

GEOPHYSICS®

Tackling cycle-skipping in full-waveform inversion with intermediate data

Journal:	<i>Geophysics</i>
Manuscript ID	GEO-2018-0096.R3
Manuscript Type:	Technical Paper
Keywords:	full-waveform inversion, tomography, nonlinear
Area of Expertise:	Seismic Inversion

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1 2 1 Tackling cycle-skipping in full-waveform inversion with intermediate data 3 4 5

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7 2 Gang Yao^{1,2}, Nuno V. da Silva³, Michael Warner³, Di Wu¹, Chenhao Yang²
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15 4 Right Running Head: intermediate data for cycle-skipping
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1 ¹ *China University of Petroleum (Beijing), The Unconventional Natural Gas Institute, State Key
2 Laboratory of Petroleum Resources and Prospecting, Beijing, China. E-mail: yaogang@cup.edu.cn;*
3
4 wudi@cup.edu.cn.

5
6 ² *Rice University, Department of Earth Science, Houston, Texas. E-mail: chenhao.yang@rice.edu.*
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10 ³ *Imperial College London, Department of Earth Science and Engineering, London, UK. E-mail:
11 n.vieira-da-silva@imperial.ac.uk; m.warner@imperial.ac.uk.*
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ABSTRACT

Full-waveform inversion (FWI) is a promising technique for recovering the Earth models for both exploration geophysics and global seismology. FWI is generally formulated as the minimization of an objective function, defined as the L2-norm of the data residuals. The non-convex nature of this objective function is one of the main obstacles for the successful application of FWI. A key manifestation of this non-convexity is cycle-skipping, which happens if the predicted data is more than half a cycle away from the recorded data. We introduce the concept of intermediate data for tackling cycle-skipping. This intermediate data set is created to sit between predicted and recorded data, and it is less than half a cycle away from the predicted data. Inverting the intermediate data rather than the cycle-skipped recorded data can then circumvent cycle-skipping. We applied this concept to invert cycle-skipped first arrivals. Firstly, we picked up the first breaks of the predicted data and the recorded data. Secondly, we linearly scaled down the time difference between the two first breaks of each shot into a series of time shifts, the maximum of which is less than half a cycle, for each trace in this shot. Thirdly, we moved the predicted data with the corresponding time shifts to create the intermediate data. Finally, we inverted the intermediate data rather than the recorded data. Since the intermediate data is not cycle-skipped and contains the traveltime information of the recorded data, FWI with intermediate data updates the background velocity model in the correct direction. Thus, it produces a background velocity model accurate enough for carrying out conventional FWI to rebuild the intermediate- and short-wavelength components of the velocity model. Our numerical examples using synthetic data validate the intermediate-data concept for tackling cycle-skipping and demonstrate its effectiveness for the application to first arrivals.

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INTRODUCTION

Full-waveform inversion (FWI) aims to recover the seismic properties of the Earth, including P-wave velocity, anisotropic parameters, and density. FWI is formulated as the optimization of an objective function defined by the misfit between recorded and predicted data, constrained by a wave equation (Tarantola, 1984). Compared to seismic imaging techniques that utilize the adjoint operator instead of the inverse operator, e.g., seismic migration (Claerbout, 1992), FWI, in which the optimization process acts as applying inverse operators to the seismic data, has the capability to recover quantitatively accurate models. Furthermore, the use of full wave equations allows for more accurate simulations of wavefield propagation, giving FWI the advantage over other techniques that employ simplified wave equations, e.g., migration with one-way wave equations and Kirchhoff migration with ray-tracing. For these two reasons, FWI is becoming one of the most promising techniques for exploration geophysics (Virieux and Operto, 2009; Warner et al., 2013; Debens et al., 2015; da Silva et al., 2016; da Silva et al., 2018; Morgan et al., 2016; Yao et al., 2018b) as well as global seismology (Zhu et al., 2012; Chen et al., 2015; Tao et al., 2017).

However, several aspects of FWI hinder its advance. Firstly, the high computational cost is a key obstacle to FWI for 3D field data applications because it requires solving wave equations repeatedly. This difficulty can be alleviated using more efficient inversion strategies, e.g., random shot selection (Herrmann et al., 2013; Warner et al., 2013), source encoding (Krebs et al., 2009), preconditioning (Baumstein et al., 2009; Burgess and Warner, 2015; Hu, 2016; Biondi et al., 2017; Yao et al., 2017), faster numerical-modeling codes (Zhang and Yao, 2012; Wang et al., 2014; Yao et al., 2016; Wang et al., 2017; Yao et al., 2018a), and faster hardware (Brown, 2007; Nemeth et al., 2008; Weiss and Shragge, 2013).

Secondly, because of the ill-posed nature of geophysical data inversion, FWI may produce a model that is very different from the true model but fits the recorded data due to incomplete acquisition coverage and the dimensionality of the model space including estimating multi-parameters (e.g., Baumstein, 2014; da Silva et al., 2016). The conditioning of the inverse problem can be ameliorated by

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3 61 developing better acquisition equipment to record more complete data (e.g., Shen et al., 2017), and
4 62 incorporating suitable mathematical constraints (e.g., Tikhonov and Arsenin, 1977; Fehmers and
5 63 Höcker, 2003; Esser et al., 2016; Trinh et al., 2017).
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10 64 Thirdly, the objective function used in conventional FWI (square of the L2-norm of the residuals) is
11 65 non-convex. Hence, convergence towards the global minimum is not guaranteed when using local
12 66 gradient-based methods. In such cases, the iterative inversion often converges towards a local minimum
13 67 if the initial model sits in a basin of attraction of the objective function that is away from the one
14 68 containing the global minimum. As a result, the recovered model can be very different from the true
15 69 model, and it is likely to be even worse than the initial model. Generally, seismic data is more nonlinear
16 70 for long-wavelength background velocity than for short-wavelength impedance contrast (Jannane et al.,
17 71 1989). Consequently, it is much easier to invert impedance contrast (reflectivity) with least-squares
18 72 reverse-time migration (LSRTM) (e.g. Dai et al., 2012; Yao and Wu, 2015; Yao and Jakubowicz, 2016)
19 73 than to recover the velocity model with FWI.
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23 74 One key manifestation of this non-convexity is cycle-skipping. It happens if the events in the
24 75 predicted data are more than half a cycle away from the corresponding events in the recorded data. As
25 76 a result, cycle-skipping generally leads FWI to converge to a local minimum. This results in an incorrect
26 77 estimation of the model parameters.
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30 78 This problem has been addressed in several ways. One is to build an initial model that is accurate
31 79 enough to produce the predicted data less than half a cycle away from the recorded data. Although this
32 80 approach can guarantee, at least to some extent, that the successive estimates of the model parameters
33 81 are carried out within the basin of attraction of the global minimum, it is usually difficult to generate
34 82 such an accurate initial model.
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38 83 The multi-scale strategy (Bunks et al., 1995) is perhaps the most widely and successfully used
39 84 approach for mitigating cycle skipping. The lower the frequency, the wider the half cycle. Consequently,
40 85 if the inversion starts from the lowest frequency in the recorded data and then the frequency is increased
41 86 sequentially, the possibility of cycle-skipping to occur is substantially reduced. This is the main reason
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3 87 why techniques aiming to create low-frequency signals by using mathematical extrapolation (Shin and
4 88 Ho Cha, 2008; Shin and Ho Cha, 2009; Li and Demanet, 2016), or acquiring low-frequency signals by
5 89 modifying conventional acquisition (Kalinicheva et al., 2017), or inventing new acquisition systems
6 90 (Baeten et al., 2013; Dellinger et al., 2016), have been investigated.

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9 91 Alternative objective functions have been proposed to resolve the cycle-skipping issue. The
10 92 fundamental mechanism of this type of method is to create a much broader convex region around the
11 93 global minimum than that of conventional FWI, the objective function of which is the square of the L₂-
12 94 norm of the data residual (Tarantola, 1984). The envelope of an oscillating seismic trace has a much
13 95 larger period than the original seismic signal. Thus, an objective function formulated by the difference
14 96 of envelopes between the predicted data and the recorded data has a very broad convex region around
15 97 the global minimum (Wu et al., 2014; Liu and Zhang, 2017). Consequently, the inversion is more likely
16 98 to converge towards the global minimum even with a poorer starting model. Similarly, the superposition
17 99 of the two single-frequency signals, the frequencies of which are close, produces a low-frequency
18 100 envelope as well; therefore, the objective function with this envelope has a wide convex region around
19 101 the global minimum. This method is referred to as the beat-tone method (Hu, 2014).

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22 102 The alternative objective functions can also be formulated in the extended domain to solve the cycle-
23 103 skipping issue. For example, by weighting the crosscorrelation function or designing a Wiener filter
24 104 between the predicted traces and the corresponding recorded traces, the objective function can generate
25 105 a large convex region around the global minimum, hence improving the convergence properties of the
26 106 inversion algorithm (van Leeuwen and Mulder, 2010; Luo and Sava, 2011; Warner and Guasch, 2016;
27 107 Zhu and Fomel, 2016). A similar mechanism has been applied to the image-domain inversion referred
28 108 to as wave-equation migration velocity analysis (WEMVA). In this method, the objective function is
29 109 formed by weighting the sub-surface offset image gather (Symes, 2008; Zhang and Biondi, 2013; Fu
30 110 and Symes, 2017). There are also some other objective functions by domain extension, for instance,
31 111 wavefield reconstruction inversion (WRI) (van Leeuwen and Herrmann, 2013; da Silva and Yao, 2018),
32 112 matched source extension (Huang et al., 2017), tomographic full-waveform inversion (Biondi and
33 113 Almomin, 2014), and differential semblance (Symes and Carazzone 1991; Plessix et al., 2000; Mulder

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3 114 and ten Kroode, 2002). The methods mentioned above have the common aspect of being based upon an
4 unphysical extension that is penalized along with the inversion.
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8 116 Alongside the extended domain methods, there are other objective functions that mitigate the cycle-
9 skipping problem, including full-traveltime inversion (FTI) (Luo and Schuster, 1991; Luo et al., 2016),
10 adjustable full-waveform inversion (AFWI) (Jiao et al., 2015), FWI with optimal transport distance
11 (Métivier et al., 2016; Yang and Engquist, 2018; Yang et al., 2018), and the scaled-Sobolev objective
12 function (Zuberi and Pratt, 2018), for example.
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15 Wang et al. (2016) applied dynamic warping to shift the recorded data set to less than half a cycle
16 away from the predicted data set to generate a series of data sets, which are used to update velocity
17 models without cycle-skipping. In this paper, we further develop the methodology of Wang et al. (2016)
18 to tackle the cycle-skipping issue using an intermediate data set, which is a generalized concept. The
19 latter is a data set that retains some of the characteristics of the recorded data set. In addition, it is also
20 sufficiently close to the predicted data set such that it is not cycle-skipped with respect to a current
21 model. The new method is described as follows. Firstly, we create an intermediate data set by shifting
22 the predicted data set towards the recorded data set but within half a cycle; secondly, we invert the
23 intermediate data instead of the recorded data; then the first two steps are repeated until the time
24 difference between the predicted data and the recorded data is less than half a cycle; finally, when this
25 condition is met, we complete the inversion with conventional FWI (Tarantola, 1984). Since the
26 intermediate data includes the missing information of the predicted data relative to the recorded data
27 but also are not cycle-skipped to the predicted data, inverting the intermediate data can produce a correct
28 model without cycle-skipping. Furthermore, since the intermediate data originates from the predicted
29 data, the modeling kernel of the inversion is more compatible with the intermediate data than the
30 recorded data, and consequently, the inversion with the intermediate data set is robust. We demonstrate
31 the application of this method in the inversion of first arrivals with numerical examples.
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METHOD

139 Conventional FWI (Tarantola, 1984) is generally formulated as the minimization of the L2-norm of
140 the data residual, expressed as

$$141 \quad \phi = \frac{1}{2} \|Ru - d\|_2^2 , \quad (1)$$

142 subject to

$$143 \quad Au = s , \quad (2)$$

144 where R represents the restriction operator to extract the wavefield at the position of receivers, u is the
145 predicted wavefield, d is the recorded data, A is the operator of the wave equation, and s is the source
146 wavelet. In equation 2, both A and u are a function of the model parameters, m . Generally, in
147 practical applications, the number of the elements of m can reach the order of several millions or even
148 higher. In addition, solving wave equations numerically is very computationally intensive.
149 Consequently, local gradient-based methods, e.g., steepest-descent and conjugate-gradient, are
150 commonly used in FWI to minimize the objective function (equation 1). However, the objective
151 function shown in equation 1 has many local minima – mainly related to the oscillatory nature of seismic
152 data. If the starting model is not close enough to the true model, then the events in the predicted data
153 can be more than half a cycle away from their corresponding events in the recorded data. This
154 phenomenon is known as ‘cycle-skipping’, and it can cause the iterative optimization to stall at a local
155 minimum.

156 To illustrate this phenomenon, a trace generated with a 5-Hz Ricker wavelet is used as the record,
157 which is shown as the solid curves in Figures 1b and 1c. The same trace with a different time shift is
158 employed as the predicted data. The objective function (equation 1) as a function of time-lag can then
159 be computed and is illustrated in Figure 1a. The nearest peaks to the global minimum are located at a
160 time lag of -86 ms and 86 ms, which are half a cycle. The dashed curves in Figures 1b and 1c represent

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3 161 the predicted data with a time shift of 50 ms and 125 ms, respectively, which are corresponding to the
4 starting point 1 and 2 in Figure 1a. As can be seen, the two starting points are positioned in different
5 basins of attraction, and the gradients point towards opposite directions. Consequently, local gradient-
6 based methods will converge towards the global minimum of the objective function when starting at
7 point 1 but to a local minimum when starting at point 2. As a result, FWI with starting point 2 produces
8 a wrong estimate of the model parameters. In addition, this wrongly estimated model is likely to be
9 worse than the starting model. Therefore, to achieve a successful inversion with FWI, cycle-skipping
10 should be avoided.

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12 169 Herein, we introduce a method for tackling cycle-skipping by generating an intermediate data set.
13 The events of this intermediate data sit between those in the predicted data and those in the recorded
14 data. In addition, the events in the intermediate data are less than half a cycle away from the events in
15 the predicted data set. The inversion then inverts intermediate data instead of the recorded data. The
16 mechanism for this method is illustrated by the sketch in Figure 2. The red curve represents the recorded
17 trace while the black curves are for the predicted traces. The initial traveltime difference between the
18 recorded trace and the predicted trace 1 is Δt_0 , which is much larger than half a cycle. To avoid cycle-
19 skipping, we can shift the predicted data, which is represented by the blue curves, by Δt_s , which is
20 smaller than half a cycle. This shifted predicted data is the intermediate data, which is closer to the
21 recorded data than the originally predicted data. Consequently, if we invert the intermediate data instead
22 of the recorded data, FWI can produce correct updates. As the correct updates are added into the initial
23 model, the new predicted data, e.g., predicted data i , is closer to the record, but still more than half a
24 cycle away from the record. We can then produce new intermediate data, e.g., intermediate data i , and
25 invert it. By repeating this process, we can gradually improve the initial model to a point such that
26 cycle-skipping is avoided. At this point, conventional FWI can be carried out without encountering
27 issues related to cycle skipping.

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29 185 In this paper, we only investigate generating an intermediate data set from first arrivals, for the sake
30 of demonstrating the potential and validity of the concept. We also postulate the possibility of using
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3 187 more events in the data in order to generate such mapping. However, we do not explore that possibility
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5 188 further herein.
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9 189 For a surface acquisition, the first arrivals are direct arrivals for short offsets but refractions for far
10 offsets. The direct arrivals only carry the information for the surface update of the model while the
11 refractions include crucial information for the background update of the deeper part of the model. For
12 a cross-well setting, the first arrivals are transmitted waves, which carry the background update
13 information of the whole model. As a result, the inversion with first arrivals in this paper can
14 successfully update the background model. The algorithm can be outlined as follows:
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- 195 1. *Calculating the time duration of half a cycle, $0.5 T$, for the inverted frequencies;*
- 196 2. *Picking the first breaks of the record, d , which are indicated by the red curve in Figure 3;*
- 197 3. *Picking the first breaks of the predicted data, d_0 , which are represented by the black curve in*
198 *Figure 3;*
- 199 4. *Scaling down the difference between the two first breaks for each shot into a series of time*
200 *shifts, the maximum of which is less than half a cycle;*
- 201 5. *Shifting the predicted data towards the record by the magnitude of the time shifts computed in*
202 *step 4; the shift produces the intermediate data, d_i , the first breaks of which are indicated by*
203 *the blue curve in Figure 3;*
- 204 6. *Producing a window function, W , the value of which is one around the first arrivals but*
205 *decreases to zero when it is away from the first arrivals;*
- 206 7. *Minimizing the new objective function,*

$$207 \quad \phi = \frac{1}{2} \|W(Ru - d_i)\|_2^2; \quad (3)$$

- 208 8. *Repeating step 3 to step 7 until the difference of the first breaks of the record and the predicted*
209 *data is smaller than half a cycle.*

In the first step, the quantity of half a cycle is obtained easily by firstly filtering the wavelet to keep the frequencies used for inversion, secondly computing the functional value between the filtered wavelet and its shifted version, which will be like the one in Figure 1a, and finally measuring the time lag of the nearest peak to the zero lag. The approach we outline herein is independent of the first-break picking method. Thus any method of picking first breaks used in exploration geophysics and global seismology can be used here to identify the arrival time of first arrivals. We used the method of Wong et al. (2009) for all the examples in this paper. In the fourth step, the scaling is defined by linearly mapping the difference of the two first breaks in the range from $-0.5 T$ to $0.5 T$. The window function, W , can be a Gaussian window or a cosine-squared window, which is used to select the data, so that the inversion is restricted to first arrivals only.

