

Quantum Field Theory

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1 Standard Model

1.1 Electroweak Standard Model

Lagrangian with a global $SU(2) \times U(1)$ symmetry:

$$\mathcal{L} = i\bar{L}^{(l)}\gamma_\mu\partial^\mu L^{(l)} + i\bar{l}_R\gamma_\mu\partial^\mu l_R + \frac{1}{2}\partial_\mu\Phi^*\partial^\mu\Phi - m^2\Phi^*\Phi - \frac{1}{4}\lambda(\Phi^*\Phi)^2 - h_e\bar{L}^{(l)}\Phi e_R - \text{h.c.}$$

where $l = e, \mu, \tau$ and $a = 1, 2$, $l_{L,R} = \frac{1}{2}(1 \mp \gamma_5)l$ and

$$L^{(l)} = \begin{pmatrix} \nu_{(l)L} \\ l_L \end{pmatrix}$$

Local $SU(2) \times U(1)$ symmetry:

This consists of two things. First changing the partial derivatives to covariant ones:

$$\partial^\mu \rightarrow D^\mu = \partial^\mu - \frac{i}{2}g\tau_k A_k^\mu - \frac{i}{2}g'Y B^\mu$$

and second adding the kinetic terms

$$-\frac{1}{4}F_{\mu\nu}^a F^{a\mu\nu} - \frac{1}{4}B_{\mu\nu}B^{\mu\nu}$$

of the vector gauge particles to the lagrangian.

$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + g\epsilon^{abc}A_\mu^b A_\nu^c$$

$$B_{\mu\nu} = \partial_\mu B_\nu - \partial_\nu B_\mu$$

$$\Phi = e^{\frac{i}{v}\pi^a(x)\tau^a} \begin{pmatrix} 0 \\ \frac{1}{\sqrt{2}}(v + H(x)) \end{pmatrix}$$

This breaks the gauge invariance. The $\partial^\mu \pi^a$ are going to be added to A_μ^a so we can set $\pi_a = 0$ now.

1.1.1 Higgs Terms

$$\mathcal{L}_{Higgs} = \frac{1}{2}\partial_\mu \Phi^* \partial^\mu \Phi - m^2 \Phi^* \Phi - \frac{1}{4}\lambda(\Phi^* \Phi)^2$$

Plugging in the covariant derivatives and Φ in U-gauge (symmetry breaking):

$$\begin{aligned} \mathcal{L}_{Higgs} &= \frac{1}{2}\Phi^+ (\overleftarrow{\partial}_\mu + igA_\mu^a \frac{\tau^a}{2} + ig'Y B_\mu) (\overrightarrow{\partial}^\mu + igA^{a\mu} \frac{\tau^a}{2} + ig'Y B^\mu) \Phi - \lambda(\Phi^+ \Phi - \frac{v^2}{2})^2 = \\ &= \Phi_U^+ (\overleftarrow{\partial}_\mu + igA_\mu^a \frac{\tau^a}{2} + ig'Y B_\mu) (\overrightarrow{\partial}^\mu + igA^{a\mu} \frac{\tau^a}{2} + ig'Y B^\mu) \Phi_U - \lambda(\Phi_U^+ \Phi_U - \frac{v^2}{2})^2 = \\ &= \frac{1}{2}\partial_\mu H \partial^\mu H - \lambda v^2 H^2 - \lambda v H^3 - \frac{1}{4}\lambda H^4 + \\ &+ \frac{1}{8}(v+H)^2 \left(2g^2 \frac{A_\mu^1 + iA_\mu^2}{\sqrt{2}} \frac{A^{1\mu} - iA^{2\mu}}{\sqrt{2}} + (g^2 + 4Y^2 g'^2) \frac{gA_\mu^3 - 2Yg'B_\mu}{\sqrt{g^2 + 4Y^2 g'^2}} \frac{gA^{3\mu} - 2Yg'B^\mu}{\sqrt{g^2 + 4Y^2 g'^2}} \right) = \\ &= \frac{1}{2}\partial_\mu H \partial^\mu H - \lambda v^2 H^2 - \lambda v H^3 - \frac{1}{4}\lambda H^4 + \frac{1}{8}(v+H)^2 \left(2g^2 W_\mu^- W^{+\mu} + \frac{g^2}{\cos^2 \theta_W} Z_\mu Z^\mu \right) = \\ &= \frac{1}{2}\partial_\mu H \partial^\mu H - \lambda v^2 H^2 + \frac{1}{4}g^2 v^2 W_\mu^- W^{+\mu} + \frac{g^2 v^2}{8 \cos^2 \theta_W} Z_\mu Z^\mu - \lambda v H^3 - \frac{1}{4}\lambda H^4 + \\ &+ \frac{1}{2}v g^2 W_\mu^- W^{+\mu} H + \frac{g^2}{4 \cos \theta_W} v Z_\mu Z^\mu H + \frac{1}{4}g^2 W_\mu^- W^{+\mu} H^2 + \frac{g^2}{8 \cos \theta_W} Z_\mu Z^\mu H^2 \end{aligned}$$

Where we put

$$W_\mu^\pm = \frac{1}{\sqrt{2}}(A_\mu^1 \mp iA_\mu^2)$$

$$Z_\mu = \frac{g}{\sqrt{g^2 + 4Y^2 g'^2}} A_\mu^3 - \frac{2Yg'}{\sqrt{g^2 + 4Y^2 g'^2}} B_\mu$$

we defined θ_W by the relation

$$\cos \theta_W = \frac{g}{\sqrt{g^2 + 4Y^2 g'^2}}$$

so that the expressions simplify a bit, e.g. we now get:

$$\begin{aligned}\sin \theta_W &= \frac{2Y g'}{\sqrt{g^2 + 4Y^2 g'^2}} \\ Z_\mu &= \cos \theta_W A_\mu^3 - \sin \theta_W B_\mu \\ g^2 + 4Y^2 g'^2 &= \frac{g^2}{\cos^2 \theta_W}\end{aligned}$$

1.1.2 Yukawa terms

$$\begin{aligned}\mathcal{L}_{Yukawa} &= -h_e \bar{L} \Phi e_R - \text{h.c.} = -h_e \bar{L} \Phi_U e_R - \text{h.c.} = \\ &= -\frac{1}{\sqrt{2}} h_e (v + H) (\bar{e}_L e_R + \bar{e}_R e_L) = -\frac{1}{\sqrt{2}} h_e (v + H) \bar{e} e = \\ &= -\frac{1}{\sqrt{2}} h_e v \bar{e} e - \frac{1}{\sqrt{2}} h_e \bar{e} e H\end{aligned}$$

