

Quantum Bootstrapping

Christopher Granade

www.cgranade.com/research/talks/msr-2014

Joint work with:

Nathan Wiebe Christopher Ferrie D. G. Cory

Institute for Quantum Computing
University of Waterloo, Ontario, Canada

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Q100

We want to build a quantum computer.

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Need to push past what a classical computer can do. How do we get to 100 qubits?

Building Large Systems: Computational Limits

Computational limits affect many aspects of building large quantum systems:

- Characterization of H

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- Verification of control

Building Large Systems: Computational Limits

Computational limits affect many aspects of building large quantum systems:

- Characterization of H
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- ~~Design of control sequences~~
- Verification of control

Here, we focus on characterization and verification.

Bootstrapping to Q100

Express challenges in terms of *simulation*, then use quantum simulators.

Use small quantum simulators to characterize and verify large devices, bootstrap up to Q100 scale.

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 - Sequential Monte Carlo: algorithm for Bayesian inference

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- Learning control distortions

Bayesian Approaches to Characterization and Control

Likelihood Function

Model data collection as a probability distribution:

$$\Pr(d|\underline{x}; \underline{e})$$

d : data, \underline{x} : model, \underline{e} : experiment

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Example

Larmor precession at an unknown ω and T_2 :

$$H(\omega) = \frac{\omega}{2} \sigma_z, \quad |\psi_{\text{in}}\rangle = |+\rangle, \quad M = \{|+\rangle\langle +|, |-\rangle\langle -|\}$$

$$\Pr(d = 0|\underline{x} = (\omega, T_2); \underline{e} = (t)) = \frac{1}{2}(1 - e^{-t/T_2}) + e^{-t/T_2} \cos^2(\omega t/2)$$

Updating Knowledge

Given a likelihood, we can reason about

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what we know having seen some data.

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By Bayes' Rule: $\Pr(\underline{x}|\underline{d}, \underline{e}) = \frac{\Pr(\underline{d}|\underline{x}; \underline{e})}{\Pr(\underline{d}|\underline{e})} \Pr(\underline{x}).$
 \implies Simulation is a resource for learning.

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Estimate $\hat{\underline{x}}$ as the expectation over \underline{x} ,

$$\hat{\underline{x}} = \mathbb{E}[\underline{x}] = \int \underline{x} \Pr(\underline{x}) \, d\underline{x}.$$

Sequential Monte Carlo

SMC (aka *particle filter*): numerical algorithm for generating samples from a distribution, using a transition kernel.

$$\text{prior} \xrightarrow{\text{Bayes' Rule}} \text{posterior}$$

Posterior samples then approximate \int / \mathbb{E} .

SMC Approximation

$$\Pr(\underline{x}) \approx \sum_i^n w_i \delta(\underline{x} - \underline{x}_i)$$

(Doucet and Johansen 2011; Huszár and Hounsby [10/s86](#); Granade et al. 2012 [10/s87](#))

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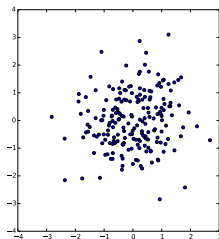
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QInfer Open-source implementation for quantum info.

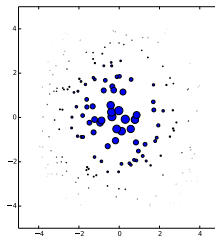
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Ambiguity and Impovrishment

Ambiguity in SMC approximation:



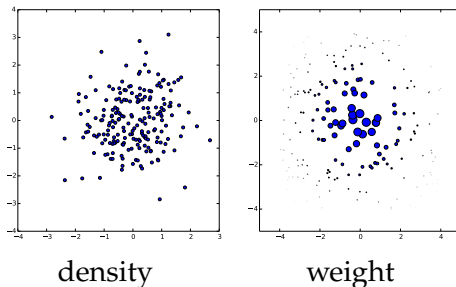
density



weight

Ambiguity and Impovrishment

Ambiguity in SMC approximation:



Using weight is less numerically stable, results in smaller *effective* number of particles.

$$n_{\text{ess}} := 1 / \sum_i w_i^2$$

Numerical Stability and Resampling

As data D is collected, $\Pr(\underline{x}_i|D) \rightarrow 0$ for most initial particles $\{x_i\}$.

■ $\Rightarrow n_{\text{ess}} \rightarrow 0$ as data is collected.

Resampling: move information from weights to the density of SMC particles.

- Resampling when $n_{\text{ess}}/n \leq 0.5$ preserves stability.
- Monitoring n_{ess} can herald some kinds of failures.

Liu and West Algorithm

Draw new particles \underline{x}' from kernel density estimate:

$$\Pr(\underline{x}') \propto \sum_i w_i \exp \left((\underline{x}' - \underline{\mu}_i)^T \underline{\underline{\Sigma}} (\underline{x}' - \underline{\mu}_i) \right)$$

$$\underline{\mu}_i := a \underline{x}_i + (1 - a) \mathbb{E}[\underline{x}] \quad \underline{\underline{\Sigma}} := h^2 \text{Cov}[\underline{x}] \quad w'_i := 1/n$$

(West 1993; Isard and Blake 1998 [10/cc76f6](#); Liu and West 2001)

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Parameters a and h can be set based on application:

- $a = 1, h = 0$: Bootstrap filter, used in state-space applications like CONDENSATION.
- $a^2 + h^2 = 1$: Ensures $\mathbb{E}[\underline{x}'] = \mathbb{E}[\underline{x}]$ and $\text{Cov}(\underline{x}') = \text{Cov}(\underline{x})$, but assumes unimodality.
- $a = 1, h \geq 0$: Allows for multimodality, emulating state-space with synthesized noise.