EXAMPLES

In this section, we show numerical examples using a model, which contains two Gaussian anomalies over a homogeneous background, and the Marmousi model to demonstrate the effectiveness of the new method using intermediate data for tackling cycle-skipping.

A synthetic model with two Gaussian anomalies

The synthetic model with two Gaussian anomalies is shown in Figure 4a. Its background velocity is 3000 m/s while the two Gaussian anomalies have a velocity difference of ± 1000 m/s to the background. In this example, a cross-well acquisition geometry is chosen. In total, 122 shots are fired at a depth of 100 m, from a distance of 160 m to 9840 m with a shot spacing of 80 m. A 10-Hz Ricker wavelet is used as the source wavelet. The receiver array with 1001 traces for each shot is fixed at a depth of 2900 m.

The initial model shown in Figure 4b has a constant velocity of 2800 m/s, comprising significant differences to the true model. These large differences cause the predicted data to be more than half a cycle away from the recorded data. This can be observed from the comparison of the shot gathers, depicted in Figure 5a. Conventional FWI (Tarantola, 1984) with this starting model will suffer from

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3 235 cycle-skipping. To demonstrate this, we carried out conventional FWI from 5 Hz to 9 Hz with a
4 frequency increment of 1 Hz. Note that one frequency in the inversion means a narrow frequency band
5 because the frequency was extracted using band-pass filtering. Each frequency is inverted for 5
6 iterations. Afterwards, 15 iterations are carried out with the full bandwidth. In total, the inversion
7 encompasses 40 iterations. The gradient of the first iteration is shown in Figure 6a. One can observe
8 that there are large areas dominated by a negative value of the gradient (in blue color), indicating that
9 the velocity should decrease in these regions. However, the velocity in the starting model should be
10 increased except in the area of the low-velocity blob. Thus, conventional FWI converges towards an
11 inaccurate velocity model as a consequence of cycle-skipping as depicted in Figure 7a.

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14 244 To deal with cycle-skipping, we generated an intermediate data set and inverted it in each iteration.
15 Firstly, we measured the size of half a cycle, which is about 43 ms for the 10-Hz Ricker wavelet.
16 Secondly, we picked the first breaks for both predicted data and recorded data. Thirdly, we linearly
17 mapped the time difference of each trace between the two picks into a time shift that ranges from -30
18 ms to 30 ms, so that the maximum shift is less than half a cycle. Fourthly, we applied the shift to the
19 predicted data to generate the intermediate data. A comparison between the predicted data and the
20 intermediate data of one shot is shown in Figure 5b. Finally, we inverted the intermediate data instead
21 of the record by minimizing the objective function shown in Equation 3. The residual is weighted by a
22 cosine-squared function shown in Figure 5c. The inversion is carried out with full-bandwidth data in
23 this example with the intermediate data. The gradient of the first iteration is shown in Figure 6b. One
24 can observe the significant differences between this gradient (Figure 6b) and that for the conventional
25 FWI (Figure 6a). The gradient with intermediate data is dominated by positive values (red color),
26 meaning an increase in velocity. Hence, the inversion with the intermediate data leads to correct updates
27 of the velocity model. The intermediate data is regenerated at the beginning of each iteration.

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30 258 In this paper, we applied a steepest-descent method to minimize the two objective functions
31 expressed in equations 1 and 3. The gradient is preconditioned by the pseudo-Hessian matrix following
32 the approach of Shin et al. (2001), which is computed by stacking the scaled source-wavefield energy

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3 261 of all shots. The step-length is computed by assuming a linear relationship between the model and data
4 262 perturbations. Its mathematical derivation is outlined in Appendix A.
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8 263 The inversion result after 40 iterations is shown in Figure 7b. By comparison, one can see that the
9 264 inversion with intermediate data properly fixed the background velocity and recovered the two Gaussian
10 265 anomalies. This high-quality inversion result can also be verified by the good match between the
11 266 predicted data with the inverted model and the originally generated record, which is shown in Figure
12 267 5d. After completing the intermediate data inversion, we then perform inversion with conventional FWI
13 268 with the same inversion setting as the one shown in Figure 7a. Figure 7c shows the final result after 40
14 269 iterations of the conventional FWI. By comparison of Figures 7b and 7c, it can be seen that the
15 270 conventional FWI further enhanced the Gaussian anomalies and converged to the global minimum.
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271 **The Marmousi model**

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40 272 The objective of this example is demonstrating the robustness of the outlined intermediate-data
41 273 approach even when there is a discrepancy between the laws of physics of the inversion algorithm and
42 274 that of the real world. This aspect is relevant as most inversion algorithms assume that the Earth is a
43 275 fluid. However, elastic effects generally affect the acquired data.

44 276 In this test, we used both a surface geometry and a surface-to-horizontal-well geometry, which is a
45 277 90-degree rotated cross-well setting. In order to mimic the real world, we generated the record with the
46 278 elastic wave equation (Virieux, 1986) in isotropic media with the true velocity models of the P-wave
47 279 (Figure 8a) and the S-wave (Figure 8b) and the true density model. The pressure record is generated by
48 280 summing the normal stress components, τ_{xx} and τ_{zz} . In each geometry, we generated the data firing
49 281 128 shots at a depth of 25 m from a distance of 0.187 km to 12.888 km with a spacing of 100 m. The
50 282 source wavelet is a 10-Hz Ricker wavelet. The receiver array has 1051 receivers and their position is
51 283 fixed at a depth of 25 m for the surface geometry and 2900 m for the surface-to-horizontal-well
52 284 geometry. An absorbing boundary is applied around the entire domain for the surface-to-horizontal-

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3 285 well geometry but a free-surface boundary condition is used on the top boundary of the model for the
4 286 surface geometry, introducing source and receiver ghosts as well as surface-related multiples.
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8 287 In our inversion scheme, we invert for P-wave velocity only. In this test, the inversion includes two
9 288 stages. In the first stage, we aim to use the proposed intermediate-data method to correct the background
10 289 P-wave velocity. An acoustic wave equation with a constant density of 1000 kg/m^3 was used as the
11 290 modeling kernel for FWI with intermediate data, which is achieved by minimizing equation 3. In the
12 291 second stage, we applied conventional FWI, which minimizes equation 1 (Tarantola, 1984), to refine
13 292 the P-wave velocity model recovered from the first stage. In order to honor the amplitude of the recorded
14 293 data, an elastic wave equation was applied in this stage, in which the density is constrained with
15 294 Gardner's relation (Gardner et al., 1974) and the S-wave velocity is formed using a constant v_p/v_s ratio
16 295 of 1.9 throughout each iteration.
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19 296 In the first stage of this test, we ran the inversion with the intermediate data from 5 Hz to 10 Hz with
20 297 an increment of 1 Hz. Each frequency means a narrow band and was inverted for 10 iterations. This
21 298 means that the inversion completed a total of 60 iterations. We used a maximum offset of 5 km and 10
22 299 km for the surface-to-horizontal-well geometry and the surface geometry, respectively. The resulting
23 300 P-wave velocity model was further used as a starting model for conventional FWI in the second stage.
24 301 This inversion was carried out starting at 5 Hz up to 24 Hz, incrementing with 1 Hz after completing a
25 302 set of 5 iterations. Overall this inversion represents a total of 100 iterations of FWI. All offsets present
26 303 in the data were used in this stage.
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29 304 *Surface-to-horizontal-well geometry*

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32 305 In the surface-to-horizontal-well test, we designed a 1D initial velocity model shown in Figure 8c,
33 306 which has substantial errors. One shot of the predicted data from the initial model (Figure 8c) is shown
34 307 in Figure 9b. Compared to the corresponding recorded shot, one can observe that the first arrivals of
35 308 predicted data are fully cycle-skipped in relation to those in recorded data, which is indicated by the
36 309 black dashed curves.
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To overcome cycle-skipping, we applied the same method as in the first example, i.e., we shifted the predicted data towards the recorded data. The shift was calculated by linearly scaling down the difference between the first breaks into the range from -30 ms to 30 ms. Consequently, the maximum shift is less than half a cycle of the 10-Hz Ricker wavelet, which is about 43 ms. By shifting the predicted data towards the recorded data, we got the intermediate data, which carries the mismatch information between the predicted data and the recorded data but is not cycle-skipped. In addition, the shift only extracts the traveltime information from the recorded data but does not extract any elastic amplitude information, which cannot be properly handled by an acoustic modeling kernel. One shot gather of the intermediate data is shown in Figure 9c.

The gradient of the conventional FWI and the gradient of FWI with the intermediate data for the first iteration are shown in Figure 10. One can observe that the gradient of the conventional FWI has a negative background gradient on the left side but positive on the right side. However, since the initial velocity is much lower than that of the true velocity, a correct gradient should be positive (in red color). In addition, this gradient includes strong high-wavenumber events generated by the reflection events in the record. These high-wavenumber components do not carry information about the background model; hence they do not contribute to its update. By contrast, the gradient of FWI with the intermediate data is smooth and is dominated by the red color, which means the inversion will decrease the slowness (increase the velocities).

Figure 12 shows the recovered model. Compared with the true P-wave velocity model (Figure 8a), one can observe that the conventional FWI inversion converged towards an inaccurate estimate of the P-wave velocity model (Figure 12a) because of the cycle-skipping while the FWI with intermediate data inverted a smooth velocity model (Figure 12b), which has a correct background trend of the true P-wave velocity. The effectiveness of the intermediate-data method can be seen from the first-break picks of predicted data moving towards that of the recorded data progressively throughout iterations (Figure 11). This can also be seen from the model and gradient evolution shown in Figures B-3 and B-4 of Appendix B. However, if the intermediate data is generated by shifting the recorded data towards the predicted data, we obtained a much less accurate P-wave velocity, which is shown in Figure 12c.

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3 337 The reason for this is that there is waveform discrepancy of the first arrivals between the predicted data
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5 338 and the recorded data due to different laws of physics applied for the data generation. One illustration
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7 339 is shown in Figures B-1 and B-2 of Appendix B. This less accurate result can also be seen from the
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9 340 vertical profile at the distance of 6.25 km and model errors shown in Figure 13. Since FWI with
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11 341 intermediate data recovered a good background P-wave velocity model (Figure 12b), subsequent
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13 342 conventional FWI with the surface record recovered a high-resolution correct velocity model, which is
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15 343 shown in Figure 12d.

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18 344 *Surface Geometry*

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21 345 For the surface acquisition geometry, we chose a simple initial velocity model shown in Figure 8d.
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24 346 As the turning waves and refractions have a good coverage on the top central part of this model, a
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26 347 Gaussian-shape low-velocity anomaly is embedded there for emphasizing the cycle-skipping.

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29 348 Figure 14a depicts a shot gather of the recorded data. The first-break picks are represented by the
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31 dashed curves. The corresponding shot gather of the predicted data from the initial velocity model is
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33 shown in Figure 14b. By comparison, one can see that the first arrivals for the direct arrivals are matched
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35 very well but the refractions are not, which means that the background velocity of the initial velocity
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37 model is inaccurate, as it does not produce the correct traveltime for diving and refracted waves. The
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39 mismatch for the far offsets is even more than half a cycle, which means that the conventional FWI
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41 converges towards a local minimum.

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44 355 To overcome cycle-skipping, we shifted the predicted data in the same way as we did for the surface-
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46 to-horizontal-well test to generate the intermediate data. Figure 14c shows the shot gather of the
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48 intermediate data set, corresponding to the shot gather in Figure 14a. As the intermediate data has less
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50 than half a cycle time shift to the predicted data, carrying out FWI with it generates a gradient that is
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52 not affected by cycle-skipping. Figure 15c shows the gradient at the first iteration. As can be seen, the
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54 gradient is dominated by positive values, which means the inversion decreases the slowness (increases
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56 the velocity). Note that the gradient shown in Figure 15c includes only low wavenumbers because the
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58 intermediate data includes only first arrivals, mainly refracted waves.

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3 363 By contrast, Figure 15a shows the corresponding gradient from the conventional FWI. The high-
4 wavenumber components of the gradient are generated by the reflection events while the low-
5 wavenumber components, which are essential for the background update, are formed mainly from
6 refracted waves. Note that during the inversion, the water layer, the bottom boundary of which is
7 indicated by the green dotted line, is set as the true water velocity, 1500 m/s. In order to highlight the
8 low-wavenumber components, we smoothed the gradient shown in Figure 15a with a Gaussian filter,
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11 369 $\frac{1}{2\pi \cdot \sigma^2} \cdot e^{-\frac{x^2+z^2}{2\sigma^2}}$, where the standard deviation parameter, σ , is 10, and x and z are in number of
12 cells (Abdullah and Schuster, 2015). The smoothed version of the gradient is shown in Figure 15b. As
13 can be seen, the gradient around the area at a distance of 6 km and a depth of 1 km is negative, which
14 indicates that the inversion performs an incorrect update in that region due to cycle-skipping. Since the
15 low-wavenumber components of the gradient have an incorrect sign in the middle of the top area due
16 to cycle-skipping, the conventional FWI does not fix the background velocity but effectively migrates
17 the reflection events in this area and below in a least-square sense. However, these migrated events are
18 at the wrong depth and defocused as the background velocity model is incorrect, and it is not updated
19 throughout the successive iterations. This can be observed in Figure 17a: the faults are obscured; an
20 artificial low-velocity anomaly, which is a typical consequence of cycle-skipping during the application
21 of FWI on the Marmousi model, appears close to the top of the faults.
22

23
24 380 By contrast, FWI with intermediate data produced a proper background update, which can be
25 observed in Figure 17b. As a result, the predicted data matches the picked first break of the recorded
26 data (Figure 14d). The effectiveness of the inversion can be further seen from the first-break picks of
27 predicted data shifting towards those of recorded data gradually throughout iterations (Figure 16).
28 Figures B-5 and B-6 in Appendix B provide further evidence of the effectiveness. However, the
29 inversion becomes much less efficient with intermediate data generated by shifting the recorded data
30 towards predicted data. Its result is shown in Figure 17c. This can be further verified by the vertical
31 profile at a distance of 6.25 km and the model errors shown in Figure 18. We can also observe one
32 interesting phenomenon from Figure 14: a smooth background model produces a simple record, where
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3 389 the first arrival is the primary and the other arrivals are predominantly multiples. Thus, the first arrivals
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5 390 may contain sufficient information for updating the background P-wave velocity. We then carried out
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7 391 the conventional FWI utilizing the P-wave velocity model estimated with the intermediate data (Figure
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9 392 17b). Figure 17d shows the resulting high-resolution P-wave velocity model.

10
11 393 **DISCUSSION**

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13 394 In this paper, we have demonstrated using intermediate data to tackle cycle-skipping of first arrivals.
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15 395 The first arrivals in cross-well geometries and at the far-offsets in the surface acquisitions contain
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17 396 information for background velocity updates. As a result, we have successfully recovered the
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19 397 background velocity, which is accurate enough for starting the conventional FWI to build the
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21 398 intermediate- and high-wavenumber components of the velocity model. In this application of the
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23 399 intermediate-data concept to first arrivals, we need to pick the first arrivals to the accuracy of half a
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25 400 period of the starting inversion frequency. Less than this accuracy may result in an inadequate result for
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27 401 the following conventional FWI. The routine first-arrival picking methods, e.g., the short-term-average
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29 402 over long-term-average ratio (STA/LTA) (Allen, 1982) and energy ratio (Wong et al., 2009), may be
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31 403 sufficient to process the predicted data. However, the first arrivals of the real data are usually
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33 404 contaminated by noise. In this case, we may need to seek other, more intelligent picking methods, for
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35 405 example, the methods based on artificial intelligence (e.g., Chen, 2018; Yuan et al. 2018), or even with
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37 406 the assistance of manual picking.

