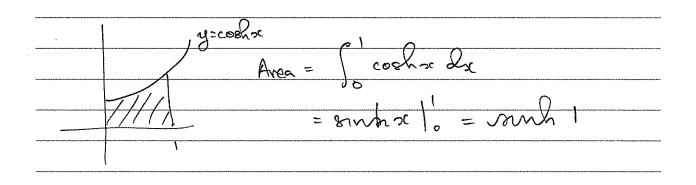
Extended Answer Section

There are four questions in this section, each with a number of parts. Write your answers in the space provided below each part. There is extra space at the end of the paper.

MARKS

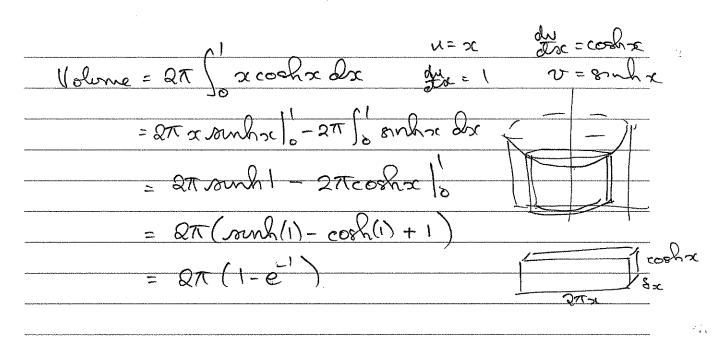
- 1. (a) Let D be the region of the plane bounded by the x-axis, the y-axis, the line x=1, and the curve $y=\cosh x$.
 - (i) Compute the area of D.

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(ii) Compute the volume of the solid obtained by rotating D about the y-axis

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(b) Let
$$I(x) = \int_0^x \sqrt{1+t^3} \, dt$$
. Calculate the integral
$$\int_0^1 x I(x) \, dx.$$

Note: The constant I(1) will appear in your answer.

L.	u = I(x)	ma = x	
	$\frac{du}{dx} = \sqrt{1+x^3}$	$19 = \frac{1}{2} x^2$	
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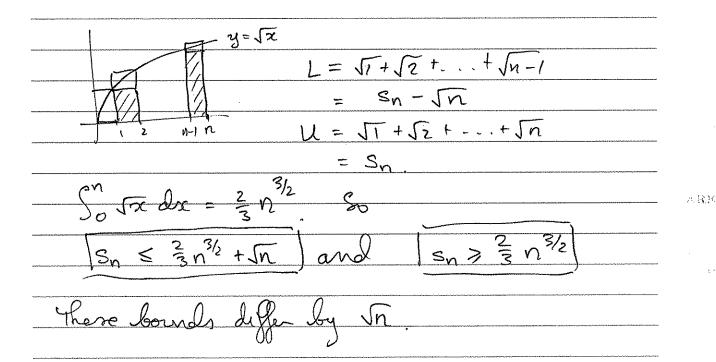
$$\int_{0}^{1} \alpha T(x) dx = \frac{1}{2} \alpha^{2} T(x) \Big|_{0}^{1} - \frac{1}{2} \int_{0}^{1} x^{2} \sqrt{1 + x^{2}} dx$$

$$U=1+\infty = \frac{1}{2}I(1)-\frac{1}{2}\cdot\frac{1}{3}\int_{0}^{0}\frac{du}{dx}\int U dx$$

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- (c) Let $s_n = \sqrt{1} + \sqrt{2} + \sqrt{3} + \dots + \sqrt{n}$.
 - (i) Let P be the partition of [0, n] into n subintervals of length 1. Use the corresponding upper and lower Riemann sums for the integral $\int_0^n \sqrt{x} \, dx$ to find upper and lower bounds for s_n , such that the bounds differ by at most \sqrt{n} .



(ii) Hence, or otherwise, calculate the limit $\lim_{n\to\infty} \frac{s_n}{n^{3/2}}$.

By above: $\frac{2}{3} \leq \frac{s_n}{n^{3/2}} \leq \frac{2}{3} + n^{-1} \Rightarrow \frac{2}{3}$ and so by the sequence slow $\frac{s_n}{n \to \infty} = \frac{2}{n^{3/2}}$ 2. (a) (i) Use a comparison test to show that $\int_0^\infty \frac{e^x}{7 + 2\cosh(2x)} dx$ converges.