The term $\bar{L} \Phi e_R$ is $U(1)$ (hypercharge) invariant, so

$$-Y_L + Y + Y_R = 0$$

1.1.3 Leptonic Terms

$$\begin{aligned}\mathcal{L} &= i \bar{L} \gamma^\mu \partial_\mu L + i \bar{e}_R \gamma^\mu \partial_\mu e_R \rightarrow \\ &\rightarrow i \bar{L} \gamma^\mu (\partial_\mu - i g A_\mu^a \frac{\tau^a}{2} - i g' Y_L B_\mu) L + i \bar{e}_R \gamma^\mu (\partial_\mu - i g' Y_R B_\mu) e_R = \\ &= i \bar{L} \gamma^\mu \partial_\mu L + i \bar{e}_R \gamma^\mu \partial_\mu e_R + g \bar{L} \gamma^\mu \frac{\tau^a}{2} L A_\mu^a + g' Y_L \bar{L} \gamma^\mu L B_\mu + g' Y_R \bar{e}_R \gamma^\mu e_R B_\mu = \\ &= i \bar{L} \gamma^\mu \partial_\mu L + i \bar{e}_R \gamma^\mu \partial_\mu e_R + \frac{g}{\sqrt{2}} (\bar{\nu}_L \gamma^\mu e_L W_\mu^+ + \text{h.c.}) + \frac{1}{2} g \bar{L} \gamma^\mu \tau^3 L A_\mu^3 + g' Y_L \bar{L} \gamma^\mu L B_\mu + g' Y_R \bar{e}_R \gamma^\mu e_R B_\mu = \\ &= i \bar{\nu}_L \gamma^\mu \partial_\mu \nu_L + i \bar{e} \gamma^\mu \partial_\mu e + \frac{g}{\sqrt{2}} (\bar{\nu}_L \gamma^\mu e_L W_\mu^+ + \text{h.c.}) + \frac{1}{2} g \bar{\nu}_L \gamma^\mu \nu_L A_\mu^3 - \frac{1}{2} g \bar{e}_L \gamma^\mu e_L A_\mu^3 \\ &\quad + g' Y_L \bar{\nu}_L \gamma^\mu \nu_L B_\mu + g' Y_L \bar{e}_L \gamma^\mu e_L B_\mu + g' Y_R \bar{e}_R \gamma^\mu e_R B_\mu = \\ &= i \bar{\nu}_L \gamma^\mu \partial_\mu \nu_L + i \bar{e} \gamma^\mu \partial_\mu e + \frac{g}{\sqrt{2}} (\bar{\nu}_L \gamma^\mu e_L W_\mu^+ + \text{h.c.}) \\ &+ [(\frac{1}{2} g \sin \theta_W + Y_L g' \cos \theta_W) \bar{\nu}_L \gamma^\mu \nu_L + (-\frac{1}{2} g \sin \theta_W + Y_L g' \cos \theta_W) \bar{e}_L \gamma^\mu e_L + Y_R g' \cos \theta_W \bar{e}_R \gamma^\mu e_R] A_\mu \\ &+ [(\frac{1}{2} g \cos \theta_W - Y_L g' \sin \theta_W) \bar{\nu}_L \gamma^\mu \nu_L + (-\frac{1}{2} g \cos \theta_W - Y_L g' \sin \theta_W) \bar{e}_L \gamma^\mu e_L - 2 Y_L g' \sin \theta_W \bar{e}_R \gamma^\mu e_R] Z_\mu\end{aligned}$$

Where we substituted new fields Z_μ and A_μ for the old ones A_μ^3 and B_μ using the relation:

$$\begin{aligned}Z_\mu &= \cos \theta_W A_\mu^3 - \sin \theta_W B_\mu \\ A_\mu &= \sin \theta_W A_\mu^3 + \cos \theta_W B_\mu\end{aligned}$$

The angle θ_W must be the same as in the Higgs sector, so that the field Z_μ is the same. We now need to make the following requirement in order to proceed further:

$$Y = -Y_L$$

This follows for example by requiring that neutrinos have zero charge, i.e. setting $\frac{1}{2}g \sin \theta_W + Y_L g' \cos \theta_W = 0$ and substituting for θ_W from the definition (see the Higgs terms), from which one gets $Y = -Y_L$. From $-Y_L + Y + Y_R = 0$ we now get

$$Y_R = 2Y_L$$

it now follows:

$$\begin{aligned} \frac{1}{2}g \sin \theta_W + Y_L g' \cos \theta_W &= 0 \\ -\frac{1}{2}g \sin \theta_W + Y_L g' \cos \theta_W &= -g \sin \theta_W \\ Y_R g' \cos \theta_W &= -g \sin \theta_W \\ \tan \theta_W &= -2Y_L \frac{g'}{g} \end{aligned}$$

and the Lagrangian can be further simplified:

$$\begin{aligned} \mathcal{L} &= i\bar{\nu}_L \gamma^\mu \partial_\mu \nu_L + i\bar{e} \gamma^\mu \partial_\mu e + \frac{g}{\sqrt{2}} (\bar{\nu}_L \gamma^\mu e_L W_\mu^+ + \text{h.c.}) \\ &\quad - g \sin \theta_W (\bar{e}_L \gamma^\mu e_L + \bar{e}_R \gamma^\mu e_R) A_\mu \\ &\quad + \frac{g}{\cos \theta_W} \left[\frac{1}{2} \bar{\nu}_L \gamma^\mu \nu_L + \left(-\frac{1}{2} + \sin^2 \theta_W\right) \bar{e}_L \gamma^\mu e_L + \sin^2 \theta_W \bar{e}_R \gamma^\mu e_R \right] Z_\mu = \\ &= i\bar{\nu}_L \gamma^\mu \partial_\mu \nu_L + i\bar{e} \gamma^\mu \partial_\mu e + \frac{g}{2\sqrt{2}} (\bar{\nu} \gamma^\mu (1 - \gamma_5) e W_\mu^+ + \text{h.c.}) - g \sin \theta_W \bar{e} \gamma^\mu e A_\mu \\ &\quad + \frac{g}{2 \cos \theta_W} \left[\bar{\nu} \gamma^\mu (1 - \gamma_5) \nu + \bar{e} \gamma^\mu \left(-\frac{1}{2} + 2 \sin^2 \theta_W + \frac{1}{2} \gamma_5\right) e \right] Z_\mu \end{aligned}$$

Where we used the relations $\bar{\nu}_L \gamma^\mu e_L = \frac{1}{2} \bar{\nu} \gamma^\mu (1 - \gamma_5) e$ and $\bar{\nu}_R \gamma^\mu e_R = \frac{1}{2} \bar{\nu} \gamma^\mu (1 + \gamma_5) e$.

1.1.4 Gauge terms

$$\begin{aligned} \mathcal{L}_{Gauge} &= -\frac{1}{4} F_{\mu\nu}^a F^{a\mu\nu} - \frac{1}{4} B_{\mu\nu} B^{\mu\nu} = \\ &= -\frac{1}{4} (\partial_\mu A_\nu^a - \partial_\nu A_\mu^a + g \epsilon^{abc} A_\mu^b A_\nu^c) (\partial^\mu A^{a\nu} - \partial^\nu A^{a\mu} + g \epsilon^{ajk} A^{j\mu} A^{k\nu}) - \frac{1}{4} B_{\mu\nu} B^{\mu\nu} = \\ &= -\frac{1}{4} \partial_\mu A_\nu^a \partial^\mu A^{a\nu} - \frac{1}{4} B_{\mu\nu} B^{\mu\nu} - \frac{1}{2} (\partial_\mu A_\nu^a - \partial_\nu A_\mu^a) g \epsilon^{abc} A^{b\mu} A^{c\nu} - \frac{1}{4} g^2 \epsilon^{abc} \epsilon^{ajk} A_\mu^b A_\nu^c A^{k\mu} A^{l\nu} = \\ &= -\frac{1}{2} W_{\mu\nu}^- W^{+\mu\nu} - \frac{1}{4} A_{\mu\nu} A^{\mu\nu} - \frac{1}{4} Z_{\mu\nu} Z^{\mu\nu} - g [(\partial_\mu A_\nu^1 - \partial_\nu A_\mu^1) A^{2\mu} A^{3\nu} + \text{cycl. perm. (123)}] \\ &\quad - \frac{1}{4} g^2 [(A_\mu^a A^{a\mu})(A_\nu^b A^{b\nu}) - (A_\mu^a A_\nu^a)(A^{b\mu} A^{b\nu})] = \\ &= -\frac{1}{2} W_{\mu\nu}^- W^{+\mu\nu} - \frac{1}{4} A_{\mu\nu} A^{\mu\nu} - \frac{1}{4} Z_{\mu\nu} Z^{\mu\nu} - g [A_\mu^1 A_\nu^2 \overleftrightarrow{\partial}^\mu A^{3\nu} + \text{cycl. perm. (123)}] \\ &\quad - \frac{1}{4} g^2 [(A_\mu^a A^{a\mu})(A_\nu^b A^{b\nu}) - (A_\mu^a A_\nu^a)(A^{b\mu} A^{b\nu})] = \end{aligned}$$