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39 407 The intermediate data approach introduced herein includes only the traveltime errors of the first
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41 arrivals. Compared with other conventional velocity-building methods relying on the first arrivals, e.g.,
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43 409 first-break travelttime tomography (e.g., Zelt and Smith 1992), one noticeable difference is that FWI
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45 410 with intermediate data estimates the velocity model with a wave equation instead of ray tracing. The
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47 411 key advantages are then twofold. First, the wave equation yields more accurate wave paths as it can
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49 412 deal more easily with multi-arrivals and is not affected by the existence of shadow areas and caustics,
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51 413 which exist in ray tracing. Second, it is a natural extension of any conventional FWI algorithm, only
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53 414 requiring the implementation of a method for travelttime picking. In fact, we can further extend the

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3 415 concept of intermediate data to much broader applications of FWI than using the traveltime information
4 416 of the first arrivals only. Algorithms, such as dynamic warping (Hale, 2013; Ma and Hale, 2013), are
5 417 potentially good alternatives when mapping multiple cycle-skipped events, including reflection events,
6 418 in the predicted data to the corresponding events in the recorded data, for producing an intermediate
7 419 data set. Inverting the intermediate data set will avoid cycle-skipping (e.g., Wang et al., 2016).
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10
11 420 We can also investigate the concept of intermediate data to deal with other complications existing
12 421 in FWI. It is well known that it is much easier to invert ‘inverse-crime’ data sets, which are generated
13 422 by using the modeling kernels for inversion, than field data sets or data sets generated with different
14 423 kernels. To mitigate this difficulty, we can generate an intermediate data set that has the information of
15 424 the recorded data set, e.g., traveltime, missed in the predicted data set but also is compatible with the
16 425 modeling kernel of FWI, e.g. the Marmousi example in this paper. Inverting the intermediate data set
17 426 is then equivalent to inverting the desired information in the recorded data set only. This approach has
18 427 the advantage of eliminating undesired events, such as records of S-wave, which cannot be handled by
19 428 the modeling kernel if the Earth is considered to be an acoustic body. The existence of events that cannot
20 429 be accounted for by the modeling kernel can introduce artifacts in the estimated models. Hence, this
21 430 approach can, in principle, mitigate these artifacts.
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24 431 CONCLUSION

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27 432 In this paper, we have proposed a generalized concept of intermediate data to tackle cycle-skipping
28 433 in full-waveform inversion (FWI). The principle of this method is to create an intermediate data set, the
29 434 events of which sit between the events of the predicted data and the recorded data and are less than half
30 435 a cycle away from the predicted data, and then invert the intermediate data rather than the recorded
31 436 data. We have successfully applied this concept to invert the first arrivals with both surface-to-
32 437 horizontal-well and surface acquisition geometries. In these applications, we picked first breaks of both
33 438 the recorded data set and the predicted data set, and then linearly scaled down the time difference
34 439 between the two first breaks into a time shift that is less than half a cycle. We then create the
35 440 intermediate data set by moving the predicted data set with the time shift, and finally invert the
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intermediate data set. The theoretical analyses and numerical examples have validated this method. Our tests also show that shifting the predicted data yields a more accurate inverted velocity model than shifting the recorded data, especially in the case of existing discrepancies between the laws of physics of the inversion algorithm and that of the real world. Moreover, this concept of intermediate data might be applied to much broader areas in FWI than just using the first arrivals to drive the inversion. With other multi-event mapping algorithms instead of first-break picking, the concept of intermediate data might be extended to overcome cycle-skipping for multi-events. There is also potential for the concept of intermediate data to deal with other complications in FWI, such as inverting data sets including S-waves with acoustic modeling kernels.

ACKNOWLEDGEMENTS

451 This research was supported by the National Key Research and Development Program of China (No.
452 2017YFC1500303), NSFC (Grant No. 41630209), NSF (Grant No. 1547228), Science Foundation of
453 China University of Petroleum, Beijing (No. 2462018BJC001, No. 2462015YJRC012), and the
454 FULLWAVE consortium. The authors would like to thank Editor, Anatoly Baumstein and Deyan
455 Draganov, and Reviewer, Ettore Biondi and two anonymous reviewers, for their comments and
456 suggestions, which helped to improve and clarify the manuscript significantly. Finally, the authors
457 would also like to show gratitude to our colleague, George Stronge, for the proofreading.

459

APPENDIX A

460

CALCULATION OF STEP LENGTH

461 To find the optimal step length, α , that minimizes the functional, we start from the current model, \mathbf{m}_0 ,
462 which gives the residual, $\delta\mathbf{d}_0$. We then perturb the current model by a small amount, $\delta\mathbf{m}$ to form a
463 trial model,

464
$$\mathbf{m}_1 = \mathbf{m}_0 + \delta\mathbf{m}, \quad (\text{A-1})$$

465 where $\delta\mathbf{m}$ can be generated by scaling the preconditioned gradient to a few percentages of \mathbf{m}_0 in
466 magnitude. Our target is to find a new model,

467
$$\mathbf{m}_\alpha = \mathbf{m}_0 + \alpha\delta\mathbf{m} = \mathbf{m}_0 + \alpha(\mathbf{m}_1 - \mathbf{m}_0) \quad (\text{A-2})$$

468 that minimizes the functional

469
$$\frac{1}{2}\|\delta\mathbf{d}_\alpha\|_2^2, \quad (\text{A-3})$$

470 where $\delta\mathbf{d}_\alpha$ denotes the data residual corresponding to \mathbf{m}_α . By assuming the data perturbation is
471 linearly dependent on the model perturbation, we then have

472
$$\delta\mathbf{d}_\alpha = \delta\mathbf{d}_0 + \alpha(\delta\mathbf{d}_1 - \delta\mathbf{d}_0). \quad (\text{A-4})$$

473 If we insert equation A-4 into equation A-3, and set the derivative of the functional with respect to α
474 to zero, then we obtain the optimal step length,

475
$$\alpha = \frac{\delta\mathbf{d}_0^T(\delta\mathbf{d}_1 - \delta\mathbf{d}_0)}{(\delta\mathbf{d}_1 - \delta\mathbf{d}_0)^T(\delta\mathbf{d}_1 - \delta\mathbf{d}_0)}. \quad (\text{A-5})$$

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APPENDIX B

478

SUPPLEMENTARY FIGURES

479 In this section, we provide additional tests and figures to the main body of this paper. The first test
480 is to compare the seismic shot gathers generated with the true P-wave, S-wave velocities, and densities
481 of the Marmousi model. Figure B-1 shows one shot gather generated with the surface-to-horizontal-
482 well geometry, which is used for the Marmousi example. The source is located at a distance of 3.1875
483 km. The source signature is a 10-Hz Ricker wavelet. Figure B-2 shows the counterpart of Figure B-1
484 with the surface geometry used for the Marmousi example. By comparison of these figures, we can see
485 clearly that the records with different laws of physics share similar overall appearance and first-break
486 times, but have a significant difference in wiggles and amplitudes.

487 Figures B-3 and B-5 show the gradients at different iterations in FWI with intermediate data while
488 Figures B-4 and B-6 depict the evolution of the velocity model during the inversion.

489

490 REFERENCES

- 491 Abdullah, A., and G. T. Schuster, 2015, Inverting reflections using full-waveform inversion with
492 inaccurate starting models: 85th Annual International Meeting, SEG, Expanded Abstracts, 1148-1153.
- 493 Allen, R., 1982, Automatic phase pickers: Their present use and future prospects: Bulletin of the
494 Seismological Society of America, **72**, no. 6B, S225-S242.
- 495 Baeten, G., J. W. de Maag, R.-E. Plessix, R. Klaassen, T. Qureshi, M. Kleemeyer, F. t. Kroode, and
496 Z. Rujie, 2013, The use of low frequencies in a full-waveform inversion and impedance inversion land
497 seismic case study: Geophysical Prospecting, **61**, 701-711, doi: 10.1111/1365-2478.12010.
- 498 Baumstein, A., J. Anderson, D. Hinkley, and J. Krebs, 2009, Scaling of the objective function
499 gradient for full wavefield inversion: 79th Annual International Meeting, SEG, Expanded Abstracts,
500 2243-2247.
- 501 Baumstein, A., 2014, Extended subspace method for attenuation of crosstalk in multi-parameter full
502 wavefield inversion: 84th Annual International Meeting, SEG, Expanded Abstracts, 1121-1125.
- 503 Biondi, B., and A. Almomin, 2014, Simultaneous inversion of full data bandwidth by tomographic
504 full-waveform inversion: Geophysics, **79**, no. 3, WA129-WA140, doi: 10.1190/geo2013-0340.1.
- 505 Biondi, E., G. Barnier, and B. Biondi, 2017, Preconditioned elastic full-waveform inversion with
506 approximated Hessian: 79th Annual International Meeting, SEG, Expanded Abstracts, 1654-1658.
- 507 Brown, S., 2007, Performance comparison of finite-difference modeling on Cell, FPGA, and
508 multicore computers: 77th SEG Annual meeting, SEG, Expanded Abstracts, 2110-2114.
- 509 Bunks, C., F. Saleck, S. Zaleski, and G. Chavent, 1995, Multiscale seismic waveform inversion:
510 Geophysics, **60**, 1457–1473, doi: 10.1190/1.1443880.
- 511 Burgess, T., and M. Warner, 2015, Preconditioning FWI with approximate receiver Green's
512 functions: 85th SEG Annual meeting, SEG, Expanded Abstracts, 1116-1121.

- 1
2
3 513 Chen, M., F. Niu, Q. Liu, J. Tromp, and X. Zheng, 2015, Multiparameter adjoint tomography of the
4 crust and upper mantle beneath East Asia: 1. Model construction and comparisons: *Journal of*
5 *Geophysical Research: Solid Earth*, **120**, 1762–1786, doi: 10.1002/2014JB011638.
6
7 516 Chen, Y., 2018, Automatic microseismic event picking via unsupervised machine learning:
8 *Geophysical Journal International*, **212**, no. 1, 88-102, doi: 10.1093/gji/ggx420.
9
10 518 Claerbout, J. F., 1992, *Earth Soundings Analysis: Processing Versus Inversion*: Blackwell Scientific
11 Publications.
12
13 520 da Silva, N. V., A. Ratcliffe, V. Vinje, and G. Conroy, 2016, A new parameter set for anisotropic
14 multiparameter full-waveform inversion and application to a North Sea data set: *Geophysics*, **81**, no. 4,
15 U25–U38, doi: 10.1190/geo2015-0349.1.
16
17 523 da Silva, N. V., and G. Yao, 2018, Wavefield reconstruction inversion with a multiplicative cost
18 function: *Inverse Problems*, **34**, 015004, doi: 10.1088/1361-6420/aa9830.
19
20 525 da Silva, N., G. Yao, and M. Warner, 2018, Semi-global visco-acoustic full waveform inversion:
21 *Geophysics*, 1-78. doi: 10.1190/geo2017-0773.1.
22
23 527 Dai, W., P. Fowler, and G. T. Schuster, 2012, Multi-source least-squares reverse time migration:
24 *Geophysical Prospecting*, **60**, 681–695.
25
26 529 Debens, H., M. Warner, A. Umpleby, and N. Silva, 2015, Global anisotropic 3D FWI: 85th Annual
27 International Meeting, SEG, Expanded Abstracts, 1193–1197.
28
29 531 Dellinger, J., A. Ross, D. Meaux, A. Brenders, G. Gesoff, J. Etgen, J. Naranjo, G. Openshaw, and
30 M. Harper, 2016, Wolfspark®, an “FWI-friendly” ultralow-frequency marine seismic source: 86th
31 Annual International Meeting, SEG, Expanded Abstracts, 4891-4895.
32
33 534 Esser, E., L. Guasch, F. Herrmann, and M. Warner, 2016, Constrained waveform inversion for
34 automatic salt flooding: *The Leading Edge*, **35**, no. 3, 235-239, doi: 10.1190/tle35030235.1.

1
2
3 536 Fehmers, G. C., and Höcker, C. F. W., 2003, Fast structural interpretation with structure-oriented
4 filtering: *Geophysics*, **68**, no. 4, 1286–1293, doi: 10.1190/1.1598121.
5
6
7
8
9
10
11
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44
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46
47
48
49
50
51
52
53
54
55
56
57
58
59
60

538 Fu, L., and W. Symes, 2017, An adaptive multiscale algorithm for efficient extended waveform
539 inversion: *Geophysics*, **82**, no. 3, R183–R197, doi: 10.1190/geo2016-0426.1.

540 Gardner, G., L. Gardner and A. Gregory, 1974, Formation velocity and density — the diagnostic
541 basics for stratigraphic traps: *Geophysics*, **39**, no. 6, 770–780, doi: 10.1190/1.1440465.

542 Hale, D., 2013, Dynamic warping of seismic images: *Geophysics*, **78**, no. 2, S105–S115, doi:
543 10.1190/geo2012-0327.1.

544 Herrmann, F. J., I. Hanlon, R. Kumar, T. van Leeuwen, X. Li, B. Smithyman, H. Wason, A. J.
545 Calvert, M. Javanmehri, and E. T. Takougang, 2013, Frugal full-waveform inversion: From theory to a
546 practical algorithm: *The Leading Edge*, **32**, no. 9, 1082–1092, doi: 10.1190/tle32091082.1

547 Hu, W., 2014, FWI without low frequency data - beat tone inversion: 84th Annual International
548 Meeting, SEG, Expanded Abstracts, 1116–1120.

549 Hu, W., 2016, Optimal Green's function estimation for FWI acceleration: 86th Annual International
550 Meeting, SEG, Expanded Abstracts, 1074–1078.

551 Huang, G., R. Nammour, and W. Symes, 2017, Full-waveform inversion via source-receiver
552 extension: *Geophysics*, **82**, no. 3, R153–R171, doi: 10.1190/geo2016-0301.1.

553 Jannane, M., W. Beydoun, E. Crase, D. Cao, Z. Koren, E. Landa, M. Mendes, A. Pica, M. Noble, G.
554 Roeth, S. Singh, R. Snieder, A. Tarantola, D. Trezeguet, and M. Xie, 1989, Wavelengths of earth
555 structures that can be resolved from seismic reflection data: *Geophysics*, **54**, 906–910, doi:
556 10.1190/1.1442719.

557 Jiao, K., D. Sun, X. Cheng, and D. Vigh, 2015, Adjustive full waveform inversion: 85th Annual
558 International Meeting, SEG, Expanded Abstracts, 1091–1095.

