Modeling 2DoF System

ARO (@_art_of_electronics_)
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1 Introduction

This document presents the modeling of a 2 degrees-of-freedom (DoF) system using Lagrangian mechanics and co-energy analysis. A Simulink block diagram is also provided for simulation purposes.

2 2DoF System Description

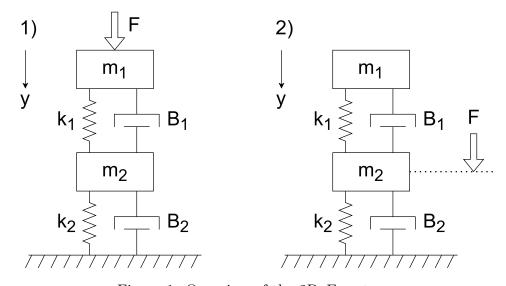


Figure 1: Overview of the 2DoF system

We consider a 2DoF system in 2 cases: first, where force is applied to outer mass m_1 ; second, where force is applied to inner mass m_2 . In this model, forcing can also be applied as displacement, given by initial conditions. In order to perform modeling, the 2DoF system must be uncoupled into 2 separate equivalent 1DoF subsystems.

3 Lagrangian Formulation

3.1 Forcing on m_1 - subsystem 1a

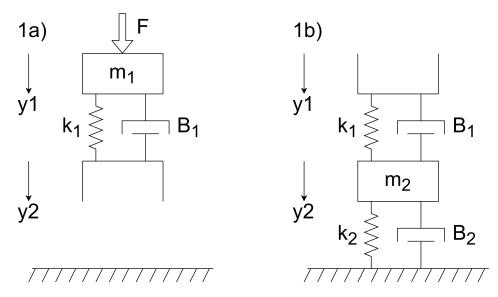


Figure 2: Uncoupled 2DoF system with forcing on m_1

The kinetic and potential energies of subsystem 1a are defined as:

$$T_1 = \frac{1}{2}m_1\dot{y_1}^2$$

$$U_1 = k_1 \int_{y_1}^{y_2} y dy = \frac{1}{2}k_1y_1^2 - \frac{1}{2}k_1y_2^2$$

The system also has a damper, which is a dissipative element:

$$D_1 = B_1 \int_{\dot{y_1}}^{\dot{y_2}} \dot{y} dy = \frac{1}{2} B_1 \dot{y}_1^2 - \frac{1}{2} B_1 \dot{y}_2^2$$

The Lagrangian of the subsystem 1a is:

$$\mathcal{L}(y, \dot{y}) = T_1 - U_1 = \frac{1}{2}m_1\dot{y_1}^2 - \left(\frac{1}{2}k_1y_1^2 - \frac{1}{2}k_1y_2^2\right)$$

The Rayleigh dissipation function can be written as:

$$\mathcal{D} = D_1 - P_1 = \left(\frac{1}{2}B_1\dot{y}_1^2 - \frac{1}{2}B_1\dot{y}_2^2\right) - F\dot{y}_1$$

Using Lagrange's equation with a non-conservative force (damper), the equation expands into the following form:

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{q}} \right) - \frac{\partial \mathcal{L}}{\partial q} + \frac{\partial \mathcal{D}}{\partial \dot{q}} = 0$$

This leads to the following differential equation:

$$\frac{d}{dt} \left(\frac{\partial \left(\frac{1}{2} m_1 \dot{y}_1^2 - \left(\frac{1}{2} k_1 y_1^2 - \frac{1}{2} k_1 y_2^2 \right) \right)}{\dot{y}} \right) - \frac{\partial \left(\frac{1}{2} m_1 \dot{y}_1^2 - \left(\frac{1}{2} k_1 y_1^2 - \frac{1}{2} k_1 y_2^2 \right) \right)}{y} + \frac{\partial \left(\left(\frac{1}{2} B_1 \dot{y}_1^2 - \frac{1}{2} B_1 \dot{y}_2^2 \right) - F \dot{y}_1 \right)}{\dot{y}} = 0$$

$$\frac{d}{dt} \left(\frac{\partial \left(\frac{1}{2} m_1 \dot{y_1}^2 - 0 \right)}{\dot{y}} \right) - \frac{\partial \left(0 - \left(\frac{1}{2} k_1 y_1^2 - \frac{1}{2} k_1 y_2^2 \right) \right)}{y} + \frac{\partial \left(\left(\frac{1}{2} B_1 \dot{y_1}^2 - \frac{1}{2} B_1 \dot{y_2}^2 \right) - F \dot{y_1} \right)}{\dot{y}} = 0$$