 $\frac{e^{x}}{7 + 2\cosh(2x)} \le \frac{e^{x}}{7 + e^{2x}}$ (since $\cosh(2x) > \frac{1}{2}e^{x}$)

 $\int_{0}^{\infty} e^{-x} dx = 1$ converges, we conclude $\int_{0}^{\infty} e^{x} dx$ converges by

(ii) Using an appropriate substitution, or otherwise, calculate the integral

 $\int_{0}^{1} \frac{xe^{\sqrt{1+x^2}}}{\sqrt{1+x^2}} \, dx.$

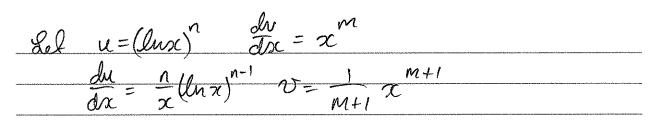
OR

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(b) (i) For integers
$$m,n\geq 0$$
 let $I_{m,n}=\int_0^1 x^m(\ln x)^n\,dx$. Show that for $n\geq 1$,
$$I_{m,n}=-\frac{n}{m+1}I_{m,n-1},$$

and hence compute $I_{m,n}$.

[You may use the fact that $\lim_{x\to 0^+} x^{\alpha} (\ln x)^{\beta} = 0$ for all $\alpha > 0$ and $\beta \geq 0$.]



$$I_{m,n} = \frac{x^{m+1}(\ln x)^n}{m+1} \frac{1}{10} = \frac{n}{m+1} \int_0^1 \frac{m}{x} (\ln x)^{n-1} dx$$

$$\frac{g_{0}}{f_{0}} = \frac{f_{0}}{f_{0}} = \frac{f_{0}}{$$

(ii)Hence show that

$$\int_0^1 x^{-x} \, dx = \sum_{k=1}^\infty n^{-n}.$$

 $\int_0^1 x^{-x}\,dx = \sum_{k=1}^\infty n^{-n}.$ You may assume that any reasonable series manipulations are valid.

$\int_{0}^{1} \frac{-x}{x} dx = \int_{0}^{1} \frac{-x \ln x}{e} dx$
$= \int_{0}^{1/2} \left(\frac{2}{k} - \frac{(-1)^{k}}{2} \alpha^{k} (\ln x)^{k}\right) dx$
(annué this $\sum_{k=0}^{\infty} \frac{(-1)^k}{k!} \int_0^1 x \left(\frac{k!}{k!} \right) dx$
(annué this $\sum_{k=0}^{\infty} \frac{(-1)^k}{k!} \int_0^1 x (lnx)^k dx$
$\frac{8}{(1)^{k}} = \frac{8}{(1)^{k}} \frac{(1)^{k}}{(1)^{k+1}} \frac{k!}{(1)^{k+1}}$
$\sum_{k=0}^{\infty} \frac{1}{(k+1)^{k+1}}$
$k = o(k+1)^{n}$ $= \sum_{k=0}^{\infty} n^{k}$

3. (a) Find the general solution of the differential equation

$$\frac{d^2y}{dx^2} + 3\frac{dy}{dx} + 2y = 3e^{-2x}$$

(1) Homogeneous soln: $\chi^2 + 3\chi + 2 = 0$ $\chi = -2, -1$

 $y_h(x) = Ae^{-2x} + Be^{-x}$; A, B constants

2) Particolar volutio: This is resonant

yp(α) = Cαe-201

 $\frac{4}{4}\frac{dy}{dx} = \frac{-2x}{ce} - 2cxe = c(1-2x)e^{-2x}$

 $\frac{d^{2}q}{dx^{2}} = -2ce^{-2x} -2x$ $= 4c(x-1)e^{-2x}$

60 + Cx - 4C + 3C - 6Cx + 2Cx = 3

& -C = 3, C = -3.

4p(x) = -3e-2x

(3) General soon:

y(x) = Ae + Be - 3xe

(b) Find the general solution of

$$\frac{dy}{dx} = \frac{2x+1}{x^2+x+1}(1-y),$$

and show that every solution converges to the equilibrium solution y=1 for $x\to\infty$.

Ossession Assoliumen emission American In Assolium emission emissi	$\frac{1}{1-y} \frac{dy}{dx} = \frac{2x+1}{x^2+x+1}$ (separable)
	$\int_{1-y}^{1-y} dy = \int_{2x+1}^{2x+1} \frac{2x+1}{x^2+x+1} dx$
	$-\ln 1-y = \ln x^2+x+1 + C$
	$\frac{-\ln x^2+x+1 -c}{e}$
underen pårere hurere er	e : [sc2+x+1]
	$= e^{-C} \frac{1}{x^2 + x + 1} \left(\frac{x^2 + x + 1}{x + 1} > 0 \right)$
akudan akasan makasan ma	$= e^{-C} \frac{1}{x^2 + x + 1} (x^2 + x + 1 > 0)$ $1 - y = \pm e^{-C}$ $x^2 + x + 1$
or water new arrangement and the second	$\Rightarrow y = 1 + \frac{A}{x^2 + x + 1} \text{for } A \text{ a constant}$
as	$x \to \infty$ we have $y \to 1$.

