

§ 11.2 杨氏双缝干涉



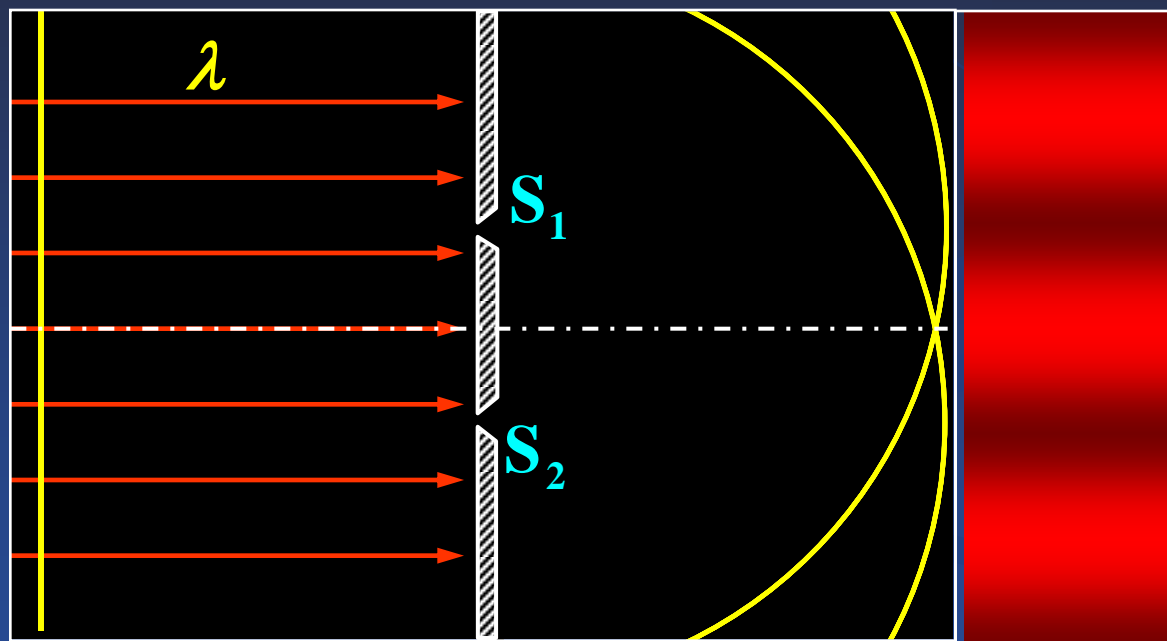


1773~1829

Thomas Young

在1801年首先用实验的方法研究了光的干涉现象，为光的波动理论确定了实验基础。

实验装置图：



一、观测屏上的光强分布

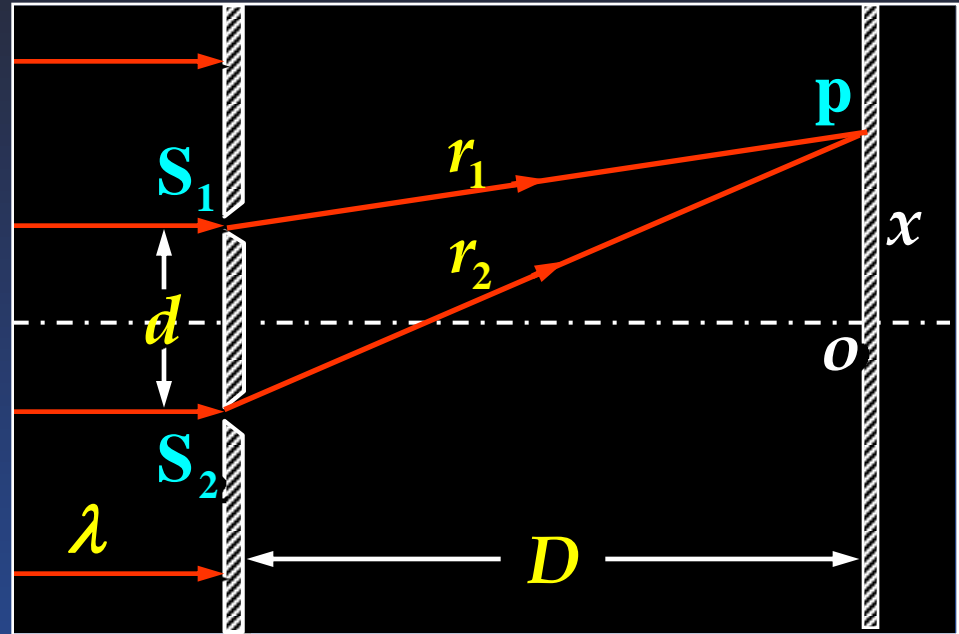
p点光振动方程:

$$E_{1p} = A_{10} \cos(\omega t - \frac{2\pi}{\lambda} r_1)$$

$$E_{2p} = A_{20} \cos(\omega t - \frac{2\pi}{\lambda} r_2)$$

$$E_p = A_p \cos(\omega t + \varphi)$$

$$A_p^2 = A_{10}^2 + A_{20}^2 + 2A_{10}A_{20} \cos(\Delta\phi) \quad \Delta\phi = \frac{2\pi}{\lambda} (r_2 - r_1)$$

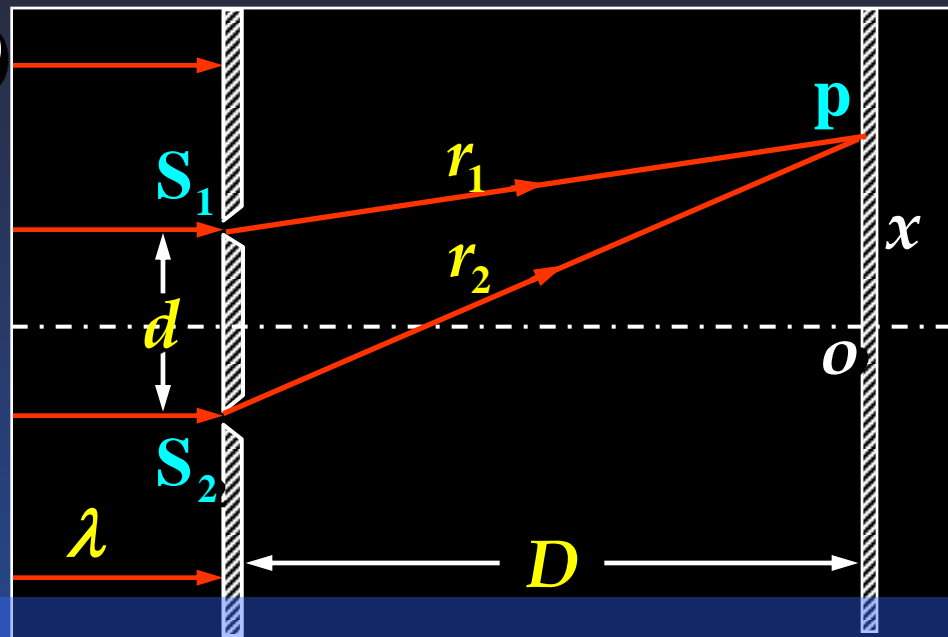


光强: $I_p \propto A_p^2$, $I_1 \propto A_{10}^2$, $I_2 \propto A_{20}^2$

$$I_p = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos(\Delta\phi)$$

设 $x \ll D$ (旁轴条件)