$$\frac{d}{dt}\left(2\cdot\frac{1}{2}m_1\dot{y_1}^{2-1}\right) - \left(-\left(2\cdot\frac{1}{2}k_1y_1^{2-1} - 2\cdot\frac{1}{2}k_1y_2^{2-1}\right)\right) + \left(\left(2\cdot\frac{1}{2}B_1\dot{y_1}^{2-1} - 2\cdot\frac{1}{2}B_1\dot{y_2}^{2-1}\right) - F\dot{y_1}^{1-1}\right) = 0$$

$$\frac{d}{dt}\left(m_1\dot{y_1}\right) - \left(-\left(k_1y_1 - k_1y_2\right)\right) + \left(\left(B_1\dot{y_1} - B_1\dot{y_2}\right) - F\right) = 0$$

$$\frac{d}{dt}\left(m_1\dot{y_1}\right) - \left(-k_1y_1 + k_1y_2\right) + \left(B_1\dot{y_1} - B_1\dot{y_2} - F\right) = 0$$

$$\frac{d}{dt}\left(m_1\dot{y_1}\right) + k_1y_1 - k_1y_2 + B_1\dot{y_1} - B_1\dot{y_2} - F = 0$$

$$m_1\ddot{y_1} + k_1y_1 - k_1y_2 + B_1\dot{y_1} - B_1\dot{y_2} - F = 0$$

Sorting yields the equation's final form:

$$m_1\ddot{y_1} + B_1\dot{y_1} + k_1y_1 - B_1\dot{y_2} - k_1y_2 = F$$

3.2 Forcing on m_1 - subsystem 1b

The kinetic and potential energies of subsystem 1b are defined as:

$$T_2 = \frac{1}{2}m_2\dot{y_2}^2$$

$$U_2 = -k_2 \int_0^{y_2} y dy + k_1 \int_{y_2}^{y_1} y dy = \frac{1}{2}k_2y_2^2 + \frac{1}{2}k_1y_2^2 - \frac{1}{2}k_1y_1^2$$

The system also has a damper, which is a dissipative element:

$$D_2 = -B_2 \int_0^{\dot{y_2}} \dot{y} dy + B_1 \int_{\dot{y_2}}^{\dot{y_1}} \dot{y} dy = \frac{1}{2} B_2 \dot{y_2}^2 + \frac{1}{2} B_1 \dot{y_2}^2 - \frac{1}{2} B_1 \dot{y_1}^2$$

The Lagrangian of the subsystem 1a is:

$$\mathcal{L}(y,\dot{y}) = T_2 - U_2 = \frac{1}{2}m_2\dot{y}_2^2 - \left(\frac{1}{2}k_2y_2^2 + \frac{1}{2}k_1y_2^2 - \frac{1}{2}k_1y_1^2\right)$$

The Rayleigh dissipation function can be written as:

$$\mathcal{D} = D_2 - P_2 = \left(\frac{1}{2}B_2\dot{y}_2^2 + \frac{1}{2}B_1\dot{y}_2^2 - \frac{1}{2}B_1\dot{y}_1^2\right) - 0$$

Using Lagrange's equation with a non-conservative force (damper), the equation expands into the following form:

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{q}} \right) - \frac{\partial \mathcal{L}}{\partial q} + \frac{\partial \mathcal{D}}{\partial \dot{q}} = 0$$

This leads to the following differential equation:

$$\frac{d}{dt} \left(\frac{\partial \left(\frac{1}{2} m_2 \dot{y}_2^2 - \left(\frac{1}{2} k_2 y_2^2 + \frac{1}{2} k_1 y_2^2 - \frac{1}{2} k_1 y_1^2 \right) \right)}{\dot{y}} \right) - \frac{\partial \left(\frac{1}{2} m_2 \dot{y}_2^2 - \left(\frac{1}{2} k_2 y_2^2 + \frac{1}{2} k_1 y_2^2 - \frac{1}{2} k_1 y_1^2 \right) \right)}{y} + \frac{\partial \left(\left(\frac{1}{2} B_2 \dot{y}_2^2 + \frac{1}{2} B_1 \dot{y}_2^2 - \frac{1}{2} B_1 \dot{y}_1^2 \right) - 0 \right)}{\dot{y}} = 0$$

