Primordial Black Holes

Brady Metheral

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Solution

Simulation Solution

The Effects of a Primordial Black Hole on the Surface of a Neutron Star

Brady Metherall

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Black Holes

Primordial Black Holes

Brady Metherall

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Analytic Solution

Simulation Solution

- A black hole is a region of space where enough mass is concentrated such that nothing, including light, can escape its gravitational pull.
- Typical black holes form from the collapse of a massive star, and have a minimum mass on the order of a few solar masses.
- In the early universe while it was very hot and dense, small perturbations could be enough to create small black holes called primordial black holes.
- The mass of a primordial black hole is probably much less than one solar mass $(10^{-18} 10^{-6} M_{\odot})$, and microscopic in size.

Primordial Black Holes as Dark Matter

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Dark matter is matter that does not emit/absorb light, and is chargeless. Dark matter accounts for about a quarter of the energy in the universe, but we don't know what it is.

Primordial black holes may be a good candidate for dark matter since: they are chargeless, do not emit light, have a very small radius, and are non-relativistic. Also, most other proposed explanations of dark matter involve creating a new particle not part of the Standard Model.

Flat Star Model

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Model

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Investigate the collision of a primordial black hole with a neutron star with the following assumptions:

- Neutron stars are flat and infinite
- Primordial black holes are point masses
- Neutron stars are incompressible fluids
- Gravitational interactions are Newtonian
- Primordial black hole has a constant velocity

Eigenfunctions of Laplacian

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Analytic Solution

Velocity Potential Deformation of

Solution

To find the velocity potential we need the eigenfunctions of the Laplacian. First, assume a product solution and solve the Laplacian.

$$\varphi = R(r)Z(z)\Theta(\theta)T(t)$$

$$\nabla^{2}\varphi = 0$$

$$\implies \varphi \propto \begin{cases} J_{\mu}(kr) \\ Y_{\mu}(kr) \end{cases} \begin{cases} e^{-kz} \\ e^{kz} \end{cases} \begin{cases} \sin(\mu\theta) \\ \cos(\mu\theta) \end{cases} T(t)$$

$$\implies \varphi \propto J_{0}(kr)e^{kz}T(t)$$

T(t) comes from boundary conditions.

Hankel Transform

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Analytic Solution

Velocity Potential Deformation or

Surface Energy The Hankel transform is needed to solve the time component of the velocity potential, and the calculation of the energy.

Definition

The Hankel transform of a function f(s) is given by

$$(\mathscr{H}_{\nu} f)(\sigma) = \int_{0}^{\infty} f(s) J_{\nu}(s\sigma) s \, ds,$$

where J_{ν} is the Bessel function of the first kind, of order $\nu \geq -\frac{1}{2}$, and σ is a non-negative real variable.

Corollary

The Hankel transform is self-reciprocal, that is, it is its own inverse.

Solving for φ

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Velocity Potential Deformation of Surface

Simulatio Solution From the pressure condition

$$\left. \left(\frac{\partial^2 \varphi}{\partial t^2} + g \frac{\partial \varphi}{\partial z} + \frac{\partial \Phi}{\partial t} \right) \, \right|_{z=0} = 0.$$

We can write the potential as a Hankel transform

$$\frac{\partial \Phi}{\partial t}\Big|_{z=0} = \int_0^\infty \left(\mathscr{H} \left. \frac{\partial \Phi}{\partial t} \right|_{z=0} \right) (k) J_0(kr) k \, dk.$$

Then,

$$\varphi = \frac{Gmv}{g} \int_0^\infty \frac{J_0(kr)e^{kz}}{1 + kv^2/g} \left(-\operatorname{sgn}(t)e^{-kv|t|} + 2\operatorname{H}(t)\cos(\omega_k t) \right) dk,$$

with $\omega_k^2 = gk$

Deformation of Surface

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We can easily solve for the shape of the surface now that we have the velocity potential, using

$$\left. \frac{\partial \varphi}{\partial z} \right|_{z=0} = \frac{\partial \eta}{\partial t}.$$

$$\eta = \frac{Gm}{g} \int_0^\infty \frac{J_0(kr)}{1 + kv^2/g} \left(e^{-kv|t|} + 2 H(t) v \sqrt{\frac{k}{g}} \sin(\omega_k t) \right) dk$$

Analytic Surface Waves

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Simulation

Figure: Waves!

Energy Transferred

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Analytic Solution Velocity Potential

Velocity Potential Deformation o Surface Energy

Simulatio Solution We can calculate the energy by taking the sum of the kinetic, and potential energies,

$$\begin{split} E(t) &= \frac{1}{2}\rho \int_{-\infty}^{0} \int_{0}^{\infty} |\nabla \varphi|^{2} \, r dr dz \int_{0}^{2\pi} d\theta \\ &+ \rho g \int_{0}^{\infty} \int_{0}^{\eta} z dz \, r dr \int_{0}^{2\pi} d\theta \\ \lim_{t \to \infty} E(t) &= 4\pi \rho \frac{G^{2} m^{2}}{q}. \end{split}$$

We take the limit as t approaches infinity, because we are interested in the total energy transferred.

Analytic Energy

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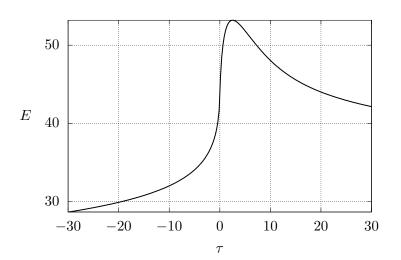
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Smooth Particle Hydrodynamics

Figure: Waves!

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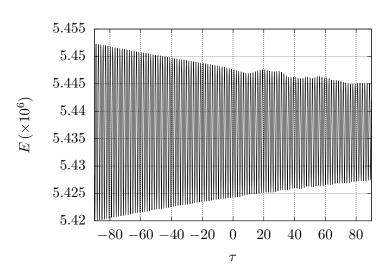
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