1
2
3 559 Kalinicheva, T., M. Warner, J. Ashley, and F. Mancini, 2017, Two- vs three-dimensional full-
4 waveform inversion in a 3D world: 87th Annual International Meeting, SEG, Expanded Abstracts,
5 560 1383–1387.
6
7

8 562 Krebs, J. R., J. E. Anderson, D. Hinkley, R. Neelamani, S. Lee, A. Baumstein, and M.-D. Lacasse,
9 563 2009, Fast full-wavefield seismic inversion using encoded sources: *Geophysics*, **74**, no. 6, WCC177-
10 564 WCC188, doi: 10.1190/1.3230502.
11
12

13 565 Li, Y., and L. Demanet, 2016, Full-waveform inversion with extrapolated low-frequency data:
14 566 *Geophysics*, **81**, no. 6, R339-R348, doi: 10.1190/geo2016-0038.1.
15
16

17 567 Liu, Z., and J. Zhang, 2017, Joint traveltimes, waveform, and waveform envelope inversion for near-
18 568 surface imaging: *Geophysics*, **82**, no. 4, R235–R244, doi: 10.1190/geo2016-0356.1.
19
20

21 569 Luo, S., and P. Sava, 2011, A deconvolution-based objective function for wave-equation inversion:
22 570 81st Annual International Meeting, SEG, Expanded Abstracts, 2788–2792.
23
24

25 571 Luo, Y., and G. T. Schuster, 1991, Wave-equation traveltimes inversion: *Geophysics*, **56**, no. 5, 645–
26 572 653, doi: 10.1190/1.1443081.
27
28

29 573 Luo, Y., Y. Ma, Y. Wu, H. Liu, and L. Cao, 2016, Full-traveltimes inversion: *Geophysics*, **81**, no. 5,
30 574 R261–R274, doi: 10.1190/geo2015-0353.1.
31
32

33 575 Ma, Y., and D. Hale, 2013, Wave-equation reflection traveltimes inversion with dynamic warping
34 576 and full-waveform inversion: *Geophysics*, **78**, no. 6, R223–R233, doi: 10.1190/geo2013-0004.1.
35
36

37 577 Métivier, L., R. Brossier, Q. Mérigot, E. Oudet, and J. Virieux, 2016, Measuring the misfit between
38 578 seismograms using an optimal transport distance: application to full waveform inversion: *Geophysical
39 579 Journal International*, **205**, 345–377, doi: 10.1093/gji/ggw014.
40
41

42 580 Morgan, J. V., S. P. S. Gulick, T. Bralower, E. Chenot, G. Christeson, P. Claeys, C. Cockell, G. S.
43 581 Collins, M. J. L. Coolen, L. Ferrière, C. Gebhardt, K. Goto, H. Jones, D. A. Kring, E. Le Ber, J. Lofi,
44
45

- 1
2
3 582 X. Long, C. Lowery, C. Mellett, R. Ocampo-Torres, G. R. Osinski, L. Perez-Cruz, A. Pickersgill, M.
4
5 583 Poelchau, A. Rae, C. Rasmussen, M. Rebollo-Vieyra, U. Riller, H. Sato, D. R. Schmitt, J. Smit, S.
6
7 584 Tikoo, N. Tomioka, J. Urrutia-Fucugauchi, M. Whalen, A. Wittmann, K. E. Yamaguchi, and W.
8
9 585 Zylberman, 2016, The formation of peak rings in large impact craters: *Science*, **354**, 878–882, doi:
10.1126/science.aah6561.
- 11
12
13
14 587 Mulder, W., and A. ten Kroode, 2002, Automatic velocity analysis by differential semblance
15 optimization: *Geophysics*, **67**, no. 4, 1184–1191, doi: 10.1190/1.1500380.
- 16
17
18 589 Nemeth, T., J. Stefani, W. Liu, R. Dimond, O. Pell, and R. Ergas, 2008, An implementation of the
19 acoustic wave equation on FPGAs: 78th Annual International Meeting, SEG, Expanded Abstracts,
20
21 590 2874–2878.
- 22
23
24 592 Plessix, R. E., W. A. Mulder, and A. P. E. ten Kroode, 2000, Automatic cross-well tomography by
25 semblance and differential semblance optimization: theory and gradient computation: *Geophysical*
26 Prospecting, **48**, no. 5, 913–935, doi: 10.1046/j.1365-2478.2000.00217.x.
- 27
28
29 595 Shen, X., I. Ahmed, A. Brenders, J. Dellinger, J. Etgen, and S. Michell, 2017, Salt model building
30 at Atlantis with full-waveform inversion: 87th Annual International Meeting, SEG, Expanded Abstracts,
31
32 596 1507–1511.
- 33
34
35 598 Shin, C., and Y. Ho Cha, 2008, Waveform inversion in the Laplace domain: *Geophysical Journal*
36 International, **173**, 922–931, doi: 10.1111/j.1365-246X.2008.03768.x.
- 37
38
39 600 Shin, C., and Y. Ho Cha, 2009, Waveform inversion in the Laplace-Fourier domain: *Geophysical*
40 Journal International, **177**, 1067–1079, doi: 10.1111/j.1365-246X.2009.04102.x.
- 41
42
43 602 Shin, C., S. Jang, and D.-J. Min, 2001, Improved amplitude preservation for prestack depth
44 migration by inverse scattering theory: *Geophysical Prospecting*, **49**, 592–606, doi: 10.1046/j.1365-
45 2478.2001.00279.x.
- 46
47
48
49
50
51
52
53
54
55
56
57
58
59
60

- 1
2
3 605 Symes, W., and J. Carazzone, 1991, Velocity inversion by differential semblance optimization:
4
5 606 Geophysics, **56**, no. 5, 654–663. doi: 10.1190/1.1443082.
6
7
8 607 Symes, W. W., 2008, Migration velocity analysis and waveform inversion: Geophysical Prospecting,
9
10 608 **56**, 765–790, doi: 10.1111/j.1365-2478.2008.00698.x.
11
12
13 609 Tao, K., S. P. Grand, and F. Niu, 2017, Full-waveform inversion of triplicated data using a
14
15 normalized-correlation-coefficient-based misfit function: Geophysical Journal International, **210**, no. 3,
16
17 611 1517–1524, doi: 10.1093/gji/ggx249.
18
19
20 612 Tarantola, A., 1984, Inversion of seismic reflection data in the acoustic approximation: Geophysics,
21
22
23 613 **49**, no. 8, 1259–1266, doi: 10.1190/1.1441754.
24
25
26 614 Tikhonov, A. N. and V. Y. Arsenin, 1977, Solution of Ill-posed Problems: John Wiley & Sons.
27
28
29 615 Trinh, P.T., R. Brossier, L. Métivier, J. Virieux, and P. Wellington, 2017, Bessel smoothing filter
30
31 for spectral-element mesh: Geophysical Journal International, **209**, no. 3, 1489–1512, doi:
32
33 617 10.1093/gji/ggx103.
34
35
36 618 van Leeuwen, T., and W. A. Mulder, 2010, A correlation-based misfit criterion for wave-equation
37
38 travelttime tomography: Geophysical Journal International, **182**, no. 3, 1383–1394, doi: 10.1111/j.1365-
39
40 620 246X.2010.04681.x.
41
42
43 621 van Leeuwen, T., and F. J. Herrmann, 2013, Mitigating local minima in full-waveform inversion by
44
45 expanding the search space: Geophysical Journal International, **195**, no. 1, 661–667, doi:
46
47 623 10.1093/gji/ggt258.
48
49
50 624 Virieux, J., 1986, P-SV wave propagation in heterogeneous media: Velocity-stress finite-difference
51
52 method: Geophysics, **51**, no. 4, 889–901, doi: 10.1190/1.1442147.
53
54
55 626 Virieux, J., and S. Operto, 2009, An overview of full-waveform inversion in exploration geophysics:
56
57
58 627 Geophysics, **74**, no. 6, WCC1–WCC26, doi: 10.1190/1.3238367.

1
2
3 628 Wang, M., Y. Xie, W. Xu, F. Loh, K. Xin, B. Chuah, T. Manning, and S. Wolfarth, 2016, Dynamic-
4 warping full-waveform inversion to overcome cycle skipping: 86th Annual International Meeting, SEG,
5
6 629 Expanded Abstracts, 1273-1277.
7
8

9
10
11 631 Wang, Y., W. Liang, Z. Nashed, X. Li, G. Liang, and C. Yang, 2014, Seismic modeling by
12 optimizing regularized staggered-grid finite-difference operators using a time-space-domain
13 dispersion-relationship-preserving method: *Geophysics*, **79**, no. 5, T277–T285, doi: 10.1190/geo2014-
14 633 0078.1.
15
16

17
18
19 635 Wang, Z., J. Li, B. Wang, Y. Xu, and X. Chen, 2017, Time-domain explicit finite-difference method
20 based on the mixed-domain function approximation for acoustic wave equation: *Geophysics*, **82**, no. 5,
21
22 636 T237–T248, doi: 10.1190/geo2017-0012.1.
23
24

25
26
27 638 Warner, M., A. Ratcliffe, T. Nangoo, J. Morgan, A. Umpleby, N. Shah, V. Vinje, I. Štekl, L. Guasch,
28
29 639 C. Win, G. Conroy, and A. Bertrand, 2013, Anisotropic 3D full-waveform inversion: *Geophysics*, **78**,
30
31 640 no. 2, R59–R80, doi: 10.1190/geo2012-0338.1.
32
33

34
35
36 641 Warner, M., and L. Guasch, 2016, Adaptive waveform inversion: Theory: *Geophysics*, **81**, no. 6,
37 R429–R445, doi: 10.1190/geo2015-0387.1.
38
39

40
41
42 643 Weiss, R. M., and J. Shragge, 2013, Solving 3D anisotropic elastic wave equations on parallel GPU
43 devices: *Geophysics*, **78**, no. 2, F7–F15, doi: 10.1190/geo2012-0063.1.
44
45

46
47
48 645 Wong, J., L. Han, J. C. Bancroft, and R. R. Stewart, 2009, Automatic time-picking of first arrivals
49 on noisy microseismic data, CSEG conference Abstract 2009,
50
51 646 https://www.crewes.org/ForOurSponsors/ConferenceAbstracts/2009/CSEG/Wong_CSEG_2009.pdf,
52
53 647 accessed 25 May 2018
54
55

56
57 649 Wu, R., J. Luo, and B. Wu, 2014, Seismic envelope inversion and modulation signal model:
58
59 650 *Geophysics*, **79**, no. 3, WA13–WA24, doi: 10.1190/geo2013-0294.1.
60

- 1
2
3 651 Yang, Y., and B. Engquist, 2018, Analysis of optimal transport and related misfit functions in FWI:
4
5 652 *Geophysics*, **83**, no. 1, A7-A12, doi: 10.1190/geo2017-0264.1.
6
7 653 Yang, Y., B. Engquist, J. Sun, and B. Froese, 2018, Application of Optimal Transport and the
8 Quadratic Wasserstein Metric to Full-Waveform Inversion: *Geophysics*, **83**, no. 1, R43-R62, doi:
9
10 655 10.1190/geo2016-0663.1.
11
12
13 656 Yao, G., and D. Wu, 2015, Least-squares reverse-time migration for reflectivity imaging: *Science
14 China Earth Sciences*, **58**, 1982–1992, doi: 10.1007/s11430-015-5143-1.
15
16
17 658 Yao, G., and H. Jakubowicz, 2016, Least-squares reverse-time migration in a matrix-based
18 formulation: *Geophysical Prospecting*, **64**, no. 3, 611-621, doi: 10.1111/1365-2478.12305.
19
20
21 660 Yao, G., D. Wu, and H. A. Debens, 2016, Adaptive finite difference for seismic wavefield modelling
22 in acoustic media: *Scientific Reports*, **6**, 30302, doi: 10.1038/srep30302.
23
24
25 662 Yao, G., N. V. da Silva, M. Warner, A. Umpleby, and D. Wu, 2017, Improved FWI convergence
26 using efficient receiver-side spatial preconditioning employing ray theory: 79th EAGE Conference and
27 Exhibition, EAGE, Extended Abstracts, We P1 14.
28
29
30 665 Yao, G., N. V. da Silva, H. A. Debens, and D. Wu, 2018a, Accurate seabed modeling using finite
31 difference methods: *Computational Geosciences*, **22**, no. 2, 469-484, doi: 10.1007/s10596-017-9705-5.
32
33
34 667 Yao, G., N. V. Silva, M. Warner, and T. Kalinicheva, 2018b, Separation of migration and
35 tomography modes of full-waveform inversion in the plane wave domain: *Journal of Geophysical
36 Research: Solid Earth*, **123**, no. 2, 1486-1501, doi: doi:10.1002/2017JB015207.
37
38
39 670 Yuan, S., J. Liu, S. Wang, T. Wang, and P. Shi, 2018, Seismic Waveform Classification and First-
40 Break Picking Using Convolution Neural Networks: *IEEE Geoscience and Remote Sensing Letters*, **15**,
41 no. 2, 272-276, doi: 10.1109/LGRS.2017.2785834.
42
43
44
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48
49
50
51
52
53
54
55
56
57
58
59
60

- 1
2
3 673 Zelt, C. A., and R. B. Smith, 1992, Seismic traveltimes inversion for 2-D crustal velocity structure:
4
5 674 Geophysical Journal International, **108**, no. 1, 16-34, doi: 10.1111/j.1365-246X.1992.tb00836.x.
6
7
8 675 Zhang, J., and Z. Yao, 2012, Optimized finite-difference operator for broadband seismic wave
9
10 modeling: Geophysics, **78**, no. 1, A13–A18, doi: 10.1190/geo2012-0277.1.
11
12
13 677 Zhang, Y., and B. Biondi, 2013, Moveout-based wave-equation migration velocity analysis:
14
15 678 Geophysics, **78**, no. 2, U31–U39, doi: 10.1190/geo2012-0082.1.
16
17
18 679 Zhu, H., E. Bozdağ, D. Peter, and J. Tromp, 2012, Structure of the European upper mantle revealed
19
20 by adjoint tomography: Nature Geoscience, **5**, 493, doi: 10.1038/ngeo1501
21
22
23 681 Zhu, H., and S. Fomel, 2016, Building good starting models for full-waveform inversion using
24
25 adaptive matching filtering misfit: Geophysics, **81**, no. 5, U61-U72, doi: 10.1190/geo2015-0596.1.
26
27
28 683 Zuberi, M. A., and R. G. Pratt, 2018, Mitigating nonlinearity in full waveform inversion using
29
30 scaled-Sobolev pre-conditioning: Geophysical Journal International, **213**, no. 1, 706-725, doi:
31
32 685 10.1093/gji/ggx549.
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List of figure caption:

688 Figure 1. Schematic illustration of cycle-skipping. (a) The value of the objective function, which is the
689 square of the L2-norm of the difference between a 5-Hz Ricker wavelet and its shifts. The two black
690 dots represent the global minimum and a local minimum. The nearest peaks to the global minimum has
691 a time lag of 86 ms, which corresponds to half a cycle, $0.5T$. (b) The 5-Hz Ricker wavelet and its shift
692 by 50 ms, which is less than half a cycle. (c) The same as (b) but the shift is 125 ms, which is greater
693 than half a cycle. The functional value for the case of (b) and (c) is indicated by the two crosses in (a).

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695 Figure 2. Schematic illustration for the method of tackling cycle-skipping with intermediate data. The
696 predicted data 1 is generated with the initial model. The recorded data indicated by the red trace has a
697 time difference of Δt_0 relative to predicted data 1. Δt_0 is greater than half a cycle. Shifting the
698 predicted data by Δt_s , which is smaller than half a cycle, generates intermediate data 1. Replacing the
699 record with intermediate data 1, FWI will update the model to make the predicted data fit intermediate
700 data 1. Then shifting the new predicted data creates the new intermediate data. By repeating the process,
701 the inversion will recover a model, which can produce the predicted data having a time difference to
702 the record less than half a cycle.

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704 Figure 3. Schematic illustration of creating intermediate data by using the first arrivals.

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706 Figure 4. A model with two Gaussian anomalies. (a) The true velocity model, which includes a constant
707 background of 3000 m/s, a high-velocity anomaly and a low-velocity anomaly. The anomalies are
708 formed with a Gaussian function of $\pm 1000 e^{-\frac{d^2}{5 \times 10^5}}$, where d (unit: m/s) is the distance to the center

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3 709 of the anomaly. The red dotted line indicates the location of sources while the blue dotted line represents
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5 710 the receiver arrays. (b) The initial velocity model, which is a constant velocity of 2800 m/s.
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712 Figure 5. Overcoming cycle-skipping with intermediate data. (a) Alternating display of a shot gather of
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719
720 Figure 6. The gradient of (a) conventional FWI and (b) FWI with intermediate data in the first iteration.
721 During the inversion, the model is updated through slowness, so that the negative gradient (blue color)
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723 velocity.

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725 Figure 7. The recovered velocity models from (a) conventional FWI starting from the initial model, (b)
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728 Figure 8. The Marmousi models and the initial models for inversion. (a) The true P-wave velocity model.
729 (b) The true S-wave velocity model. (c) The simple 1D initial P-wave velocity model for the surface-
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731 anomaly is embedded around the fault area, for the test with a surface acquisition. In (a), the red dotted

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3 732 line indicates the locations of the sources and the receiver array of the surface acquisition while the blue
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5 733 dotted line represents the locations of the receiver array of the surface-to-horizontal-well setting.
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35 769 and x and z are in number of cells. The green dotted lines indicate the sea bottom in the Marmousi
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43 772 Figure 16. Evolution of the first breaks throughout the iterations. The red curve indicates the first break
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57 778 Figure 17. Recovered P-wave velocity models. (a) The recovered P-wave velocity model by using
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7 782 generated by shifting the observed data towards the predicted data. (d) The recovered P-wave velocity
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9 783 model with conventional FWI starting from the velocity model shown in (b).

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14 785 Figure 18. (a) Velocity profile at a distance of 6.25 km. (b) Normalized L2-norm model errors. The
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16 black solid curves indicate the true velocity model while the dashed black curves are for the initial
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18 velocity model. The red and blue curves are for the recovered velocity model shown in Figures 17b and
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20 786 c, respectively.