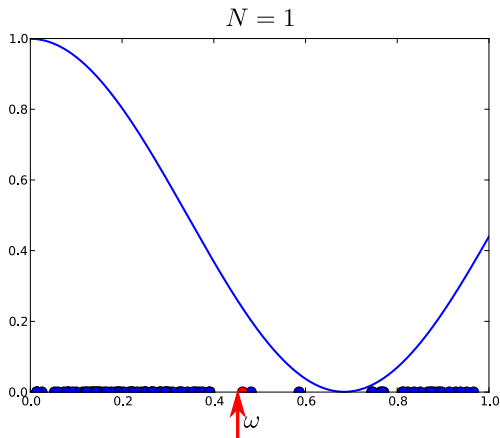
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Putting it All Together: The SMC Algorithm

- 1 Draw $\{\underline{x}_i\} \sim \pi$, set $\{w_i\} = 1/n$.
- 2 For each datum $d_j \in D$:
 - 1 $w_i \leftarrow w_i \times \Pr(d_j | \underline{x}_i; \underline{e}_j)$.
 - 2 Renormalize $\{w_i\}$.
 - 3 If $n_{\text{ess}}/n \leq 0.5$, resample.
- 3 Report $\hat{\underline{x}} := \mathbb{E}[\underline{x}] \approx \sum_i w_i \underline{x}_i$.

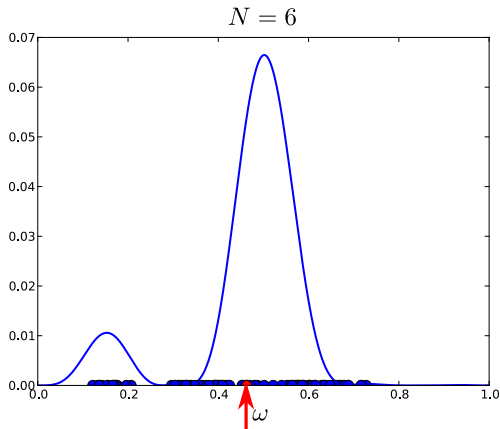
Sequential Monte Carlo

With SMC and resampling, particles move towards the true model as data is collected.



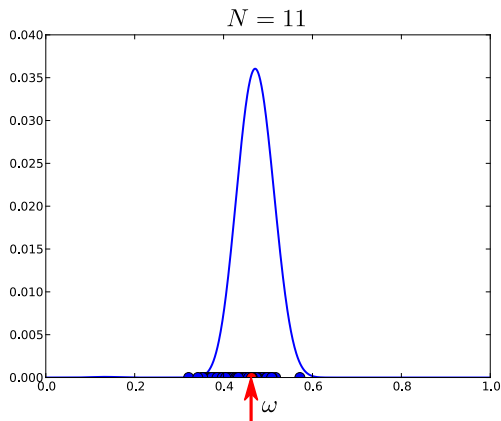
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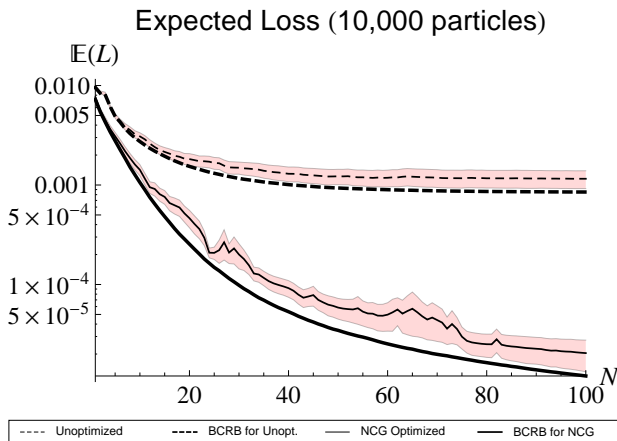
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Before bootstrapping, a few examples of SMC w/ classical resources:

Near-Optimality for \cos^2

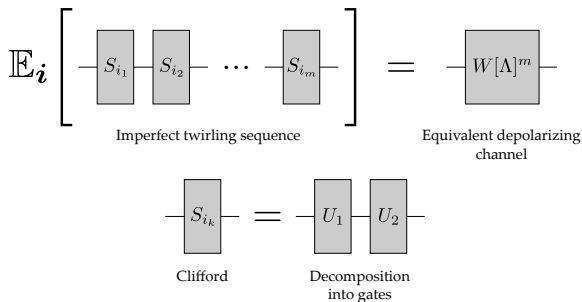
Adaptive experiment design w/ Newton Conjugate-Gradient,
vs. optimum from Bayesian Cramér-Rao Bound:



(Granade et al. 2012 [10/s87](#))

Randomized Benchmarking Example

Applying sequences of random Clifford gates *twirls* errors in a gateset, such that they can be simulated using depolarizing channels.



(Knill et al. 2008 [10/cxz9vm](#); Magesan et al. 2012 [10/tfz](#); Magesan et al. 2012 [10/s8j](#))

Randomized Benchmarking Example

Interpret survival probability as likelihood. For interleaved case, the lowest-order model is:

$$\Pr(\text{survival}|A, B, \tilde{p}, p_{\text{ref}}; m, \text{mode}) = \begin{cases} Ap_{\text{ref}}^m + B & \text{reference} \\ A(\tilde{p}p_{\text{ref}})^m + B & \text{interleaved} \end{cases}$$

A, B state preparation and measurement

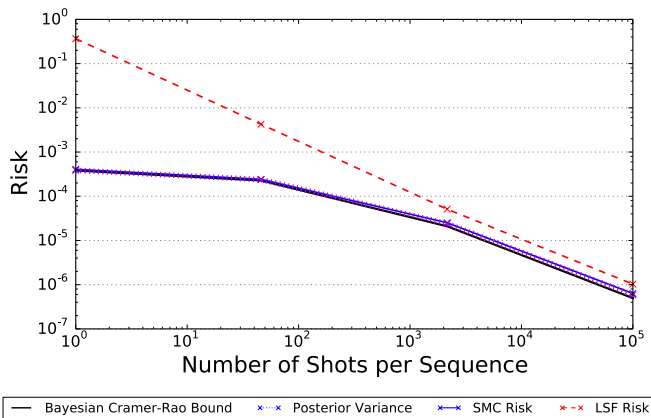
m sequence length

p_{ref} reference depolarizing parameter

\tilde{p} depolarizing parameter for gate of interest

Randomized Benchmarking Example

Using SMC, useful conclusions can be reached with significantly less data than with least-squares fitting.