$$\frac{d}{dt} \left(\frac{\partial \left(\frac{1}{2} m_2 \dot{y}_2^2 - 0 \right)}{\dot{y}} \right) - \frac{\partial \left(0 - \left(\frac{1}{2} k_2 y_2^2 + \frac{1}{2} k_1 y_2^2 - \frac{1}{2} k_1 y_1^2 \right) \right)}{y} + \frac{\partial \left(\left(\frac{1}{2} B_2 \dot{y}_2^2 + \frac{1}{2} B_1 \dot{y}_2^2 - \frac{1}{2} B_1 \dot{y}_1^2 \right) - 0 \right)}{\dot{y}} = 0$$

$$\frac{d}{dt}\left(\left(2\cdot\frac{1}{2}m_2\dot{y_2}^{2-1}\right)\right) - \left(-\left(2\cdot\frac{1}{2}k_2y_2^{2-1} + 2\cdot\frac{1}{2}k_1y_2^{2-1} - 2\cdot\frac{1}{2}k_1y_1^{2-1}\right)\right) + \left(\left(2\cdot\frac{1}{2}B_2\dot{y_2}^{2-1} + 2\cdot\frac{1}{2}B_1\dot{y_2}^{2-1} - 2\cdot\frac{1}{2}B_1\dot{y_1}^{2-1}\right)\right) = 0$$

$$\frac{d}{dt}\left(m_2\dot{y_2}\right) - \left(-\left(k_2y_2 + k_1y_2 - k_1y_1\right)\right) + \left(B_2\dot{y_2} + B_1\dot{y_2} - B_1\dot{y_1}\right) = 0$$

$$\frac{d}{dt}(m_2\dot{y_2}) + k_2\dot{y_2} + k_1\dot{y_2} - k_1\dot{y_1} + B_2\dot{y_2} + B_1\dot{y_2} - B_1\dot{y_1} = 0$$

$$m_2\ddot{y}_2 + k_2y_2 + k_1y_2 - k_1y_1 + B_2\dot{y}_2 + B_1\dot{y}_2 - B_1\dot{y}_1 = 0$$

Sorting and grouping yield the final form of the equation:

$$m_2\ddot{y_2} + (B_2 + B_1)\dot{y_2} + (k_2 + k_1)y_2 - B_1\dot{y_1} - k_1y_1 = 0$$

Equation Results

Ultimately, the equation for the 2DoF system with forcing on m_1 can be written as:

$$\begin{cases}
m_1 \ddot{y}_1 + B_1 \dot{y}_1 + k_1 y_1 - B_1 \dot{y}_2 - k_1 y_2 = F \\
m_2 \ddot{y}_2 + (B_2 + B_1) \dot{y}_2 + (k_2 + k_1) y_2 - B_1 \dot{y}_1 - k_1 y_1 = 0
\end{cases}$$

3.3 Forcing on m_2 - subsystem 2a

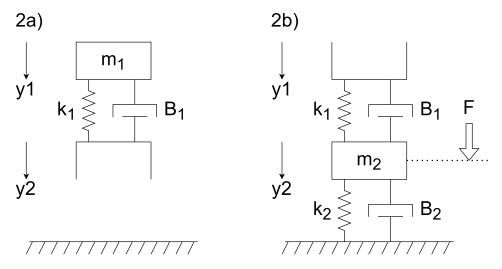


Figure 3: Uncoupled 2DoF system with forcing on m_2

The kinetic and potential energies of subsystem 2a are defined as:

$$T_1 = \frac{1}{2}m_1\dot{y}_1^2$$

$$U_1 = k_1 \int_{y_1}^{y_2} y \, dy = \frac{1}{2}k_1y_1^2 - \frac{1}{2}k_1y_2^2$$

The system also has a damper, which is a dissipative element:

$$D_1 = B_1 \int_{\dot{y_1}}^{y_2} \dot{y} dy = \frac{1}{2} B_1 \dot{y}_1^2 - \frac{1}{2} B_1 \dot{y}_2^2$$

The Lagrangian of subsystem 2a is:

$$\mathcal{L}(y, \dot{y}) = T_1 - U_1 = \frac{1}{2}m_1\dot{y_1}^2 - \left(\frac{1}{2}k_1y_1^2 - \frac{1}{2}k_1y_2^2\right)$$

The Rayleigh dissipation function can be written as:

$$\mathcal{D} = D_1 - P_1 = \left(\frac{1}{2}B_1\dot{y}_1^2 - \frac{1}{2}B_1\dot{y}_2^2\right) - 0$$

