1 Operators and Integral Transforms

Definition 1.1 (Operator). Let A, and B be vector spaces with respective subspaces, X, and Y. An operator \mathcal{T} , maps any $x \in X$ to Y, and is denoted by $\mathcal{T}(x)$.

Common examples of operators are the Sturm-Liouville operator, the Laplacian, or Hamiltonian. Our focus will be on the integral operator, or transform. Let the domain, and co-domain of the transform be C[a,b] and $K: \mathbb{R}^2 \to \mathbb{R}$, then we can define our operator $\mathcal{T}: C[a,b] \to C[a,b]$ as

$$(\mathcal{T}f)(x) = \int_{a}^{b} f(y)K(x,y)dy,$$

where K is called the kernel function.

Theorem 1.2. Tf is continuous iff $\int_a^b |f(y)| dy < \infty$, and K(x,y) is uniformly continuous on [a,b].

Proof. For all $\epsilon > 0$, choose $\delta : |x - x_0| < \delta$, so that $|K(x, y) - K(x_0, y)| < \epsilon/M$, with $M = \int_a^b |f(y)| dy$. $\mathcal{T}f$ is continuous iff

$$|(\mathcal{T}f)(x) - (\mathcal{T}f)(x_0)| = \left| \int_a^b K(x, y) f(y) dy - \int_a^b K(x_0, y) f(y) dy \right|$$

$$\leq \int_a^b |K(x, y) - K(x_0, y)| |f(y)| dy$$

$$< \int_a^b \frac{\epsilon}{M} |f(y)| dy$$

$$< \epsilon$$

Corollary 1.3. The conditions for Theorem 1.2 are satisfied if a and b are finite, as well as if f and K are bounded and continuous.

Proof. A bounded continuous function is integrable over a compact domain.

1.1 Hankel Transform

Definition 1.4 (Hankel Transform). The Hankel transform of a function f(s) is given by

$$(\mathscr{H}_{\nu} f)(\sigma) = \int_{0}^{\infty} f(s) J_{\nu}(s\sigma) s \, ds,$$

where J_{ν} is the Bessel function of the first kind, of order $\nu \geq -\frac{1}{2}$, and σ is a non-negative real variable.

Corollary 1.5 (Inverse Hankel Transform). The Hankel transform is self-reciprocal, that is, the inverse Hankel transform is also given by Definition 1.4.

Proof. The Hankel transform is self-reciprocal

$$\iff f(s) = \int_0^\infty (\mathcal{H}_{\nu} f)(\sigma) J_{\nu}(s\sigma) \sigma \, d\sigma$$

$$\iff f(s) = \int_0^\infty \int_0^\infty f(s') J_{\nu}(s\sigma) s' \, ds' J_{\nu}(s\sigma) \sigma \, d\sigma$$

$$= \int_0^\infty f(s') s' \int_0^\infty J_{\nu}(s'\sigma) J_{\nu}(s\sigma) \sigma \, d\sigma \, ds'$$

$$= f(s),$$

by the orthogonality of the Bessel functions.

$$\varphi = \int_0^\infty J_0(kr)e^{kz}T(t)dk$$

$$\left(\frac{\partial \varphi}{\partial t} + (g\eta + \Phi)\right) \Big|_{z=0} = 0$$

$$\left(\frac{\partial^2 \varphi}{\partial t^2} + g\frac{\partial \varphi}{\partial z} + \frac{\partial \Phi}{\partial t}\right) \Big|_{z=0} = 0$$
where $\Phi = \frac{-Gm}{\sqrt{r^2 + (z - vt)^2}}$

$$\varphi = \frac{Gmv}{g} \int_0^\infty \frac{J_0(kr)e^{kz}}{1 + kv^2/g} \left(-\operatorname{sgn}(t)e^{-kv|t|} + 2\operatorname{H}(t)\operatorname{cos}(\omega_k t)\right) dk$$
with $\omega_k^2 = gk$

We must write the gravitational potential as an infinite sum of Bessel functions,

$$\begin{split} \frac{\partial \Phi}{\partial t} \bigg|_{z=0} &= \int_0^\infty a(k;t) J_0(kr) k \, dk \\ a(k;t) &= \int_0^\infty \frac{\partial \Phi}{\partial t} \bigg|_{z=0} J_0(kr) r \, dr \, ^*\mathrm{HT}^* \\ a(k;t) &= Gmv^2 tk \int_0^\infty \frac{J_0(kr) r}{(r^2 + v^2 t^2)^{3/2}} dr \\ &= Gmv^2 t \frac{1}{|vt|} e^{-k|vt|} \, ^*\mathrm{HT}^* \\ &= Gmv \, \mathrm{sgn}(t) e^{-kv|t|} \end{split}$$

$$\begin{split} \int_0^\infty J_0(kr)\ddot{T}(t)dk + g \int_0^\infty kJ_0(kr)T(t)dk + Gmv \int_0^\infty \operatorname{sgn}(t)e^{-kv|t|}J_0(kr)k\,dk &= 0 \\ \int_0^\infty \left[\frac{\ddot{T}(t)}{k} + gT(t) + Gmv \operatorname{sgn}(t)e^{-kv|t|}\right] J_0(kr)k\,dk &= 0 \\ \frac{\ddot{T}(t)}{k} + gT(t) + Gmv \operatorname{sgn}(t)e^{-kv|t|} &= \int_0^\infty 0 \times J_0(kr)r\,dr = 0 \text{ *HT*} \\ T(t) &= A\cos(\omega_k t) + B\sin(\omega_k t) \text{ homogenous solution } t \geq 0 \\ \operatorname{Assume} T(t) &= Ce^{\xi|t|} \text{ for the particular solution} \\ &\Rightarrow C\left(\xi^2 \operatorname{sgn}^2(t)e^{\xi|t|} + gke^{\xi|t|}\right) + Gmvk \operatorname{sgn}(t)e^{-kv|t|} &= 0 \\ &\Rightarrow \xi = -kv, \quad C = \frac{-Gmvk \operatorname{sgn}(t)}{k^2v^2 + gk} = \frac{Gmv}{g} \frac{-\operatorname{sgn}(t)}{1 + kv^2/g} \\ \Rightarrow T(t) &= \frac{Gmv}{g} \frac{1}{1 + kv^2/g} \left(-\operatorname{sgn}(t)e^{-kv|t|} + \operatorname{H}(t) \left(\tilde{A}\cos(\omega_k t) + \tilde{B}\sin(\omega_k t)\right)\right) \\ T(t) \text{ must be at least of class } C^1 \Rightarrow \tilde{A} = 2, \quad \tilde{B} = 0 \\ \Rightarrow \varphi &= \frac{Gmv}{g} \int_0^\infty \frac{J_0(kr)e^{kz}}{1 + kv^2/g} \left(-\operatorname{sgn}(t)e^{-kv|t|} + 2H(t)\cos(\omega_k t)\right) dk \end{split}$$