



University of Michigan – Shanghai Jiao Tong University Joint Institute
Center of Optics and Optoelectronics

VE 320 – Summer 2012 Introduction to Semiconductor Device

Carrier Statistics

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NANO ENERGY LAB

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1

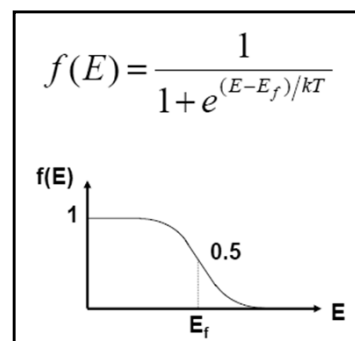
Carrier Statistics

Density of States

$$g_c(E) = \frac{m_n^* \sqrt{2m_n^* (E - E_c)}}{\pi^2 \hbar^3}$$

$$g_v(E) = \frac{m_p^* \sqrt{2m_p^* (E_v - E)}}{\pi^2 \hbar^3}$$

Fermi-Dirac distribution



Now we are ready to calculate how many electrons are there !

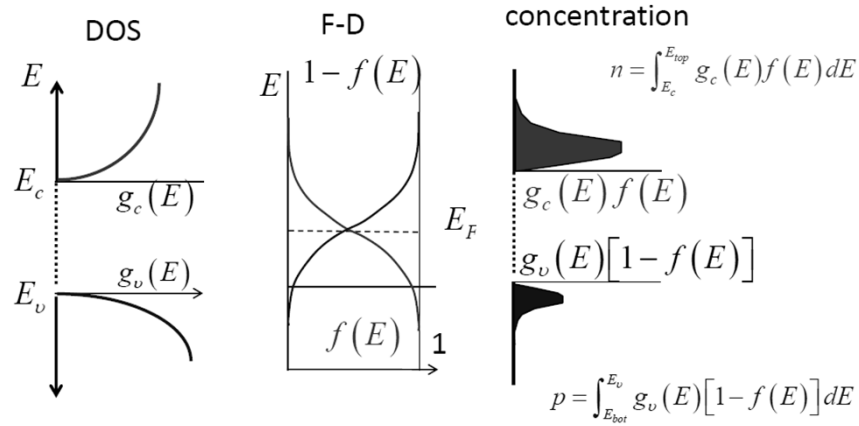


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2

Carrier Distribution



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3

Electron Concentration in 3D Solid

$$\begin{aligned}
 n &= \int_{E_c}^{E_{top}} g_c(E) f(E) dE \\
 &= \int_{E_c}^{E_{top}} 2 \times \frac{m_n^* \sqrt{2m_n^* (E - E_c)}}{2\pi^2 \hbar^3} \frac{1}{1 + e^{\beta(E - E_F)}} dE \\
 &\approx \int_{E_c}^{\infty} \frac{m_n^* \sqrt{2m_n^* (E - E_c)}}{\pi^2 \hbar^3} \frac{1}{1 + e^{\beta(E - E_c)} e^{\beta(E_c - E_F)}} dE \\
 &= N_c \frac{2}{\sqrt{\pi}} F_{1/2}(\eta_c) \quad \eta_c \equiv \beta(E_F - E_c) \\
 N_c &\equiv 2 \left(\frac{2\pi m_n^* \beta}{h^2} \right)^{3/2} \quad F_{1/2}(\eta) = \int_0^{\infty} \frac{\sqrt{z} dz}{1 + e^{z - \eta}}
 \end{aligned}$$



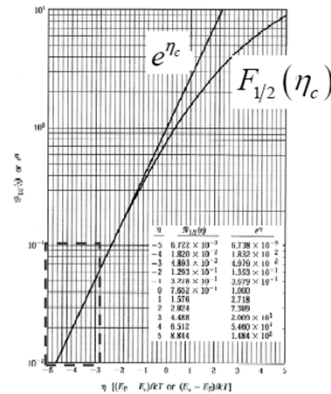
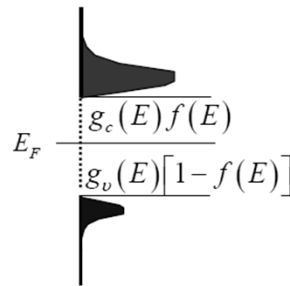
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4