Using Lagrange's equation with a non-conservative force (damper), the equation expands into the following form:

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{q}} \right) - \frac{\partial \mathcal{L}}{\partial q} + \frac{\partial \mathcal{D}}{\partial \dot{q}} = 0$$

The solving process in this model is identical to the previous case. The step-by-step process will not be explained, resulting in the following solution of subsystem 2a:

$$m_1\ddot{y_1} + B_1\dot{y_1} + k_1y_1 - B_1\dot{y_2} - k_1y_2 = 0$$

3.4 Forcing on m_2 - subsystem 2b

The kinetic and potential energies of subsystem 1b are defined as:

$$T_2 = \frac{1}{2}m_2\dot{y}_2^2$$

$$U_2 = -k_2 \int_0^{y_2} y dy + k_1 \int_{y_2}^{y_1} y dy = \frac{1}{2}k_2y_2^2 + \frac{1}{2}k_1y_2^2 - \frac{1}{2}k_1y_1^2$$

The system also has a damper, which is a dissipative element:

$$D_2 = -B_2 \int_0^{\dot{y_2}} \dot{y} dy + B_1 \int_{\dot{y_2}}^{\dot{y_1}} \dot{y} dy = \frac{1}{2} B_2 \dot{y_2}^2 + \frac{1}{2} B_1 \dot{y_2}^2 - \frac{1}{2} B_1 \dot{y_1}^2$$

The Lagrangian of subsystem 2a is:

$$\mathcal{L}(y, \dot{y}) = T_2 - U_2 = \frac{1}{2}m_2\dot{y_2}^2 - \left(\frac{1}{2}k_2y_2^2 + \frac{1}{2}k_1y_2^2 - \frac{1}{2}k_1y_1^2\right)$$

The Rayleigh dissipation function can be written as:

$$\mathcal{D} = D_2 - P_2 = \left(\frac{1}{2}B_2\dot{y}_2^2 + \frac{1}{2}B_1\dot{y}_2^2 - \frac{1}{2}B_1\dot{y}_1^2\right) - F\dot{y}_2$$

Using Lagrange's equation with a non-conservative force (damper), the equation expands into the following form:

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{q}} \right) - \frac{\partial \mathcal{L}}{\partial q} + \frac{\partial \mathcal{D}}{\partial \dot{q}} = 0$$

As mentioned above, the step-by-step process will not be explained. The resulting solution of subsystem 2b is:

$$m_2\ddot{y_2} + (B_2 + B_1)\dot{y_2} + (k_2 + k_1)y_2 - B_1\dot{y_1} - k_1y_1 = F$$

Equation Results

Ultimately, the equation for the 2DoF system with forcing on m_2 can be written as:

$$\begin{cases}
m_1 \ddot{y}_1 + B_1 \dot{y}_1 + k_1 y_1 - B_1 \dot{y}_2 - k_1 y_2 = 0 \\
m_2 \ddot{y}_2 + (B_2 + B_1) \dot{y}_2 + (k_2 + k_1) y_2 - B_1 \dot{y}_1 - k_1 y_1 = F
\end{cases}$$

4 Co-energy and Energy Expressions

Mass

The general inertia co-energy for mass is:

$$W'_m = W_m = \int_0^v mv' \, dv' = \frac{1}{2} mv^2 = \frac{1}{2} m\dot{y}^2$$

In linear systems, energy and co-energy are equal.

Spring

The general energy stored in the spring k_1 is:

$$W_s' = W_s = \int_0^y ky' \, dy' = \frac{1}{2}ky^2$$

Since the spring is conservative (no energy loss), co-energy equals energy.

Damper

The general damper dissipates energy as friction:

$$W_d = \int_0^t Bv^2(\tau) \, d\tau = \frac{1}{2}Bv^2 = \frac{1}{2}B\dot{y}^2$$

Dampers do not store energy, so they do not have co-energy.