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24 787 Figure B-1. Comparison of one shot gather generated from the surface-to-horizontal-well geometry with
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50 comparison at trace (d) 600, (e) 700 and (f) (800). Blue, red and black curves represent the traces from
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60 803 Figure B-3. Gradients of FWI with the intermediate data for the surface-to-horizontal-well geometry at
804 different iterations. The green dotted lines indicate the sea bottom in the Marmousi model.

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807 horizontal-well geometry at different iterations.

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Figure B-5. Gradients of FWI with the intermediate data for the surface geometry at different iterations.
The green dotted lines indicate the sea bottom in the Marmousi model.

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Figure B-6. Recovered P-wave velocity model using FWI with the intermediate data for the surface geometry at different iterations.

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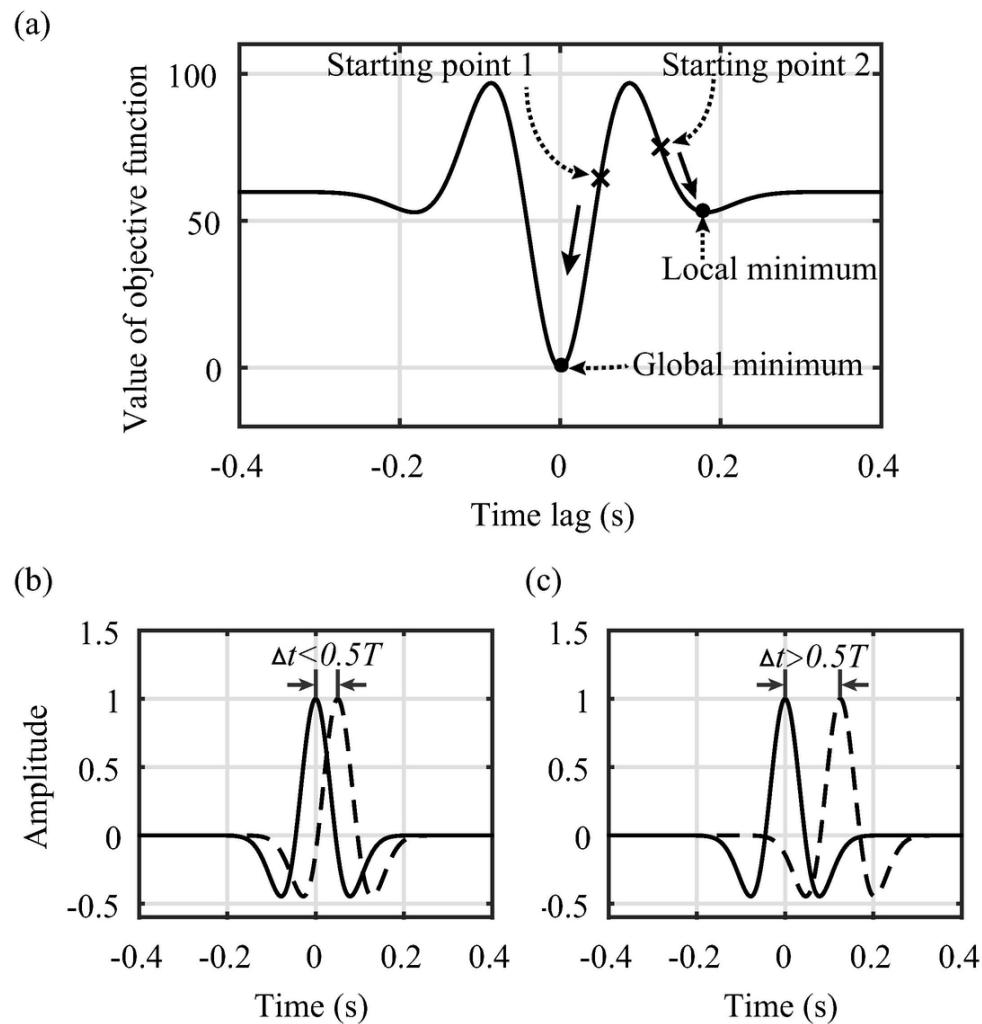


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94x96mm (300 x 300 DPI)

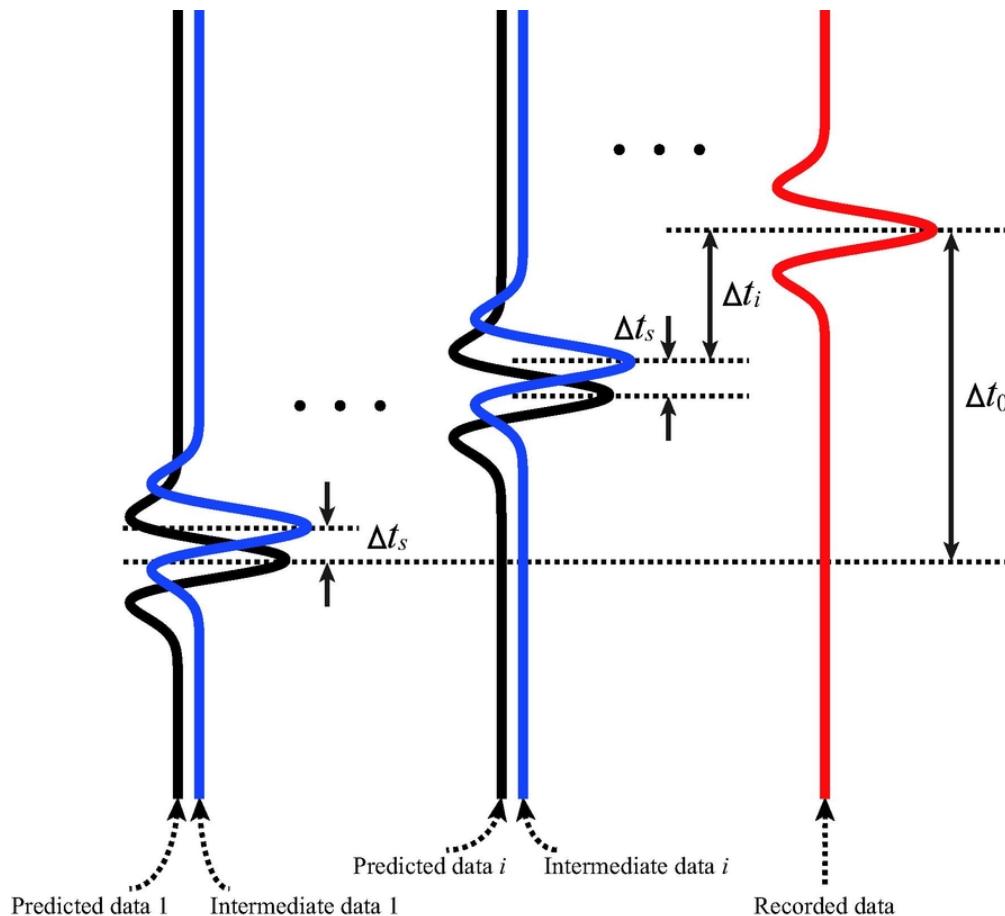


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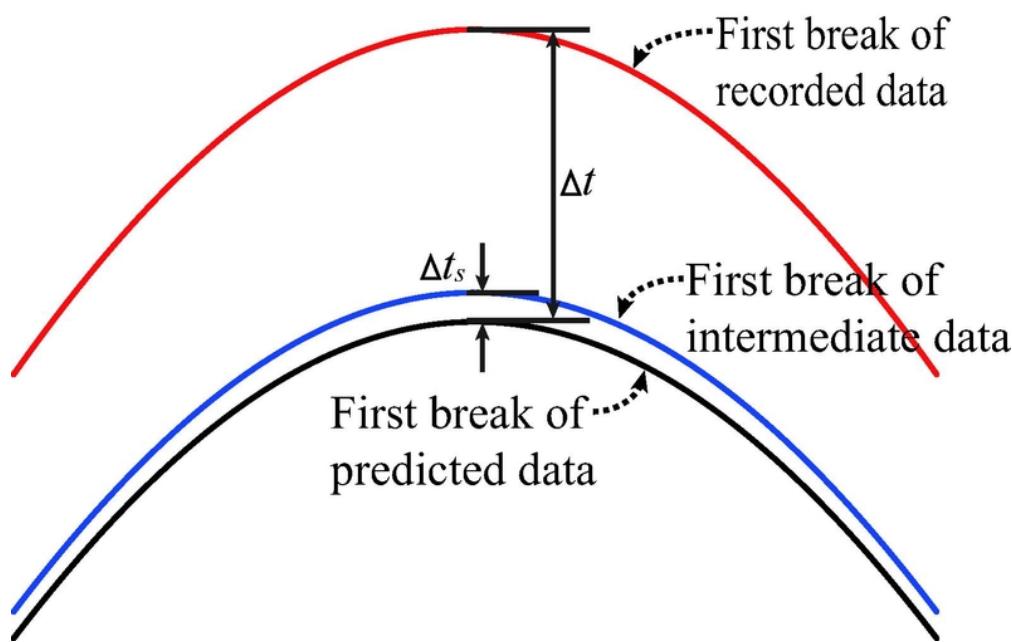


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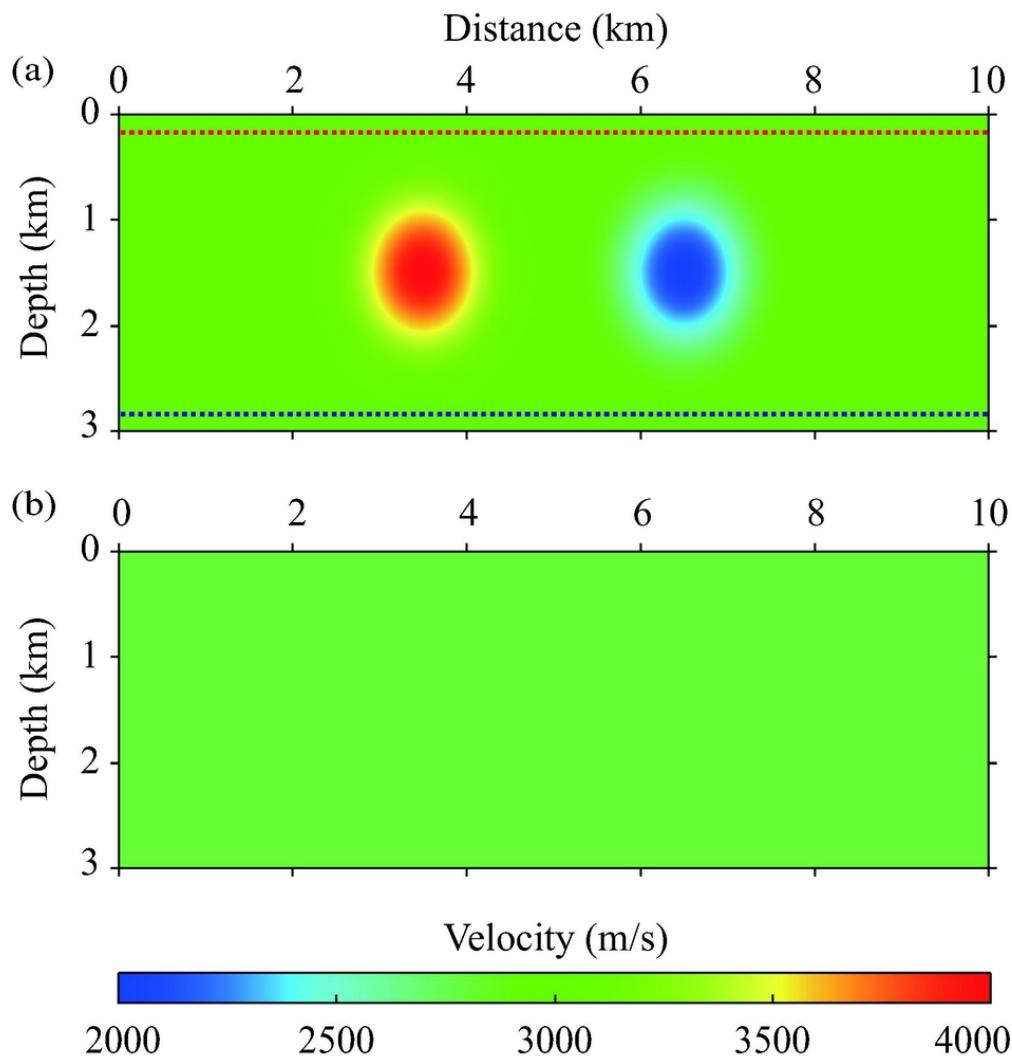


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80x84mm (300 x 300 DPI)

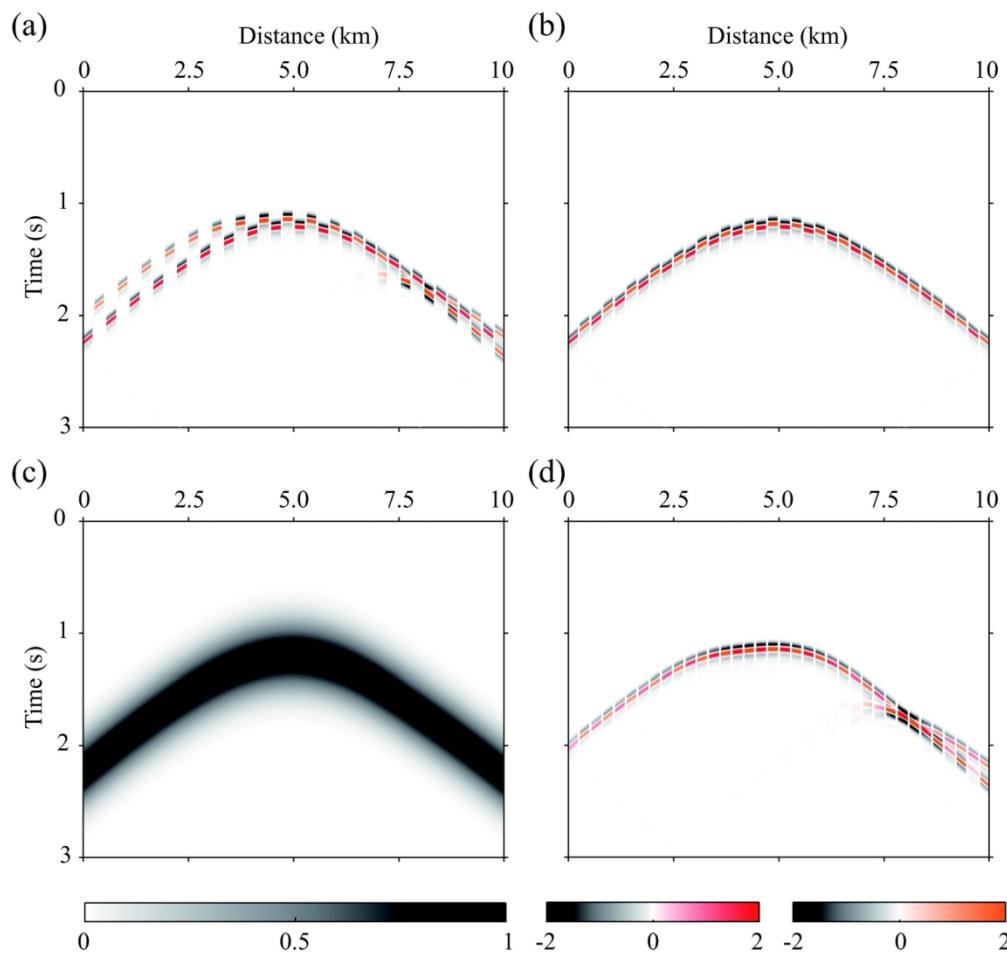


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128x122mm (300 x 300 DPI)