$$\longrightarrow I_1 \approx I_2 = I_0$$



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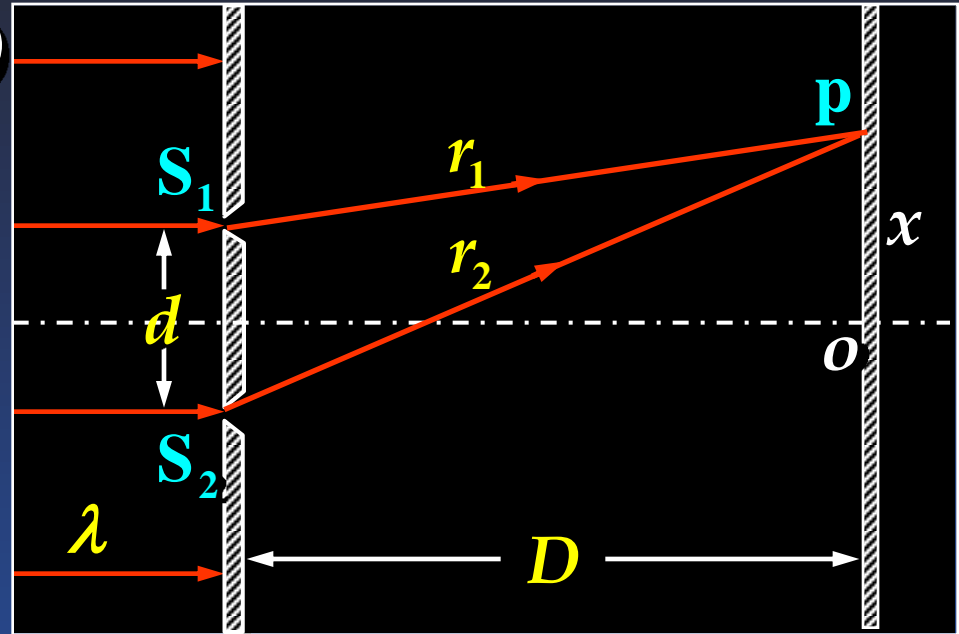
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$$I_p = 4I_0 \cos^2\left(\frac{\Delta\phi}{2}\right)$$



$$\Delta\phi = \frac{2\pi}{\lambda} (r_2 - r_1)$$

$$I_{p\min} = 0 \leq I_p \leq I_{p\max} = 4I_0$$

$\Delta\phi$ 与 p 点位置有关!

$$\Delta\phi = \begin{cases} \pm 2k\pi & I_p = I_{p\max} \quad \text{干涉加强} \rightarrow \text{明纹} \\ \pm (2k+1)\pi & I_p = I_{p\min} \quad \text{干涉减弱} \rightarrow \text{暗纹} \end{cases}$$

$$(k=0, 1, 2, \dots)$$

$$\Delta\phi = \frac{2\pi}{\lambda}(r_2 - r_1) = \frac{2\pi}{\lambda}\delta \quad (\delta \text{ 为光程差})$$

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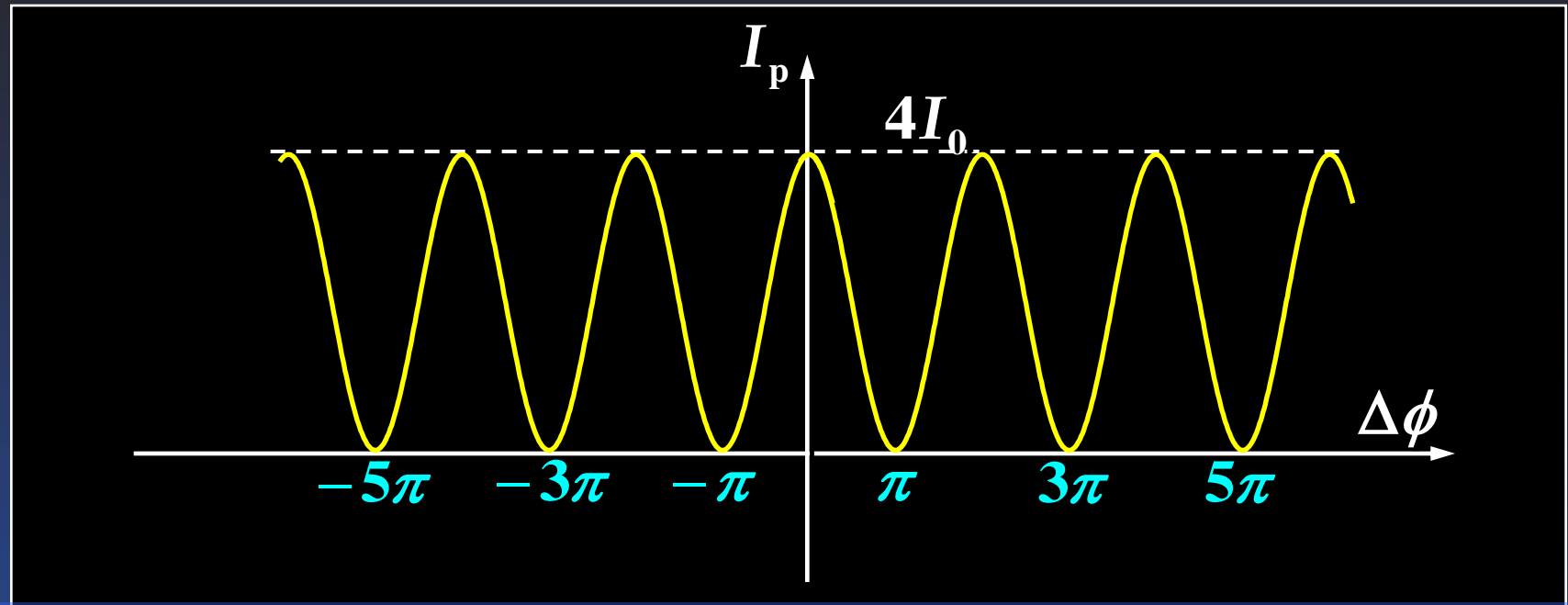
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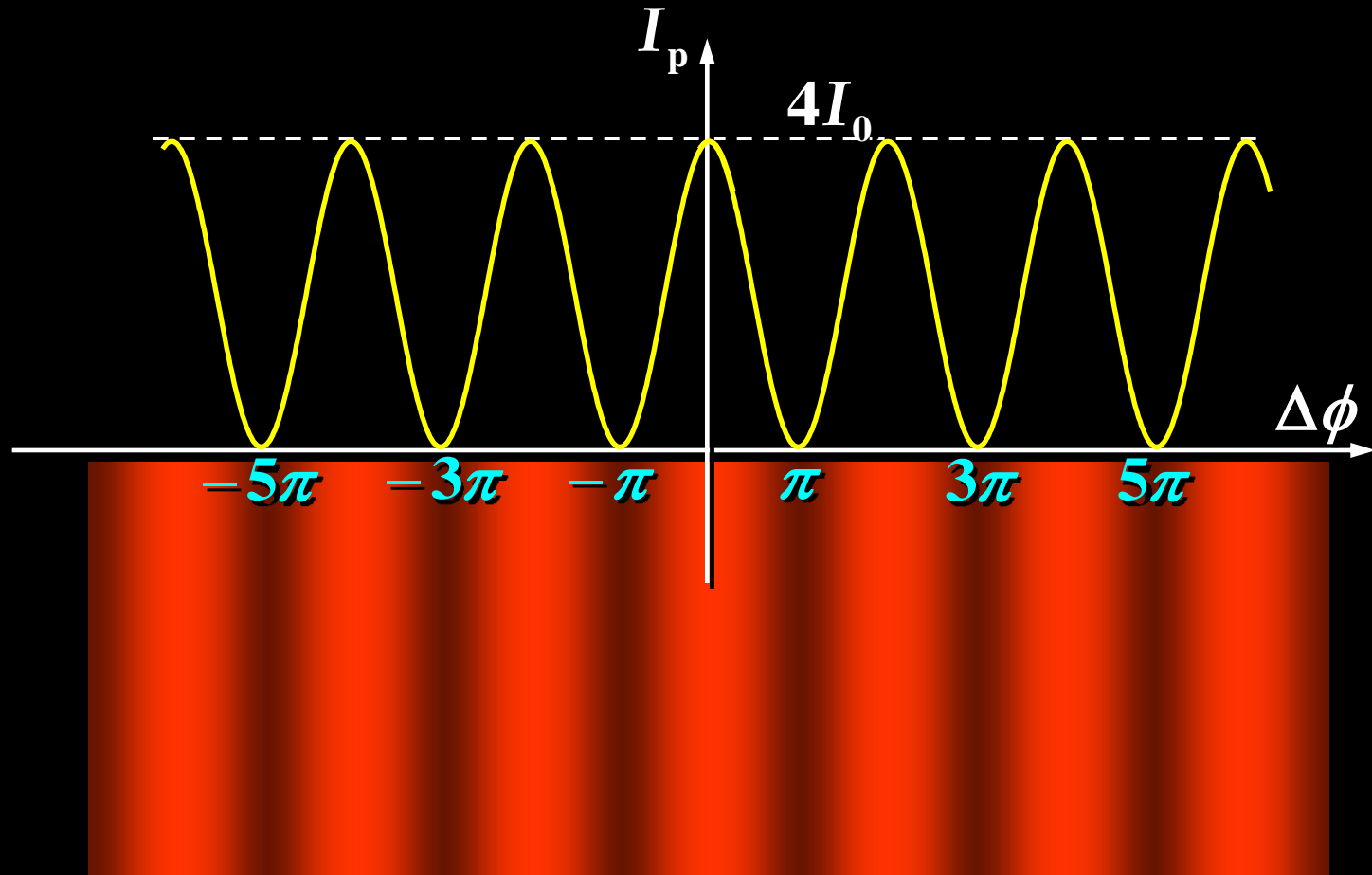
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观测屏上光强分布曲线:



$$\delta = \begin{cases} \pm 2k \frac{\lambda}{2} & I_p = I_{p\max} \quad \text{干涉加强} \rightarrow \text{明纹} \\ \pm (2k+1) \frac{\lambda}{2} & I_p = I_{p\min} \quad \text{干涉减弱} \rightarrow \text{暗纹} \end{cases}$$

观测屏上光强分布曲线:



干涉图样：平行等间距的直条纹！

二、干涉条纹位置分布

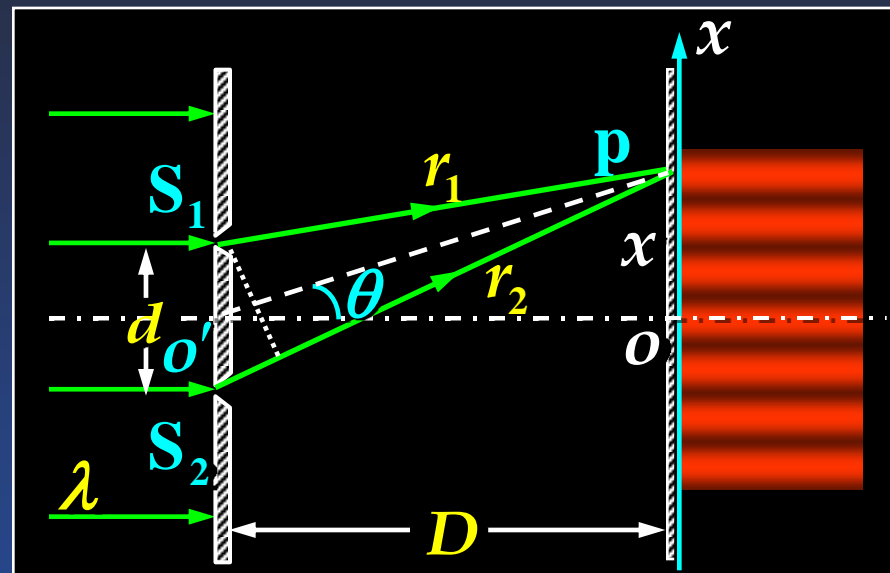
设入射光垂直入射, $d \ll D$, $x \ll D$

$$\delta = r_2 - r_1 \approx d \cdot \sin \theta$$

(θ 称为衍射角)

$$\sin \theta \approx \tan \theta = \frac{x}{D}$$

$$\delta = r_2 - r_1 \approx d \cdot \frac{x}{D} = \begin{cases} \pm 2k \frac{\lambda}{2} & \text{明纹} \\ \pm (2k + 1) \frac{\lambda}{2} & \text{暗纹} \end{cases}$$



