrow \delta^{ij}\delta_{\alpha\delta}\langle 0|\,\bar q^i_\alpha(0)\Gamma_{1,\alpha\beta}P_{+,\beta\gamma}\Gamma_{2,\gamma\delta}q^j_\delta(0)\,|0\rangle = 4\cdot AN_c \tag{3}$$

Notice that we implicitly sum over the flavors i,j, hence N_c also appears in the RHS and cancels the one on the LHS. So we can solve for A:

$$A = \frac{1}{4}\underbrace{\Gamma_{1,\alpha\beta}P_{+,\beta\gamma}\Gamma_{2,\gamma\delta}\delta_{\alpha\delta}}_{=\,\mathrm{Tr}[\Gamma_1 P_+\Gamma_2]}\langle 0|\,\bar q q\,|0\rangle \tag{4}$$