Using the co-energy equation with a non-conservative force, we can write:

$$\frac{d}{dt} \left[\frac{\partial T}{\dot{y}} \right] - \frac{\partial T}{\dot{y}} + \frac{\partial U}{\dot{y}} + \frac{\partial D}{\dot{y}} = f_i$$

4.1 Forcing on m_1 - subsystem 1a

We need to adjust the general co-energy expressions to 2DoF subsystems speciffic configuration:

$$\begin{split} W_m &= \frac{1}{2} m_1 \dot{y_1}^2 \\ W_s &= \frac{1}{2} k_1 y_1^2 - \frac{1}{2} k_1 y_2^2 \\ W_d &= \frac{1}{2} B_1 \dot{y}_1^2 - \frac{1}{2} B_1 \dot{y}_2^2 \end{split}$$

Finally, we can assign following assumptions:

$$T_1 = W_m$$

$$U_1 = W_s$$

$$D_1 = W_d$$

$$f_i = F$$

By substituting the coefficients into the equation, we obtain the following:

$$\frac{d}{dt} \left[\frac{\partial \left(\frac{1}{2} m_1 \dot{y}_1^2 \right)}{\dot{y}} \right] - \frac{\partial \left(\frac{1}{2} m_1 \dot{y}_1^2 \right)}{y} + \frac{\partial \left(\frac{1}{2} k_1 y_1^2 - \frac{1}{2} k_1 y_2^2 \right)}{y} + \frac{\partial \left(\frac{1}{2} B_1 \dot{y}_1^2 - \frac{1}{2} B_1 \dot{y}_2^2 \right)}{\dot{y}} = F$$

$$\frac{d}{dt} \left[\frac{\partial \left(\frac{1}{2} m_1 \dot{y}_1^2 \right)}{\dot{y}} \right] - 0 + \frac{\partial \left(\frac{1}{2} k_1 y_1^2 - \frac{1}{2} k_1 y_2^2 \right)}{y} + \frac{\partial \left(\frac{1}{2} B_1 \dot{y}_1^2 - \frac{1}{2} B_1 \dot{y}_2^2 \right)}{\dot{y}} = F$$

$$\frac{d}{dt} \left[m_1 \dot{y}_1 \right] + k_1 y_1 - k_1 y_2 + B_1 \dot{y}_1 - B_1 \dot{y}_2 = F$$

$$m_1 \ddot{y}_1 + B_1 \dot{y}_1 + k_1 y_1 - B_1 \dot{y}_2 - k_1 y_2 = F$$

4.2 Forcing on m_1 - subsystem 1b

In this case, we also need to adjust the general co-energy expressions to 2DoF 2nd subsystems speciffic configuration:

$$W_{m} = \frac{1}{2}m_{2}\dot{y}_{2}^{2}$$

$$W_{s} = \frac{1}{2}k_{2}y_{2}^{2}$$

$$W_{d} = \frac{1}{2}B_{2}\dot{y}_{2}^{2}$$

Finally, we can assign following assumptions:

$$T_2 = W_m$$

$$U_2 = W_s$$

$$D_2 = W_d$$

$$f_i = B_1 \Delta \dot{y} + k_1 \Delta y$$

By substituting the coefficients into the equation, we obtain the following:

$$\frac{d}{dt} \left[\frac{\partial \left(\frac{1}{2} m_2 \dot{y}_2^2 \right)}{\dot{y}} \right] - \frac{\partial \left(\frac{1}{2} m_2 \dot{y}_2^2 \right)}{y} + \frac{\partial \left(\frac{1}{2} k_2 y_2^2 \right)}{y} + \frac{\partial \left(\frac{1}{2} B_2 \dot{y}_2^2 \right)}{\dot{y}} = B_1 \Delta \dot{y} + k_1 \Delta y$$

$$\frac{d}{dt} \left[\frac{\partial \left(\frac{1}{2} m_2 \dot{y}_2^2 \right)}{\dot{y}} \right] - 0 + \frac{\partial \left(\frac{1}{2} k_2 y_2^2 \right)}{\dot{y}} + \frac{\partial \left(\frac{1}{2} B_2 \dot{y}_2^2 \right)}{\dot{y}} = B_1 (\dot{y}_1 - \dot{y}_2) + k_1 (y_1 - y_2)$$

$$m_2 \ddot{y}_2 + k_2 y_2 + B_2 \dot{y}_2 = B_1 \dot{y}_1 - B_1 \dot{y}_2 + k_1 y_1 - k_1 y_2$$

$$m_2 \ddot{y}_2 + k_2 y_2 + B_2 \dot{y}_2 = B_1 \dot{y}_1 - B_1 \dot{y}_2 + k_1 y_1 - k_1 y_2$$

$$m_2 \ddot{y}_2 + (B_2 + B_1) \dot{y}_2 + (k_2 + k_1) y_2 - B_1 \dot{y}_1 - k_1 y_1 = 0$$