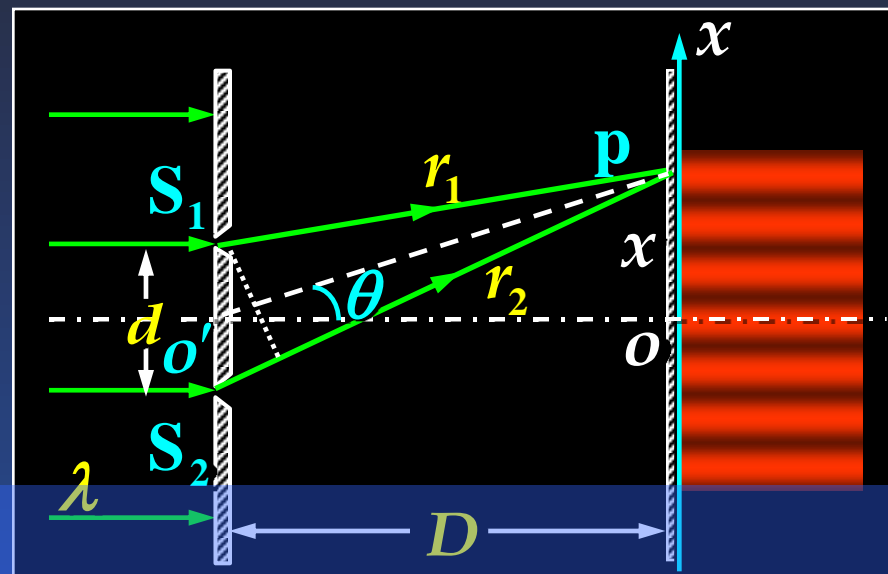
明纹位置: $x_k = \pm \frac{D}{d} \cdot 2k \cdot \frac{\lambda}{2} \quad (k=0, 1, 2, \dots)$

中央明纹: $k=0, \delta=0$

暗纹位置:

$$x_k = \pm \frac{D}{d} \cdot (2k+1) \cdot \frac{\lambda}{2}$$

$$\sin \theta \approx \tan \theta = \frac{x}{D}$$



明纹

暗纹

$$\delta = r_2 - r_1 \approx d \cdot \frac{x}{D} = \begin{cases} \pm 2k \frac{\lambda}{2} & \text{明纹} \\ \pm (2k+1) \frac{\lambda}{2} & \text{暗纹} \end{cases}$$

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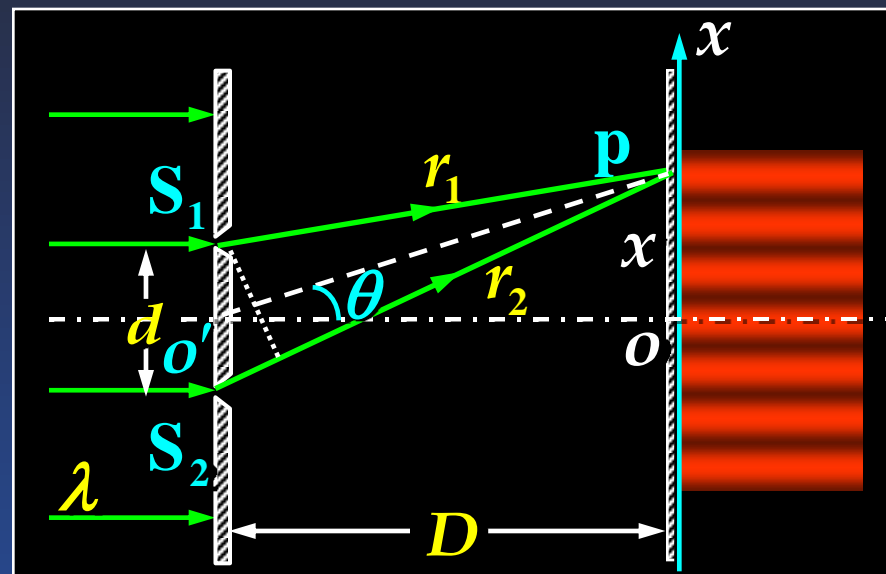
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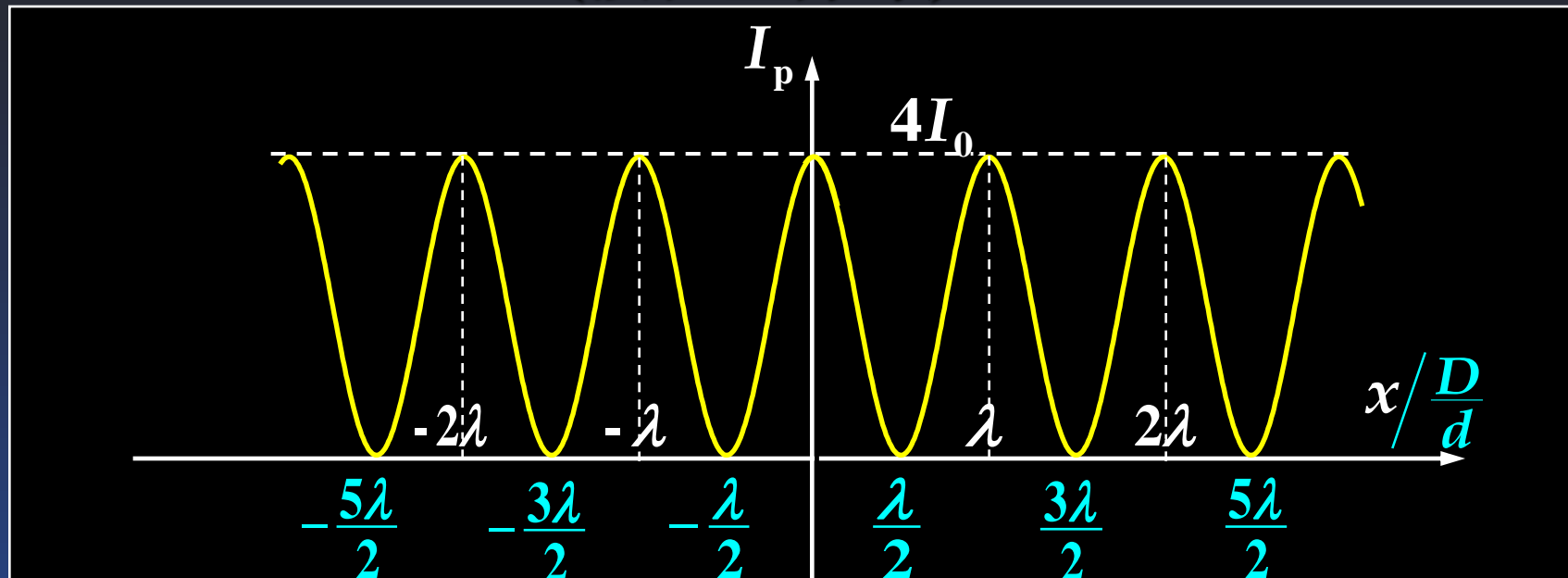
关于中央明纹对称!

