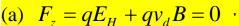
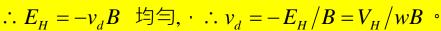
Hall Effect (Oct. 1879)

Maxwell 的書上說: \vec{B} 作用在導体上而非電流上,有無 \vec{B} 對電流分布無影響。 Hall 在 1877 年開始當 H. Rowland 的研究生。他懷疑上述說法,他想電流應被推到 導線的一側而降低有效截面積使電阻增加,但他無法測得此效應。Rowland 早先已 發現電線兩側間有微弱電壓,便建議 Hall 用金箔重作此實驗,而發現了 Hall effect。

 $\vec{v}_d = v_d \hat{x}, \quad \vec{B} = B \hat{y},$ $\vec{E}_H = E_H \hat{z}.$

 \vec{v}_d (& \vec{J} , \vec{E})均勻。





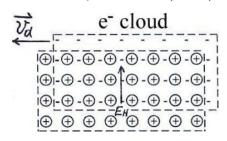
(b)
$$nq = J/v_d = J(-B/E_H) = (I/wt)(Bw/V_H) = IB/V_H t$$
.

 \therefore 量 V_H 可知 v_d & nq · 若 $V_H > 0$ · 則q > 0 。

Most metals: e^- . Metals $Co, Zn, Pb, Fe & semiconductors <math>Si, Ge: h^+$ (hole).

Force on a wire: e^- cloud 受磁力向上而拉正離子, 或說正離子受到向上 $\vec{E}_{_H}$ 的作用。

H.W.: Ex. 54; Prob. 1, 2, 4, 5, 6.



Ch. 30 Sources of the Magnetic Field

Oersted described his work ($I \Rightarrow \vec{B}$) to Paris Academy of Science in Sept. 1820. Biot-Savart announced in Oct. 1820 that $B \propto I/r$ for a long straight wire. (They measured the period of oscillation of a magnetized needle in \vec{B} field.)

Force between parallel wires

 $F_{21} = I_2 L B_1 = I_2 L (\mu_0 I_1 / 2\pi D) = (\mu_0 / 2\pi) I_1 I_2 L / D$.

 \therefore force per unit length $F/L = (\mu_0/2\pi)(I_1I_2/D)$.

Definition of 1 A: Let D = 1m & L = 1m,

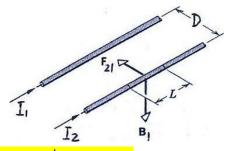
when $I_1 = I_2 = 1A$, $F = 2 \times 10^{-7} N$.

 $2 \times 10^{-7} N = (\mu_0/2\pi)(1A)^2(1m)/(1m)$, $\mu_0 = 4\pi \times 10^{-7} N/A^2$ exact.

(If D = 1 cm, L = 1 m, $I_1 = I_2 = 1A$, then $F = 2 \times 10^{-5} N = 2 \text{ dyne}$.)

 $1C \equiv 1A \cdot 1 \text{ sec}$, and the charge of electron was found to be $1.6 \times 10^{-19} C$.

 $1/4\pi \in 0$ was measured to be $9.0 \times 10^9 N \cdot m^2/C^2$.



$$1/\mu_0 \in_0 = (1/4\pi \in_0)/(\mu_0/4\pi) = (9.0 \times 10^9 \, N \cdot m^2/C^2)/(10^{-7} \, N \cdot \sec^2/C^2)$$
$$= 9.0 \times 10^{16} \, m^2/\sec^2 = \text{ (speed of light)}^2.$$

$d\vec{B}$ of a current element

For a long straight wire $B = \mu_0 I/2\pi r \leftrightarrow E = \lambda/2\pi \in_0 r$. $dB \propto 1/r^2$? Biot-Savart (Dec. 1820, after Laplace's hint): $dB = (\mu_0/4\pi) I d\ell \sin\theta/r^2$, $d\vec{B} = (\mu_0/4\pi) I d\vec{\ell} \times \hat{r}/r^2$.

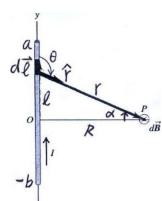
\vec{B} of a moving charge

$$Id\vec{\ell} = JAd\vec{\ell} = \vec{J}Ad\ell = nq\vec{v}_dAd\ell = (nqAd\ell)\vec{v}_d = (dQ)\vec{v}_d,$$

∴ $d\vec{B} = (\mu_0/4\pi)(dQ)\vec{v}_d \times \hat{r}/r^2$ · 但只適用於:

(a)緩慢且非加速;或(b)電流圈中的穩定電流(雖電荷有加速)。





$$B = (\mu_0 I/4\pi) \int d\ell \sin\theta / r^2$$
. $\sin\theta = \sin(\alpha + \pi/2) = \cos\alpha$,

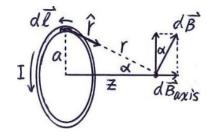
$$\ell = R \tan \alpha , \quad d\ell = R d\alpha / \cos^2 \alpha , \quad 1/r^2 = \cos^2 \alpha / R^2 .$$

$$B = (\mu_0 I / 4\pi) \int_{\alpha_1}^{\alpha_2} \cos \alpha d\alpha / R = (\mu_0 I / 4\pi R) (\sin \alpha_2 - \sin \alpha_1)$$