Equation Results

Ultimately, the equation for the 2DoF system with forcing on m_1 , solved via co-energy expression can be written as:

$$\begin{cases}
m_1 \ddot{y_1} + B_1 \dot{y_1} + k_1 y_1 - B_1 \dot{y_2} - k_1 y_2 = F \\
m_2 \ddot{y_2} + (B_2 + B_1) \dot{y_2} + (k_2 + k_1) y_2 - B_1 \dot{y_1} - k_1 y_1 = 0
\end{cases}$$

4.3 Forcing on m_2 - subsystem 2a

We need to adjust the general co-energy expressions to 2DoF subsystems speciffic configuration:

$$W_m = \frac{1}{2}m_1\dot{y_1}^2$$

$$W_s = \frac{1}{2}k_1y_1^2 - \frac{1}{2}k_1y_2^2$$

$$W_d = \frac{1}{2}B_1\dot{y}_1^2 - \frac{1}{2}B_1\dot{y}_2^2$$

Finally, we can assign following assumptions:

$$T_1 = W_m$$

$$U_1 = W_s$$

$$D_1 = W_d$$

$$f_i = 0$$

By substituting the coefficients into the equation, we obtain the following:

$$\frac{d}{dt} \left[\frac{\partial \left(\frac{1}{2} m_1 \dot{y}_1^2 \right)}{\dot{y}} \right] - \frac{\partial \left(\frac{1}{2} m_1 \dot{y}_1^2 \right)}{y} + \frac{\partial \left(\frac{1}{2} k_1 y_1^2 - \frac{1}{2} k_1 y_2^2 \right)}{y} + \frac{\partial \left(\frac{1}{2} B_1 \dot{y}_1^2 - \frac{1}{2} B_1 \dot{y}_2^2 \right)}{\dot{y}} = 0$$

$$\frac{d}{dt} \left[\frac{\partial \left(\frac{1}{2} m_1 \dot{y}_1^2 \right)}{\dot{y}} \right] - 0 + \frac{\partial \left(\frac{1}{2} k_1 y_1^2 - \frac{1}{2} k_1 y_2^2 \right)}{y} + \frac{\partial \left(\frac{1}{2} B_1 \dot{y}_1^2 - \frac{1}{2} B_1 \dot{y}_2^2 \right)}{\dot{y}} = 0$$

$$\frac{d}{dt} \left[m_1 \dot{y}_1 \right] + k_1 y_1 - k_1 y_2 + B_1 \dot{y}_1 - B_1 \dot{y}_2 = 0$$

$$m_1 \ddot{y}_1 + B_1 \dot{y}_1 + k_1 y_1 - B_1 \dot{y}_2 - k_1 y_2 = 0$$

4.4 Forcing on m_2 - subsystem 2b

In this case, we also need to adjust the general co-energy expressions to 2DoF 2nd subsystems speciffic configuration:

$$W_{m} = \frac{1}{2}m_{2}\dot{y}_{2}^{2}$$

$$W_{s} = \frac{1}{2}k_{2}y_{2}^{2}$$

$$W_{d} = \frac{1}{2}B_{2}\dot{y}_{2}^{2}$$

Finally, we can assign following assumptions:

$$T_2 = W_m$$

$$U_2 = W_s$$

$$D_2 = W_d$$

$$f_i = F + B_1 \Delta \dot{y} + k_1 \Delta y$$

By substituting the coefficients into the equation, we obtain the following:

$$\frac{d}{dt} \left[\frac{\partial \left(\frac{1}{2} m_2 \dot{y}_2^2 \right)}{\dot{y}} \right] - \frac{\partial \left(\frac{1}{2} m_2 \dot{y}_2^2 \right)}{y} + \frac{\partial \left(\frac{1}{2} k_2 y_2^2 \right)}{y} + \frac{\partial \left(\frac{1}{2} B_2 \dot{y}_2^2 \right)}{\dot{y}} = F + B_1 \Delta \dot{y} + k_1 \Delta y$$

$$\frac{d}{dt} \left[\frac{\partial \left(\frac{1}{2} m_2 \dot{y}_2^2 \right)}{\dot{y}} \right] - 0 + \frac{\partial \left(\frac{1}{2} k_2 y_2^2 \right)}{y} + \frac{\partial \left(\frac{1}{2} B_2 \dot{y}_2^2 \right)}{\dot{y}} = F + B_1 (\dot{y}_1 - \dot{y}_2) + k_1 (y_1 - y_2)$$

