



University of Michigan – Shanghai Jiao Tong University Joint Institute  
Center of Optics and Optoelectronics

## VE 320 – Summer 2012 Introduction to Semiconductor Device

### Non-ideal Effects

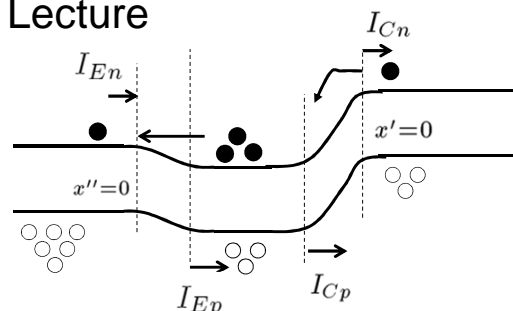
Instructor: Professor Hua Bao

**NANO ENERGY LAB**

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### Previous Lecture



$$I_{En} = -qAD_E \frac{d\Delta n_E}{dx''} \Big|_{x''=0} = qA \frac{D_E}{L_E} n_{E0} (e^{qV_{EB}/kT} - 1)$$

$$I_{Cn} = qAD_C \frac{d\Delta n_C}{dx'} \Big|_{x'=0} = -qA \frac{D_C}{L_C} n_{C0} (e^{qV_{CB}/kT} - 1)$$

$$I_{Ep} = qA \frac{D_B}{L_B} p_{B0} \left[ \frac{\cosh(W/L_B)}{\sinh(W/L_B)} (e^{qV_{EB}/kT} - 1) - \frac{1}{\sinh(W/L_B)} (e^{qV_{CB}/kT} - 1) \right]$$

$$I_{Cp} = qA \frac{D_B}{L_B} p_{B0} \left[ \frac{1}{\sinh(W/L_B)} (e^{qV_{EB}/kT} - 1) - \frac{\cosh(W/L_B)}{\sinh(W/L_B)} (e^{qV_{CB}/kT} - 1) \right]$$



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## Terminal Currents

$$I_E = qA \left[ \left( \frac{D_E}{L_E} n_{E0} + \frac{D_B}{L_B} p_{B0} \frac{\cosh(W/L_B)}{\sinh(W/L_B)} \right) (e^{qV_{EB}/kT} - 1) - \left( \frac{D_B}{L_B} p_{B0} \frac{1}{\sinh(W/L_B)} \right) (e^{qV_{CB}/kT} - 1) \right]$$

$$I_C = qA \left[ \left( \frac{D_B}{L_B} p_{B0} \frac{1}{\sinh(W/L_B)} \right) (e^{qV_{EB}/kT} - 1) - \left( \frac{D_C}{L_C} n_{C0} + \frac{D_B}{L_B} p_{B0} \frac{\cosh(W/L_B)}{\sinh(W/L_B)} \right) (e^{qV_{CB}/kT} - 1) \right]$$

- Why BJT is not equivalent to two pn junction diode?
- What happens if  $W \gg L_B$  ?



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## Performance Parameters

$$\left. \begin{aligned} \gamma &= \frac{I_{Ep}}{I_E} = \frac{I_{Ep}}{I_{Ep} + I_{En}} \\ \alpha_T &= \frac{I_{Cp}}{I_{Ep}} \\ \alpha_{dc} &= \gamma \alpha_T = \frac{I_{Cp}}{I_E} \approx \frac{I_C}{I_E} \\ \beta &= \frac{I_C}{I_B} = \frac{\alpha_{dc}}{1 - \alpha_{dc}} \end{aligned} \right\}$$

Without narrow base assumption.

$$\gamma = \left[ 1 + \left( \frac{D_E L_B N_B}{D_B L_E N_E} \right) \frac{\sinh(W/L_B)}{\cosh(W/L_B)} \right]^{-1}$$

$$\alpha_T = [\cosh(W/L_B)]^{-1}$$

$$\alpha_{dc} = \left[ \cosh(W/L_B) + \left( \frac{D_E L_B N_B}{D_B L_E N_E} \right) \sinh(W/L_B) \right]^{-1}$$

$$\beta_{dc} = \left[ \cosh(W/L_B) + \left( \frac{D_E L_B N_B}{D_B L_E N_E} \right) \sinh(W/L_B) - 1 \right]^{-1}$$

Generally one want large common emitter current gain. Then the doping concentration in the base should be as small as possible. However ...