相邻两条明(暗)纹间距:

$$\Delta x = |x_{k+1} - x_k| = \frac{D}{d} \lambda$$



观测屏上光强曲线(按位置分布):

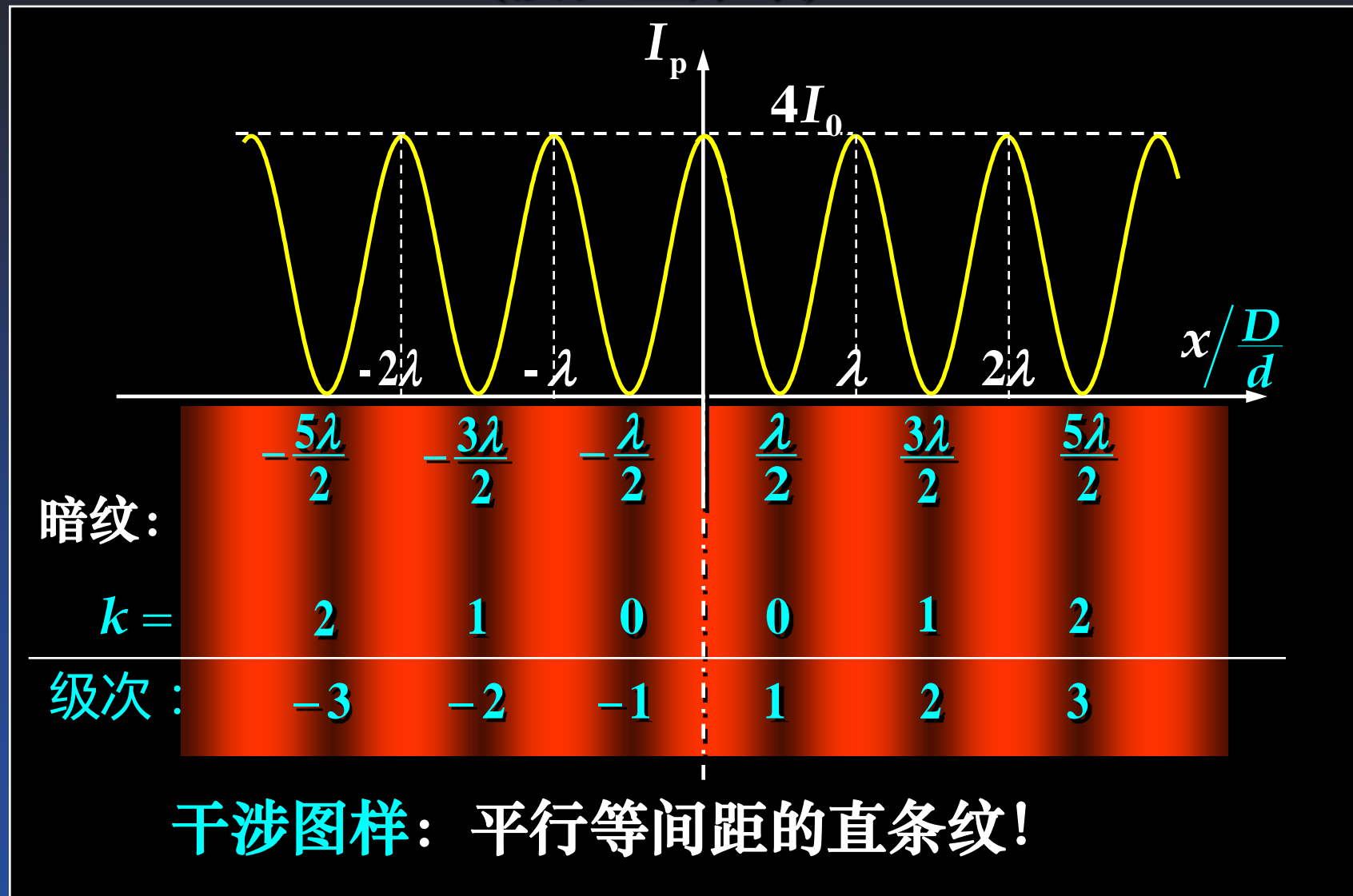


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明纹:

$k =$ 3 2 1 0 1 2 3

级次: -3 -2 -1 中央明纹 1 2 3

x

暗纹:

$k =$ 2 1 0 0 1 2

级次: -3 -2 -1 1 2 3

干涉图样: 平行等间距的直条纹!

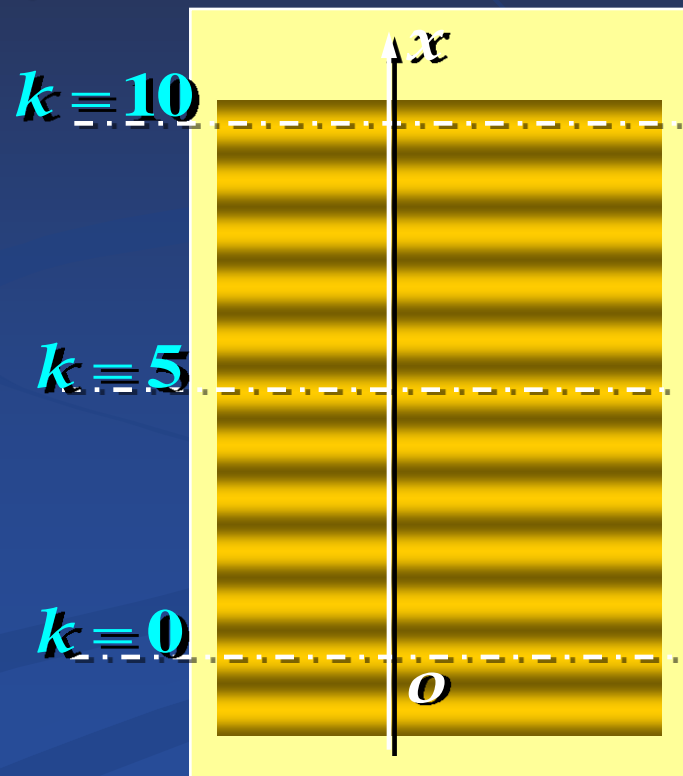
例 如图, $\lambda = 5500\text{\AA}$ 垂直入射, $d = 2 \times 10^{-4}\text{ m}$, $D = 2\text{ m}$
 求 (1) 第10级明纹间距; (2) 用厚度 $e = 6.7 \times 10^{-6}\text{ m}$ 、 $n = 1.58$ 的云母片覆盖一缝后, 中央明纹移到原来的某级明纹处, 原来该处明纹的级次为多少?

解 求第10级明纹间距:

