

# An Intro to the coarse-grained ME and universal Lindblad ME

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## 0.1 Model

In this tutorial, we consider a standard single qubit annealing Hamiltonian

$$H(s) = -\frac{1}{2}(1-s)\sigma_x - \frac{1}{2}s\sigma_z$$

coupling to an Ohmic bath via  $\sigma_z$  operator. We solve the open system dynamics via three different MEs: Redfield equation, coarse-grained ME(CGME), and universal Lindblad equation(ULE). Unlike the Redfield equation, CGME and ULE generate CP maps.

## 0.2 Coarse-grained ME

Coarse-grained ME is a completely positive ME obtained by applying an additional time coarse-graining approximate to the Redfield equation. More details of CGME can be found in [Mozgunov and Lidar](#). We first solve the original Redfield equation and CGME and compare both cases' instantaneous ground state population.

```
using OrdinaryDiffEq, Plots, LaTeXStrings
using OpenQuantumTools

# Hamiltonian
H = DenseHamiltonian([(s)->1-s, (s)->s], -[σx, σz]/2, unit=:ħ)
# initial state
u0 = PauliVec[1][1]
# coupling
coupling = ConstantCouplings(["Z"], unit=:ħ)
# bath
bath = Ohmic(1e-4, 4, 16)
annealing = Annealing(H, u0; coupling=coupling, bath=bath)

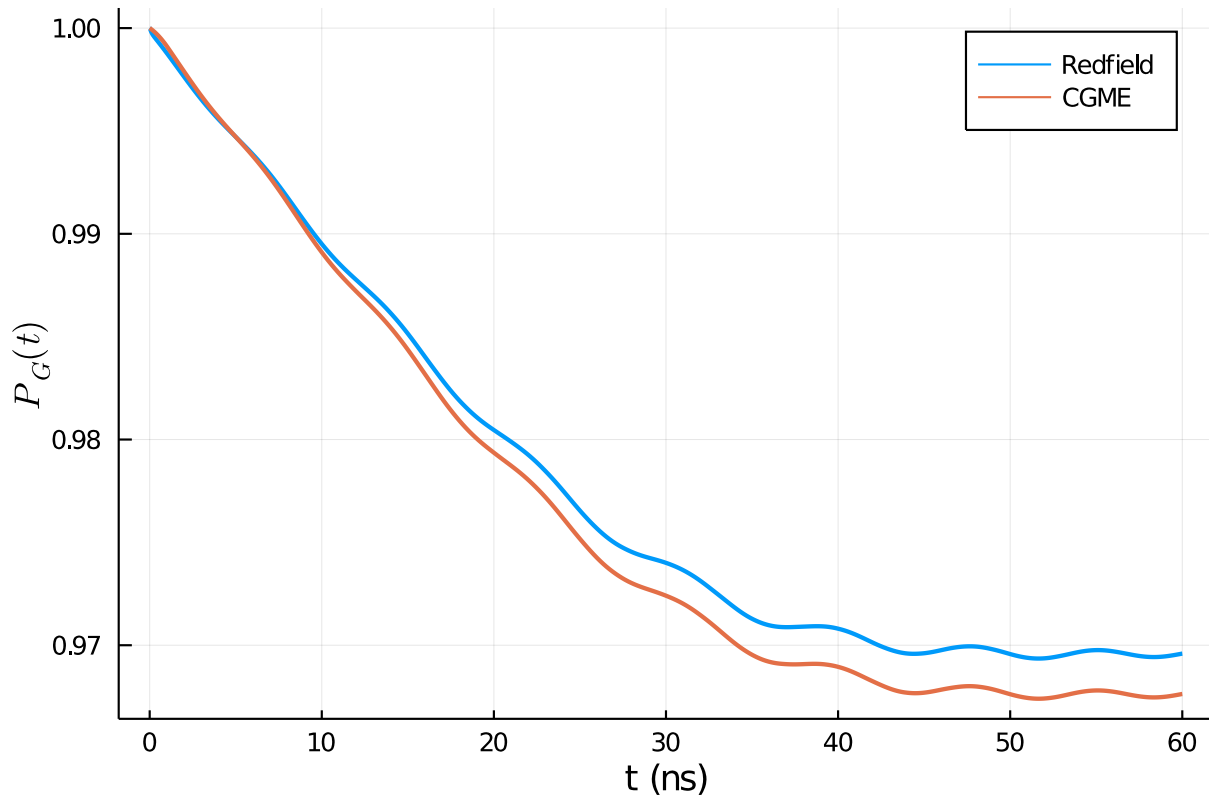
tf = 60
U = solve_unitary(annealing, tf, alg=Tsit5(), abstol=1e-8, reltol=1e-8)
U = InplaceUnitary(U)

@time solr = solve_redfield(annealing, tf, U, alg=Tsit5())
# we set the integration error tolerance to 1e-5 for speed
@time solc = solve_cgme(annealing, tf, U, alg=Tsit5(), int_atol=1e-5, int_rtol=1e-5)
plot(solr, H, [0], 0:0.01:tf, linewidth=2, xlabel="t (ns)", ylabel="\$P_G(t)\$",
label="Redfield")
```

```
plot!(solc, H, [0], 0:0.01:tf, linewidth=2, label="CGME")
```