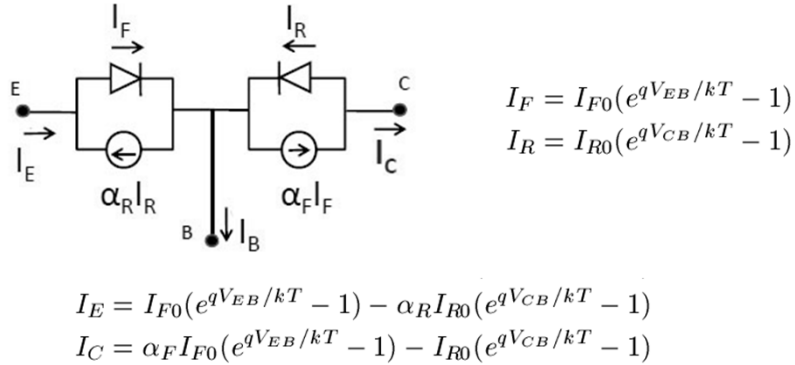
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## Ebers-Moll Equations

- Node equation for BJT terminals



The four parameters can be treated as empirical parameters.



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## More about Ebers-Moll Equations

### Exercise 11.5 Forward Current Gain

$$\alpha_F = \frac{qA \frac{D_B}{L_B} \frac{p_{B0}}{\sinh(W/L_B)}}{qA \left( \frac{D_E}{L_E} n_{E0} + \frac{D_B}{L_B} p_{B0} \frac{\cosh(W/L_B)}{\sinh(W/L_B)} \right)}$$

$$\alpha_F = \left[ \cosh(W/L_B) + \left( \frac{D_E L_B n_{E0}}{D_B L_E p_{B0}} \right) \sinh(W/L_B) \right]^{-1}$$

$$n_{E0}/p_{B0} = N_B/N_E$$



$$\alpha_F = \alpha_{dc}$$



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## Input and Output

(1) Common base input [ $I_E = I_E(V_{EB}, V_{CB})$ ]

$$I_E = I_{F0}(e^{qV_{EB}/kT} - 1) - \alpha_R I_{R0}(e^{qV_{CB}/kT} - 1)$$

(2) Common base output [ $I_C = I_C(V_{CB}, I_E)$ ]

$$I_C = \alpha_F I_E - (1 - \alpha_F \alpha_R) I_{R0}(e^{qV_{CB}/kT} - 1)$$

(3) Common emitter input [ $I_B = I_B(V_{EB}, V_{EC})$ ]

$$I_B = [(1 - \alpha_F)I_{F0} + (1 - \alpha_R)I_{R0}e^{-qV_{EC}/kT}] e^{qV_{EB}/kT} - [(1 - \alpha_F)I_{F0} + (1 - \alpha_R)I_{R0}]$$

(4) Common emitter output [ $I_C = I_C(V_{EC}, I_B)$ ]

$$I_C = \frac{(\alpha_F I_{F0} - I_{R0} e^{-qV_{EC}/kT})[I_B + (1 - \alpha_F)I_{F0} + (1 - \alpha_R)I_{R0}]}{(1 - \alpha_F)I_{F0} + (1 - \alpha_R)I_{R0} e^{-qV_{EC}/kT}} + I_{R0} - \alpha_F I_{F0}$$

