



Chaos in the  
Electric  
Curtain

Owen Myers

# Chaos in the Electric Curtain

Owen Myers   Junru Wu   Jeff Marshall

University of Vermont  
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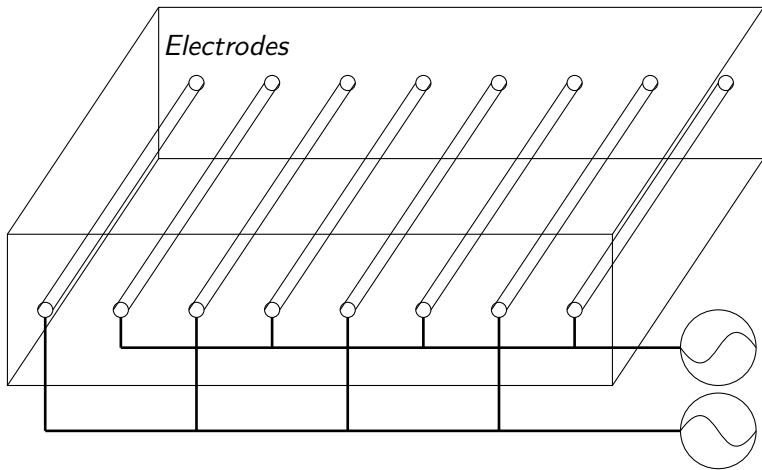


# The Electric Curtain

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## *Dielectric Surface*





# Why Study the Electric Curtain's Chaotic Dynamics

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A better understanding of the sensitivity of particle motion to physical parameters will improve applications such as:

- Dust mitigation in extraterrestrial environments (Mazumder et al., 2007)
- Separation of cells in solution (Masuda et al., 1988)
- Separation of by-products from agricultural processes (Weiss et al., 1984)
- Separation of particles with different charge to mass ratios (Kawamoto, 2008)



# Simplifications

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- 1 Primarily investigated single particle motion on the surface of the EC (1D) model.
- 2 Only preliminary work on particle motion above the surface.
- 3 In both cases dissipative forces are considered.
  - 1D Dissipative forces: Rolling resistance and viscous fluid force. Both are proportional to velocity (Brillanttov and Poschell, 1999).
  - 2D: Viscous fluid force (non-dissipative collisions)
- 4 For the center line charge approximation the electrodes are considered as one dimensional wires. To make this approximation valid we need the surface to be an adequate distance above the plane of electrodes.



# Center Line Charge Approximation

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Approximate electric field equations (Masuda and Kamimura, 1975):

$$E_x = \frac{kQ}{4\pi\epsilon_0} \sum_{n=0}^1 \frac{\sin(kx - n\pi) \cos(\omega t - n\pi)}{\cosh(ky) - \cos(kx - n\pi)}$$

$$E_y = \frac{kQ}{4\pi\epsilon_0} \sum_{n=0}^1 \frac{\sinh(ky) \cos(\omega t - n\pi)}{\cosh(ky) - \cos(kx - n\pi)}$$



# Non-Dimensionalized Parameters

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$$t' = \omega t$$

$$x' = kx$$

$$y' = ky$$

$$j = \frac{k^2 q Q}{4\pi\epsilon_0 m^2 \omega^2}$$

$$g' = \frac{gk}{\omega^2}$$

$$\beta' = \frac{\beta}{m\omega}$$

$j$  is the dimensionless interaction amplitude.

Dimensionless Equations of Motion:

$$\frac{d^2 x'}{dt'^2} + \beta' \frac{dx'}{dt'} = \sum_{n=0}^1 j \frac{\sin(x' - n\pi) \cos(t' - n\pi)}{\cosh(y') - \cos(x' - n\pi)}$$

$$\frac{d^2 y'}{dt'^2} + \beta' \frac{dy'}{dt'} = \sum_{n=0}^1 j \frac{\sinh(y') \cos(t' - n\pi)}{\cosh(y') - \cos(x' - n\pi)} - g'$$



# Interaction Amplitude Regimes

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Place particle directly above electrode to solve for when the electrostatic force equals the gravitational force.

$$\sum_{n=0}^1 j \frac{\sinh(y') \cos(t' - n\pi)}{\cosh(y') - \cos(x' - n\pi)} = g'$$