(Granade, Ferrie and Cory 2014 [1404.5275](#))

SMC in Nitrogen Vacancy Centers

Would like to learn hyperfine coupling $\underline{\underline{A}}$ between e^- spin \underline{S} and ^{13}C spin \underline{I} .

$$H(\underline{x}) = \Delta_{\text{zfs}} S_z^2 + \gamma(\underline{B} + \underline{\delta B}) \cdot \underline{S} + \underline{S} \cdot \underline{\underline{A}} \cdot \underline{I}$$

$$\underline{x} = (\Delta_{\text{zfs}}, \underline{\delta B}, \underline{\underline{A}}, \alpha, \beta, T_{2,e}^{-1}, T_{2,C}^{-1})$$

α, β : visibility parameters

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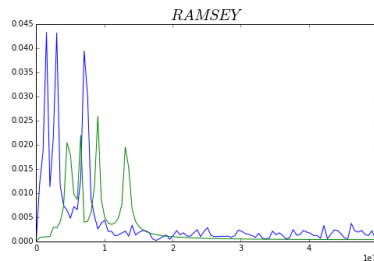
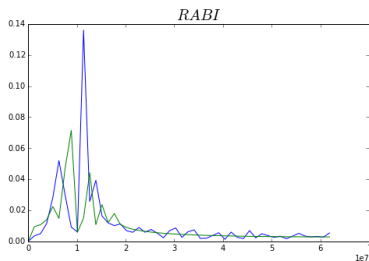
α, β : visibility parameters

- Analytic estimate sensitive to error $\underline{\delta B}$ in static field.
- Use multiple \underline{B} settings to decorrelate $\underline{\delta B}, \underline{\underline{A}}$.
- Each experiment informs about multiple parameters

Preliminary Results from Rabi Experiment

As a test, attempt to learn $\underline{\delta B}$, Δ_{zfs} , $\delta\omega_{\text{Rabi}}$ and A_N (coupling to nitrogen spin).

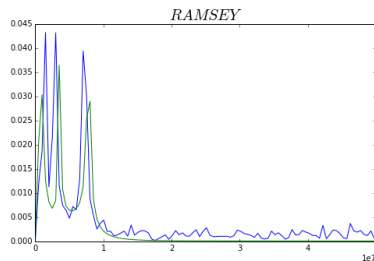
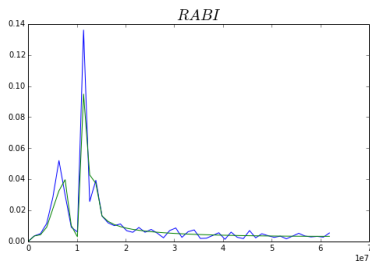
Preliminary Results from Rabi Experiment



Simulation with prior mean

Each average: 30k shots per point, 100 Rabi points + 200 Ramsey points

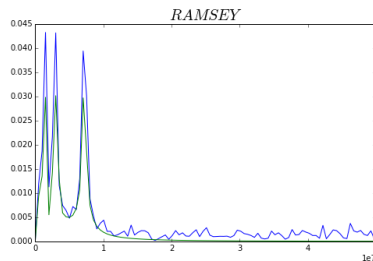
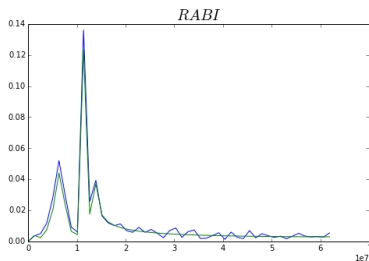
Preliminary Results from Rabi Experiment



Simulation with posterior mean, 20 averages

Each average: 30k shots per point, 100 Rabi points + 200 Ramsey points

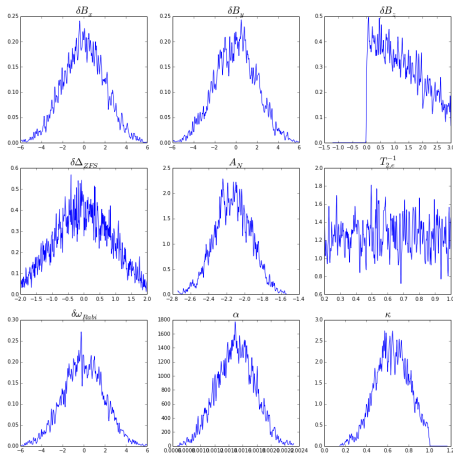
Preliminary Results from Rabi Experiment



Simulation with posterior mean, 486 averages

Each average: 30k shots per point, 100 Rabi points + 200 Ramsey points

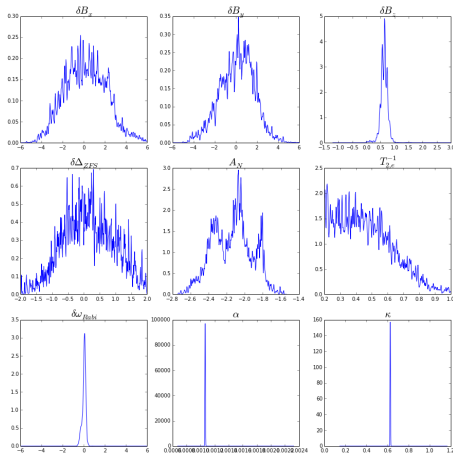
Preliminary Results from Rabi Experiment