$$\begin{aligned}
&= -\frac{1}{2}W_{\mu\nu}^-W^{+\mu\nu}-\frac{1}{4}A_{\mu\nu}A^{\mu\nu}-\frac{1}{4}Z_{\mu\nu}Z^{\mu\nu}-ig(W_\mu^0W_\nu^-\overleftrightarrow{\partial}^\mu W^{+\nu}+\text{cycl. perm. (0-+)}) \\
&-g^2[\frac{1}{2}(W_\mu^+W^{-\mu})^2-\frac{1}{2}(W_\mu^+W^{+\mu})(W_\nu^-W^{-\nu})+(W_\mu^0W^{0\mu})(W_\nu^+W^{-\nu})-(W_\mu^-W_\nu^+)(W^{0\mu}W^{0\nu})] = \\
&= -\frac{1}{2}W_{\mu\nu}^-W^{+\mu\nu}-\frac{1}{4}A_{\mu\nu}A^{\mu\nu}-\frac{1}{4}Z_{\mu\nu}Z^{\mu\nu}+\mathcal{L}_{WW\gamma}+L_{WWZ}+L_{WW\gamma\gamma}+L_{WWWW}+L_{WWZZ}+L_{WWZ\gamma}
\end{aligned}$$

Where $W_\mu^0 = A_\mu^3 = \cos\theta_W Z_\mu + \sin\theta_W A_\mu$ and:

$$\begin{aligned}
\mathcal{L}_{WW\gamma} &= -ig\sin\theta_W(A_\mu W_\nu^-\overleftrightarrow{\partial}^\mu W^{+\nu} + \text{cycl. perm. (A W}^- \text{W}^+)) \\
\mathcal{L}_{WWZ} &= -ig\cos\theta_W(Z_\mu W_\nu^-\overleftrightarrow{\partial}^\mu W^{+\nu} + \text{cycl. perm. (Z W}^- \text{W}^+)) \\
\mathcal{L}_{WW\gamma\gamma} &= -g^2\sin^2\theta_W(W_\mu^-W^{+\mu}A_\nu A^\nu - W_\mu^-A^\mu W_\nu^+A^\nu) \\
\mathcal{L}_{WWWW} &= \frac{1}{2}g^2(W_\mu^-W^{-\mu}W_\nu^+W^{+\nu} - W_\mu^-W^{+\mu}W_\nu^-W^{+\nu}) \\
\mathcal{L}_{WWZZ} &= -g^2\cos^2\theta_W(W_\mu^-W^{+\mu}Z_\nu Z^\nu - W_\mu^-Z^\mu W_\nu^+Z^\nu) \\
\mathcal{L}_{WWZ\gamma} &= g^2\sin\theta_W\cos\theta_W(-2W_\mu^-W^{+\mu}A_\nu Z^\nu + W_\mu^-Z^\mu W_\nu^+A^\nu + W_\mu^-A^\mu W_\nu^+Z^\nu)
\end{aligned}$$

1.1.5 GWS Lagrangian

Plugging everything together we get the GWS Lagrangian:

$$\begin{aligned}
\mathcal{L} &= \frac{1}{2}\partial_\mu H\partial^\mu H - \lambda v^2 H^2 + \frac{1}{4}g^2v^2W_\mu^-W^{+\mu} + \frac{g^2v^2}{8\cos^2\theta_W}Z_\mu Z^\mu - \lambda v H^3 - \frac{1}{4}\lambda H^4 + \\
&+ \frac{1}{2}vg^2W_\mu^-W^{+\mu}H + \frac{g^2}{4\cos\theta_W}vZ_\mu Z^\mu H + \frac{1}{4}g^2W_\mu^-W^{+\mu}H^2 + \frac{g^2}{8\cos\theta_W}Z_\mu Z^\mu H^2 \\
&- \frac{1}{\sqrt{2}}h_e v\bar{e}e - \frac{1}{\sqrt{2}}h_e \bar{e}e H \\
&- \frac{1}{2}W_{\mu\nu}^-W^{+\mu\nu} - \frac{1}{4}A_{\mu\nu}A^{\mu\nu} - \frac{1}{4}Z_{\mu\nu}Z^{\mu\nu} + \mathcal{L}_{WW\gamma} + L_{WWZ} + L_{WW\gamma\gamma} + L_{WWWW} + L_{WWZZ} + L_{WWZ\gamma} \\
&+ i\bar{\nu}_L\gamma^\mu\partial_\mu\nu_L + i\bar{e}\gamma^\mu\partial_\mu e + \frac{g}{2\sqrt{2}}(\bar{\nu}\gamma^\mu(1-\gamma_5)eW_\mu^+ + \text{h.c.}) - g\sin\theta_W\bar{e}\gamma^\mu eA_\mu \\
&+ \frac{g}{2\cos\theta_W}[\bar{\nu}\gamma^\mu(1-\gamma_5)\nu + \bar{e}\gamma^\mu(-\frac{1}{2} + 2\sin^2\theta_W + \frac{1}{2}\gamma_5)e]Z_\mu \\
&+ (e, \nu_e, h_e \leftrightarrow \mu, \nu_\mu, h_\mu) + (e, \nu_e, h_e \leftrightarrow \tau, \nu_\tau, h_\tau)
\end{aligned}$$

The free parameters are $g, \theta_W, v, \lambda, h_e, h_\mu, h_\tau$.

1.1.6 Particle Masses

The particle masses are deduced from the terms

$$\mathcal{L} = -\frac{1}{2}m_H^2 H^2 + m_W^2 W_\mu^- W^{+\mu} + \frac{1}{2}m_Z^2 Z_\mu Z^\mu - m_e \bar{e}e + \dots$$

comparing to the above:

$$\mathcal{L} = -\lambda v^2 H^2 + \frac{1}{4}g^2 v^2 W_\mu^- W^{+\mu} + \frac{g^2 v^2}{8 \cos^2 \theta_W} Z_\mu Z^\mu - \frac{1}{\sqrt{2}} h_e v \bar{e}e + \dots$$

we get

$$\begin{aligned} m_W &= \frac{1}{2}gv \\ m_Z &= \frac{gv}{2 \cos \theta_W} = \frac{m_W}{\cos \theta_W} \\ m_H &= v\sqrt{2\lambda} \\ m_e &= \frac{1}{\sqrt{2}}h_e v \end{aligned}$$

1.1.7 Dimensional Analysis

The evolution operator is dimensionless:

$$U(-\infty, \infty) = T \exp \left(\frac{i}{\hbar} \int_{-\infty}^{\infty} d^4x \mathcal{L}(x) \right)$$