The inequality:

$$\frac{j}{g'} \leq \frac{\cosh^2(y') - 1}{2 \cos(t') \cosh(y')}$$

needs to be satisfied for the particle to remain on the surface.  
When we make  $\cos(t')$  one and put it  $y' = 1$  we find that:

$$\frac{j}{g'} \approx .448$$

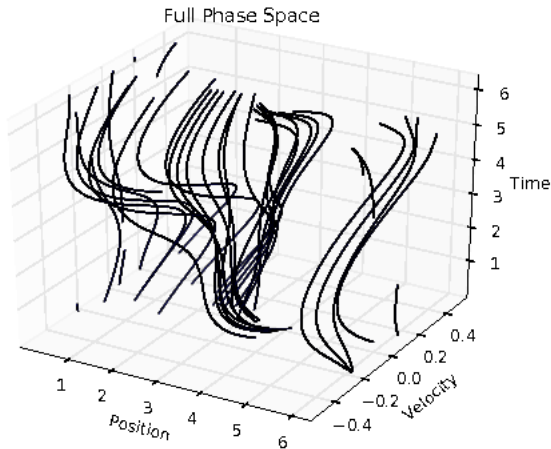
is the critical ratio that needs to be considered.



# The full periodic phase space

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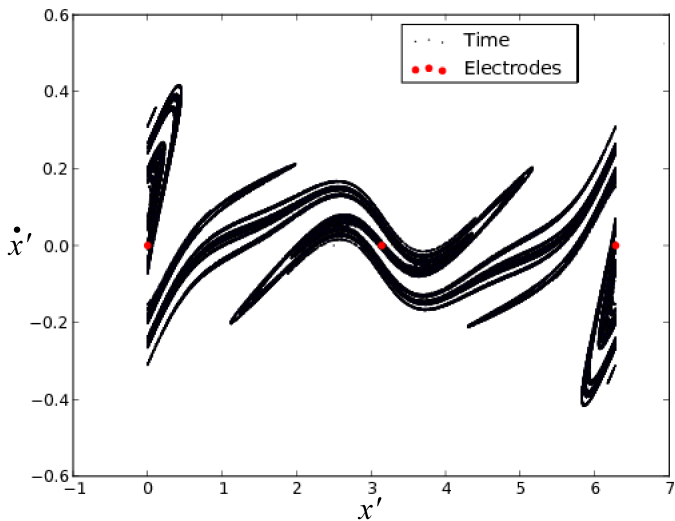




# Full phase space time sliced

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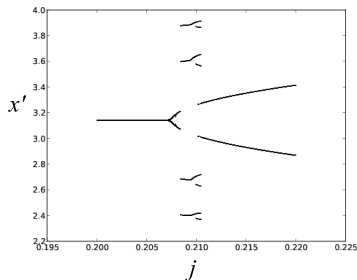


# Period doubling?

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Period 2 Orbit



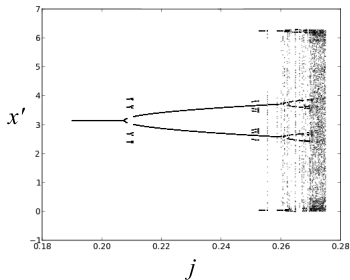
Period 4 Orbit



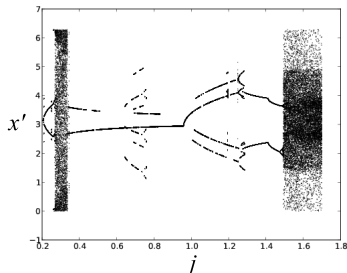
# A unique period doubling route to chaos

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Chaotic Orbit

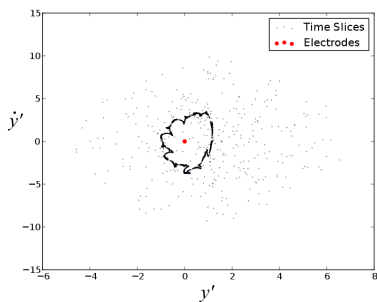
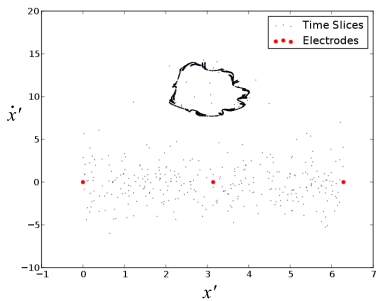




# Preliminary 2D Work

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# Conclusion

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- We have analyzed a simplified model of the electric curtain by using Poincaré sections and bifurcation diagrams.
- Particle trajectories in a real Electric curtain have the potential to be chaotic but may also exhibit order. Understanding the chaotic dynamics may help improve the electric curtain for its current applications and may suggest the possibility of new applications.
- **Future work:**
  - Improve our computational model and start to include particle-particle interactions.
  - Experimentally verify our computational findings.
  - Quantify the strength of chaos (find the Lyapunov spectrum)



# Bibliography

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