Marginals of prior

Each average: 30k shots per point, 100 Rabi points + 200 Ramsey points

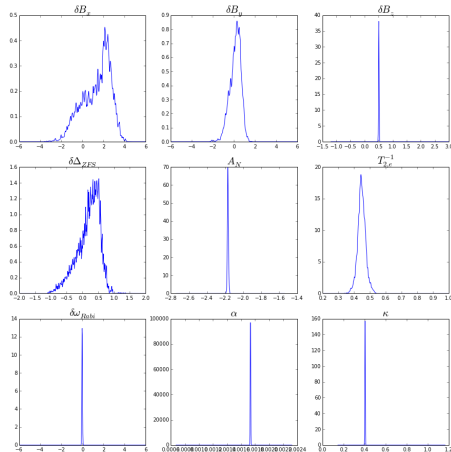
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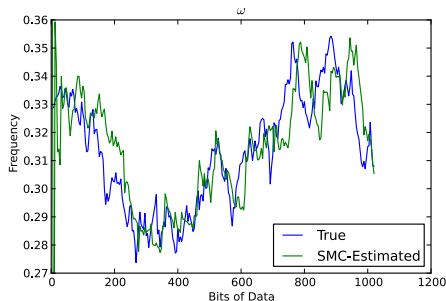
Each average: 30k shots per point, 100 Rabi points + 200 Ramsey points

We can do a few more things with SMC, some of which will be very useful in the semiquantum case.

State-Space SMC

Can move particles at each timestep $\underline{x}(t_k) \sim \Pr(\underline{x}(t_k) | \underline{x}(t_{k-1}))$.

This represents *tracking* of a stochastic process.



Confidence and Credible Regions

Characterizing uncertainty of estimates is critical for many applications:

Definition (Confidence Region)

X_α is an α -confidence region if $\Pr_D(\underline{x}_0 \in X_\alpha(D)) \geq \alpha$.

Confidence and Credible Regions

Characterizing uncertainty of estimates is critical for many applications:

Definition (Credible Region)

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(Granade et al. 2012 [10/s87](#); Ferrie 2014 [10/tb4](#))

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Credible regions can be calculated from posterior $\Pr(\underline{x}|D)$ by demanding

$$\int_{X_\alpha} d\Pr(\underline{x}|D) \geq \alpha.$$

High Posterior Density

Want credible regions that are *small* (most powerful).

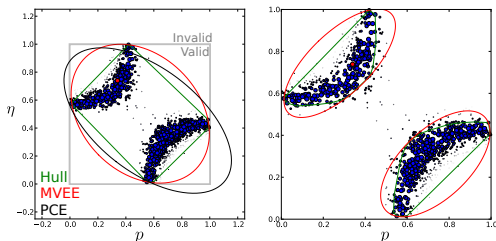
- Posterior covariance ellipses (PCE)— good for approximately normal posteriors
- Convex hull— very general, but verbose description
- Minimum volume enclosing ellipses (MVEE)— good approximation to hull

(Granade et al. 2012 [10/s87](#); Ferrie 2014 [10/tb4](#))

Comparison of HPD Estimators

For multimodal distributions, clustering can be used to exclude regions of small support.

For a noisy coin model (heads probability p , visibility η):



Left, no clustering. Right, DBSCAN.

Plot courtesy of Chris Ferrie. (Ferrie 2014 [10/tb4](#))

Bayes Factors and Model Selection

Drunk Under the Streetlights

In SMC update $w_i \mapsto w_i \times \Pr(d|\underline{x}; \underline{e})/\mathcal{N}$,

$$\mathcal{N} = \mathcal{N}(d) \approx \Pr(d|\underline{e}).$$

Is this useful?

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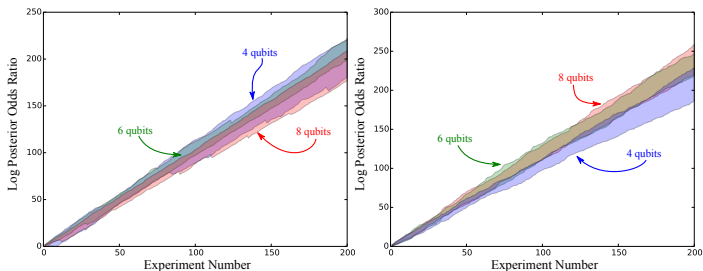
Is this useful?

Collecting normalizations \mathcal{N}_A and \mathcal{N}_B for models A, B at each step gives (assuming uniform prior over A, B)

$$\text{Bayes factor} = \frac{\Pr(D|A; \underline{e})}{\Pr(D|B; \underline{e})} \approx \frac{\prod_{d \in D} \mathcal{N}_A(d)}{\prod_{d \in D} \mathcal{N}_B(d)}$$

For full data record, can multiply normalization records to select A versus B .

For example, deciding between linear- (left) and complete-graph (right) Ising models:



(Wiebe, Granade, Ferrie and Cory 2014 [10/tdk](#))

Towards Bootstrapping

SMC uses *simulation* as a resource for *learning*.

Simulation calls: main cost to SMC (n each Bayes update).

Towards Bootstrapping

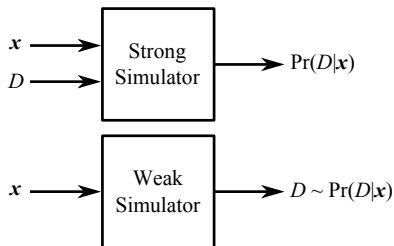
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Big Idea

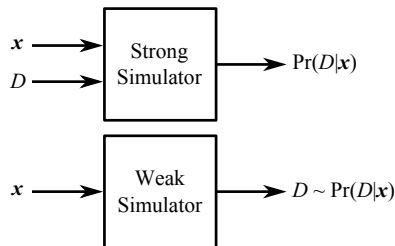
Use quantum simulation to extend SMC past classical resources.

Weak and Strong Simulation



(Ferrie and Granade 2014 [10/tdj](#))

Weak and Strong Simulation



Quantum simulation produces data, not likelihoods. Must sample to estimate likelihood.

Adaptive Likelihood Estimation

Solution

Treat estimating the likelihood as a secondary estimation problem:

Learn likelihood of untrusted system from frequencies of trusted system.

