# An tutorial on polaron transformed Redfield equation

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## 0.1 Correlation function in polaron frame

This is a tutorial for using the polaron transformed Redfield equation (PTRE) in HOQST. For more details on PTRE, Xu and Cao is a good reference.

In this example, we solve both the Redfield equation and PTRE for a single qubit model with system Hamiltonian

$$H_{\rm S} = \epsilon \sigma_z + \Delta \sigma_x$$

coupling to an Ohmic bath via  $\sigma_z$  interaction

$$H = H_{\rm S} + \sigma_z \otimes B + H_{\rm B}$$
.

Loosely speaking, the main difference between the Redfield equation and PTRE is that they have different bath correlation functions. For the Redfield equation, the bath correlation function is

$$C(t_1, t_2) = \langle B(t_1)B(t_2) \rangle$$
.

In the polaron frame, however, the bath correlation function becomes

$$K(t_1, t_2) = \exp \left\{ -4 \int_0^t \int_{-\infty}^0 C(t_1, t_2) dt_1 dt_2 \right\}.$$

Again, interested readers can refer to [Amin and Averin] and Leggett et al. for more details.

#### 0.1.1 Error bound on the second-order master equation

The simplest analysis is to compare the error bounds given in Mozgunov and Lidar between the Redfield equation and PTRE. We define the error scaling parameter as

$$error = \frac{\tau_{\rm B}}{\tau_{\rm SB}}$$
.

Then we plot the error ratio between the Redfield equation and PTRE

$$R = \frac{error_{\text{Redfield}}}{error_{\text{PTRE}}} \ ,$$

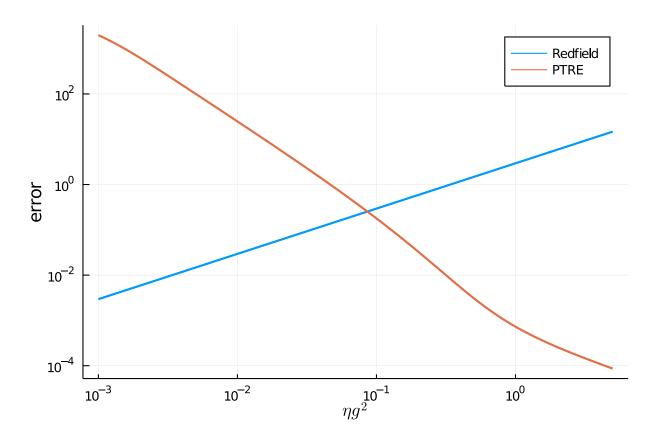
vs. the system bath coupling strength  $\eta g^2$  while fixing other parameters in the Ohmic bath.

```
using OrdinaryDiffEq, OpenQuantumTools, Plots
using LaTeXStrings
function err_bound(tf, cfun)
    tsb, esb = \tau_{SB}(cfun)
    tb, eb = \tau_B(cfun, tf, tsb)
    tb / tsb
end
fc = 4; T = 12; tf = 1000;
\etalist = log_uniform(1e-3, 5, 1000)
err_ratio = []
err_clist = []
err_klist = []
for \eta in \etalist
    bath = Ohmic(\eta, fc, T)
    cfun = (x)->correlation(x, bath)
    pfun = (x)->polaron_correlation(x, bath)
    err_c = err_bound(tf, cfun)
    err_k = err_bound(tf, pfun)
    push!(err_clist, err_c)
    push!(err_klist, err_k)
    push!(err_ratio, err_c/err_k)
idx = findfirst((x)->x>=1, err_ratio)
plot(\etalist, err_ratio, xscale=:log10, yscale=:log10, label="", linewidth=2)
vline!([\etalist[idx]], label="", linestyle=:dash, linewidth=2)
annotate!([(0.5, 1.0, Plots.text("polaron")), (0.01, 1.0, Plots.text("Redfield"))])
xlabel!(L"\eta g^2")
ylabel!("R")
     10<sup>5.0</sup>
     10<sup>2.5</sup>
    10<sup>0.0</sup>
                              Redfield
                                                                     polaron
    10<sup>-2.5</sup>
    10<sup>-5.0</sup>
                                                     \eta g^{10^{-1}}
                                  10-2
            10<sup>-3</sup>
                                                                                 10<sup>0</sup>
```

From above figure we can see that, as the system-bath coupling strength is bigger than  $10^{-1}$ ,

PTRE should have better error scaling than the usual form of Redfield equation. We also plot the corresponding error values for both Redfield and PTRE:

```
plot(\etalist, err_clist, xscale=:log10, yscale=:log10, label="Redfield", linewidth=2) plot!(\etalist, err_klist, xscale=:log10, yscale=:log10, label="PTRE", linewidth=2) xlabel!(L"\eta g^2") ylabel!("error")
```



The above figure confirms that the range of applicability of the Redfield equation and PTRE are weak and strong coupling regimes, respectively.