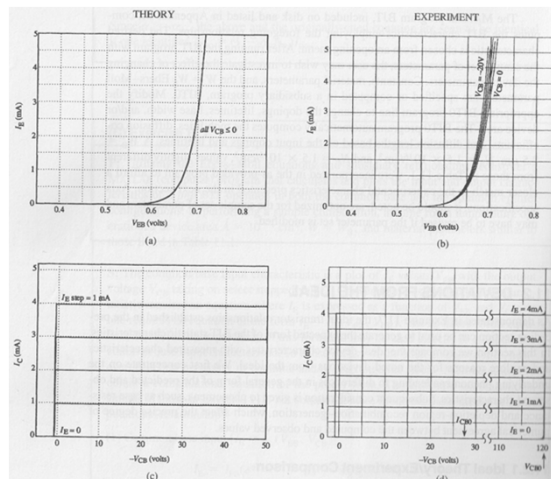


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## Ideal and Non-ideal



Input current increases with increasing negative values of  $V_{CB}$

Break down at large  $-V_{CB}$

Common base input and output

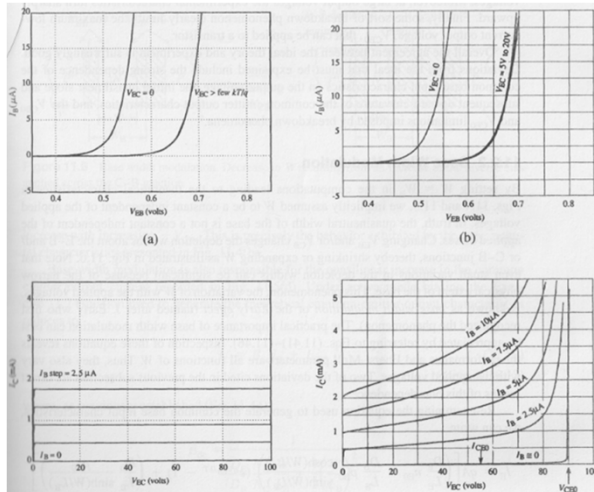


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## Common Emitter



Theory and experiment are almost consistent

Breakdown voltage is lower than common base breakdown.

Output current increases with  $V_{EC}$



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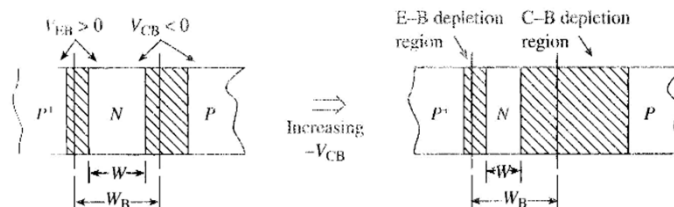
## Base Width Modulation

The plots are generated with the assumption of  $W = W_B$

$$I_E = qA \left[ \left( \frac{D_E}{L_E} n_{E0} + \frac{D_B}{L_B} p_{B0} \frac{\cosh(W/L_B)}{\sinh(W/L_B)} \right) (e^{qV_{EB}/kT} - 1) - \left( \frac{D_B}{L_B} p_{B0} \frac{1}{\sinh(W/L_B)} \right) (e^{qV_{CB}/kT} - 1) \right]$$

$$I_E \approx qA \frac{D_B}{W} p_{B0} e^{qV_{EB}/kT}$$

Common base input is dependent on the base width.



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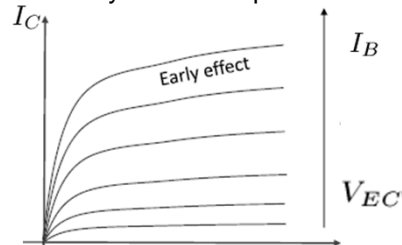
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## Base Width Modulation (Early Effect)

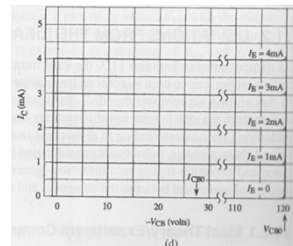
$$\beta_{dc} = \left[ \frac{D_E N_B W}{D_B N_E L_E} + \frac{1}{2} \left( \frac{W}{L_B} \right)^2 \right]^{-1}$$

$$\alpha_{dc} = \left[ 1 + \frac{D_E N_B W}{D_B N_E L_E} + \frac{1}{2} \left( \frac{W}{L_B} \right)^2 \right]^{-1}$$

Do you see the problem here?



Common emitter current gain is sensitive to base width but common base current gain is not.

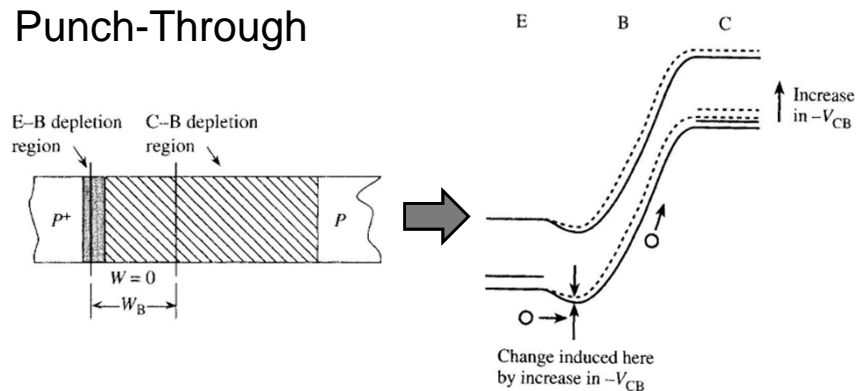


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## Punch-Through



- Base width modulation results in zero quasi-neutral width in the base region.
- The barrier can be reduced by C-B bias.
- **If it occurs**, it limits the maximum voltage can be applied on the C-B diode.