Adaptive Likelihood Estimation

Solution

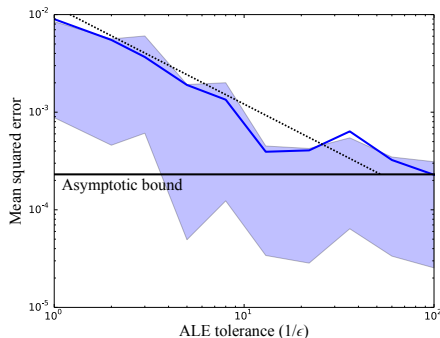
Treat estimating the likelihood as a secondary estimation problem:

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SMC is robust to likelihood estimation errors.

Performance of SMC+ALE

Ex: Simple ‘photodetector’ model $\Pr(0|p) = \alpha p + (1 - p)\beta$



α, β known bright, dark references

(Ferrie and Granade 2014 [10/t dj](#))

ALE Example: Two-Outcome Models

Given:

d result of measurement

D' set of samples from weak simulator

Hedged binomial estimate of likelihood ℓ from frequency k/K :

$$\hat{\ell} = \frac{k + \beta}{K + 2\beta},$$

where $\beta \approx 0.509$, $k := |\{d' \in D' | d' = d\}|$, $K = |D'|$.

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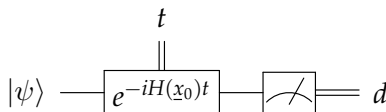
where $\beta \approx 0.509$, $k := |\{d' \in D' | d' = d\}|$, $K = |D'|$.

Variance well-known, so collect until a fixed *tolerance* is reached.

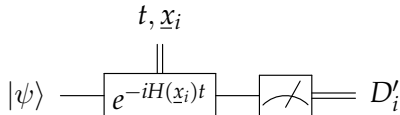
Quantum Likelihood Evaluation

Compare *classical* outcomes of unknown and trusted systems.

Unknown System



Simulator



For each \underline{x}_i :

- repeatedly sample from quantum simulation of $e^{-it\underline{x}_i}$, getting D'_i .
- estimate $\hat{\ell}_i$ from D'_i .

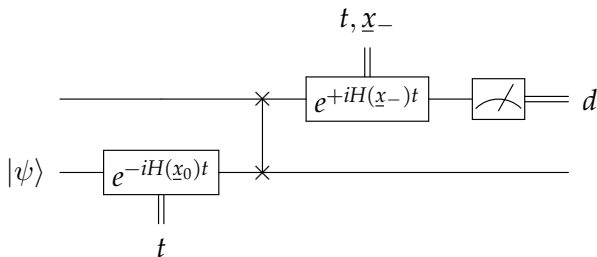
SMC update: $w_i \mapsto w_i \hat{\ell}_i / \sum_i w_i \hat{\ell}_i$.

QLE can work, but as $t \rightarrow \infty$, $\Pr(d|\underline{x}; t) \rightsquigarrow 1/\dim \mathcal{H}$.
Thus, $t \geq t_{\text{eq}}$ is uninformative.

By CRB, error then scales as $O(1/Nt_{\text{eq}}^2)$.

Interactive QLE

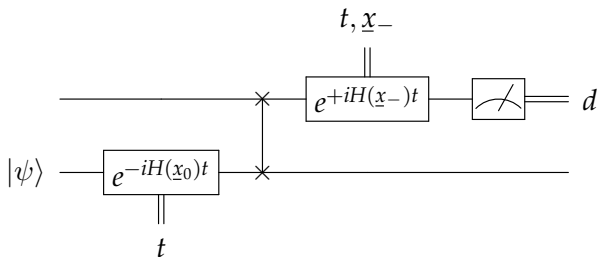
Solution: couple unknown system to a quantum simulator, then invert evolution by hypothesis \underline{x}_- .



(Wiebe, *Granade*, Ferrie and Cory 2014 [10/tf3](#))

Interactive QLE

Solution: couple unknown system to a quantum simulator, then invert evolution by hypothesis \underline{x}_- .



Echo

If $\underline{x}_- \approx \underline{x}_0$, then $|\langle \psi | e^{-it(H(\underline{x}_0) - H(\underline{x}_-))} | \psi \rangle|^2 \approx 1$.

(Wiebe, Granade, Ferrie and Cory 2014 10/13)

Posterior Guess Heuristic

Inversion connects the model and experiment spaces.
Use this duality as a heuristic for experiment design.

- Choose $\underline{x}_-, \underline{x}'_- \sim \Pr(\underline{x})$, the most recent posterior.
- Choose $t = 1/\|\underline{x}_- - \underline{x}'_-\|$.
- Return $\underline{e} = (\underline{x}_-, t)$.

Alternate Interpretation

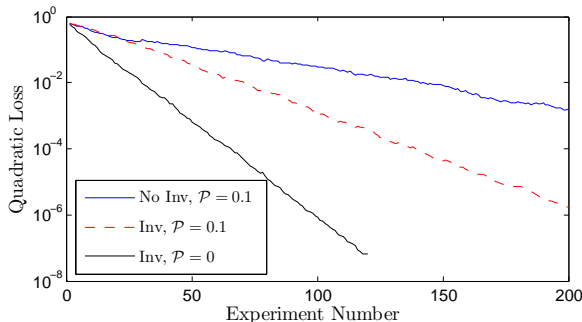
QHL finds $\hat{\underline{x}}$ such that $H(\hat{\underline{x}})$ most closely approximates “unknown” system H_0 .

Gives an α -credible bound on error introduced by replacing $H_0 \rightarrow H(\hat{\underline{x}})$.

Ising Model on Spin Chains

Hamiltonian: nearest-neighbor Ising models on a chain of nine qubits.

Interactivity allows for dramatic improvements over QLE.

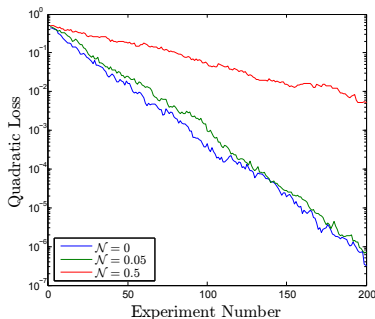


\mathcal{P} : adaptive likelihood estimation tolerance.