### 0.1.2 Solving PTRE

Since PTRE and the Redfield equation have identical forms, solve\_redfield can also be used for PTRE. To see this, let's first write down the PTRE for our example.

$$\dot{\rho}_{\rm S} = \epsilon \sigma_z + [\sigma_i, \Lambda_i(t) \rho_{\rm S}(t)] + h.c.$$

where  $i, j \in [+, -], i \neq j$  and

$$\Lambda_i(t) = \Delta^2 \int_0^t K(t-\tau)U(t,\tau)\sigma_j U^{\dagger}(t,\tau)d\tau .$$

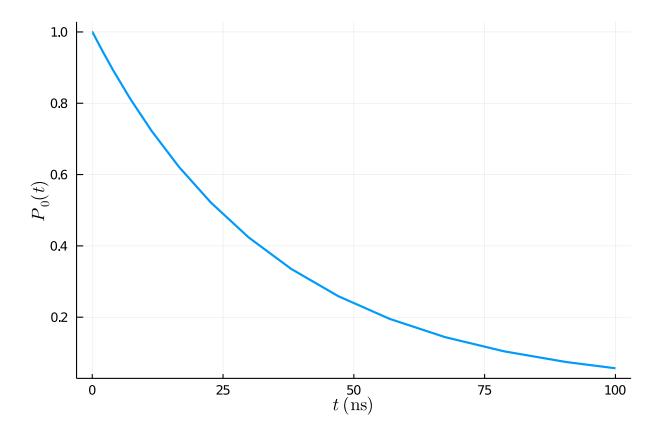
From the above equations, it is clear that the following steps are needed to define an annealing process in the polaron frame:

- 1. define a new Hamiltonian  $H = \epsilon \sigma_z$ ;
- 2. define new coupling operators  $\sigma_{-}$  and  $\sigma_{+}$ ;

3. define new correlated bath with two-point correlation  $K_{i,j}(t_1,t_2)$ ;

The following code block illustrates how these can be done in HOQST:

```
# assume \epsilon = 1
    const \Delta = 0.1
    # define the Ohmic bath in polaron transformed frame
    \eta = 0.5; bath = Ohmic(\eta, fc, T)
    K(t1, t2) = \Delta^2 * polaron_correlation(t1-t2, bath)
    cfun = [nothing K; K nothing]
    pbath = CorrelatedBath(((1,2),(2,1)), correlation=cfun)
    # define coupling as \sigma\text{+} and \sigma\text{-} operators
    \sigma p = [0 \ 1; 0 \ 0.0im]; \ \sigma m = [0 \ 0; 1 \ 0.0im]
    coupling = ConstantCouplings([\sigmap, \sigmam])
    # manually define the unitary operator
    U(t) = \exp(-2.0im * \pi * \sigma z * t)
    H = DenseHamiltonian([(s)->1.0], [\sigma z])
    u0 = PauliVec[3][1]
    annealing = Annealing(H, u0, coupling = coupling, bath = pbath)
    tf = 100
    sol_ptre = solve_redfield(annealing, tf, U, alg=Tsit5(), Ta=2, reltol=1e-5)
    pop_e = [real(s[1,1]) for s in sol_ptre.u]
    plot(sol_ptre.t, pop_e, xlabel=L"t\ (\mathrm{ns})", ylabel=L"P_0(t)", label="",
linewidth = 2)
```

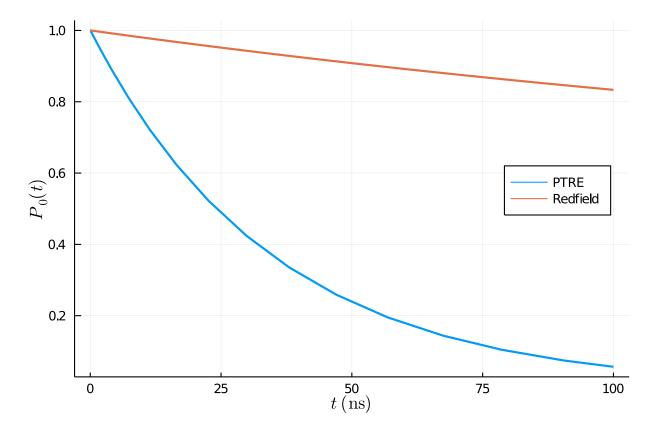


For historical reasons, this is known as an example of "incoherent tunneling". The off-diagonal elements of the density matrix in computational bases(Z-bases) during the entire evolution are 0(shown in the next section).

#### 0.1.3 Redfield equation

What would happen to the Redfield equation in this regime? We can also try:

```
H = DenseHamiltonian([(s)->1.0], [\sigma z+0.1*\sigma x])
coupling = ConstantCouplings(["Z"])
annealing = Annealing(H, u0, coupling = coupling, bath = bath)
tf = 100
sol_redfield = solve_redfield(annealing, tf, U, alg=Tsit5(), Ta=40, reltol=1e-5,
callback=PositivityCheckCallback())
pop_e_redfield = [real(s[1,1]) for s in sol_redfield.u]
plot(sol_ptre.t, pop_e, xlabel=L"t\ (\mathrm{ns})", ylabel=L"P_0(t)", label="PTRE",
linewidth = 2, legend = :right)
plot!(sol_redfield.t, pop_e_redfield, xlabel=L"t\ (\mathrm{ns})", ylabel=L"P_0(t)",
label="Redfield", linewidth = 2)
```



PTRE gives a much stronger decay than the Redfield equation for the parameters chosen in this example. One can also verify the amplitude of the off-diagonal elements during the evolution. Unlike PTRE, the solution of the Redfield equation has non-vanishing off-diagonal elements of the density matrix.

```
t_axis = range(0, 5, length=100)
off_diag_ptre = [abs(sol_ptre(t)[1,2]) for t in t_axis]
off_diag_redfield = [abs(sol_redfield(t)[1,2]) for t in t_axis]
plot(t_axis, off_diag_ptre, xlabel=L"t\ (\mathrm{ns})", ylabel=L"\lvert \rho_{12}\
\rvert|(t)", label="PTRE", linewidth = 2, legend=:right)
plot!(t_axis, off_diag_redfield, xlabel=L"t\ (\mathrm{ns})", ylabel=L"|\rho_{12}(t)|",
label="Redfield", linewidth = 2)
```