$$m_2 \ddot{y}_2 + k_2 y_2 + B_2 \dot{y}_2 = F + B_1 \dot{y}_1 - B_1 \dot{y}_2 + k_1 y_1 - k_1 y_2$$

$$m_2 \ddot{y}_2 + k_2 y_2 + B_2 \dot{y}_2 = F + B_1 \dot{y}_1 - B_1 \dot{y}_2 + k_1 y_1 - k_1 y_2$$

$$m_2 \ddot{y}_2 + k_2 y_2 + B_2 \dot{y}_2 = F + B_1 \dot{y}_1 - B_1 \dot{y}_2 + k_1 y_1 - k_1 y_2$$

$$m_2 \ddot{y}_2 + (B_2 + B_1) \dot{y}_2 + (k_2 + k_1) y_2 - B_1 \dot{y}_1 - k_1 y_1 = F$$

Equation Results

Ultimately, the equation for the 2DoF system with forcing on m_1 , solved via co-energy expression can be written as:

$$\begin{cases}
m_1 \ddot{y}_1 + B_1 \dot{y}_1 + k_1 y_1 - B_1 \dot{y}_2 - k_1 y_2 = 0 \\
m_2 \ddot{y}_2 + (B_2 + B_1) \dot{y}_2 + (k_2 + k_1) y_2 - B_1 \dot{y}_1 - k_1 y_1 = F
\end{cases}$$

5 Simulink Model for Lagrangian and Co-energy Expression - Forcing on m_1

In order to model the equation properly, we must translate the dynamic equation into a form understandable for Simulink:

$$\begin{cases} m_1 \ddot{y_1} + B_1 \dot{y_1} + k_1 y_1 - B_1 \dot{y_2} - k_1 y_2 = F \\ m_2 \ddot{y_2} + (B_2 + B_1) \dot{y_2} + (k_2 + k_1) y_2 - B_1 \dot{y_1} - k_1 y_1 = 0 \end{cases}$$

$$\begin{cases} \ddot{y_1} = \frac{1}{m_1} \left(F - B_1 \dot{y_1} - k_1 y_1 + B_1 \dot{y_2} + k_1 y_2 \right) \\ \ddot{y_2} = \frac{1}{m_2} \left(-(B_2 + B_1) \dot{y_2} - (k_2 + k_1) y_2 + B_1 \dot{y_1} + k_1 y_1 \right) \end{cases}$$

Below, the Simulink model is presented.

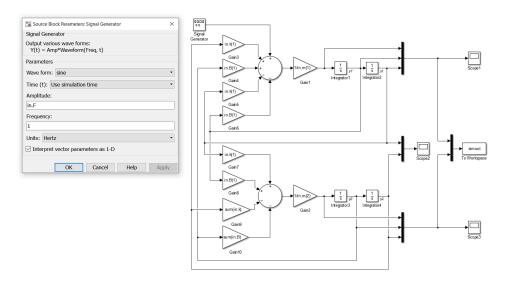


Figure 4: Simulink model of 2DoF system - Forcing on m_1

6 Simulink Model for Lagrangian and Co-energy Expression - Forcing on m_2

Same as above, we must translate the dynamic equation into a form understandable for Simulink:

$$\begin{cases} m_1 \ddot{y_1} + B_1 \dot{y_1} + k_1 y_1 - B_1 \dot{y_2} - k_1 y_2 = 0 \\ m_2 \ddot{y_2} + (B_2 + B_1) \dot{y_2} + (k_2 + k_1) y_2 - B_1 \dot{y_1} - k_1 y_1 = F \end{cases}$$

$$\begin{cases} \ddot{y_1} = \frac{1}{m_1} \left(-B_1 \dot{y_1} - k_1 y_1 + B_1 \dot{y_2} + k_1 y_2 \right) \\ \ddot{y_2} = \frac{1}{m_2} \left(F - (B_2 + B_1) \dot{y_2} - (k_2 + k_1) y_2 + B_1 \dot{y_1} + k_1 y_1 \right) \end{cases}$$

Below, the Simulink model is presented.