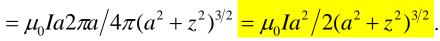
=
$$(\mu_0 I/4\pi R)[a/(R^2+a^2)^{1/2}+b/(R^2+b^2)^{1/2}].$$

- (1) a = b, $B = (\mu_0 I / 4\pi) 2a / R(R^2 + a^2)^{1/2}$.
- (2) $a,b \rightarrow \infty$, $B = \mu_0 I / 2\pi R$.
- (3) $b = 0, a \to \infty, B = \mu_0 I / 4\pi R$.

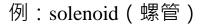
例: A circular loop



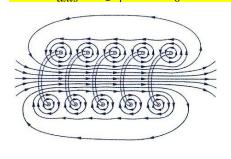
 $dec{eta} \quad dec{\ell} \perp ec{r} \;,\; \therefore dB = (\mu_0/4\pi)Id\ell/r^2 \;, \ dB_{axis} = dB\sinlpha \ = (\mu_0/4\pi)(Id\ell/r^2)(a/r) \;. \ B_{axis} = (\mu_0Ia/4\pi r^3)\int_0^{2\pi a}d\ell \;.$

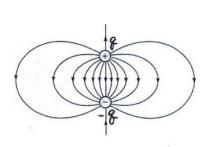


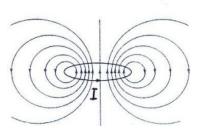
- (1) At z = 0, $B_{axis} = \mu_0 I/2a$. \leftrightarrow 直導線 $B = \mu_0 I/2a\pi$.
- (2) When $z \gg a$, $B_{axis} \approx \mu_0 I a^2 \pi / 2\pi z^3 = \mu_0 \mu / 2\pi z^3$ $\leftrightarrow E_{axis} \approx p / 2\pi \in_0 z^3$. (右圖)

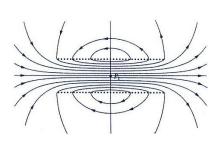


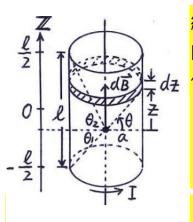












繞得很緊密,# of turns per unit length $n=N/\ell$ 。

由圓形電流圈的 B_{axis} 可知 $dB = \mu_0 (ndzI)a^2/2(a^2+z^2)^{3/2}$ 。

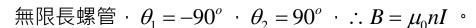
 $d \neq$ 代 $z = a \tan \theta + dz = a d\theta / \cos^2 \theta +$ 得

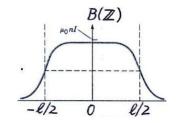
 $dB = [\mu_0 n I a^3 d\theta / \cos^2 \theta] / 2(a^2 + a^2 \tan^2 \theta)^{3/2}$

 $= [\mu_0 n I d\theta / \cos^2 \theta] / [2/\cos^3 \theta] = \mu_0 n I d\theta \cos \theta / 2.$

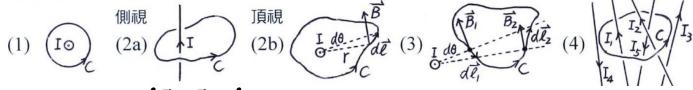
 $\therefore B = (\mu_0 nI/2) \int_{\theta}^{\theta_2} \cos \theta d\theta$

 $= (\mu_0 nI/2)[\sin \theta_2(Z) - \sin \theta_1(Z)].$





Ampere's law (Ampere 不喜歡 Biot-Savart 的工作,因為 (a) 實驗沒精確到足以宣稱 $\sin \theta$; (b) 不存在 isolated current element,它永遠是線路的一部份) 先考慮無限長直導線電流 I 與積分迴路 C。



- (1) C 正圓形: $\oint_C \vec{B} \cdot d\vec{\ell} = \oint_C B d\ell = (\mu_0 I/2\pi r)2\pi r = \mu_0 I$ 。
- (2) C 任意形狀: $\vec{B} \cdot d\vec{\ell} = Bd\ell_{\parallel} = (\mu_0 I/2\pi r)rd\theta = (\mu_0 I/2\pi)d\theta$ · ind. of r · $\therefore \oint_C \vec{B} \cdot d\vec{\ell} = \mu_0 I$ 與正圓形相同(或= $(\mu_0 I/2\pi)\int_0^{2\pi} d\theta = \mu_0 I$)。
- (3) I 在 C 外部:::成對抵消 · :: $\oint_C \vec{B} \cdot d\vec{\ell} = 0$ (或 = $(\mu_0 I/2\pi) \int_{\theta_0}^{\theta_0} d\theta = 0$) °
- (4) 多電流: $\oint_C \vec{B} \cdot d\vec{\ell} = \oint_C (\sum \vec{B}_i) \cdot d\vec{\ell} = \sum_{i \notin C \setminus D} \mu_0 I_i = \mu_0 I_{encl} \cdot I_{encl}$ 是在 C 內的總電流。