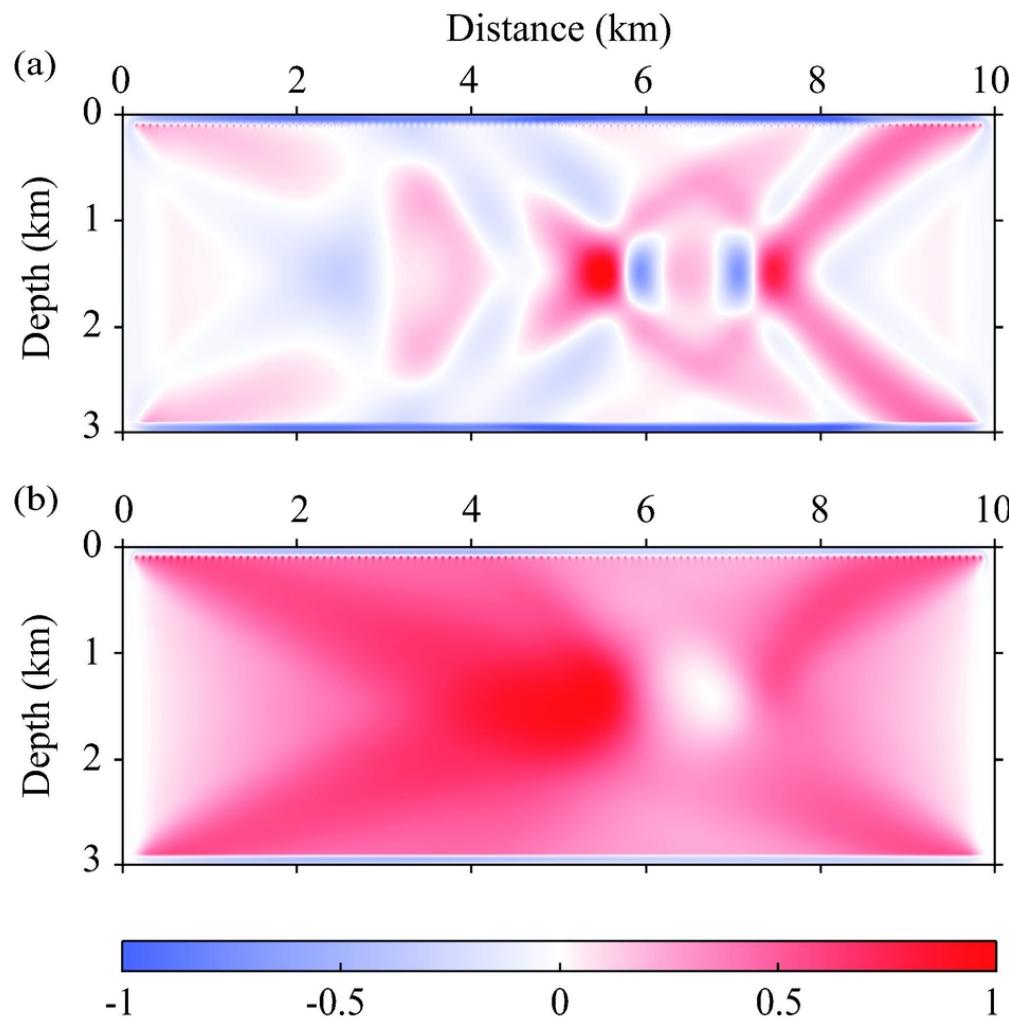
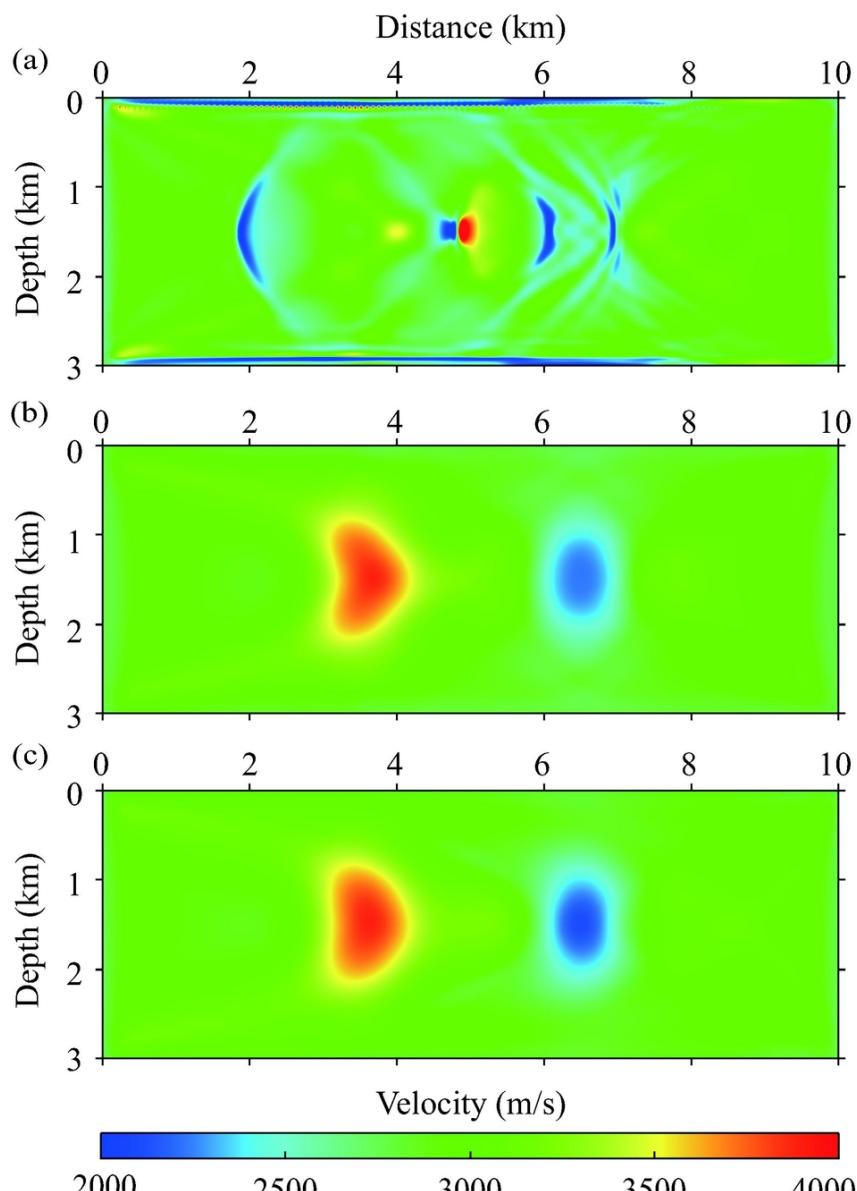


Figure 6. The gradient of (a) conventional FWI and (b) FWI with intermediate data in the first iteration. During the inversion, the model is updated through slowness, so that the negative gradient (blue color) indicates FWI will decrease velocity while the positive gradient (red color) means FWI will increase velocity.

80x81mm (300 x 300 DPI)



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Figure 7. The recovered velocity models from (a) conventional FWI starting from the initial model, (b) FWI with intermediate data starting from the initial model, and (c) conventional FWI starting from (b).

80x113mm (300 x 300 DPI)

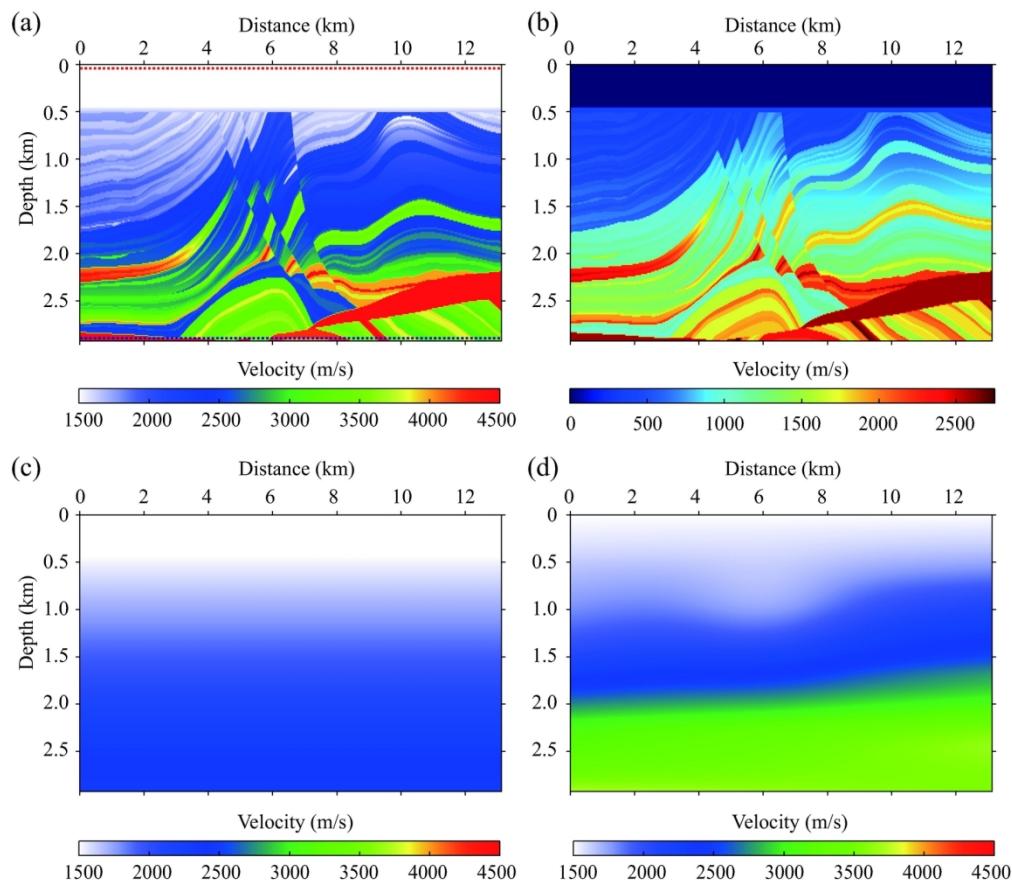


Figure 8. The Marmousi models and the initial models for inversion. (a) The true P-wave velocity model. (b) The true S-wave velocity model. (c) The simple 1D initial P-wave velocity model for the surface-to-horizontal-well test. (d) The simple initial P-wave velocity model, in which a Gaussian low-velocity anomaly is embedded around the fault area, for the test with a surface acquisition. In (a), the red dotted line indicates the locations of the sources and the receiver array of the surface acquisition while the blue dotted line represents the locations of the receiver array of the surface-to-horizontal-well setting.

155x136mm (300 x 300 DPI)

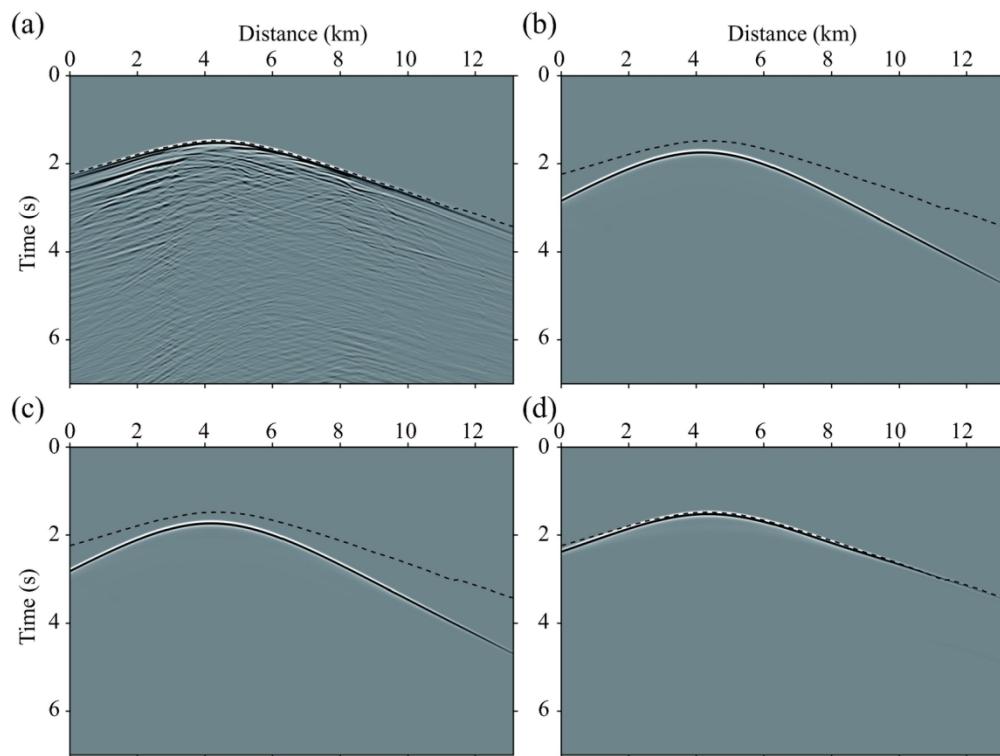


Figure 9. One shot gather in FWI with intermediate data for the surface-to-horizontal-well test. (a) Recorded data. (b) Predicted data with the initial velocity model shown in Figure 8c. (c) The intermediate data generated by shifting the predicted data shown in (b) within half a cycle. (d) The predicted data with the recovered model shown in Figure 12b. The dashed black curves indicate the first-break picks of the record shown in (a).

139x105mm (300 x 300 DPI)

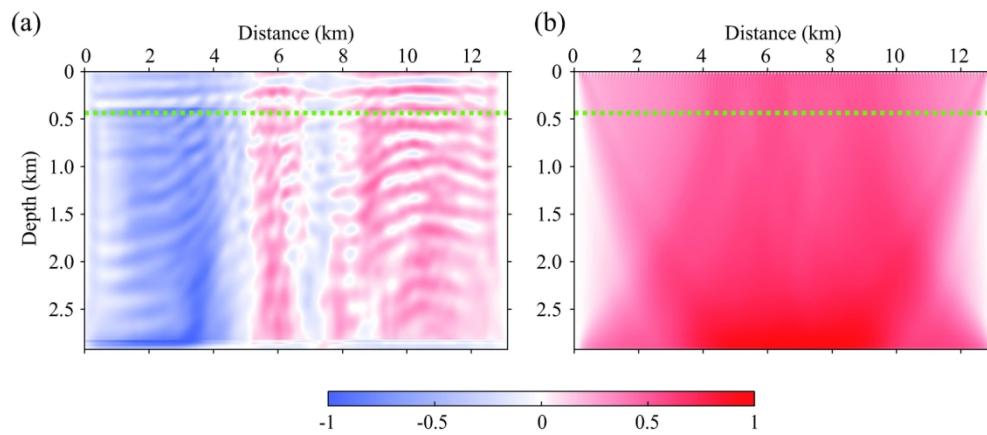


Figure 10. The gradient of (a) conventional FWI and (b) FWI with intermediate data in the first iteration. The green dotted lines indicate the sea bottom in the Marmousi model.

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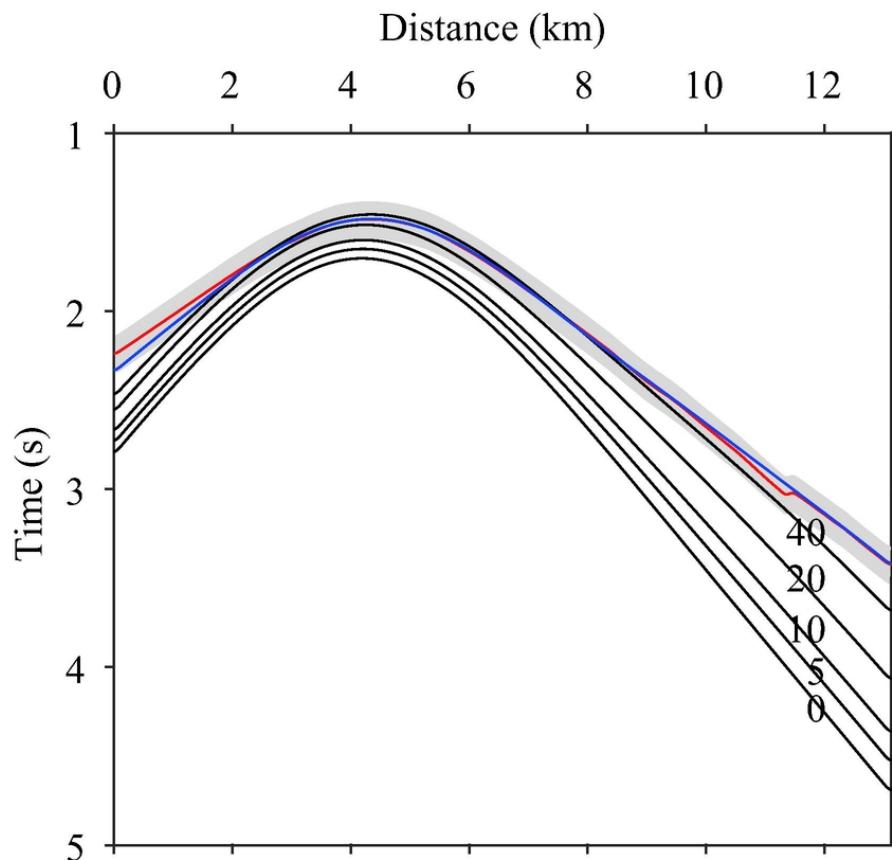


Figure 11. Evolution of first breaks throughout the iterations. The red curve indicates the first break of the recorded data while the blue curve represents the first break of the predicted data after 60 iterations. The black curves denote the first breaks for early iterations. The iteration numbers are drawn on top of these curves. The light gray strip around the red curve indicates the half period of the 5-Hz Ricker wavelet.