Why the base doping cannot be too low?



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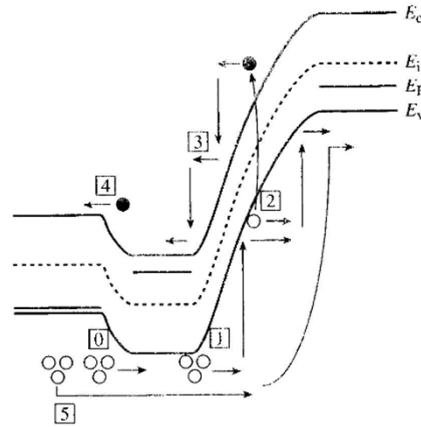
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## Avalanche Breakdown

- Common base:  
breakdown occurs when  
the C-B diode breakdown
- Common emitter:  
breakdown occurs at a  
lower  $V_{CE}$  voltage.

Base current is held constant so that the extra electrons have to be injected into the emitter, which creates extra holes injected into the collector.



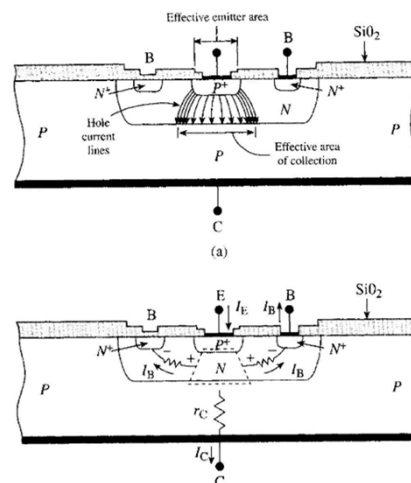
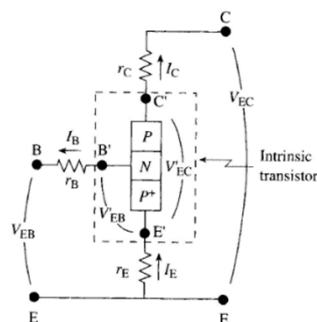
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## Geometric Effect

- Emitter area and collector area are not the same
- Series resistance
- Current crowding



Can you see another reason why base doping cannot be too low?



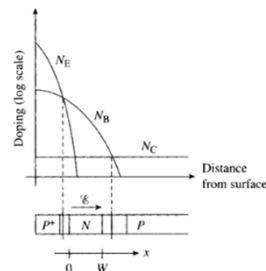
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## R-G Current / Graded Base

- In most operation conditions the R-G current is generally not important. For small forward bias of the E-B junction, the R-G current is dominant.
- The doping profile in the base is “graded” and therefore induces an electric field in the base.



$$J_N = q\mu_n n \mathcal{E} + qD_N \frac{dn}{dx} = 0$$

$$\mathcal{E}(x) = -\frac{D_N}{\mu_n} \frac{dn/dx}{n} \cong -\frac{kT}{q} \frac{dN_B(x)/dx}{N_B(x)}$$

The electric field assists the transport of minority carriers in the base, and thus reduces the R-G in the base. This is beneficial.

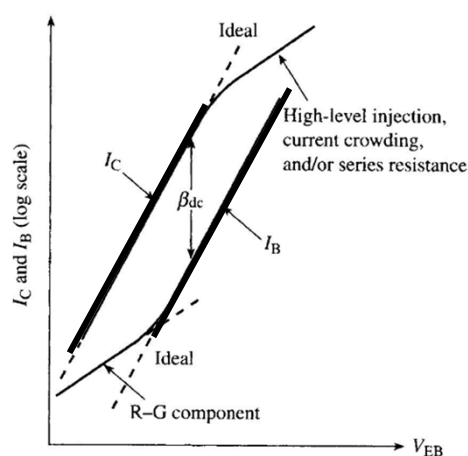


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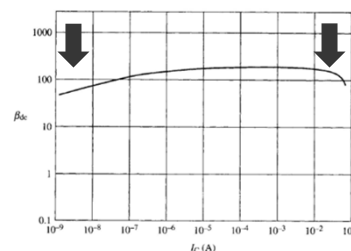
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## Gummel Plot



In a certain region,  $I_C$  is proportional to  $I_B$  and the signal is amplified. Outside this region the non-ideality is due to similar reasons as a non-ideal diode.



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## Self-test

- How to obtain the ideal BJT current-voltage relationship?
- How can BJT amplify a signal?
- What are the reasons for the nonideality of BJT?



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