$$x_{10} = 10 \cdot \frac{D}{d} \lambda$$

$$x_{-10} = -10 \cdot \frac{D}{d} \lambda$$

$$x_{10} - x_{-10} = 20 \frac{D}{d} \lambda = 110\text{ (mm)}$$



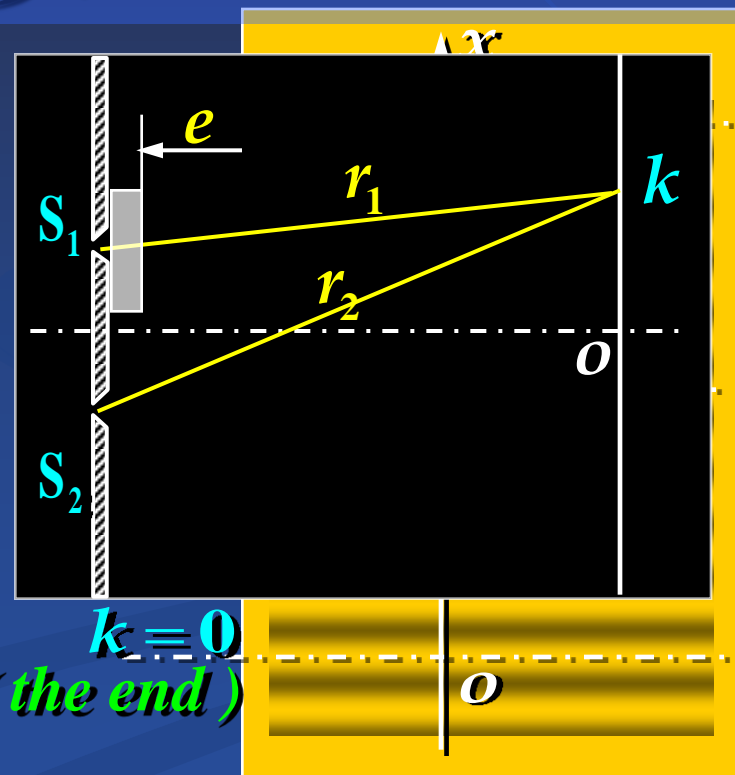
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前: $\delta = r_2 - r_1 = 2k \frac{\lambda}{2}$

后: $\delta' = r_2 - (r_1 - e + ne) = 0$

$k = \frac{n-1}{\lambda} e \approx 7$

即移到原来的第 **7** 级明纹上。



三、几点讨论

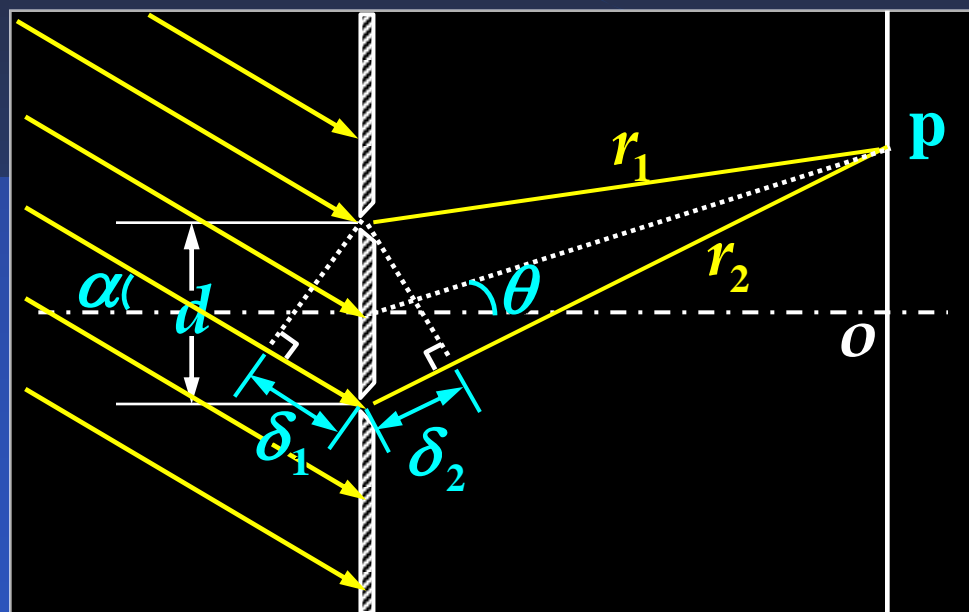
1、斜入射: $\delta = \delta_1 + \delta_2 = d \cdot \sin \theta + d \cdot \sin \alpha$

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即移到原来的第7级明纹上。 (the end)

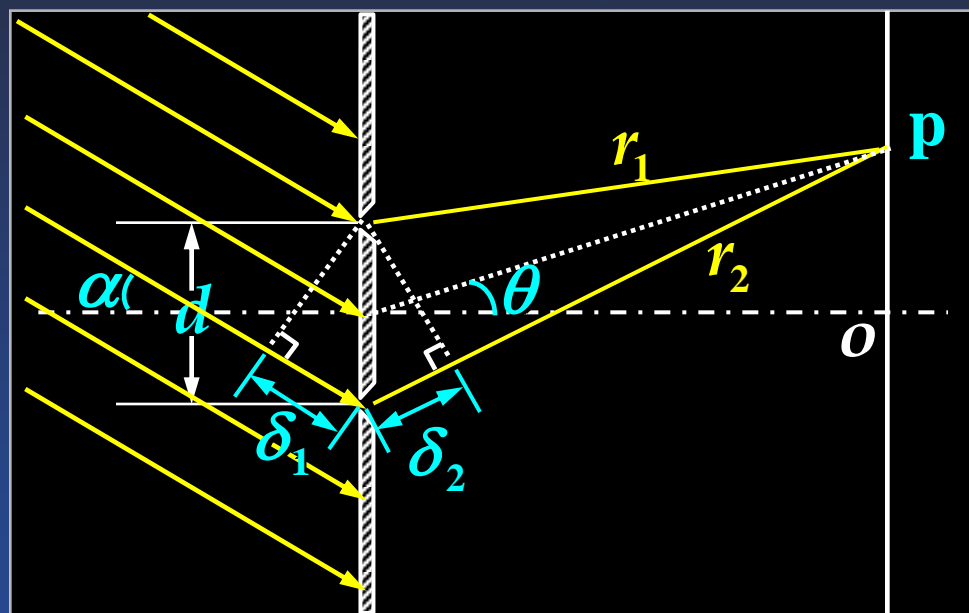
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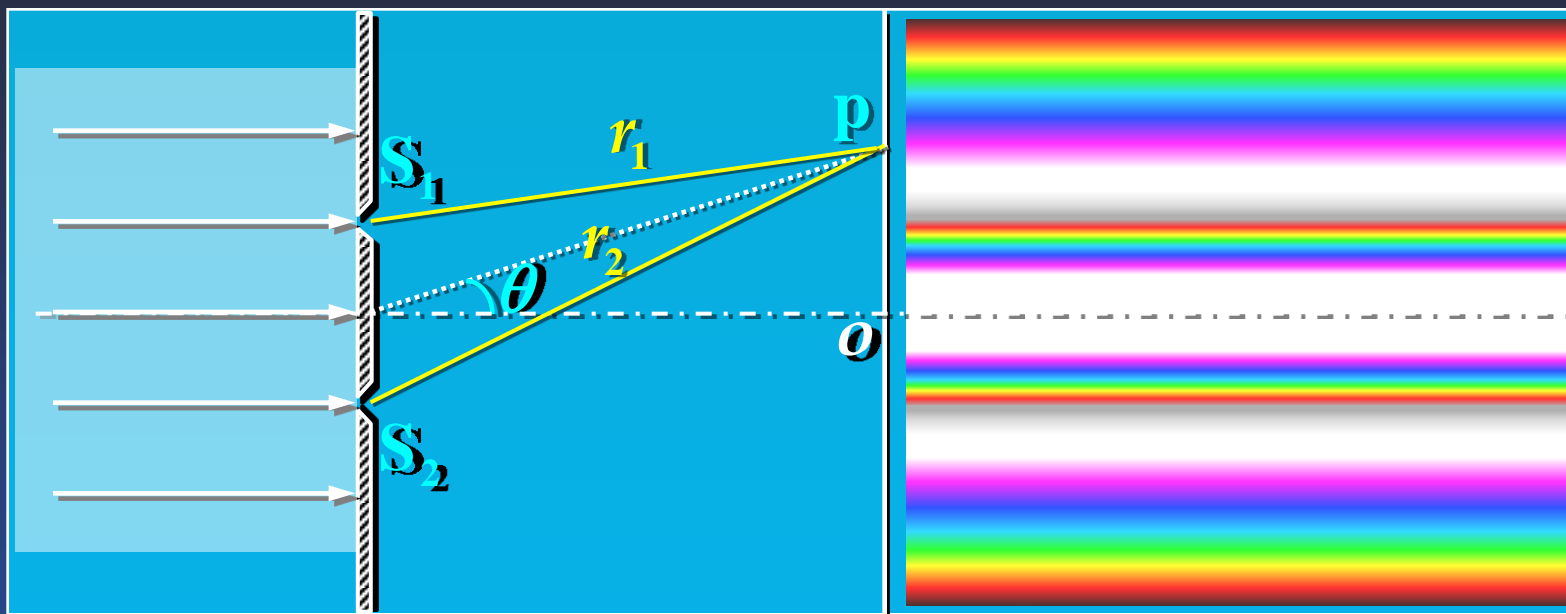
(注意 θ 、 α 的正负)