So:

$$\left[\int_{-\infty}^{\infty} d^4x \mathcal{L}(x) \right] = [\hbar] = M^0$$

where M is an arbitrary mass scale. Length unit is M^{-1} , so then

$$[\mathcal{L}(x)] = M^4$$

For the particular forms of the Lagrangians above we get:

$$[m\bar{e}e] = [m^2 Z_\mu Z^\mu] = [m^2 H^2] = [i\bar{e}\gamma^\mu \partial_\mu e] = [\mathcal{L}] = M^4$$

so $[\bar{e}e] = M^3$, $[Z_\mu Z^\mu] = [H^2] = M^2$ and we get

$$[e] = [\bar{e}] = M^{\frac{3}{2}}$$

$$[Z_\mu] = [Z^\mu] = [H] = [\partial_\mu] = [\partial^\mu] = M^1$$

Example: what is the dimension of G_μ in $\mathcal{L} = -\frac{G_\mu}{\sqrt{2}}[\bar{\psi}_{\nu\mu}\gamma^\mu(1-\gamma_5)\psi_\mu][\bar{\psi}_e\gamma^\mu(1-\gamma_5)\psi_{\nu_e}]$? Answer:

$$\begin{aligned} [\mathcal{L}] &= [G_\mu \bar{\psi}\psi\bar{\psi}\psi] \\ M^4 &= [G_\mu]M^{\frac{3}{2}}M^{\frac{3}{2}}M^{\frac{3}{2}}M^{\frac{3}{2}} \\ [G_\mu] &= M^{-2} \end{aligned}$$

1.1.8 Quarks

$$\begin{aligned}\mathcal{L}_{fermion} += & \sum_{q=d,s,b} i\bar{L}_0^{(q)} \gamma^\mu \partial_\mu L_0^{(q)} + \sum_{q=d,u,s,c,b,t} i\bar{q}_{0R} \gamma^\mu \partial_\mu q_{0R} \\ \mathcal{L}_{Yukawa} += & - \sum_{\substack{q=d,s,b \\ q'=d,s,b}} h_{qq'} i\bar{L}_0^{(q)} \Phi q'_{0R} + \text{h.c.} - \sum_{\substack{q=d,s,b \\ q'=u,c,t}} \tilde{h}_{qq'} i\bar{L}_0^{(q)} \tilde{\Phi} q'_{0R} + \text{h.c.}\end{aligned}$$

More to be added here...

1.2 QFT

1.2.1 Evolution Operator, S-Matrix Elements

The evolution operator U is defined by the equations:

$$\begin{aligned}|\phi(t_2)\rangle &= U(t_2, t_1) |\phi(t_1)\rangle \\ i\hbar \frac{\partial U(t, t_1)}{\partial t} &= H(t) U(t, t_1) \\ U(t_1, t_1) &= 1\end{aligned}$$

We are interested in calculating the S matrix elements:

$$\langle f|U(-\infty, \infty)|i\rangle = \langle f|S|i\rangle = S_{fi}$$

so we first calculate $U(-\infty, \infty)$. Integrating the equation for the evolution operator:

$$U(t_2, t_1) = U(t_1, t_1) - \frac{i}{\hbar} \int_{t_1}^{t_2} H(t) U(t, t_1) dt = 1 - \frac{i}{\hbar} \int_{t_1}^{t_2} H(t) U(t, t_1) dt$$

Now:

$$\begin{aligned}S &= U(-\infty, \infty) = 1 - \frac{i}{\hbar} \int_{-\infty}^{\infty} H(t') U(t', -\infty) dt' = \\ &= 1 + \left(-\frac{i}{\hbar}\right) \int_{-\infty}^{\infty} H(t') U(t', -\infty) dt' + \left(-\frac{i}{\hbar}\right)^2 \int_{-\infty}^{\infty} \int_{-\infty}^{t'} H(t') H(t'') U(t'', -\infty) dt' dt'' = \\ &= \dots = \sum_{n=0}^{\infty} \left(-\frac{i}{\hbar}\right)^n \frac{1}{n!} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \dots T\{H(t_1) H(t_2) \dots\} dt_1 dt_2 \dots = \\ &= T \exp \left(-\frac{i}{\hbar} \int_{-\infty}^{\infty} H(t) dt \right) = T \exp \left(-\frac{i}{\hbar} \int_{-\infty}^{\infty} d^4x \mathcal{H}(x) \right)\end{aligned}$$

If \mathcal{L} doesn't contain derivatives of the fields, then $\mathcal{H} = -\mathcal{L}$ so:

$$U(-\infty, \infty) = T \exp \left(\frac{i}{\hbar} \int_{-\infty}^{\infty} d^4x \mathcal{L}(x) \right)$$

Let's write $S = 1 + iT$ and $|i\rangle = |k_1 \dots k_m\rangle$, $|f\rangle = |p_1 \dots p_n\rangle$. As a first step now, let's investigate a scalar field, e.g. $\mathcal{L} = -\frac{\lambda}{4} \phi^4$ (e.g. a Higgs self interaction term above), we'll look at other fields later:

$$\langle f|S|i\rangle = \langle f|iT|i\rangle = \langle p_1 \dots p_n | iT | k_1 \dots k_m \rangle = \frac{1}{\tilde{D}(k_1) \dots \tilde{D}(k_m)} \frac{1}{\tilde{D}(p_1) \dots \tilde{D}(p_n)}$$

$$\int d^4x_1 \cdots d^4x_m e^{-i(k_1x_1 + \cdots + k_mx_m)} \int d^4y_1 \cdots d^4y_n e^{+i(p_1y_1 + \cdots + p_ny_n)} G(x_1, \dots, x_m, y_1, \dots, y_n)$$

where

$$G(x_1, \dots, x_n) = \langle 0|T\{\phi(x_1) \cdots \phi(x_n)\}|0\rangle = \frac{\langle 0|T\{\phi_I(x_1) \cdots \phi_I(x_n) \exp\left(\frac{i}{\hbar} \int_{-\infty}^{\infty} d^4x \mathcal{L}(x)\right)\}|0\rangle}{\langle 0|T \exp\left(\frac{i}{\hbar} \int_{-\infty}^{\infty} d^4x \mathcal{L}(x)\right)|0\rangle}$$

This is called the LSZ formula. Now we use the Wick contraction, get some terms like $D_{23}D_{34}$ integrate things out, this will give the delta function and $\tilde{D}(p)$'s and that's it.

Let's see how it goes for $\mathcal{L} = -\frac{\lambda}{4}\phi^4$ for the process $k_1 + k_2 \rightarrow p_1 + p_2$:

$$\begin{aligned} \langle p_1 p_2 | S | k_1 k_2 \rangle &= \frac{\int d^4x_1 d^4x_2 e^{-i(k_1x_1 + k_2x_2)} \int d^4y_1 d^4y_2 e^{-i(p_1y_1 + p_2y_2)}}{\tilde{D}(k_1)\tilde{D}(k_2)\tilde{D}(p_1)\tilde{D}(p_2)} \\ &= \frac{\langle 0|T\{\phi_I(x_1)\phi_I(x_2)\phi_I(y_1)\phi_I(y_2) \exp\left(-\frac{i\lambda}{4\hbar} \int d^4x \phi_I^4(x)\right)\}|0\rangle}{\langle 0|T \exp\left(-\frac{i\lambda}{4\hbar} \int d^4x \phi_I^4(x)\right)|0\rangle} = \\ &= \frac{\int d^4x_1 d^4x_2 e^{-i(k_1x_1 + k_2x_2)} \int d^4y_1 d^4y_2 e^{-i(p_1y_1 + p_2y_2)}}{\tilde{D}(k_1)\tilde{D}(k_2)\tilde{D}(p_1)\tilde{D}(p_2)} \\ &\quad \left[\frac{\langle 0|T\{\phi_I(x_1)\phi_I(x_2)\phi_I(y_1)\phi_I(y_2)\}|0\rangle}{\langle 0|T \exp\left(-\frac{i\lambda}{4\hbar} \int d^4x \phi_I^4(x)\right)|0\rangle} + \right. \\ &\quad + \frac{\left(-\frac{i\lambda}{4\hbar}\right) \int d^4x \langle 0|T\{\phi_I(x_1)\phi_I(x_2)\phi_I(y_1)\phi_I(y_2)\phi_I^4(x)\}|0\rangle}{\langle 0|T \exp\left(-\frac{i\lambda}{4\hbar} \int d^4x \phi_I^4(x)\right)|0\rangle} + \\ &\quad \left. + \frac{\left(-\frac{i\lambda}{4\hbar}\right)^2 \int d^4x d^4y \langle 0|T\{\phi_I(x_1)\phi_I(x_2)\phi_I(y_1)\phi_I(y_2)\phi_I^4(x)\phi_I^4(y)\}|0\rangle}{\langle 0|T \exp\left(-\frac{i\lambda}{4\hbar} \int d^4x \phi_I^4(x)\right)|0\rangle} + \dots \right] = \\ &= \frac{1}{\tilde{D}(k_1)\tilde{D}(k_2)\tilde{D}(p_1)\tilde{D}(p_2)} \\ &\quad \left[(2\pi)^4 \delta^{(4)}(p_1 + p_2) (2\pi)^4 \delta^{(4)}(k_1 + k_2) \tilde{D}(p_1)\tilde{D}(k_1) + \right. \\ &\quad \left. (-i\lambda) 6(2\pi)^4 \delta^{(4)}(p_1 + p_2 - k_1 - k_2) \tilde{D}(k_1)\tilde{D}(k_2)\tilde{D}(p_1)\tilde{D}(p_2) + \right. \\ &\quad \left. (-i\lambda)(\text{disconnected terms with not enough } \tilde{D}(\dots)\text{'s}) + (-i\lambda)^2(\dots) + \dots \right] = \\ &= (2\pi)^4 \delta^{(4)}(p_1 + p_2 - k_1 - k_2) \left[6(-i\lambda) + 3(-i\lambda)^2 \int \frac{d^4k}{(2\pi)^4} \tilde{D}(k)\tilde{D}(p_1 + p_2 - k) + (-i\lambda)^3(\dots) + \dots \right] \end{aligned}$$

The denominator cancels with the disconnected terms. We used the Wick contractions (see below for a thorough explanation+derivation):

$$\begin{aligned} \langle 0|T\{\phi_I(x_1)\phi_I(x_2)\phi_I(y_1)\phi_I(y_2)\}|0\rangle &= D(x_1-x_2)D(y_1-y_2) + D(x_2-y_1)D(x_1-y_2) + D(x_2-y_2)D(x_1-y_1) \\ \langle 0|T\{\phi_I(x_1)\phi_I(x_2)\phi_I(y_1)\phi_I(y_2)\phi_I^4(x)\}|0\rangle &= D(x_1-x)D(x_2-x)D(y_1-x)D(y_2-x) + \text{disconnected} \\ \langle 0|T\{\phi_I(x_1)\phi_I(x_2)\phi_I(y_1)\phi_I(y_2)\phi_I^4(x)\phi_I^4(y)\}|0\rangle &= D(x_1-x)D(x_2-x)D(y_1-y)D(y_2-y)D(x-y)D(x-y) \\ &\quad + \text{disconnected} \end{aligned}$$

Where the "disconnected" terms are $D(x_1 - y_1)D(x_2 - y_2)D(x - x)D(x - x)$ and similar. When they are integrated over, they do not generate enough $\tilde{D}(p_1)$ propagators to cancel the propagators from the LSZ formula, which will cause the terms to vanish.

For the $\mathcal{L} = \phi^3(x)$ theory, one also needs the following contractions:

$$\langle 0|T\{\phi_I(x_1)\phi_I(x_2)\phi_I(y_1)\phi_I(y_2)\phi_I^3(x)\}|0\rangle = 0$$

$$\langle 0|T\{\phi_I(x_1)\phi_I(x_2)\phi_I(y_1)\phi_I(y_2)\phi_I^3(x)\phi_I^3(y)\}|0\rangle = D(x_1 - x)D(x_2 - x)D(x - y)D(y_1 - y)D(y_2 - y)$$

Thus it is clear that the only difference from the above is the factor $D(x - y)$ which after integrating changes to $\tilde{D}(p_1 + p_2)$ and this ends up in the final result.

One always gets the delta function in the result, so we define the matrix element \mathcal{M}_{fi} by:

$$S_{fi} = (2\pi)^4 \delta^{(4)}(p_1 + p_2 + \dots - k_1 - k_2 - \dots) i\mathcal{M}_{fi}$$

1.2.2 Wick Theorem

As seen above, we need to be able to calculate

$$\langle 0|T\{\phi_I(x_1) \dots \phi_I(x_n)\}|0\rangle$$

The Wick theorem says, that this is equal to all possible contractions of fields (all fields need to be contracted), where a contraction is defined as:

$$\langle 0|T\{\phi_I(x)\phi_I(y)\}|0\rangle \equiv D(x - y) = \int \frac{d^4p}{(2\pi)^4} \tilde{D}(p) e^{-ip(x-y)}$$

with

$$\tilde{D}(p) = \frac{i}{p^2 - m^2 + i\epsilon}$$

A few lowest possibilities:

$$\langle 0|T\{\phi_I(x_1)\}|0\rangle = 0$$

$$\langle 0|T\{\phi_I(x_1)\phi_I(x_2)\}|0\rangle = D_{12}$$

$$\langle 0|T\{\phi_I(x_1)\phi_I(x_2)\phi_I(x_3)\}|0\rangle = 0$$

$$\langle 0|T\{\phi_I(x_1)\phi_I(x_2)\phi_I(x_3)\phi_I(x_4)\}|0\rangle = \text{disconnected}$$

$$\langle 0|T\{\phi_I(x_1)\phi_I(x_2)\phi_I(x_3)\phi_I(x_4)\phi_I(x)\}|0\rangle = 0$$

$$\langle 0|T\{\phi_I(x_1)\phi_I(x_2)\phi_I(x_3)\phi_I(x_4)\phi_I^2(x)\}|0\rangle = \text{disconnected}$$

$$\langle 0|T\{\phi_I(x_1)\phi_I(x_2)\phi_I(x_3)\phi_I(x_4)\phi_I^3(x)\}|0\rangle = 0$$

$$\langle 0|T\{\phi_I(x_1)\phi_I(x_2)\phi_I(x_3)\phi_I(x_4)\phi_I^4(x)\}|0\rangle = 4! D(x_1 - x)D(x_2 - x)D(x_3 - x)D(x_4 - x) + \text{disconnected}$$

$$\langle 0|T\{\phi_I(x_1)\phi_I(x_2)\phi_I(x_3)\phi_I(x_4)\phi_I^3(x)\phi_I^3(y)\}|0\rangle =$$

$$= D(x_1 - x)D(x_2 - x)D(x - y)D(x_3 - y)D(x_4 - y) + \text{disconnected}$$

$$\langle 0|T\{\phi_I(x_1)\phi_I(x_2)\phi_I(x_3)\phi_I(x_4)\phi_I^4(x)\phi_I^4(y)\}|0\rangle =$$

$$= D(x_1 - x)D(x_2 - x)D(x - y)D(x - y)D(x_3 - y)D(x_4 - y) + \text{disconnected}$$

For the last two equations, not all possibilities of the connected graphs are listed (and also the combinatorial factor is omitted).