(Wiebe, Granade, Ferrie and Cory 2014 10/tf3)

Ising Model on the Complete Graph

With IQLE, can also learn on complete interaction graphs. We show the performance as a function of the depolarization strength \mathcal{N} .

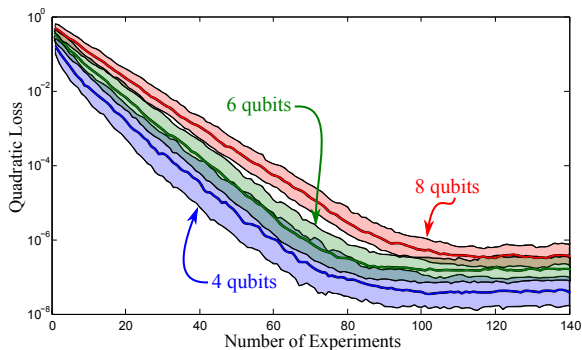


\mathcal{N} : depolarizing noise following SWAP gate.

(Wiebe, Granade, Ferrie and Cory 2014 [10/tdk](#))

Ising Model with the Wrong Graph

Simulate with spin chains, suppose “true” system is complete, with non-NN couplings $O(10^{-4})$.



(Wiebe, Granade, Ferrie and Cory 2014 [10/tdk](#))

Scaling Parameter

$\dim \underline{x}$, not $\dim \mathcal{H}$, determines scaling of IQLE.

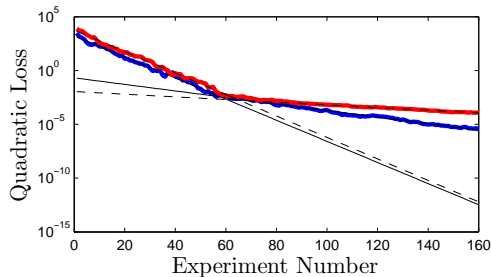


Figure : 4 qubit (red) and 6 qubit (blue) complete graph IQLE

(Wiebe, Granade, Ferrie and Cory 2014 [10/tf3](#))

Scaling and Dimensionality

In spin-chain and complete graph, average error decays exponentially,

$$L(N) \propto e^{-\gamma N}$$

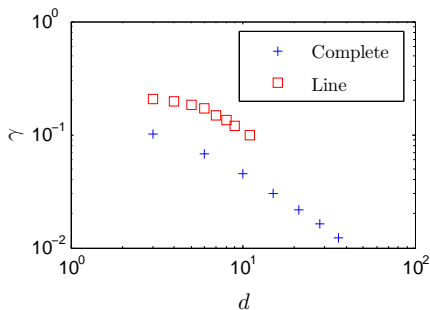
(Wiebe, *Granade*, Ferrie and Cory 2014 [10/tf3](#))

Scaling and Dimensionality

In spin-chain and complete graph, average error decays exponentially,

$$L(N) \propto e^{-\gamma N}$$

Assess scaling by finding $\gamma = \gamma(\dim \underline{x})$:



With quantum simulation, learning *may* scale efficiently.

(Wiebe, Granade, Ferrie and Cory 2014 [10/tf3](#))

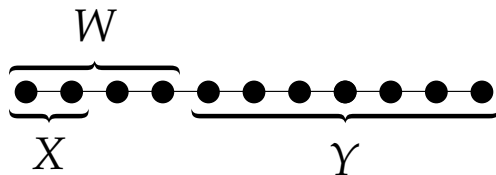
SMC + IQLE:

- Possibly scalable with quantum resources.
- Robust to finite sampling.
- Robust to approximate models.

Still requires simulator be at least as large as system of interest.

Information Locality

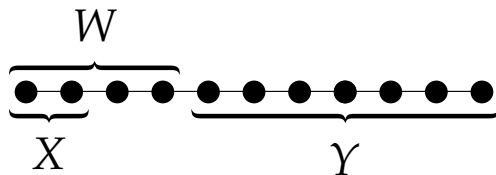
To go further, we want to *localize* our experiment, such that we can simulate on a smaller system.



Measure on X , simulate on W , and ignore all terms with support over Y .

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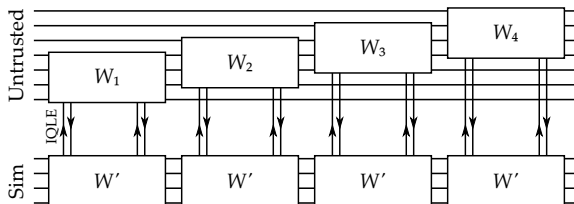


Measure on X , simulate on W , and ignore all terms with support over Y .

Gives *approximate* model that can be used to learn Hamiltonian restricted to X .

Local and Global Particle Clouds

To reconstruct the entire system, we need to combine data from different partitions.



Separate out one partition W_k at a time, maintain a *global* cloud of particles.

Local and Global Particle Clouds

Initialize $\{\underline{x}_i\}$ over entire system. Then, for each simulated subregister W_k :

- 1 Make “local” particle cloud $\{\underline{x}_i|_{W_k}\}$ by slicing $\{\underline{x}_i\}$.
- 2 Run SMC+IQLE with $\{\underline{x}_i|_{W_k}\}$ as a prior.
- 3 Ensure that the final “local” cloud has been resampled (has equal weights).
- 4 Overwrite parameters in “global” cloud $\{\underline{x}_i\}$ corresponding to post-resampling $\{\underline{x}_i|_{W_k}\}$.

In this way, all parameters are updated by an SMC run.

Q50 Example

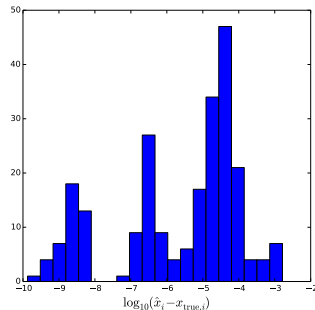
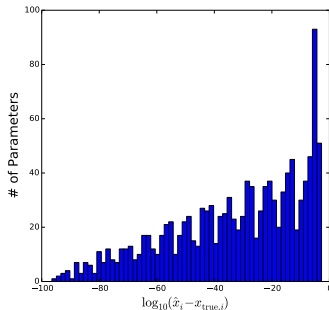
Goal: characterize a 50-qubit Ising model (complete graph) with unknown ZZ couplings.