Boltzmann vs. Fermi-Dirac Statistics

$$n = N_c \frac{2}{\sqrt{\pi}} F_{1/2}(\eta_c) \rightarrow N_c e^{\eta_c} \quad \text{if} \quad -\eta_c \equiv \beta(E_c - E_F) > 3$$



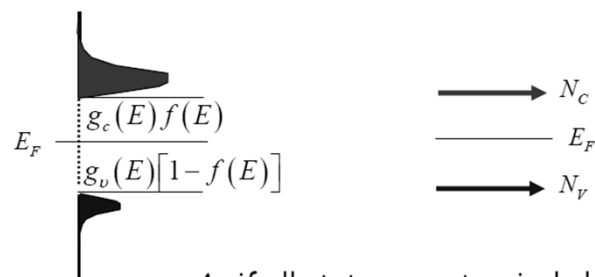
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5

Effective Density of States

$$n = N_c \frac{2}{\sqrt{\pi}} F_{1/2}(\eta_c) \rightarrow N_c e^{-\beta(E_c - E_F)} \quad \text{if} \quad E_c - E_F > 3\beta$$



As if all states are at a single level E_c



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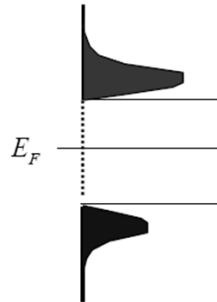
6

Law of Mass-Action

$$n = N_C e^{-\beta(E_C - E_F)}$$

$$p = N_V e^{+\beta(E_V - E_F)}$$

$$\begin{aligned} n \times p &= N_C N_V e^{-\beta(E_C - E_V)} \\ &= N_C N_V e^{-\beta E_g} \end{aligned}$$



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Fermi-Level for Intrinsic Semiconductors

$$n = p = n_i$$

$$n_i^2 = N_C N_V e^{-\beta E_g}$$

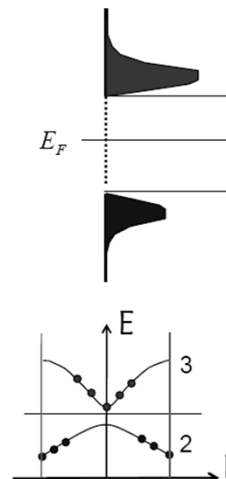
$$n_i = \sqrt{N_C N_V} e^{-\beta E_g/2}$$

$$E_F \equiv E_i$$

$$n = p \Rightarrow N_C e^{-\beta(E_C - E_i)} = N_V e^{+\beta(E_V - E_i)}$$

$$E_i = \frac{E_G}{2} + \frac{1}{2\beta} \ln \frac{N_V}{N_C}$$

$$np = n_i^2$$



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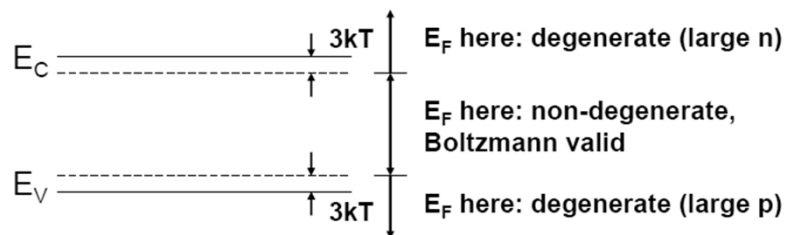
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8

Degenerate and Nondegenerate

For $E_V + 3kT < E_F < E_C - 3kT$, Boltzmann approximation is accurate

Semiconductor is said to be “non-degenerate”



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9

Doping

- We can see that the number of carrier of intrinsic semiconductor is very small.
- Is there a way to change the carrier concentration?