1.2.3 Fermions and Vector Bosons

For fermions:

$$\langle 0|T\{\psi_I(x)\bar{\psi}_I(y)\}|0\rangle \equiv S(x-y) = \int \frac{d^4p}{(2\pi)^4} \tilde{S}(p) e^{-ip(x-y)}$$

with

$$\tilde{S}(p) = \frac{i}{\not{p} - m + i\epsilon} = \frac{i(\not{p} + m)}{p^2 - m^2 + i\epsilon}$$

For vector bosons:

$$\langle 0|T\{A_\mu(x)A_\nu(y)\}|0\rangle \equiv D_{\mu\nu}(x-y) = \int \frac{d^4p}{(2\pi)^4} \tilde{D}_{\mu\nu}(p) e^{-ip(x-y)}$$

with

$$\tilde{D}_{\mu\nu}(p) = i \frac{-g_{\mu\nu} + \frac{p_\mu p_\nu}{m^2}}{p^2 - m^2 + i\epsilon}$$

For massless bosons:

$$\tilde{D}_{\mu\nu}(p) = i \frac{-g_{\mu\nu}}{p^2 + i\epsilon}$$

1.2.4 Feynman Rules

We can deduce a set of rules, so that one doesn't have to repeat the whole calculation each time. For a scalar field we derived the rules above, for fermion and vector boson fields it's more difficult.

1.2.5 ZZH interaction

Let's calculate the $\mathcal{L}_{ZZH} = \lambda Z_\mu Z^\mu H$ interaction in the SM, where $\lambda = \frac{g^2}{4 \cos \theta_W}$. Consider $H(p) \rightarrow Z(k) + Z(l)$:

$$\begin{aligned} \langle f|S|i\rangle &= \langle f|iT|i\rangle = \langle kl|iT|p\rangle = \frac{\varepsilon_\alpha(k)\varepsilon^\alpha(l)}{\tilde{D}_{\mu\nu}(k)\tilde{D}^{\mu\nu}(l)} \frac{1}{\tilde{D}(p)} \\ &= \frac{\varepsilon_\alpha(k)\varepsilon^\alpha(l)}{\tilde{D}_{\mu\nu}(k)\tilde{D}^{\mu\nu}(l)} \frac{1}{\tilde{D}(p)} \\ &= \int d^4x_1 e^{-ipx_1} \int d^4y_1 d^4y_2 e^{+i(ky_1+ly_2)} \langle 0|T\{Z_\mu(y_1)Z^\mu(y_2)H(x_1)\}|0\rangle = \\ &= \frac{\varepsilon_\alpha(k)\varepsilon^\alpha(l)}{\tilde{D}_{\mu\nu}(k)\tilde{D}^{\mu\nu}(l)} \frac{1}{\tilde{D}(p)} \\ &= \int d^4x_1 e^{-ipx_1} \int d^4y_1 d^4y_2 e^{+i(ky_1+ly_2)} \int d^4x i\lambda D_{\alpha\mu}(y_1-x) D^{\alpha\mu}(y_2-x) D(x_1-x) = \\ &= i\lambda(2\pi)^4 \delta^{(4)}(p-k-l) \varepsilon_\alpha(k)\varepsilon^\alpha(l) \end{aligned}$$

where we used the fact, that the only nonzero element of the Green function is

$$\int d^4x \langle 0|T\{Z_\alpha(y_1)Z^\alpha(y_2)H(x_1)Z_\mu(x)Z^\mu(x)H(x)\}|0\rangle$$

1.2.6 eeH interaction

Let's calculate the $\mathcal{L}_{eeH} = -\lambda \bar{e}eH$ interaction in the SM, where $\lambda = \frac{h_e}{\sqrt{2}}$. Consider $H(p) \rightarrow e^-(k) + e^+(l)$:

$$\begin{aligned}\langle f|S|i\rangle &= \langle f|iT|i\rangle = \langle kl|iT|p\rangle = \frac{\bar{u}(k)v(l)}{\tilde{S}(k)\tilde{S}(l)} \frac{1}{\tilde{D}(p)} \\ &\int d^4x_1 e^{-ipx_1} \int d^4y_1 d^4y_2 e^{+i(ky_1+ly_2)} \langle 0|T\{\bar{e}(y_1)e(y_2)H(x_1)\}|0\rangle = \\ &= \frac{\bar{u}(k)v(l)}{\tilde{S}(k)\tilde{S}(l)} \frac{1}{\tilde{D}(p)} \\ &\int d^4x_1 e^{-ipx_1} \int d^4y_1 d^4y_2 e^{+i(ky_1+ly_2)} \int d^4x (-i\lambda) S(y_1-x) S(y_2-x) D(x_1-x) = \\ &= (-i\lambda)(2\pi)^4 \delta^{(4)}(p-k-l) \bar{u}(k)v(l)\end{aligned}$$

where we used the fact, that the only nonzero element of the Green function is

$$\int d^4x \langle 0|T\{\bar{e}(y_1)e(y_2)H(x_1)\bar{e}(x)e(x)H(x)\}|0\rangle$$

1.2.7 ee gamma interaction

Let's calculate the $\mathcal{L}_{ee\gamma} = -\lambda \bar{e}\gamma^\mu e A_\mu$ interaction in the SM, where $\lambda = g \sin \theta_W$. Consider $\gamma(p) \rightarrow e^-(k) + e^+(l)$:

$$\begin{aligned}\langle f|S|i\rangle &= \langle f|iT|i\rangle = \langle kl|iT|p\rangle = \frac{\bar{u}(k)v(l)}{\tilde{S}(k)\tilde{S}(l)} \frac{\varepsilon_\mu(p)}{\tilde{D}_{\alpha\beta}(p)} \\ &\int d^4x_1 e^{-ipx_1} \int d^4y_1 d^4y_2 e^{+i(ky_1+ly_2)} \langle 0|T\{\bar{e}(y_1)e(y_2)A^\mu(x_1)\}|0\rangle = \\ &= \frac{\bar{u}(k)v(l)}{\tilde{S}(k)\tilde{S}(l)} \frac{\varepsilon_\mu(p)}{\tilde{D}_{\alpha\beta}(p)} \\ &\int d^4x_1 e^{-ipx_1} \int d^4y_1 d^4y_2 e^{+i(ky_1+ly_2)} \int d^4x (-i\lambda) S(y_2-x) \gamma^\mu S(y_1-x) D_\mu^\alpha(x_1-x) = \\ &= (2\pi)^4 \delta^{(4)}(p-k-l) \bar{u}(k)(-i\lambda)\gamma^\mu v(l)\varepsilon_\mu(p)\end{aligned}$$

where we used the fact, that the only nonzero element of the Green function is

$$\begin{aligned}\int d^4x \langle 0|T\{\bar{e}(y_1)e(y_2)A^\alpha(x_1)\bar{e}(x)\gamma^\mu e(x)A_\mu(x)\}|0\rangle = \\ = \pm S(y_2-x) \gamma^\mu S(y_1-x) D_\mu^\alpha(x_1-x)\end{aligned}$$