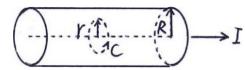
Ampere's law: $\oint_C \vec{B} \cdot d\vec{\ell} = \mu_0 I_{encl}$ for current loop & C of any shapes $\cdot I_{encl}$ 是穿過以C 為邊界的任何面的總電流 $\cdot I$ & C 的方向由右手決定 \cdot

此 law 可由 Biot-Savart $\vec{B}=(\mu_0/4\pi)q\vec{v}\times\hat{r}/r^2$ 導得。但後來發現此 \vec{B} 只適用於電荷緩慢且不加速時,或雖電荷有加速但是在電流圈的穩定電流中時,而 Ampere's law 卻適用於任何情況,因此應是由 Ampere's law 在特殊條件下導出 Biot-Savart law。



例:長圓柱中的均勻電流

$$\oint \vec{B} \cdot d\vec{\ell} = \mu_0 I_{encl}$$



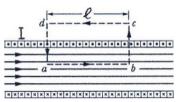
$$\Rightarrow B2\pi r = \begin{cases} \mu_0(Ir^2/R^2) \text{ for } r < R \\ \mu_0 I \text{ for } r > R \end{cases} \Rightarrow B(r) = \begin{cases} (\mu_0/2\pi)(I/R^2)r \text{ for } r < R \\ \mu_0 I/2\pi r \text{ for } r > R \end{cases}$$

例:極緊密的無窮長螺管(電流I、每單位長度繞n圈)

內部 \vec{B} 不能有徑向分量(否則有 $\oint_{S} \vec{B} \cdot d\vec{A} \neq 0$)·也不能有圓

切線分量(否則有 $\oint_C \vec{B} \cdot d\vec{\ell} \neq 0$ 但 C' 內無電流)。

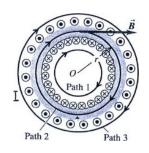
$$\oint_{C} \vec{B} \cdot d\vec{\ell} = \int_{a}^{b} \vec{B} \cdot d\vec{\ell} = B\ell = \mu_{0}(n\ell I) \implies B = \mu_{0}nI \text{ (與前同)}$$



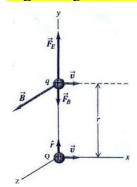
例:Toroid(電流I、共繞N圈)

$$\oint_C \vec{B} \cdot d\vec{\ell} = \mu_0 I_{encl} \Rightarrow B2\pi \ r = \mu_0(NI) \Rightarrow B(r) = \mu_0 NI/2\pi \ r.$$

當
$$r \to \infty$$
 · $N/2\pi$ $r = n$ 時 · $B = \mu_0 nI$ 。



 $\vec{F}_{\scriptscriptstyle F} \& \vec{F}_{\scriptscriptstyle B}$ between charged particles



Q在q處建立磁場 $\vec{B} = (\mu_0/4\pi)(Qv/r^2)\hat{z}$.

故 q 受磁力 $\vec{F}_B = qv\hat{x} \times \vec{B} = -(\mu_0/4\pi)(qQv^2/r^2)\hat{y}$ 。

而 q 受電力 $\vec{F}_E = q\vec{E} \approx (1/4\pi \in 0)(qQ/r^2)\hat{y}$

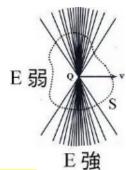
 $\therefore F_B/F_E = -\mu_0 \in_0 v^2 = -v^2/c^2 \quad \cdot$

但總力 $\vec{F}_E + \vec{F}_B = md^2\vec{r}/dt^2$ 與慣性座標無關,而在

它們的靜止座標中 $F_E = qQ/4\pi \in r^2$ · $F_B = 0$ ·

 $F_E + F_B = F_E + F_B = qQ/4\pi \in {}_{0} r^2$

故 $F_E = qQ/4\pi \in {}_0 r^2 - F_B \approx (1 + v^2/c^2) qQ/4\pi \in {}_0 r^2$ (右上圖)



H.W.: Ex. 9, 14, 21, 22; Prob. 1, 3, 4, 5, 8, 9.

Ch. 31 Electromagnetic Induction

Oersted found $I \Rightarrow B$ in 1820. Within weeks, electromagnet was found.

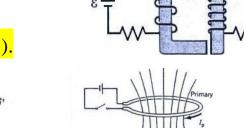
那磁場 B 能產生電流 I 嗎? Joseph Henry found in Aug. 1830 that a current was induced

by a changing magnetic field. But he did not publish immediately.

Magnetic flux $\phi_B \equiv \int_S \vec{B} \cdot d\vec{A}$ for any $\vec{B}(\vec{r},t)$ & surface S,

unit: 1 weber (W) $\equiv 1T \cdot 1m^2$.

Gauss's law for \vec{B} : $\oint_S \vec{B} \cdot d\vec{A} = 0$ (= q_M , but $q_M = 0$).



Faraday found (1831):

(1) Fixed coil, changing $\vec{B}(\vec{r},t)$ (右圖)