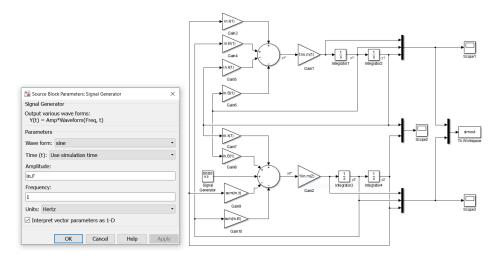


Figure 5: Simulink model of 2DoF system - Forcing on m_2

7 Matlab and Python script - forcing on m_1

For script-based system modeling we must take a slightly different approach to translating mathematical equation into code. Let $y_1(t) = y(1)$, $y'_1(t) = y(2)$, $y_2(t) = y(3)$, $y'_2(t) = y(4)$, in that case the matrix of the 2DoF equation will be written as:

$$\frac{d}{dt} \begin{bmatrix} y(1) \\ y(2) \\ y(3) \\ y(4) \end{bmatrix} = \begin{bmatrix} y(2) \\ (F - B(1) \cdot y(2) - k(1) \cdot y(1) + B(1) \cdot y(4) + k(1) \cdot y(3)) / m(1) \\ y(4) \\ (-(B(2) + B(1)) \cdot y(4) - (k(2) + k(1)) \cdot y(3) + B(1) \cdot y(2) + k(1) \cdot y(1)) / m(2) \end{bmatrix}$$

Now, the equation can be coded into Matlab:

```
Dy = [y(2); subsystem(1); y(4); subsystem(2)];
```

where:

```
subsystem1 = (F - B(1) * y(2) - k(1) * y(1) + B(1) * y(4) + k(1) * y(3))/m(1)

subsystem2 = (-(sum(B)) * y(4) - (sum(k)) * y(3) + B(1) * y(2) + k(1) * y(1))/m(2)
```

Python counts matrix elements from 0, so the equivalent Python code can be written as:

8 Matlab and Python script - forcing on m_2

For script-based system modeling we must take a slightly different approach to translating mathematical equation into code. Let $y_1(t) = y(1)$, $y'_1(t) = y(2)$, $y_2(t) = y(3)$, $y'_2(t) = y(3)$

y(4), in that case the matrix of the 2DoF equation will be written as:

$$\frac{d}{dt} \begin{bmatrix} y(1) \\ y(2) \\ y(3) \\ y(4) \end{bmatrix} = \begin{bmatrix} y(2) \\ (-B(1) \cdot y(2) - k(1) \cdot y(1) + B(1) \cdot y(4) + k(1) \cdot y(3)) / m(1) \\ y(4) \\ (F - (B(2) + B(1)) \cdot y(4) - (k(2) + k(1)) \cdot y(3) + B(1) \cdot y(2) + k(1) \cdot y(1)) / m(2) \end{bmatrix}$$

Now, the equation can be coded into Matlab:

```
Dy = [y(2); subsystem(1); y(4); subsystem(2)];
```

where:

```
subsystem1 = (-B(1) * y(2) - k(1) * y(1) + B(1) * y(4) + k

(1) * y(3))/m(1)

subsystem2 = (F-(sum(B)) * y(4) - (sum(k)) * y(3) + B(1) * y

(2) + k(1) * y(1))/m(2)
```

Python counts matrix elements from 0, so the equivalent Python code can be written as:

9 Simulation Results

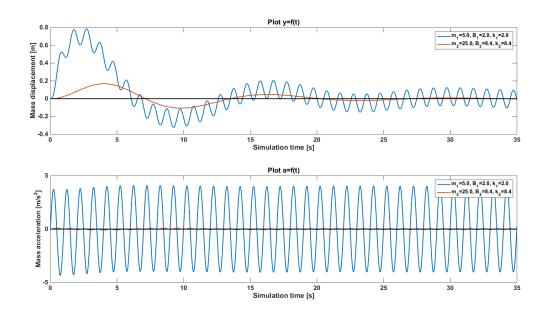


Figure 6: Forcing on m_1

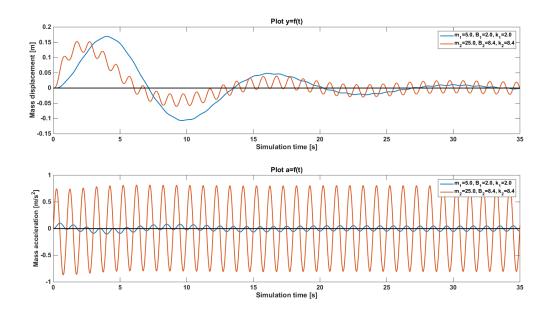


Figure 7: Forcing on m_2

10 Conclusion

This document demonstrates the Lagrangian and co-energy approach to modeling a 2DoF system. It also provides a visual block diagram suitable for simulation in Simulink and basic code for scripting.