All Hamiltonian terms commute, but initial state doesn't. Let A_X be observable, $A_{X'}$ be sim. observable.

$$\begin{aligned}\|A_X(t) - A_{X'}(t)\| &\leq \|A_X(t)\| (e^{2\|H|_Y\|t} - 1) \\ \Rightarrow t &\leq \ln \left(\frac{\delta}{\|A_X(t)\|} + 1 \right) (2\|H|_Y\|)^{-1},\end{aligned}$$

where δ is the tolerable likelihood error.

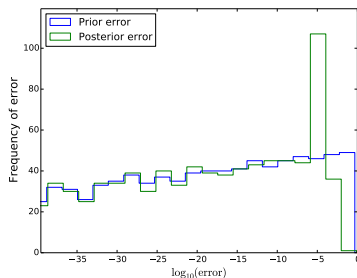
Example Q50 Run



$|X_k| = 4$, $|W_k| = 8$, $n = 20,000$, $N = 500$, exp. decaying interactions.

NB: 1225 parameter model, L_2 error of 0.3%.

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NB: 1225 parameter model, L_2 error of 0.3%.

Lieb-Robinson Bounds

More generally, for $[H|_W, H_Y] \neq 0$, use *Lieb-Robinson bound*.
 If interactions between X and Y decay sufficiently quickly, then there exists C, μ and v s. t. for any observables $A_X(t), B_Y$:

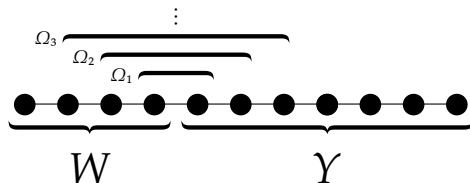
$$\|[A_X(t), B_Y]\| \leq C\|A_X(t)\|\|B_Y\|\|X\|\|Y\|(e^{v|t|} - 1)e^{-\mu d(X,Y)}$$

This *guarantees* that error due to truncation is bounded if we choose small t .

(Hastings and Koma 2006 [10/cddqgz](#); Nachtergale and Sims 2006 [10/d9xwfg](#))

Lieb-Robinson Bounds

Can find bound in terms of Hamiltonian by considering H site-by-site.

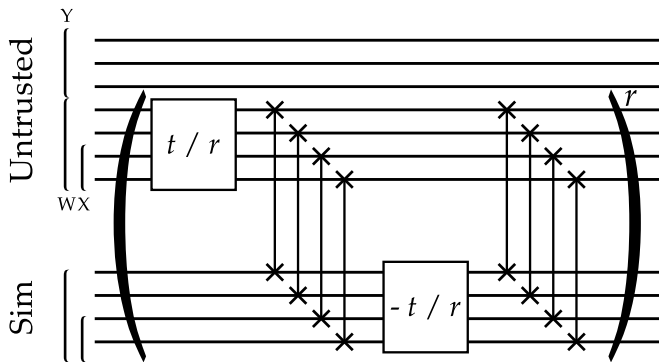


Let H_j be the Hamiltonian term containing distance- j interactions between W and Y , acting on sites Ω_j .

$$\|A(t) - e^{iH|W|t} A e^{-iH|W|t}\| \leq \sum_j C \|A\| \|H_j\| |X| |\Omega_j| e^{-\mu j} (e^{v|t|} - 1)$$

“Shaking”

Can improve the Lieb-Robinson bound by alternating between simulator and system. Using $r \approx vt$ SWAP gates, error is $O(t)$.



Distortion Operators

Control affected by classical system, *distorts* controls from intended pulse.

$$H(t) = H(t; g[\underline{p}]),$$

where \underline{p} is a pulse, $\underline{q} = g[\underline{p}]$ is the distorted pulse.

(Hincks, Granade, Borneman and Cory *forthcoming*)

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- Resonator transfer function
- Cross-talk of control lines
- Phase calibration / reference frame

(Hincks, Granade, Borneman and Cory *forthcoming*)

Learning Controls with IQLE

Learning g is also important to bootstrapping. Suppose:

- $H \neq H(t)$
- $g : \underline{y} \mapsto \underline{x}$ distorts static Hamiltonians
- g parameterized by \underline{z}

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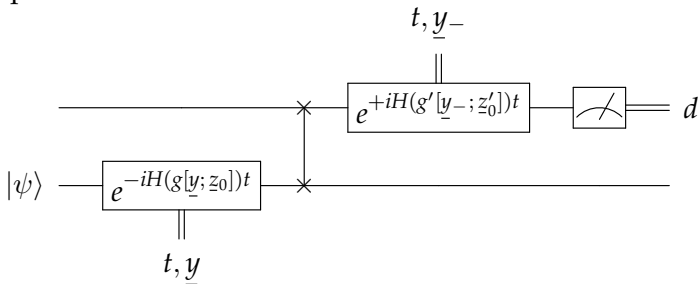
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Doesn't help with all distortions, but covers a range of important cases.

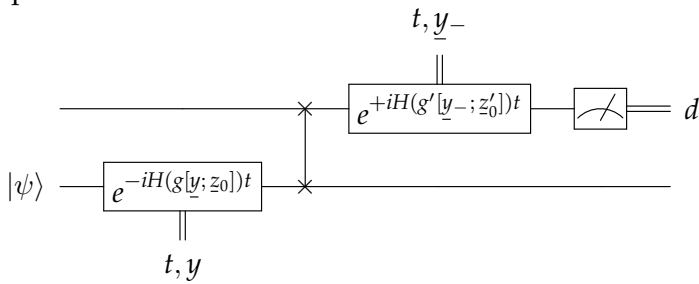
IQLE Setup

Experiment:

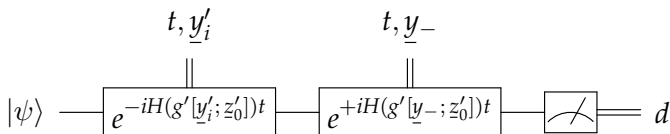


IQLE Setup

Experiment:



Simulator:



where $g'[\underline{y}'_i; \underline{z}'_0] = g[\underline{y}; \underline{z}_i]$.

