) N).	(1). no assumption -> general (2) incompressible flow 15) v. v = 0 (3) inviscid flow 11=0.4	form. (, M = const. (Newtonian)
		/()

(konvective momentum flux) due to motion!

a) momentum flux through control volume,
$$\frac{M}{Lt^2} = \frac{kg}{ms^2}$$

b) external forces acting on the control volume due to normal and shear stress,
$$\frac{M}{L^2t^2} = \frac{kg}{m^2s^2} \left(\frac{momentum}{vate} \frac{ft}{momentum} \right)$$

per unit volume

C) local and convective acceleration, $\frac{L}{t^2} = \frac{m}{s^2}$

c) local and convective acceleration,
$$\frac{L}{t^2} = \frac{m}{s^2}$$

4) Vavg =
$$\frac{100}{\text{Tr}(0.175)^2} (\frac{1}{60})$$
, $M = \frac{1}{8} \times \frac{1.15 \times 10^6}{30} \times (0.175)^2 \times \frac{\text{Tr}(0.175)^2}{100} (\times \frac{60}{1})$
= 0,0878 dyne·min·cm²

Re=
$$\frac{0.75\times100\times1.261}{5.168\times\pi(0.175)^2\times60}$$
 = 1.03 \longrightarrow It's a liminar flow. (3)

According to incompressible, larminar flow \Rightarrow $\begin{cases} V_x = f(y) \\ V_y = 0 \end{cases}$ where only exist T_{yx} .

$$= \frac{dP}{dx} + \frac{dTyx}{dy} = 0 , Tyx = u(\frac{dVx}{dy}) , \frac{dP}{dx} = \frac{P_2 - P_1}{D}$$

$$= \begin{cases} y=0, & \forall x=0 \end{cases}$$

$$= \begin{cases} y=0, & \forall x=0 \end{cases}$$

$$= \begin{cases} y=\lambda L, & \forall x=0 \end{cases}$$

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We can also easily find that when y = L, $(R-P)I^{2}$

$$V_{\text{max}} = -\frac{(R_2 - R_1)L^2}{2MD}$$

$$Vavg = \frac{1}{A} \iint_{A} Vx dA = \frac{\int_{0}^{L} Ux dZ dY}{2 L UV} = \frac{(P_{s} - P_{s})L^{2}}{3 UP} = \frac{1}{3} V_{max}$$

•

which is due to incompressible laminar flow.

We can deal with Vr=0, Vz=0, $V_0 \neq 0$ in laminar flow, and $V_0 = f(r)$. Elimating the navier - Stokes Eq into 6 direction only, we get:

$$\frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial}{\partial r} (r V_{\theta}) \right) = 0 \quad \text{, know that } \begin{cases} D_{1} = \frac{R_{1}}{2} \\ D_{2} = \frac{R_{2}}{2} \end{cases}$$

$$=) V_0 = \frac{wr}{(1-\frac{R^2}{R^2})} - \frac{wR^2}{4(1-\frac{R^2}{R^2})} + \frac{4(1-\frac{R^2}{R^2})}{4(1-\frac{R^2}{R^2})} + \frac{4}{4}$$

b)
$$T_{r\theta} = M[r \frac{\partial}{\partial r} (\frac{V_{\theta}}{r})]$$

$$F_{z} = \iint_{0}^{2\pi} (T_{r\theta}|R_{z}) \frac{R_{z}}{2} d\theta dz = \frac{2\pi L W R_{z} M}{(1 - \frac{R_{z}^{2}}{R_{i}^{2}})} + \frac{1}{2} \frac{R_{z}^{2}}{R_{i}^{2}}$$

.



(1) $M\vec{v}\vec{v} = 0$ due to $V_{r=0}$, $V_{z=0}$, $V_{0} = \frac{\omega \vec{k}}{r} = f(r)$ in incompressible fluid.

$$(2) \begin{cases} V_{r} = 0 \\ V_{\theta} = \frac{WR^{2}}{r} = -\frac{2\Psi}{2r} = \frac{1}{r} \frac{2\Psi}{2\theta}$$

$$\Rightarrow \begin{cases} \overline{\Psi} = -\omega R^2 \ln r & \text{if } \\ \phi = \omega R^2 \theta & \text{if } \end{cases}$$



velocity potential