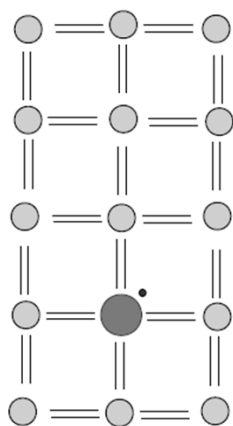


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10

Doping



II	III	IV	V	VI
4 Be	5 B	6 C	7 N	8 O
12 Mg	13 Al	14 Si	15 P	16 S
30 Zn	31 Ga	32 Ge	33 As	34 Se
48 Cd	49 In	50 Sn	51 Sb	52 Te
80 Hg	81 Tl	82 Pb	83 Bi	84 Po

Even with donors, material
is charge neutral

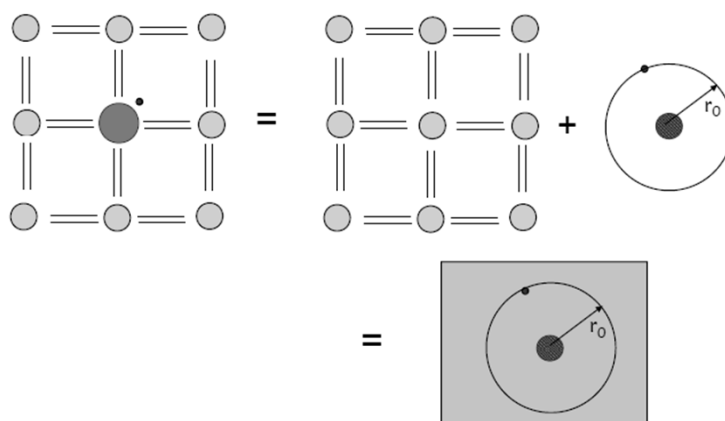


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11

Donor Atom with Hydrogen Analogy

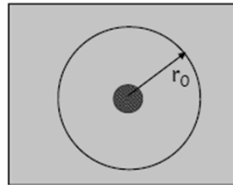


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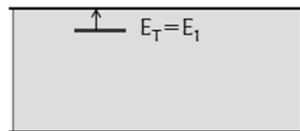
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12

Donor Atom in Real and Energy Space



~10s meV



$$\begin{aligned}
 E_1 &= -\frac{m_{host}^* q^4}{2(4\pi\epsilon_0 K_{s,host} \hbar)^2} \\
 &= -\frac{m_0 q^4}{2(4\pi\epsilon_0 \hbar)^2} \frac{m_{host}^*}{m_0} \frac{1}{K_{s,host}^2} \\
 &= -13.6 \times \frac{m_{host}^*}{m_0} \frac{1}{K_{s,host}^2}
 \end{aligned}$$

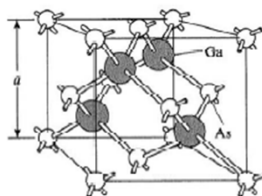
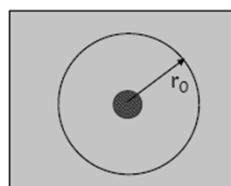


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13

Assumption of Large Radius



$$\begin{aligned}
 r_{i,P} &= \frac{4\pi\epsilon_0 K_{s,host} \hbar^2}{m_{host}^* q^2} \\
 &= \frac{4\pi\epsilon_0 \hbar^2}{m_0 q^2} \frac{m_0 K_{s,host}}{m_{host}^*} \\
 &= r_{i,H} \frac{K_{s,host}}{m_{host}^* / m_0}
 \end{aligned}$$

$$r_{i,P} = 0.53 \text{ \AA} \times \frac{12.9}{0.53} = 12.9 \text{ \AA}$$

(see tables 1.6 and 4.1)



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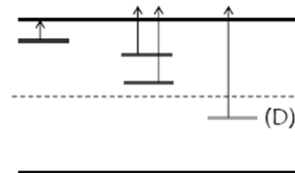
14

Characteristic of Donor Atoms

The number of donor atoms is much smaller compared to host atoms. Therefore, the electrons from one donor atom can go to the other donor atoms only via the conduction /valence bands of the host crystal.

Just like a Hydrogen atom, it is possible to have multiple localized level for a given atom (see the blue levels).

Good donors live close to the conduction band, so that they can offer electrons easily. However, if they are below the midgap, the donor levels are marked with (D) to differentiate them from acceptor atoms (which live close to the valence band).