中央明纹出现在: $\delta = d \cdot (\sin \theta + \sin \alpha) = 0 \rightarrow \theta = -\alpha$

2、白光入射:

相邻条纹间距: $\Delta x = \frac{D}{d} \lambda \propto \lambda$

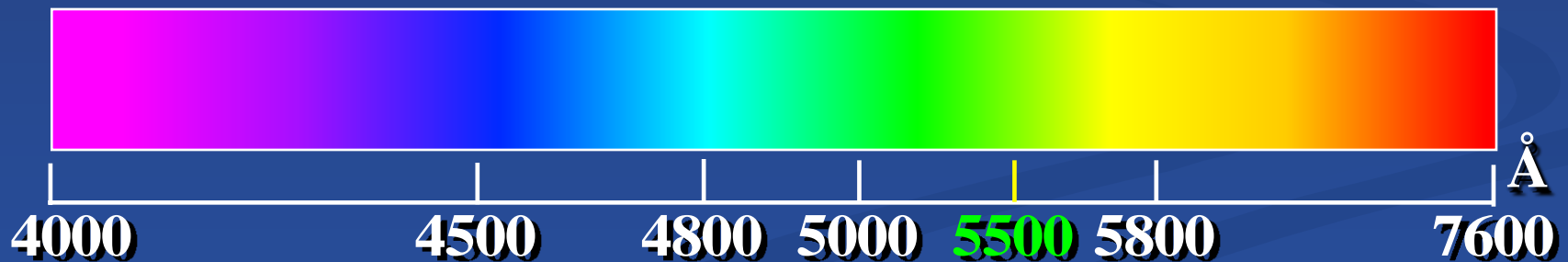
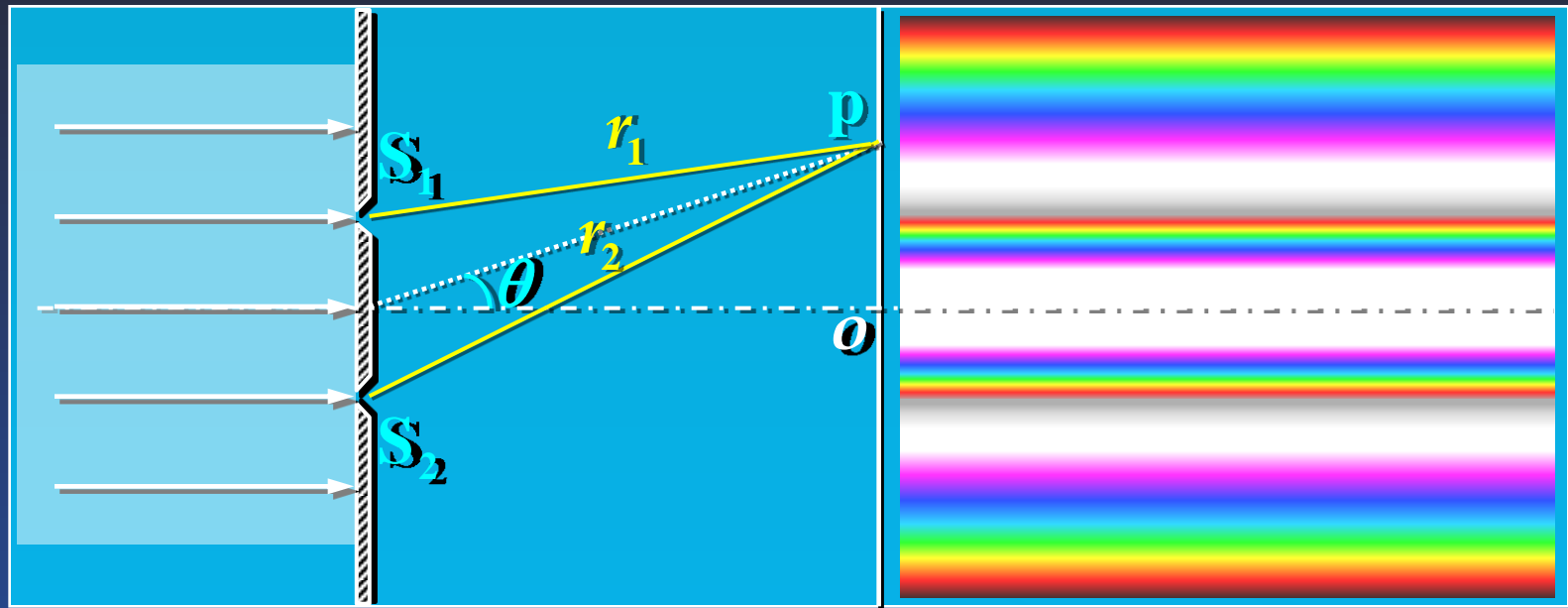