. .

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(c) Consider the differential equation of the form

$$\frac{dy}{dx} - e^{-x-y} + 1 = 0.$$

Introduce a new dependent variable u given by u = x + y, and hence find the general solution of the original equation.

u=α	$dy \Rightarrow dy = 1 + dy \Rightarrow dy = dx - 1$
Hence	•
www.vw.enhalenter.eseemanhaarannihaaranseehalababastaaran	$\frac{du}{dx} - 1 - e + 1 = 0$
	$\frac{du}{dsc} = e^{-u} \implies e^{u} \frac{du}{dsc} = 1.$
	dsc = 0 = 0 0 0 0 0 0 0 0 0 0
	u=a Sedu=x+C
	Je du = 2 + C
So	$e^{x} = x + C$, so $u = \ln x + C $
endre ndeskrudussuu uuska Audanudundanda sudassa	e - x + c, $m = x + c$
sserinnosuuduksuunikuuduunimakante	$y = \ln x+C - x$
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4. (a) A spherical raindrop evaporates at a rate proportional to its surface area, retaining the spherical shape. Derive a differential equation for the radius r(t) of the raindrop and solve it for a raindrop with initial radius r_0 to show that

$$r(t) = r_0 - \alpha t$$

for a constant $\alpha > 0$.

[Note that the volume of a sphere of radius r is $V=4\pi r^3/3$, and that the surface area is $A=4\pi r^2$, and assume that the density of water is 1.]

Set $m(t) = \frac{4}{3} \pi r(t)^3$ be the man of the drop Then $dm = -\infty$. LTC(L)²

giving 4TT (E) r(E) = - X 4TT r(E)

 $dr = -\alpha$

Thus r(t) = - at+C

at t=0, $r(0)=r_0$, so $C=r_0$.

Hence r(t)=10-xt as required

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(b) The evaporating raindrop is falling towards the ground. For this type of problem with time-dependent mass the appropriate form of Newton's second law states that the rate of change of the product of mass m with velocity v is equal to the force. The force is given by mg (with positive direction down), where g is the constant gravitational acceleration, with an additional air friction force proportional to the area πr^2 times the velocity. The friction force opposes the velocity. Show that the differential equation for the velocity v of the falling raindrop can be written as

$$\frac{dv}{dt} - \frac{k\alpha}{r(t)}v = g$$

for some constant k.

		M =	4 3 Tr3
$\frac{4}{3}\frac{d}{dt}\left(r^{3}v\right) =$	3gr - 2v		
So $2r^2v+r^3d$	W = gr3 - 3/2 r2		
So dr + (7)	$\frac{(x+2)}{\Gamma} v = g$		
So dr + (===================================	$\frac{102 = -2 + 36}{2} = \frac{-2 - 34 \cdot 36}{2}$		
		and the second s	

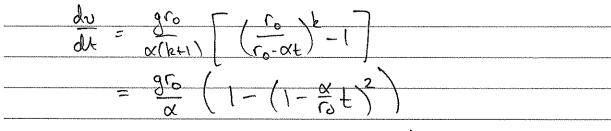
(c) Find the particular solution of the differential equation for the falling raindrop for which initially the raindrop is at rest. Assume that $k \neq -1$.

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$\frac{dv}{dt} = \frac{k\alpha}{r(t)}v = g$	V()) = U
$I(t) = e^{-\int r_0 - \alpha t} dt$	e e	(10-0xf 20)
	r(t)k	
Thus v= -kgr(t) dt		
$= g^{-k} \left(-\frac{1}{\alpha} \cdot \frac{1}{k+1} \right)$	rk+1 + C)	
g ~	n'-k	
= - g ~ (k+1) (+	C'r-k	
9_ ^	01-k	de Names and Carlos an
$0 = \frac{9}{\alpha(k+1)} r_{0} +$	C Po	ahda mananan mahambahaha ayar aya manahar yaman manana marana ayar ayar ayar ayar ayar ayar ayar
9 k+1		
$C' = \frac{9}{\alpha(k+1)} \sqrt{6}$		
$\Rightarrow v(t) = \frac{gr_0}{\alpha(k+1)}$		

(d) Assume that k=-2. Compute the distance the drop falls from rest until it is completely evaporated.

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$$v = \frac{gr_0}{\alpha} \left(t + \frac{2\alpha}{r_0} \left(1 - \frac{\alpha}{r_0} t \right) \right)$$

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V	\					

Then
$$v(t_{\infty}) = \frac{gr_0^2}{\alpha^2}$$
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