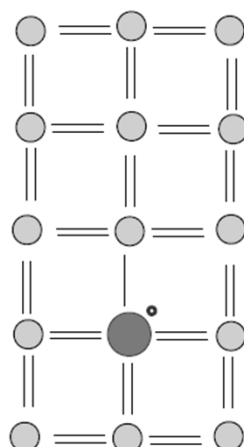


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15

Acceptor Atoms



II	III	IV	V	VI
4 Be	5 B	6 C	7 N	8 O
12 Mg	13 Al	14 Si	15 P	16 S
30 Zn	31 Ga	32 Ge	33 As	34 Se
48 Cd	49 In	50 Sn	51 Sb	52 Te
80 Hg	81 Tl	82 Pb	83 Bi	84 Po

Even with acceptor, material
is charge neutral

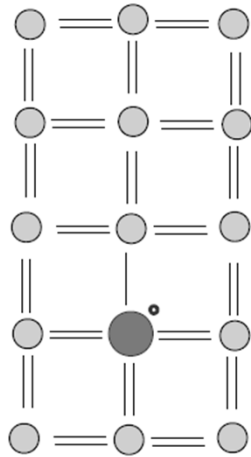


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16

Characteristic of Acceptor Atoms



II	III	IV	V	VI
4 Be	5 B	6 C	7 N	8 O
12 Mg	13 Al	14 Si	15 P	16 S
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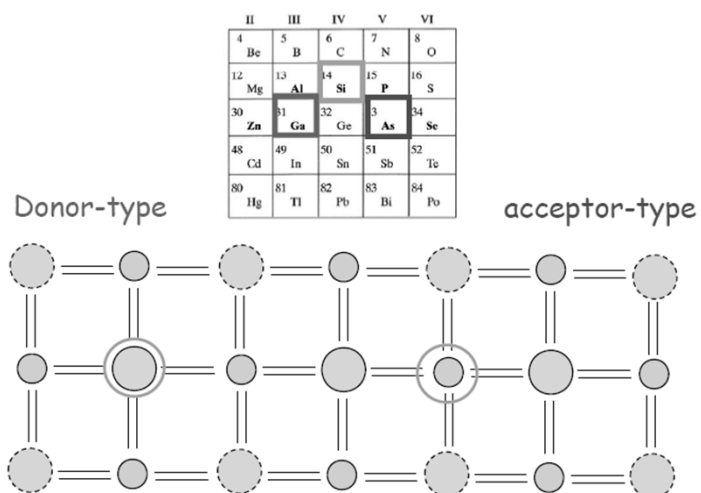


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17

Amphoteric Dopants



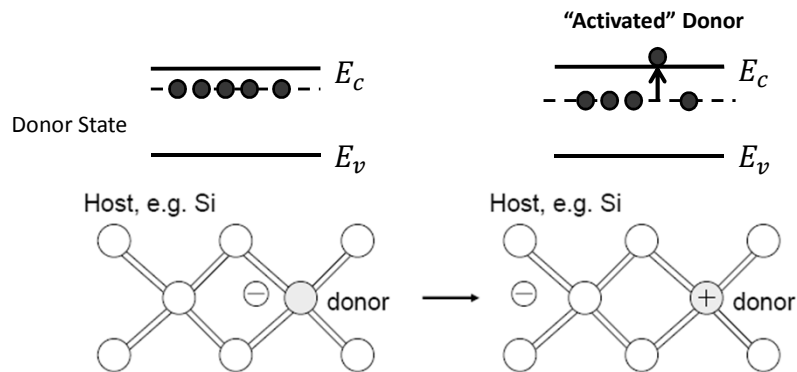
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18

Donor Action

- A net concentration of donors provides a higher electron concentration than intrinsic material
- n-type semiconductor



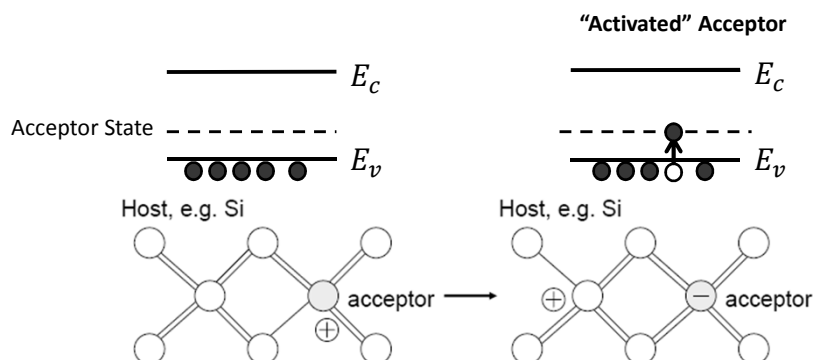
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19