Control/Distortion Duality

Unknown rescaling example: $g[\underline{y}; \underline{z}] = \underline{y} \odot \underline{z}$.

Straightforward to invert:

$$\underline{y}'_i = \underline{y} \odot \underline{z}_i \oslash \underline{z}'_0 \implies g'[\underline{y}'_i; \underline{z}'_0] = g[\underline{y}; \underline{z}_i]$$

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Shows duality between *hypothesis* and *simulation*.

Experiment		Simulation
hypothesis \underline{z}_i	\longleftrightarrow	fixed \underline{z}'_0
fixed \underline{y}		control \underline{y}'_i

QHL: special case $\underline{y} = \underline{1}$.

What must we invert by to affect $\mathbb{1}$?

Future: Bootstrapping for Control

Same approach can be used to model cross-talk:

$$\underline{z} = (\text{vec } \underline{\underline{G}}, \underline{\epsilon}) \quad g[\underline{y}; \underline{z}] = \underline{\underline{G}}\underline{y} + \underline{\epsilon}$$

Parameter count can be reduced by restricting $\underline{\underline{G}}$ (i.e.: tridiagonal).

Extend to learning time-dependent distortions?

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Extend to learning time-dependent distortions?

Can quantum resources be applied to *design* control?

- Bayesian inference: simulation as a characterization/validation resource.

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- Can use quantum simulation to offer potential scaling.
- Using robustness of SMC, can truncate simulation → bootstrapping.

Further Information

Slides, a journal reference for this work, a full bibliography and a software implementation can be found at [*http://www.cgranade.com/research/talks/lfqis-2014/*](http://www.cgranade.com/research/talks/lfqis-2014/).



Thank you for your kind attention!

Decision Theory

A few definitions help us evaluate estimates $\hat{\underline{x}}$ of \underline{x} :

Loss How well have we learned?

$$L_{\underline{\underline{Q}}}(\hat{\underline{x}}, \underline{x}) := (\hat{\underline{x}} - \underline{x})^T \underline{\underline{Q}} (\hat{\underline{x}} - \underline{x})$$

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Cramér-Rao Bound On average, how well *can* we learn?

Cramér-Rao Bound

Fisher Information

How much information about \underline{x} is obtained by sampling data?

$$\underline{\underline{I}}(\underline{x}) = \mathbb{E}_D[(\underline{\nabla}_{\underline{x}} \log \Pr(D|\underline{x}))(\underline{\nabla}_{\underline{x}} \log \Pr(D|\underline{x}))^T]$$

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The Cramér-Rao Bound tells how well any unbiased estimator can do. If $\underline{\underline{Q}} = \mathbb{1}$, then

$$R(\hat{\underline{x}}, \underline{x}) = \text{Tr}(\text{Cov}(\hat{\underline{x}})) \geq \text{Tr}(\underline{\underline{I}}(\underline{x})^{-1}).$$

Bayesian Cramér-Rao Bound

Expectation of Fisher information over prior π : the *Bayesian* Cramér-Rao bound.

$$\underline{\underline{B}} := \mathbb{E}_{\underline{x} \sim \pi} [\underline{\underline{I}}(\underline{x})], \quad r(\pi) \geq \underline{\underline{B}}^{-1}$$

For adaptive experiments, the posterior is used instead of the prior.

The BCRB can be computed iteratively: useful for tracking optimality online.

$$\underline{\underline{B}}_{k+1} = \underline{\underline{B}}_k + \begin{cases} \mathbb{E}_{\underline{x} \sim \pi} [\underline{\underline{I}}(\underline{x}; \underline{e}_{k+1})] & \text{(non-adaptive)} \\ \mathbb{E}_{\underline{x} | d_1, \dots, d_k} [\underline{\underline{I}}(\underline{x}; \underline{e}_{k+1})] & \text{(adaptive)} \end{cases}$$

(Gill and Levit 1995; Ferrie, *Granade* et al. 2012 [10/s87](#))

Method of Hyperparameters

If “true” model $\underline{x} \sim \text{Pr}(\underline{x}|\underline{y})$, for some *hyperparameters* \underline{y} , can est. \underline{y} directly:

$$\text{Pr}(d|\underline{y}; \underline{e}) = \int \text{Pr}(d|\underline{x}, \underline{y}; \underline{e}) \text{Pr}(\underline{x}|\underline{y}; \underline{e}) \, d\underline{x}.$$

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Example

For Larmor precession with $\omega \sim \text{Cauchy}(\omega_0, T_2^{-1})$,

$$\text{Pr}(d|(\omega_0, T_2^{-1}); t) = e^{-tT_2^{-1}} \cos^2(\omega_0 t/2) + (1 - e^{-tT_2^{-1}})/2.$$

Let $\underline{y} = (\omega_0, T_2^{-1})$.

Hyperparameters and Region Estimation

In some hyperparameter models, can also express as region estimator on underlying parameters.

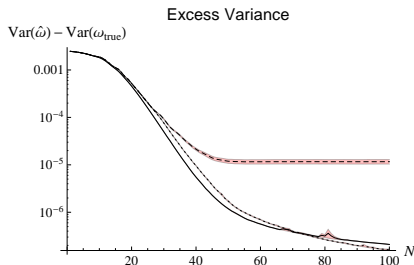


Figure : Larmor precession model w/ $\omega \sim N(\mu, \sigma^2)$, three exp. design strategies

Critically, the covariance region for ω is not smaller than the true covariance given by the hyperparameter σ^2 .

(Granade et al. 2012 [10/s87](#))