

Newtonian approximation in Causal Dynamical Triangulations

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1 Motivation

1.1 Newton's Law of Gravitation from General Relativity

Starting from the most general cylindrically symmetric (Weyl) metric [1]:

$$ds^2 = e^{2\lambda} dt^2 - e^{2(\nu-\lambda)} (dr^2 + dz^2) - r^2 e^{-2\lambda} d\phi^2 \quad (1)$$

$$g_{\mu\nu} = \begin{pmatrix} e^{2\lambda} dt^2 & 0 & 0 & 0 \\ 0 & -e^{2(\nu-\lambda)} dr^2 & 0 & 0 \\ 0 & 0 & -e^{2(\nu-\lambda)} dz^2 & 0 \\ 0 & 0 & 0 & -\frac{r^2}{e^{2\lambda}} d\phi^2 \end{pmatrix} \quad (2)$$

The definition of the Christoffel connection is: [2]

$$\Gamma_{\mu\nu}^{\lambda} = \frac{1}{2} g^{\lambda\sigma} (\partial_{\mu} g_{\nu\sigma} + \partial_{\nu} g_{\sigma\mu} - \partial_{\sigma} g_{\mu\nu}) \quad (3)$$

With the assumption of zero torsion:

$$\Gamma_{\mu\nu}^{\lambda} = \Gamma_{\nu\mu}^{\lambda} \quad (4)$$

The non-zero Christoffel connections are:

$$\begin{aligned}
\Gamma_{tr}^t &= \partial_r \lambda \\
\Gamma_{tz}^t &= \partial_z \lambda \\
\Gamma_{tt}^r &= e^{4\lambda-2v} \partial_r \lambda \\
\Gamma_{rr}^r &= \partial_r v - \partial_r \lambda \\
\Gamma_{rz}^r &= \partial_z v - \partial_z \lambda \\
\Gamma_{zz}^r &= \partial_z \lambda - \partial_z v \\
\Gamma_{\phi\phi}^r &= r e^{-2v} (r \partial_r \lambda - 1) \\
\Gamma_{tt}^z &= e^{4\lambda-2v} \partial_z \lambda \\
\Gamma_{rr}^z &= \partial_z \lambda - \partial_z v \\
\Gamma_{rz}^z &= \partial_r v - \partial_r \lambda \\
\Gamma_{zz}^z &= \partial_r v - \partial_r \lambda \\
\Gamma_{\phi\phi}^z &= r^2 e^{-2v} \partial_z \lambda \\
\Gamma_{r\phi}^\phi &= \frac{1}{r} - \partial_r \lambda \\
\Gamma_{z\phi}^\phi &= -\partial_z \lambda
\end{aligned} \tag{5}$$

The components of the Riemann tensor are given by:

$$R_{\sigma\mu\nu}^\rho = \partial_\mu \Gamma_{\nu\sigma}^\rho - \partial_\nu \Gamma_{\mu\sigma}^\rho + \Gamma_{\mu\lambda}^\rho \Gamma_{\nu\sigma}^\lambda - \Gamma_{\nu\lambda}^\rho \Gamma_{\mu\sigma}^\lambda \tag{6}$$

Using the properties of the Riemann tensor:

$$\begin{aligned}
R_{\rho\sigma\mu\nu} &= -R_{\rho\sigma\nu\mu} \\
R_{\rho\sigma\mu\nu} &= -R_{\sigma\rho\mu\nu} \\
R_{\rho\sigma\mu\nu} &= R_{\mu\nu\rho\sigma} \\
R_{\rho[\sigma\mu\nu]} &= 0
\end{aligned} \tag{7}$$

The non-zero components of the Riemann tensor are:

$$\begin{aligned}
R_{tr}^t &= -\partial_r^2 \lambda + (\partial_z \lambda)^2 - 2(\partial_r \lambda)^2 + \partial_r \lambda \partial_r v - \partial_z \lambda \partial_z v \\
R_{tz}^t &= -\partial_r \partial_z \lambda - 3\partial_r \lambda \partial_z \lambda + \partial_r \lambda \partial_z v + \partial_r v \partial_z \lambda \\
R_{rz}^t &= -\partial_z^2 \lambda - 2(\partial_z \lambda)^2 + (\partial_r \lambda)^2 - \partial_r \lambda \partial_r v + \partial_z \lambda \partial_z v \\
R_{\phi t}^t &= r e^{-2v} (r(\partial_r \lambda)^2 - \partial_r \lambda + r(\partial_z \lambda)^2) \\
R_{rz}^r &= \partial_r^2 \lambda - \partial_r^2 v + \partial_z^2 \lambda - \partial_z^2 v \\
R_{\phi z}^r &= r e^{-2v} (r \partial_z^2 \lambda - r \partial_z \lambda \partial_z v + r \partial_r \lambda \partial_r v - r(\partial_r \lambda)^2 + \partial_r \lambda - \partial_r v) \\
R_{\phi\phi}^r &= r e^{-2v} (-r \partial_r \partial_z \lambda + r \partial_r v \partial_z \lambda - r \partial_r \lambda \partial_z \lambda + r \partial_r \lambda \partial_z v - \partial_z v) \\
R_{r\phi}^\phi &= \partial_r^2 \lambda + \frac{1}{r} \partial_r v - \partial_r \lambda \partial_r v - (\partial_z \lambda)^2 + \partial_z \lambda \partial_z v + \frac{1}{r} \partial_r \lambda
\end{aligned} \tag{8}$$

The Ricci tensor is given by:

$$R_{\mu\nu} = R_{\mu\lambda\nu}^\lambda \tag{9}$$

The non-zero components of the Ricci tensor are:

$$\begin{aligned}
R_{tt} &= \frac{e^{4\lambda-2\nu}}{r} \left(r\partial_r^2\lambda + r\partial_z^2\lambda + \partial_r\lambda \right) \\
R_{rr} &= \partial_r^2\lambda - \partial_r^2\nu + \partial_z^2\lambda - \partial_z^2\nu - 2(\partial_r\lambda)^2 + \frac{1}{r}\partial_r\lambda + \frac{1}{r}\partial_r\nu \\
R_{rz} &= \frac{1}{r}\partial_z\nu - 2\partial_r\lambda\partial_z\lambda \\
R_{zz} &= \partial_r^2\lambda - \partial_r^2\nu + \partial_z^2\lambda - \partial_z^2\nu - 2(\partial_z\lambda)^2 + \frac{1}{r}\partial_r\lambda - \frac{1}{r}\partial_r\nu \\
R_{\phi\phi} &= re^{-2\nu} \left(r\partial_r^2\lambda + r\partial_z^2\lambda + \partial_r\lambda \right)
\end{aligned} \tag{10}$$

Einstein's equation in a vacuum is:

$$R_{\mu\nu} = 0 \tag{11}$$

Applying this complete set of relations to Equation (10) gives the following:

$$\partial_r^2\lambda + \frac{1}{r}\partial_r\lambda + \partial_z^2\lambda = 0 \tag{12}$$

$$\partial_r\nu = r \left(\partial_r^2\nu + \partial_z^2\nu + 2(\partial_r\lambda)^2 \right) \tag{13}$$

$$\partial_z\nu = 2r\partial_r\lambda\partial_z\lambda \tag{14}$$

$$\partial_r^2\nu + \partial_z^2\nu + (\partial_r\lambda)^2 + (\partial_z\lambda)^2 = 0 \tag{15}$$

Equation (12) is the two-dimensional Laplace equation in cylindrical coordinates, for which the known solutions are:

$$\lambda(r, z) = \sum_{m=0}^{\infty} [A_m J_m(kr) + B_m N_m(kr)] [C_m \sinh(kr) + D_m \cosh(kr)] \tag{16}$$

Plugging Equation (15) into Equation (13) gives:

$$\partial_r\nu = r \left((\partial_r\lambda)^2 - (\partial_z\lambda)^2 \right) \tag{17}$$

Using Equations (14), (16) and (17) we find solutions for ν given by:

$$\nu = \int r \left[\left((\partial_r\lambda)^2 - (\partial_z\lambda)^2 \right) dr + (2\partial_r\lambda\partial_z\lambda) dz \right] \tag{18}$$

However, before we can consider this to be a complete solution we must consider elementary flatness. This condition requires that, for any infinitesimal spacelike circle, the ratio of circumference to radius is 2π . The most likely place to run into issues is along the z-axis, for which $r = 0$. Looking back at Equation (1) we see that the necessary condition is:

$$\nu = 0 \quad \text{for} \quad r = 0 \tag{19}$$

References

- [1] J. L. Synge, *Relativity: the general theory*. North-Holland Pub. Co., 1960.
- [2] S. Carroll, *Spacetime and Geometry: An Introduction to General Relativity*. Benjamin Cummings, Sept. 2003.