80x76mm (300 x 300 DPI)

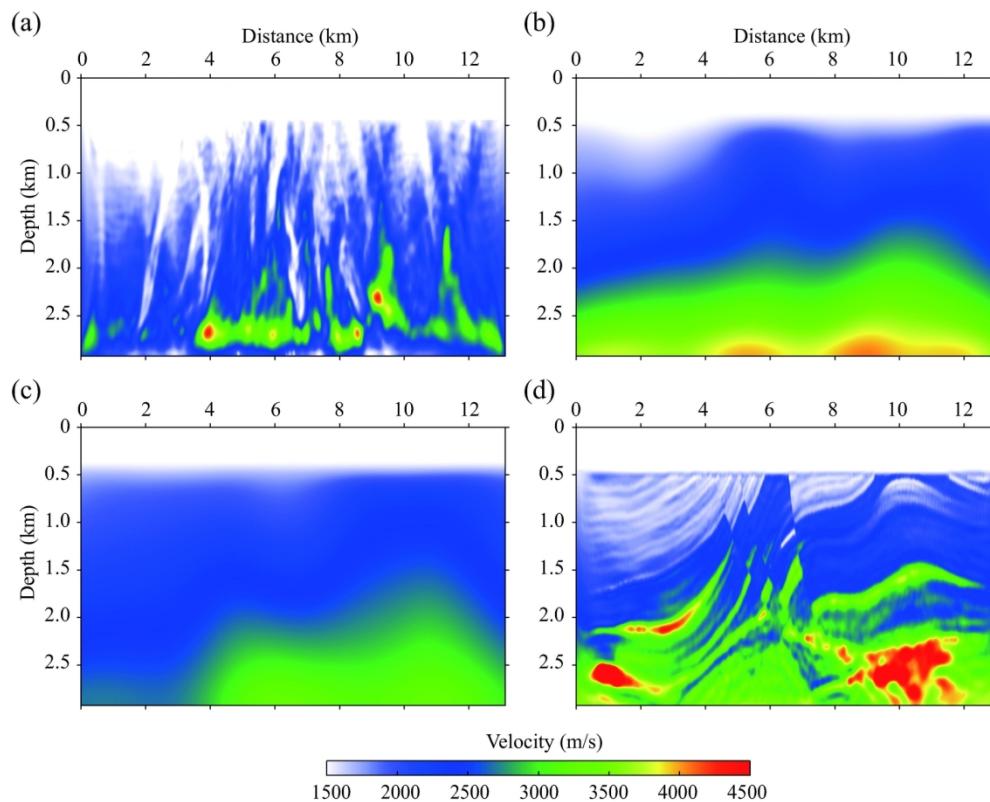


Figure 12. Recovered P-wave velocity models. (a) The recovered P-wave velocity model by using conventional FWI starting from the initial model shown in Figure 8c. (b) The recovered P-wave velocity model from the initial velocity model shown in Figure 8c by using FWI with intermediate data shifting of the predicted data towards the observed data. (c) The same as (b) but the intermediate data are generated by shifting the observed data towards the predicted data. (d) The recovered P-wave velocity with the conventional FWI starting from the velocity shown in (b).

154x123mm (300 x 300 DPI)

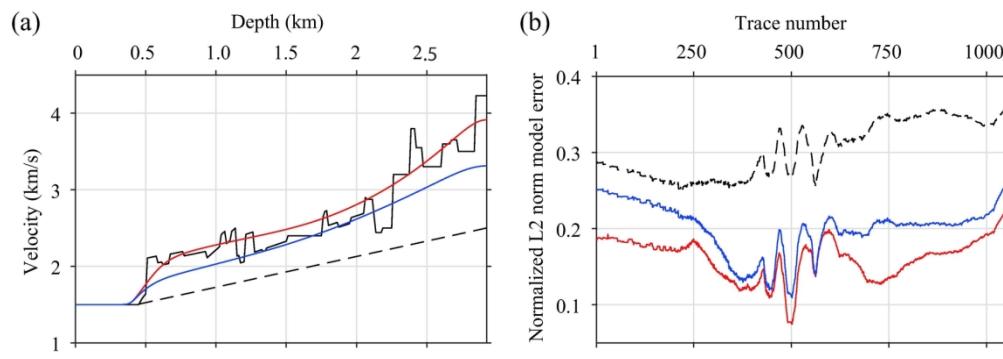


Figure 13. (a) Velocity profile at a distance of 6.25 km. (b) Normalized L2-norm model errors. The black solid curves indicate the true velocity model while the dashed black curves are for the initial velocity model. The red and blue curves are for the recovered velocity model shown in Figures 12b and 12c, respectively.

158x55mm (300 x 300 DPI)

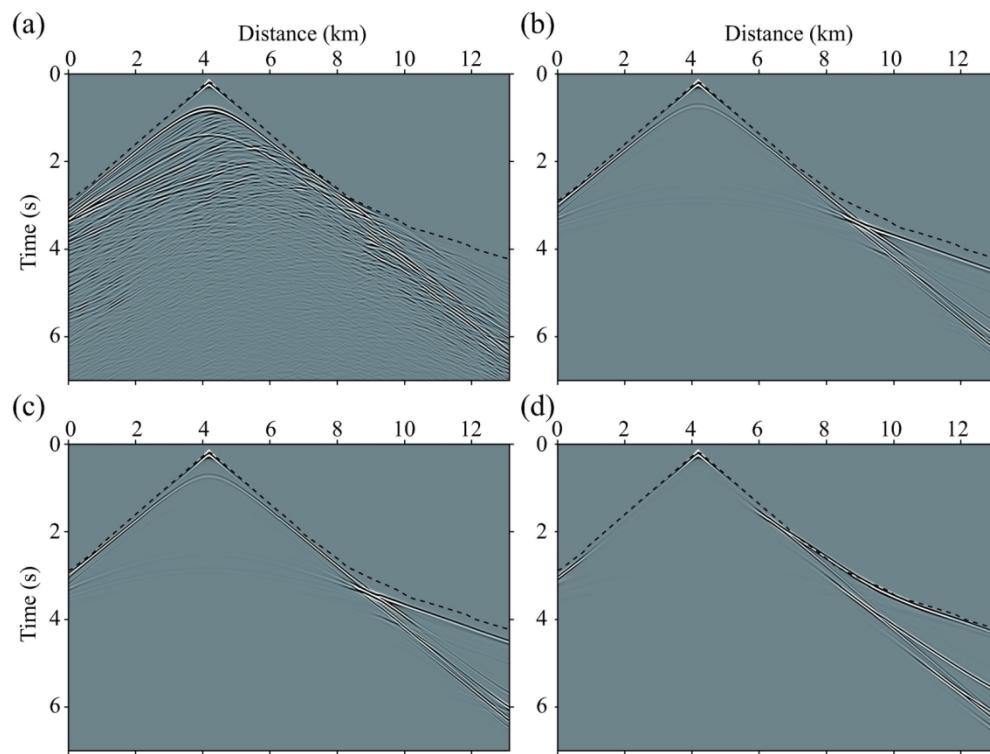
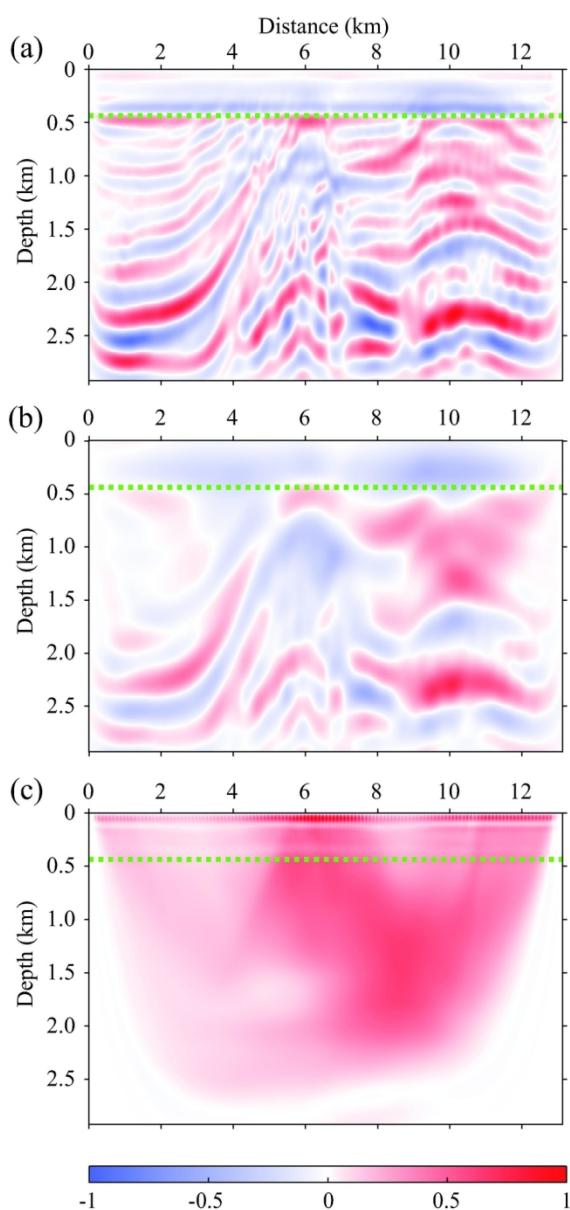


Figure 14. One shot gather in FWI with intermediate data for the test of the surface acquisition geometry.

(a) Recorded data. (b) Predicted data with the initial velocity model shown in Figure 8d. (c) The intermediate data generated by shifting the predicted data shown in (b) within half a cycle. (d) The predicted data with the recovered model shown in Figure 17b. The dashed black curves indicate the first-break picks of the record shown in (a).

140x105mm (300 x 300 DPI)



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Figure 15. The gradient of (a) conventional FWI and (c) FWI with the intermediate data in the first iteration. (b) A smoothed version of (a) using a Gaussian filter, i.e. $1/(2\pi\sigma)e^{-(x^2+z^2)/2\sigma^2}$, where $\sigma=10$, and x and z are in number of cells. The green dotted lines indicate the sea bottom in the Marmousi model.

78x165mm (300 x 300 DPI)

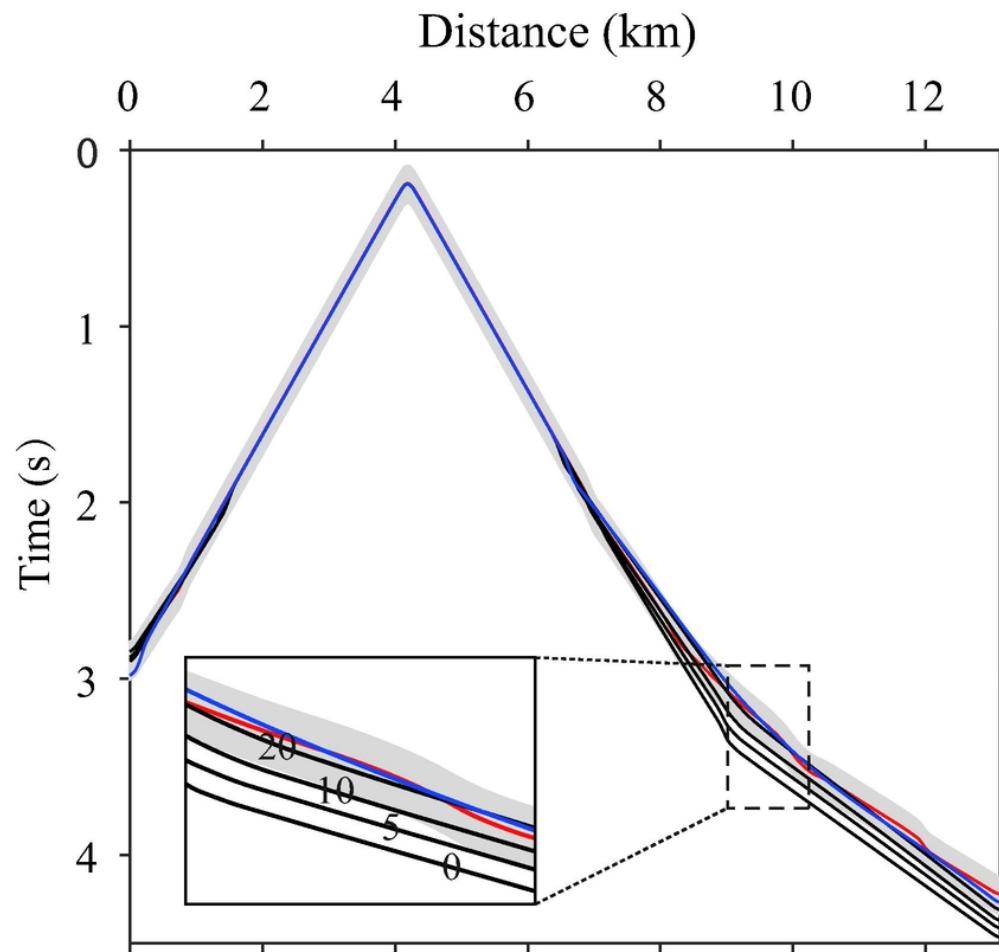


Figure 16. Evolution of the first breaks throughout the iterations. The red curve indicates the first break of the recorded data while the blue curve represents the first break of the predicted data after 60 iterations. The black curves denote the first breaks for early iterations. The iteration numbers are drawn on top of these curves. The light gray strip around the red curve indicates the half period of the 5-Hz Ricker wavelet.

71x67mm (300 x 300 DPI)

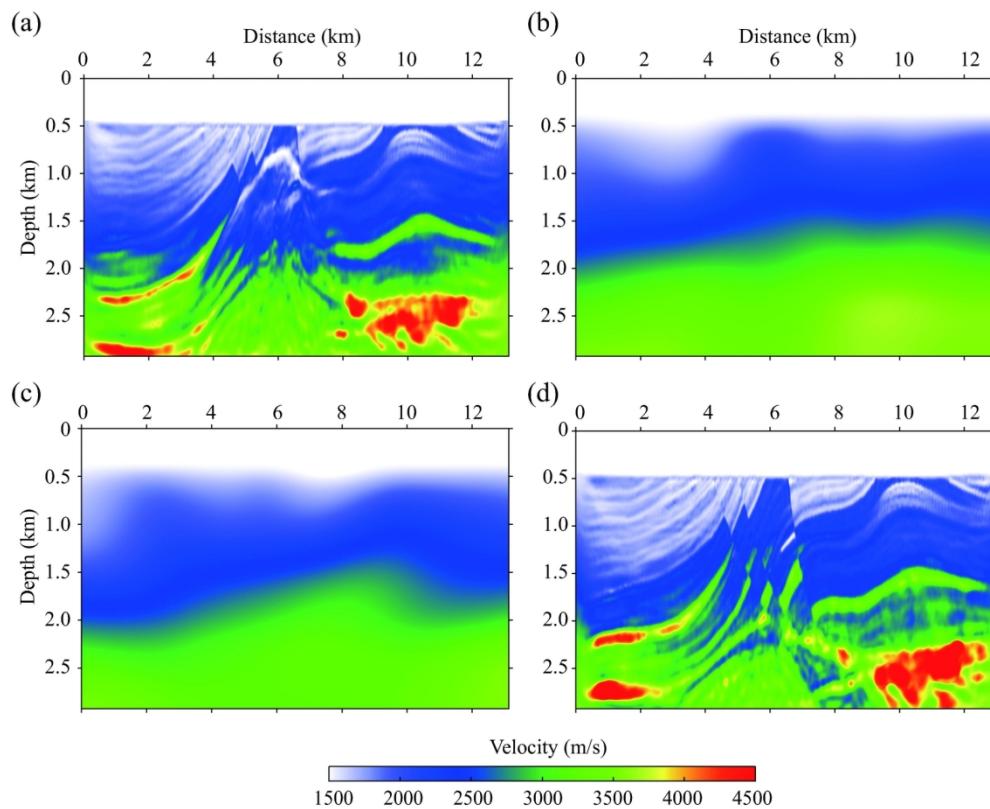


Figure 17. Recovered P-wave velocity models. (a) The recovered P-wave velocity model by using conventional FWI starting from the initial model shown in Figure 8d. (b) The recovered P-wave velocity model from the initial velocity model shown in Figure 8d by using FWI with intermediate data shifting of the predicted data towards the observed data. (c) The same as (b) but the intermediate data are generated by shifting the observed data towards the predicted data. (d) The recovered P-wave velocity model with conventional FWI starting from the velocity model shown in (b).