1.2.8 eeee interaction

Let's calculate the $\mathcal{L}_{ee\gamma} = -\lambda \bar{e}\gamma^\mu e A_\mu$ interaction in the SM, where $\lambda = g \sin \theta_W$. Consider $e^-(p_1) + e^+(p_2) \rightarrow \gamma(q) \rightarrow e^-(k_1) + e^+(k_2)$:

$$\begin{aligned}
\langle f|S|i\rangle &= \langle f|iT|i\rangle = \langle k_1 k_2 | iT | p_1 p_2 \rangle = \frac{\bar{u}(k_1)v(k_2)}{\tilde{S}(k_1)\tilde{S}(k_2)} \frac{\bar{v}(p_2)u(p_1)}{\tilde{S}(p_2)\tilde{S}(p_1)} \\
&\int d^4x_1 d^4x_2 e^{-i(p_1x_1+p_2x_2)} \int d^4y_1 d^4y_2 e^{+i(k_1y_1+k_2y_2)} \langle 0|T\{\bar{e}(y_1)e(y_2)\bar{e}(x_1)e(x_2)\}|0\rangle = \\
&= \frac{\bar{u}(k_1)v(k_2)}{\tilde{S}(k_1)\tilde{S}(k_2)} \frac{\bar{v}(p_2)u(p_1)}{\tilde{S}(p_2)\tilde{S}(p_1)} \\
&\int d^4x_1 d^4x_2 e^{-i(p_1x_1+p_2x_2)} \int d^4y_1 d^4y_2 e^{+i(k_1y_1+k_2y_2)} \int d^4x \int d^4y \\
&\quad (-i\lambda)^2 S(x_1-x)\gamma^\mu S(x_2-x)D_{\mu\nu}(x-y)S(y_1-y)\gamma^\nu S(y_2-y) = \\
&= (2\pi)^4 \delta^{(4)}(p_1+p_2-k_1-k_2) \bar{v}(p_2)(-i\lambda)\gamma^\mu u(p_1)\tilde{D}_{\mu\nu}(q=p_1+p_2=k_1+k_2)\bar{u}(k_1)(-i\lambda)\gamma^\nu v(k_2)
\end{aligned}$$

where we used the fact, that the only nonzero element of the Green function is

$$\begin{aligned}
&\int d^4x \int d^4y \langle 0|T\{\bar{e}(y_1)e(y_2)\bar{e}(x_1)e(x_2)\bar{e}(x)\gamma^\mu e(x)A_\mu(x)\bar{e}(y)\gamma^\nu e(y)A_\nu(y)\}|0\rangle = \\
&= \pm S(x_1-x)\gamma^\mu S(x_2-x)D_{\mu\nu}(x-y)S(y_1-y)\gamma^\nu S(y_2-y)
\end{aligned}$$

1.3 Low energy theories

1.3.1 Fermi-type theory

This is a low energy ($m_W^2 \gg m_\mu m_e$) model for the EW interactions, that can be derived for example from the muon decay:

$$\mu^- \rightarrow e^- + \nu_\mu + \bar{\nu}_e$$

From the SM the relevant Lagrangian is

$$\mathcal{L} = \frac{g}{2\sqrt{2}}(\bar{e}\gamma^\mu(1-\gamma_5)\nu_e W_\mu^-) + \frac{g}{2\sqrt{2}}(\bar{\mu}\gamma^\mu(1-\gamma_5)\nu_\mu W_\mu^-)$$

and one gets the diagram $\mu^- + \bar{\nu}_e \rightarrow e^- + \bar{\nu}_\mu$ and the corresponding matrix element:

$$iM = -i\frac{g^2}{8}[\bar{u}\gamma_\mu(1-\gamma_5)u]\frac{-g^{\mu\nu} + \frac{q^\mu q^\nu}{m_W^2}}{q^2 - m_W^2}[\bar{u}\gamma_\nu(1-\gamma_5)v]$$

which when the momentum transfer q is much less than m_w becomes

$$iM = -i\frac{g^2}{8m_W^2}[\bar{u}\gamma^\mu(1-\gamma_5)u][\bar{u}\gamma_\mu(1-\gamma_5)v]$$

but this element can be derived directly from the Lagrangian:

$$\mathcal{L} = -\frac{G_\mu}{\sqrt{2}}[\bar{\psi}_{\nu_\mu}\gamma^\mu(1-\gamma_5)\psi_\mu][\bar{\psi}_e\gamma^\mu(1-\gamma_5)\psi_{\nu_e}]$$

with

$$\frac{G_\mu}{\sqrt{2}} = \frac{g^2}{8m_W^2}$$

This is the universal V-A theory Lagrangian (after adding the h.c. term).

2 QED

The QED Lagrangian density is

$$\mathcal{L} = \bar{\psi}(ic\gamma^\mu D_\mu - mc^2)\psi - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}$$

where

$$\psi = (\psi_1\psi_2\psi_3\psi_4)$$

and

$$D_\mu = \partial_\mu + ieA_\mu$$

is the gauge covariant derivative and (e is the elementary charge, which is 1 in atomic units)

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

is the electromagnetic field tensor. It's astonishing, that this simple Lagrangian can account for all phenomena from macroscopic scales down to something like 10^{-13} cm. So of course Feynman, Schwinger and Tomonaga received the 1965 Nobel Prize in Physics for such a fantastic achievement.

Plugging this Lagrangian into the Euler-Lagrange equation of motion for a field, we get:

$$\begin{aligned}(ic\gamma^\mu D_\mu - mc^2)\psi &= 0 \\ \partial_\nu F^{\nu\mu} &= -ec\bar{\psi}\gamma^\mu\psi\end{aligned}$$

The first equation is the Dirac equation in the electromagnetic field and the second equation is a set of Maxwell equations ($\partial_\nu F^{\nu\mu} = -ej^\mu$) with a source $j^\mu = c\bar{\psi}\gamma^\mu\psi$, which is a 4-current coming from the Dirac equation.

The fields ψ and A^μ are quantized. The first approximation is that we take ψ as a wavefunction, that is, it is a classical 4-component field. It can be shown that this corresponds to taking three orders in the perturbation theory.

The first component A_0 of the 4-potential is the electric potential, and because this is the potential that (as we show in a moment) is in the Schrödinger equation, we denote it by V :

$$A_\mu = \left(\frac{V}{ec}, A_1, A_2, A_3\right)$$

So in the non-relativistic limit, the $\frac{V}{e}$ corresponds to the electric potential. We multiply the Dirac equation by γ^0 from left to get:

$$\begin{aligned}0 &= \gamma^0(ic\gamma^\mu D_\mu - mc^2)\psi = \gamma^0(ic\gamma^0(\partial_0 + i\frac{V}{c}) + ic\gamma^i(\partial_i + ieA_i) - mc^2)\psi = \\ &= (ic\partial_0 + ic\gamma^0\gamma^i\partial_i - \gamma^0mc^2 - V - ce\gamma^0\gamma^iA_i)\psi\end{aligned}$$

and we make the following substitutions (it's just a formalism, nothing more): $\beta = \gamma^0$, $\alpha^i = \gamma^0\gamma^i$, $p_j = -i\partial_j$, $\partial_0 = \frac{1}{c}\frac{\partial}{\partial t}$ to get

$$(i\frac{\partial}{\partial t} - c\alpha^ip_i - \beta mc^2 - V - ce\alpha^iA_i)\psi = 0.$$

This, in most solid state physics texts, is usually written as

$$i\frac{\partial\psi}{\partial t} = H\psi,$$

where the Hamiltonian is given by

$$H = c\alpha^i(p_i + eA_i) + \beta mc^2 + V.$$

The right hand side of the Maxwell equations is the 4-current, so it's given by:

$$j^\mu = c\bar{\psi}\gamma^\mu\psi$$

Now we make the substitution $\psi = e^{-imc^2t}\varphi$, which states, that we separate the largest oscillations of the wavefunction and we get

$$j^0 = c\bar{\psi}\gamma^0\psi = c\psi^\dagger\psi = c\varphi^\dagger\varphi$$

$$j^i = c\bar{\psi}\gamma^i\psi = c\psi^\dagger\alpha^i\psi = c\varphi^\dagger\alpha^i\varphi$$