0.290015 seconds (3.08 M allocations: 80.058 MiB)

51.352793 seconds (489.73 M allocations: 19.392 GiB, 5.07% gc time)



### 0.3 Universal Lindblad equation

Universal Lindblad equation (ULE) is a different CP ME proposed in [Nathan and Rudner](#). Unlike the Redfield and CGME, it depends on the jump correlator, which is the inverse Fourier transform of the square root of the noise spectrum:

$$g(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \sqrt{\gamma(\omega)} e^{i\omega t} d\omega .$$

Let's first see how it looks compared with the two-point correlation function  $C(t)$ :

```
using QuadGK
```

```
g(t) = quadgk((w)->sqrt(γ(w, bath))*exp(1.0im*w*t)/2/π, -Inf, Inf)[1]
```

```
t = range(-0.5,0.5,length=500)
```

```
g_value = g.(t)
```

```
c_value = [correlation(x, bath) for x in t];
```

```
plot(t, real.(g_value), label="Re[g(t)]", linewidth=2)
```

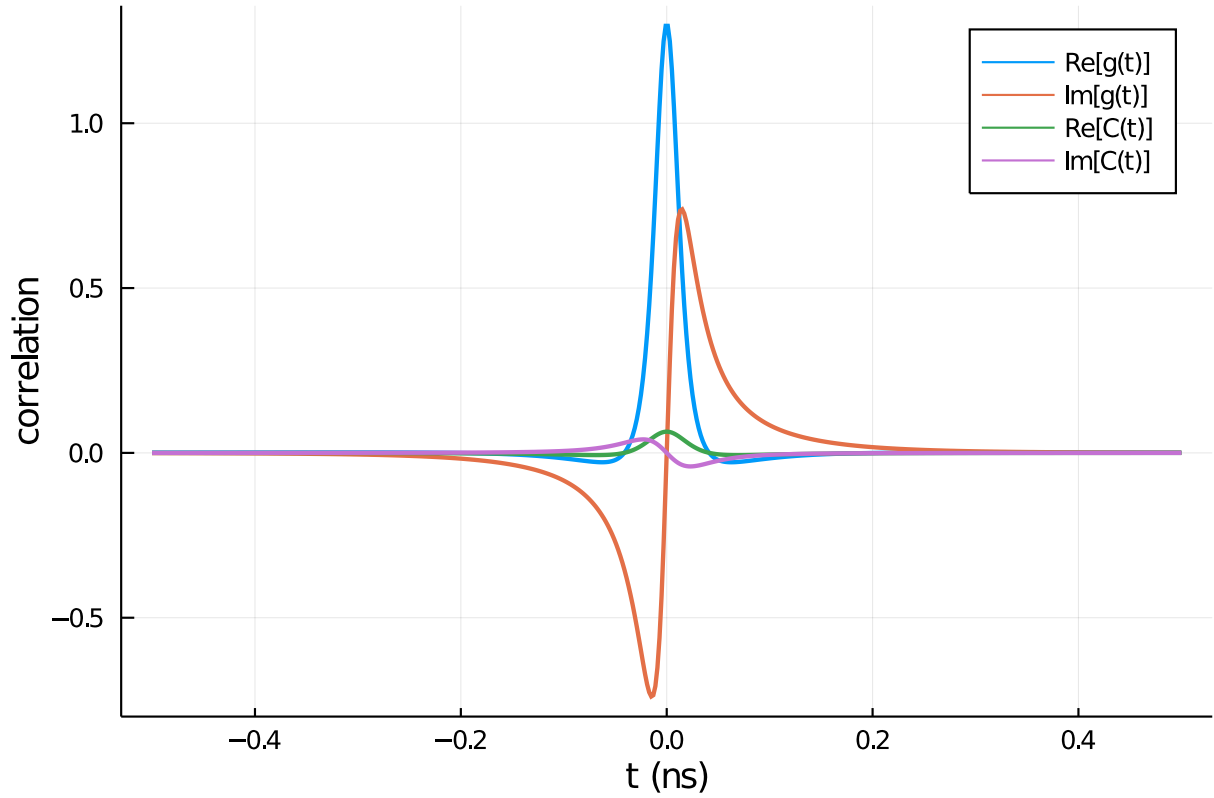
```
plot!(t, imag.(g_value), label="Im[g(t)]", linewidth=2)
```

```
plot!(t, real.(c_value), label="Re[C(t)]", linewidth=2)
```

```
plot!(t, imag.(c_value), label="Im[C(t)]", linewidth=2)
```

```
xlabel!("t (ns)")
```