153x123mm (300 x 300 DPI)

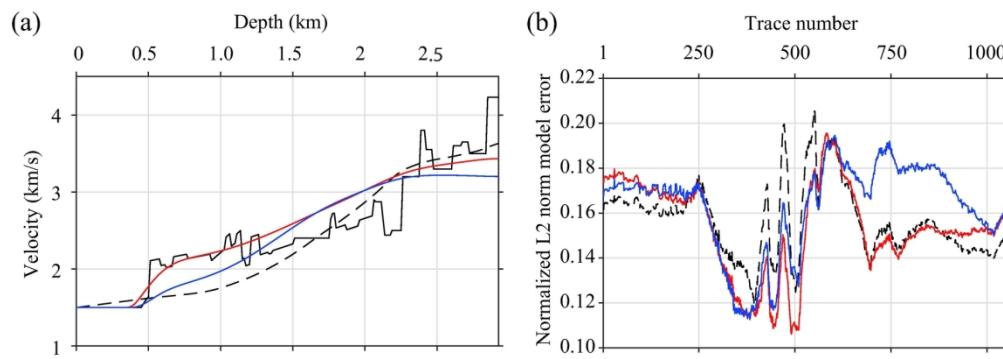


Figure 18. (a) Velocity profile at a distance of 6.25 km. (b) Normalized L2-norm model errors. The black solid curves indicate the true velocity model while the dashed black curves are for the initial velocity model. The red and blue curves are for the recovered velocity model shown in Figures 17b and c, respectively.

156x64mm (300 x 300 DPI)

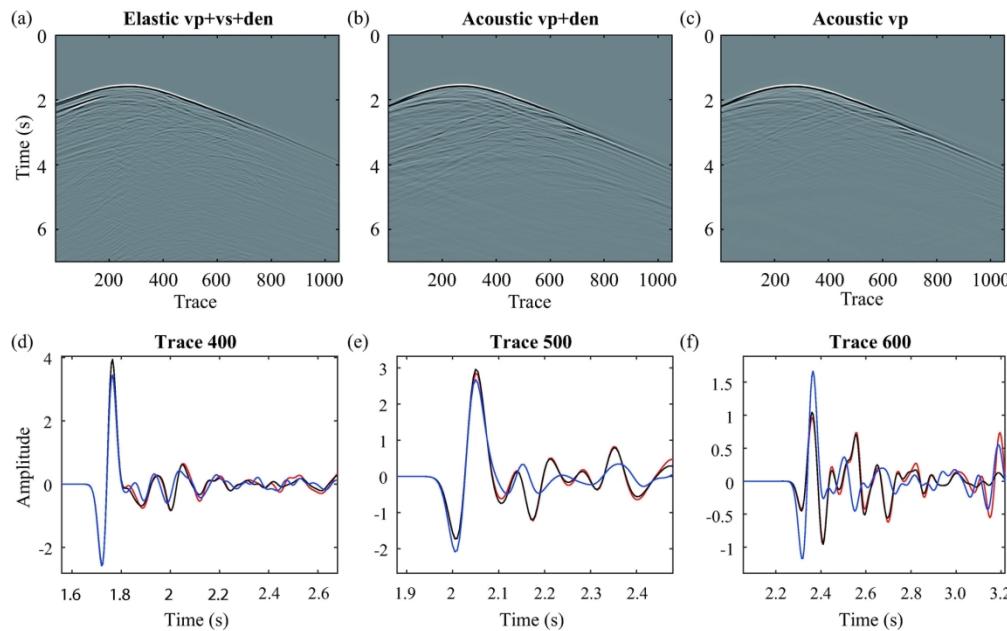


Figure B-1. Comparison of one shot gather generated from the surface-to-horizontal-well geometry with different wave equations from the Marmousi models. (a) Using the elastic wave equation with the P-wave and S-wave velocities and density models. (b) Using the acoustic wave equation with the P-wave velocity and density models. (c) Using the acoustic wave equation with the P-wave velocity model. The trace comparison at trace (d) 600, (e) 700 and (f) (800). Blue, red and black curves represent the traces from (a), (b) and (c), respectively.

163x108mm (300 x 300 DPI)

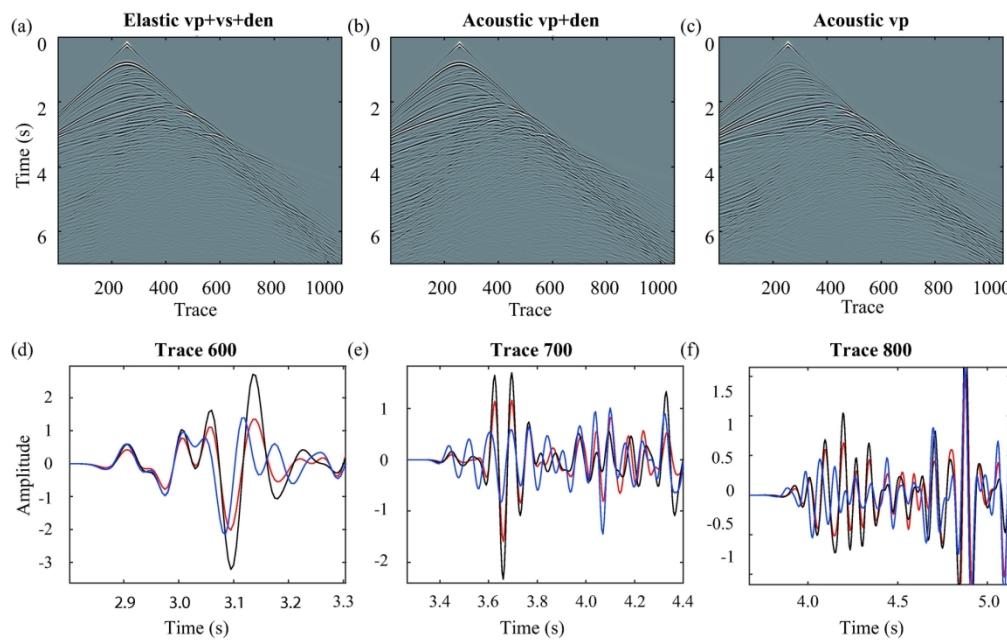


Figure B-2. Comparison of one shot gather generated from the surface geometry with different wave equations from the Marmousi models. (a) Using the elastic wave equation with the P-wave and S-wave velocities and density models. (b) Using the acoustic wave equation with the P-wave velocity and density models. (c) Using the acoustic wave equation with the P-wave velocity model. The trace comparison at trace (d) 600, (e) 700 and (f) (800). Blue, red and black curves represent the traces from (a), (b) and (c), respectively.

163x109mm (300 x 300 DPI)

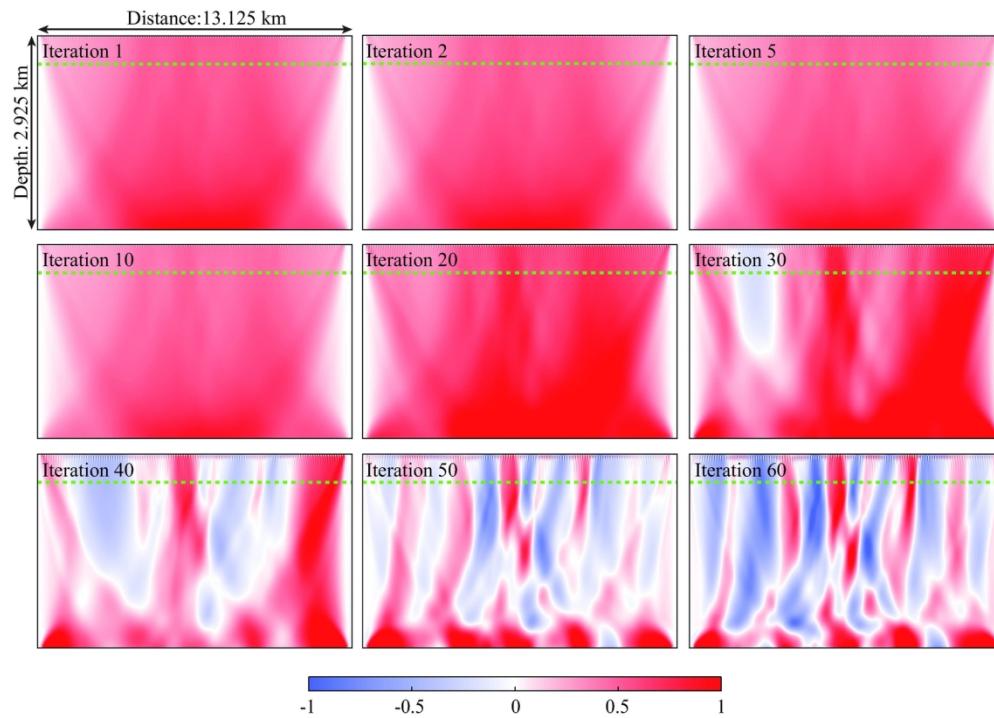


Figure B-3. Gradients of FWI with the intermediate data for the surface-to-horizontal-well geometry at different iterations. The green dotted lines indicate the sea bottom in the Marmousi model.

160x113mm (300 x 300 DPI)

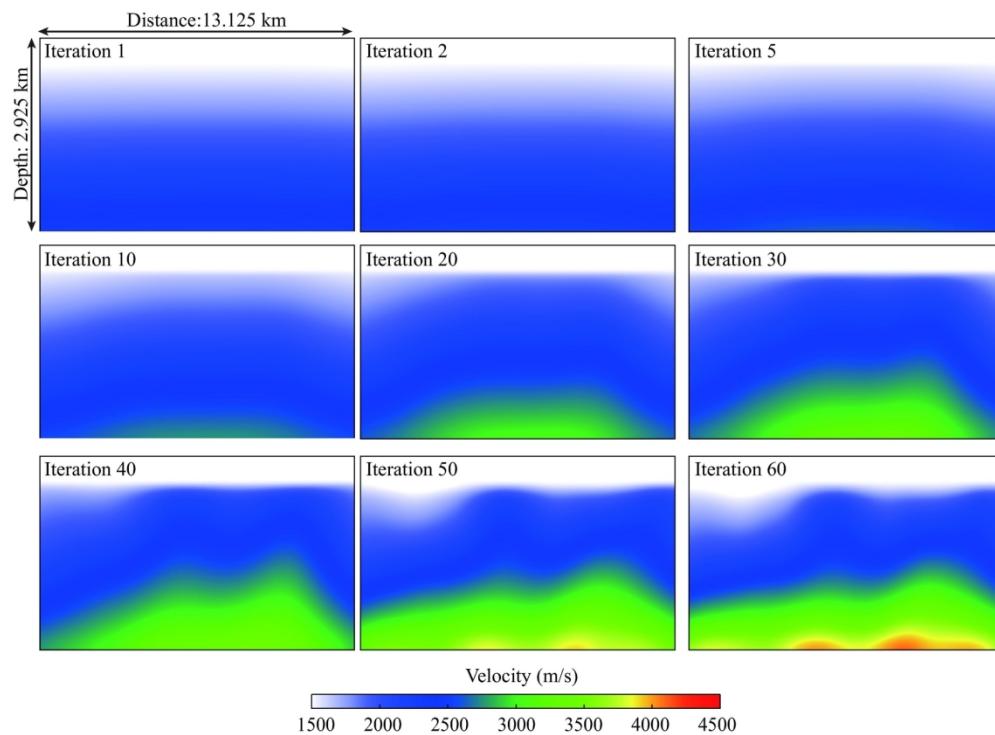


Figure B-4. Recovered P-wave velocity model using FWI with the intermediate data for the surface-to-horizontal-well geometry at different iterations.

160x116mm (300 x 300 DPI)

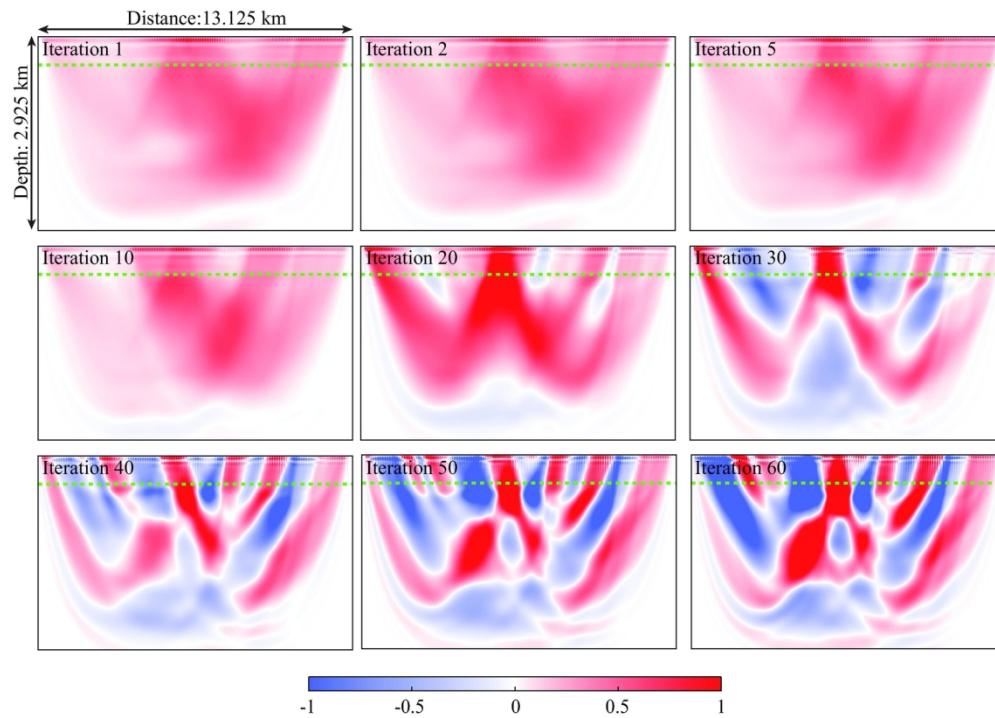


Figure B-5. Gradients of FWI with the intermediate data for the surface geometry at different iterations. The green dotted lines indicate the sea bottom in the Marmousi model.

160x113mm (300 x 300 DPI)

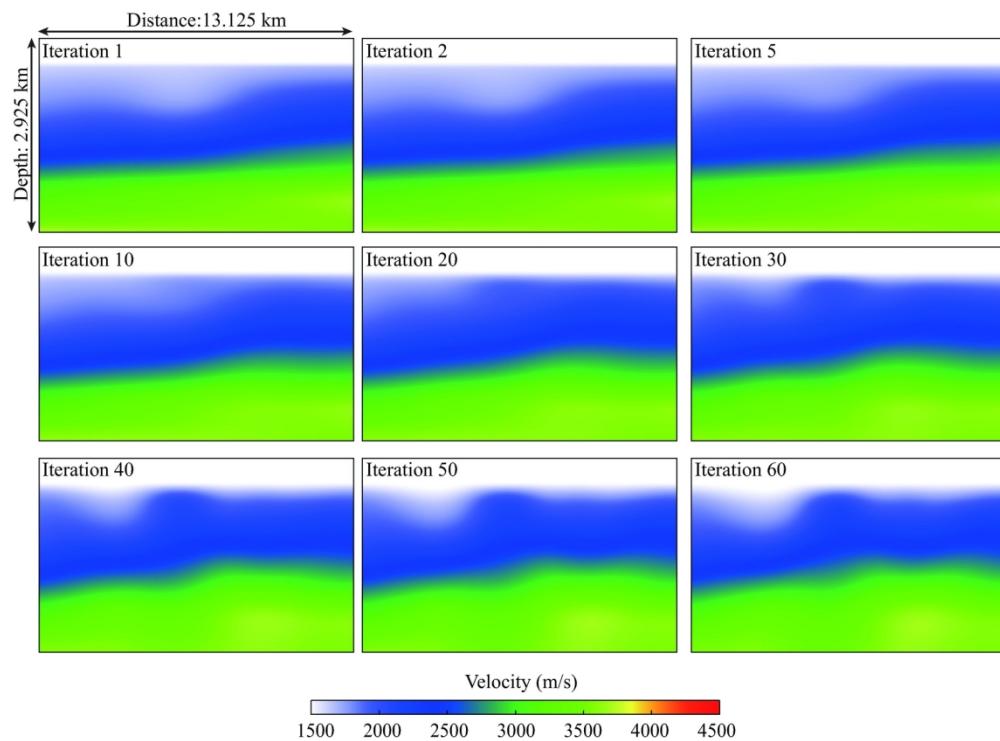


Figure B-6. Recovered P-wave velocity model using FWI with the intermediate data for the surface geometry at different iterations.

160x117mm (300 x 300 DPI)

DATA AND MATERIALS AVAILABILITY

Data associated with this research are available and can be obtained by contacting the corresponding author.