Acceptor Action

- A net concentration of acceptors provides a higher hole concentration than intrinsic material
- p-type semiconductor

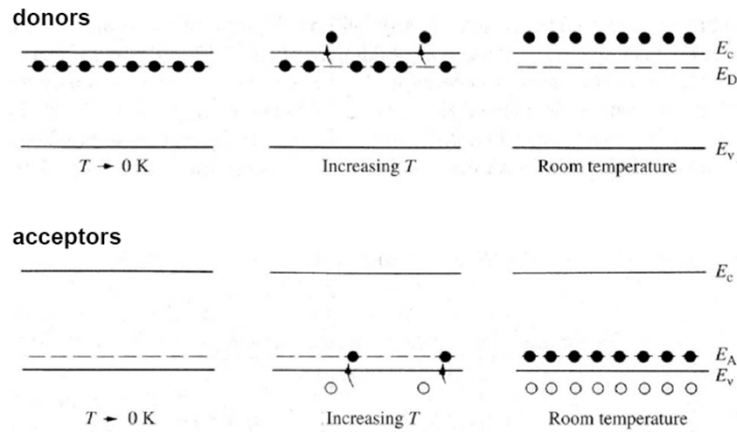


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20

Temperature Activation of Dopants



For "good dopant" majority of dopants ionized at room temperature



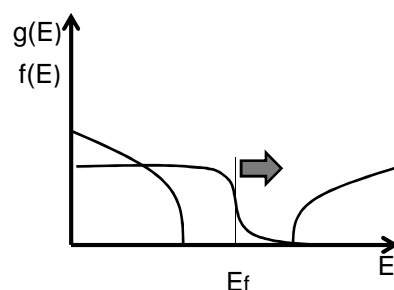
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21

Fermi Level for Doped Semiconductor

- n-type semiconductors have excess electrons in the conduction band, therefore the probability of finding an electron in the conduction band increases, so the Fermi level moves to a higher energy.



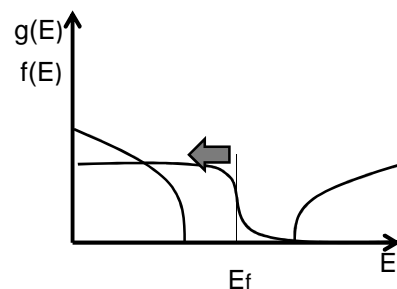
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22

Fermi Level for Doped Semiconductor

- p-type semiconductors have excess holes in the valence band, therefore the probability of finding an electron in the valence band decreases, so the Fermi level moves to a lower energy.



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23

Occupation of Donor/Acceptor Levels

- The number of activated dopants are dependent on binding energy and temperature

$$\begin{array}{l}
 \begin{array}{|l|} \hline \text{\# of Activated} \\ \text{Donor/Acceptor} \\ \hline \end{array} \rightarrow \frac{N_D^+}{N_D} = \frac{1}{1 + g_D e^{(E_f - E_D)/kT}} \quad g_D = 2 \\
 \begin{array}{|l|} \hline \text{\# of Donor/Acceptor} \\ \hline \end{array} \rightarrow \frac{N_A^-}{N_A} = \frac{1}{1 + g_A e^{(E_A - E_f)/kT}} \quad g_A = 4
 \end{array}$$

Degeneracy factors

- E_D , E_A are Donor and Acceptor levels
- Spin up, spin down for donors
- Spin up, spin down, heavy hole, light hole for acceptors



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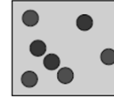
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24

Charge Neutral Relationship

A bulk material must be charge neutral over all ...

$$\int [p - n + N_D^+ - N_A^-] dV = 0$$



Further if the material is *spatially homogeneous*

$$p - n + N_D^+ - N_A^- = 0$$

$$N_V e^{-(E_F - E_V)/k_B T} - N_C e^{-(E_C - E_F)/k_B T} + \frac{N_D}{1 + 2e^{(E_F - E_D)/k_B T}} - \frac{N_A}{1 + 4e^{(E_A - E_F)/k_B T}} = 0$$

Note that with dopants, it can **no longer** be assumed that $n=p$.