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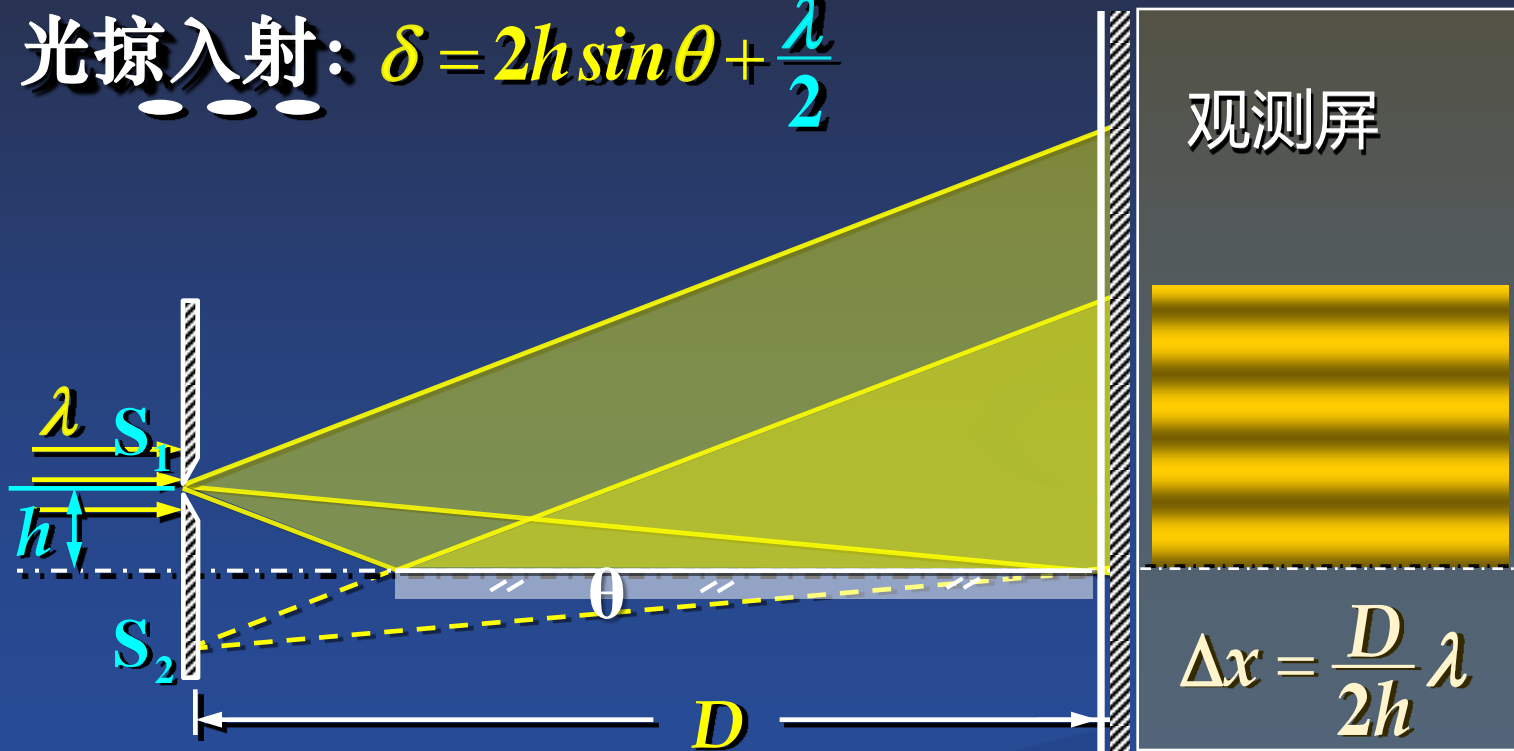
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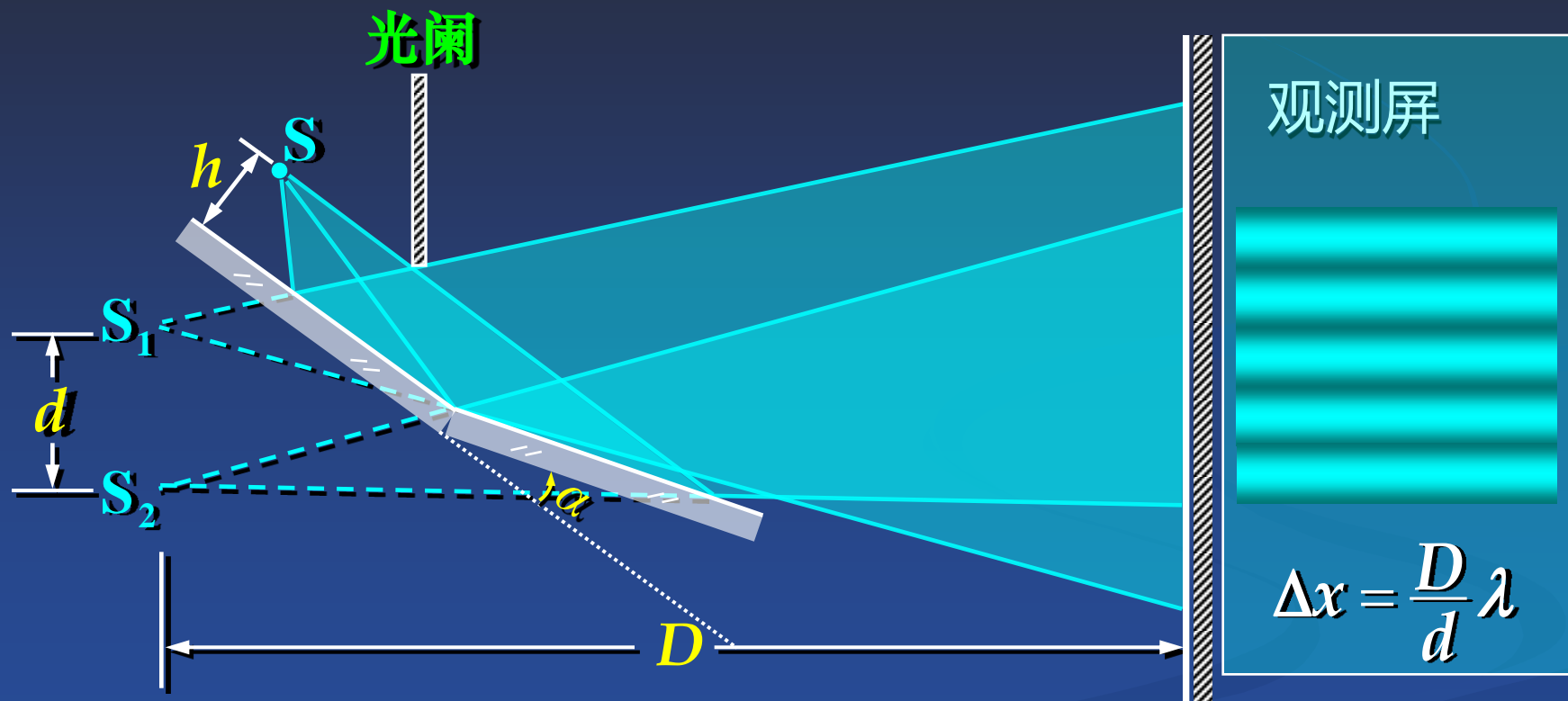
四、另外几种分波阵面法干涉实验

1、洛埃镜

光掠入射: $\delta = 2h \sin \theta + \frac{\lambda}{2}$



2、菲涅尔双镜



归纳:

1. 干涉规律:

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2. 相邻明纹(暗纹)间距: $\Delta x = \frac{D}{d} \lambda$

3. 双缝干涉条纹特点: 平行、等间距、等亮度的条纹!

4. 条纹位置分布:

end