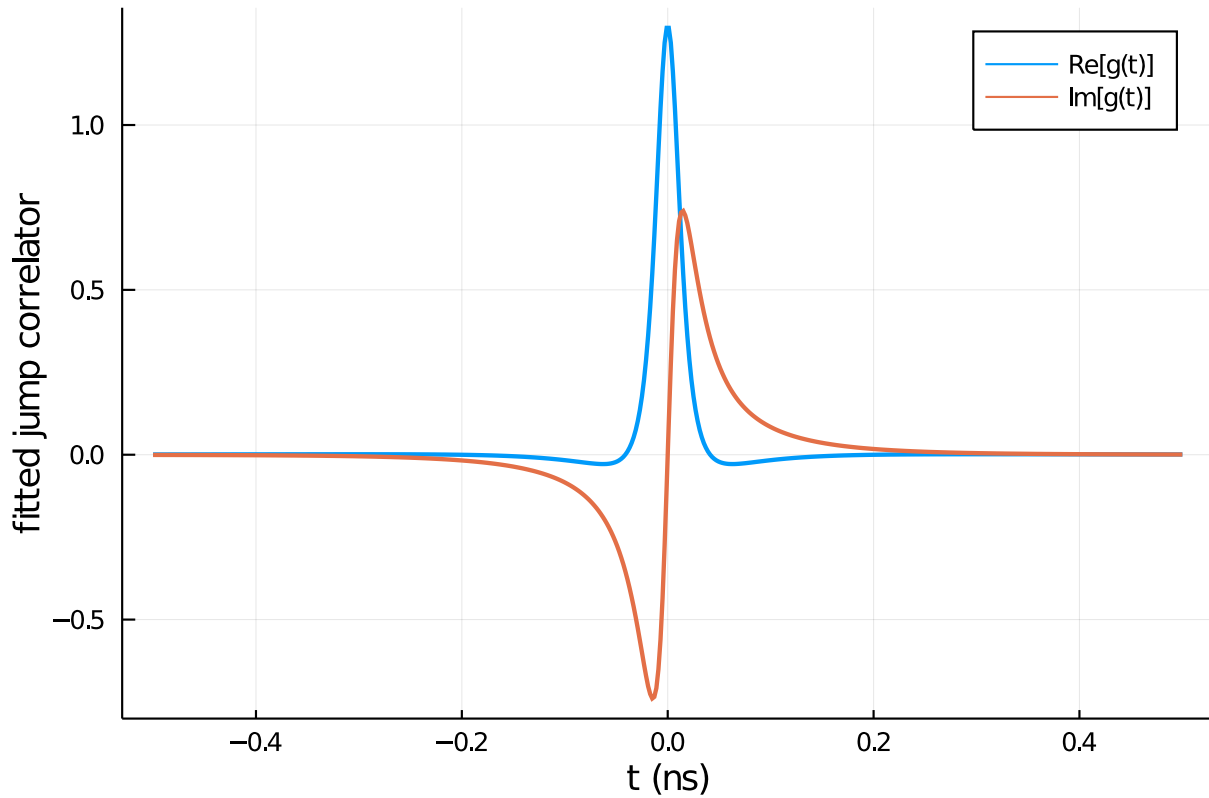
```
ylabel!("correlation")
```



From the above picture, we can see that the jump correlator and two-point correlation function roughly have the same time scale. To avoid recalculating the inverse Fourier transform within the solver, we can precalculate  $g(t)$  and construct interpolation from these pre-computed values. This procedure can be done by the following code block:

```
t = range(-4,4,length=2000)
g_value = g.(t)
gf = construct_interpolations(t, g_value, extrapolation = "flat")

t = range(-0.5,0.5,length=500)
g_value = gf.(t)
plot(t, real.(g_value), label="Re[g(t)]", linewidth=2)
plot!(t, imag.(g_value), label="Im[g(t)]", linewidth=2)
xlabel!("t (ns)")
ylabel!("fitted jump correlator")
```



Finally we solve ULE and compare the result with the Redfield equation and CGME:

```
ubath = ULEBath(gf)
annealing = Annealing(H, u0; coupling=coupling, bath=ubath)
@time solu = solve_ule(annealing, tf, U, alg=Tsit5(), int_atol=1e-5, int_rtol=1e-5)
plot(solr, H, [0], 0:0.01:tf, linewidth=2, xlabel="t (ns)", ylabel="\$P_G(t)\$",
label="Redfield")
plot!(solc, H, [0], 0:0.01:tf, linewidth=2, label="CGME")
plot!(solu, H, [0], 0:0.01:tf, linewidth=2, label="ULE")
```

0.506539 seconds (5.09 M allocations: 133.383 MiB, 2.78% gc time)