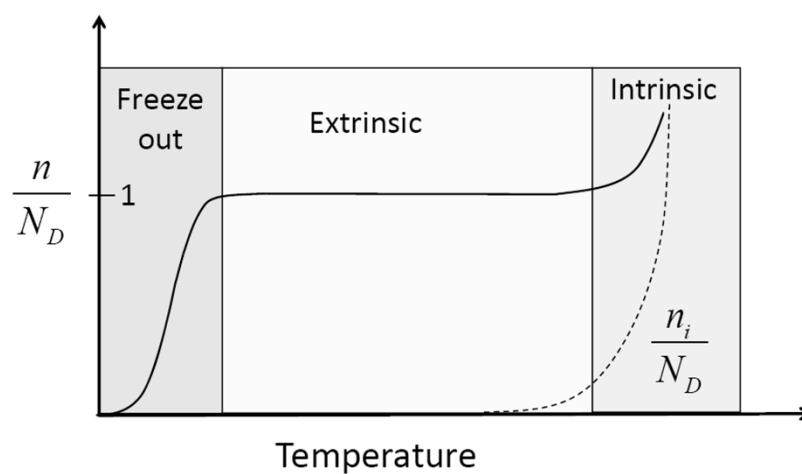


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25

Temperature Activation of Dopants

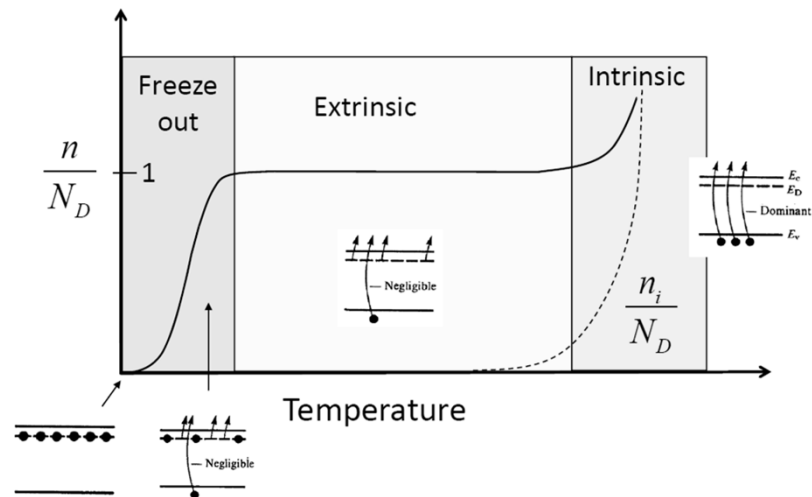


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26

Physical Interpretation



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27

Charge Neutral Relationship

In the extrinsic region, one can assume that all donors and acceptors are ionized (activated). Therefore,

$$n - p + N_D - N_A = 0$$

Law of mass action $np = n_i^2$

Carrier density

$$n = \frac{N_D - N_A}{2} + \left[\left(\frac{N_D - N_A}{2} \right)^2 + n_i^2 \right]^{1/2}$$

$$p = \frac{N_A - N_D}{2} + \left[\left(\frac{N_A - N_D}{2} \right)^2 + n_i^2 \right]^{1/2}$$



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28

Common Case

In doped semiconductor where

$$N_D - N_A \approx N_D \gg n_i$$

$$n \approx N_D \quad \text{Majority carriers}$$

$$p \approx n_i^2 / N_D \quad \text{Minority carriers}$$

$$\text{or } N_A - N_D \approx N_A \gg n_i$$

$$p \approx N_A$$

$$n \approx n_i^2 / N_A$$



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29

Carrier Concentration

$$\text{Si: } n_i = 1.6 \times 10^{10} \text{ cm}^{-3}$$

$$\text{GaAs: } n_i = 2 \times 10^6 \text{ cm}^{-3}$$

$$\text{Ge: } n_i = 2 \times 10^{13} \text{ cm}^{-3}$$

$$\text{Doping level for silicon} \sim 10^{13} \sim 10^{18} \text{ cm}^{-3}$$

($> 10^{18}$ for heavily doped silicon)

$$\text{Atomic number density of silicon} \sim 10^{22} \text{ cm}^{-3}$$



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30

Fermi-Level for Doped Semiconductors

$$n = n_i e^{(E_F - E_i)/k_B T}$$
$$p = n_i e^{(E_i - E_F)/k_B T}$$



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31

Summary



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32