The Dirac equation implies the Klein-Gordon equation:

$$\begin{aligned} (-ic\gamma^\mu D_\mu - mc^2)(ic\gamma^\nu D_\nu - mc^2)\psi &= (c^2\gamma^\mu\gamma^\nu D_\mu D_\nu + m^2c^4)\psi = \\ &= (c^2D^\mu D_\mu - ic^2[\gamma^\mu, \gamma^\nu]D_\mu D_\nu + m^2c^4)\psi = 0 \end{aligned}$$

Note however, the ψ in the true Klein-Gordon equation is just a scalar, but here we get a 4-component spinor. Now:

$$D_\mu D_\nu = (\partial_\mu + ieA_\mu)(\partial_\nu + ieA_\nu) = \partial_\mu\partial_\nu + ie(A_\mu\partial_\nu + A_\nu\partial_\mu + (\partial_\mu A_\nu)) - e^2A_\mu A_\nu$$

$$[D_\mu, D_\nu] = D_\mu D_\nu - D_\nu D_\mu = ie(\partial_\mu A_\nu) - ie(\partial_\nu A_\mu)$$

We rewrite $D^\mu D_\mu$:

$$\begin{aligned} D^\mu D_\mu &= g^{\mu\nu}D_\mu D_\nu = \partial^\mu\partial_\mu + ie((\partial^\mu A_\mu) + 2A^\mu\partial_\mu) - e^2A^\mu A_\mu = \\ &= \partial^\mu\partial_\mu + ie((\partial^0 A_0) + 2A^0\partial_0 + (\partial^i A_i) + 2A^i\partial_i) - e^2(A^0 A_0 + A^i A_i) = \\ &= \partial^\mu\partial_\mu + i\frac{1}{c^2}\frac{\partial V}{\partial t} + 2i\frac{V}{c^2}\frac{\partial}{\partial t} + ie(\partial^i A_i) + 2ieA^i\partial_i - \frac{V^2}{c^2} - e^2A^i A_i \end{aligned}$$

We use the identity $\frac{\partial}{\partial t}\left(e^{-imc^2t}f(t)\right) = e^{-imc^2t}(-imc^2 + \frac{\partial}{\partial t})f(t)$ to get:

$$\begin{aligned} L &= c^2\partial^\mu\psi^*\partial_\mu\psi - m^2c^4\psi^*\psi = \frac{\partial}{\partial t}\psi^*\frac{\partial}{\partial t}\psi - c^2\partial^i\psi^*\partial_i\psi - m^2c^4\psi^*\psi = \\ &= (imc^2 + \frac{\partial}{\partial t})\varphi^*(-imc^2 + \frac{\partial}{\partial t})\varphi - c^2\partial^i\varphi^*\partial_i\varphi - m^2c^4\varphi^*\varphi = \\ &= 2mc^2\left[\frac{1}{2}i(\varphi^*\frac{\partial\varphi}{\partial t} - \varphi\frac{\partial\varphi^*}{\partial t}) - \frac{1}{2m}\partial^i\varphi^*\partial_i\varphi + \frac{1}{2mc^2}\frac{\partial\varphi^*}{\partial t}\frac{\partial\varphi}{\partial t}\right] \end{aligned}$$

The constant factor $2mc^2$ in front of the Lagrangian is of course irrelevant, so we drop it and then we take the limit $c \rightarrow \infty$ (neglecting the last term) and we get

$$L = \frac{1}{2}i(\varphi^*\frac{\partial\varphi}{\partial t} - \varphi\frac{\partial\varphi^*}{\partial t}) - \frac{1}{2m}\partial^i\varphi^*\partial_i\varphi$$

After integration by parts we arrive at

$$L = i\varphi^*\frac{\partial\varphi}{\partial t} - \frac{1}{2m}\partial^i\varphi^*\partial_i\varphi$$

The nonrelativistic limit can also be applied directly to the Klein-Gordon equation:

$$\begin{aligned}
0 &= (c^2 D^\mu D_\mu + m^2 c^4) \psi = \\
&= \left(c^2 \partial^\mu \partial_\mu + i \frac{\partial V}{\partial t} + 2iV \frac{\partial}{\partial t} + iec^2 (\partial^i A_i) + 2iec^2 A^i \partial_i - V^2 - e^2 c^2 A^i A_i + m^2 c^4 \right) e^{-imc^2 t} \varphi = \\
&= \left(\frac{\partial^2}{\partial t^2} - c^2 \nabla^2 + 2iV \frac{\partial}{\partial t} + i \frac{\partial V}{\partial t} + iec^2 (\partial^i A_i) + 2iec^2 A^i \partial_i - V^2 - e^2 c^2 A^i A_i + m^2 c^4 \right) e^{-imc^2 t} \varphi = \\
&= e^{-imc^2 t} \left((-imc^2 + \frac{\partial}{\partial t})^2 - c^2 \nabla^2 + 2iV(-imc^2 + \frac{\partial}{\partial t}) + i \frac{\partial V}{\partial t} + iec^2 (\partial^i A_i) + 2iec^2 A^i \partial_i - V^2 + \right. \\
&\quad \left. - e^2 c^2 A^i A_i + m^2 c^4 \right) \varphi = \\
&= e^{-imc^2 t} \left(-2imc^2 \frac{\partial}{\partial t} + \frac{\partial^2}{\partial t^2} - c^2 \nabla^2 + 2Vmc^2 + 2iV \frac{\partial}{\partial t} + i \frac{\partial V}{\partial t} + iec^2 (\partial^i A_i) + 2iec^2 A^i \partial_i - V^2 + \right. \\
&\quad \left. - e^2 c^2 A^i A_i \right) \varphi = \\
&= -2mc^2 e^{-imc^2 t} \left(i \frac{\partial}{\partial t} + \frac{\nabla^2}{2m} - V - \frac{1}{2mc^2} \frac{\partial^2}{\partial t^2} - \frac{i}{2mc^2} \frac{\partial V}{\partial t} + \frac{V^2}{2mc^2} - \frac{iV}{mc^2} \frac{\partial}{\partial t} + \right. \\
&\quad \left. - \frac{ie}{2m} \partial^i A_i - \frac{ie}{m} A^i \partial_i + \frac{e^2}{2m} A^i A_i \right) \varphi
\end{aligned}$$

Taking the limit $c \rightarrow \infty$ we again recover the Schrödinger equation:

$$i \frac{\partial}{\partial t} \varphi = \left(-\frac{\nabla^2}{2m} + V + \frac{ie}{2m} \partial^i A_i + \frac{ie}{m} A^i \partial_i - \frac{e^2}{2m} A^i A_i \right) \varphi,$$

we rewrite the right hand side a little bit:

$$\begin{aligned}
i \frac{\partial}{\partial t} \varphi &= \left(\frac{1}{2m} (\partial^i \partial_i + ie \partial^i A_i + 2ie A^i \partial_i - e^2 A^i A_i) + V \right) \varphi, \\
i \frac{\partial}{\partial t} \varphi &= \left(\frac{1}{2m} (\partial^i + ie A^i) (\partial_i + ie A_i) + V \right) \varphi,
\end{aligned}$$

And we get the usual form of the Schrödinger equation for the vector potential $\mathbf{A} = (A_1, A_2, A_3)$:

$$i \frac{\partial}{\partial t} \varphi = \left(-\frac{(\nabla + ie\mathbf{A})^2}{2m} + V \right) \varphi.$$