

# Warm Molecular Gas in M51: Mapping H<sub>2</sub> Excitation and Mass with the Spitzer Infrared Spectrograph

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## ABSTRACT

We have mapped the molecular gas in a radial strip across the disk of M51 using  $\text{H}_2$  S(0) –  $\text{H}_2$  S(5) pure rotational mid-infrared emission lines using the Infrared Spectrograph (IRS) on the Spitzer Space Telescope. The morphology of the molecular gas changes significantly in the different line maps. We find that the  $\text{H}_2$  excitation temperature varies from 100 – 600 K across the galactic disk. The nucleus contains the hottest gas ( $T \sim 600$  K) and the interarm regions have the coolest ( $T \sim ??$ ) gas. We measured the  $\text{H}_2$  mass surface density by modeling the gas distribution with a two temperature model: a warm ( $T = 100 - 300$  K) phase using the low  $J$  lines and a hot phase ( $T = 400 - 600$  K) using the high  $J$  lines. The hot phase gas surface density peaks at  $0.24 \text{ M}_\odot/\text{pc}^2$  at the nucleus, whereas the warm gas phase surface density reaches a maximum of  $xx \text{ M}_\odot/\text{pc}^2$  in the north-eastern spiral arm. The spatial variation in the hot and warm phases indicates that the molecular gas is being heated by different excitation mechanisms in the nucleus and in the spiral arms. There are significant spatial offsets between the warm  $\text{H}_2$  phase and the colder ( $\sim 10$  K)  $\text{H}_2$  gas traced by CO ( $J = 1 - 0$ ) emission, such that the warm  $\text{H}_2$  is offset towards the downstream side of the cold gas. The warm  $\text{H}_2$  emission is spatially coincident with  $\text{H}\alpha$  emission, indicating that the warm  $\text{H}_2$  gas phase is primarily heated by the HII regions. In contrast the hotter  $\text{H}_2$  emission is spatially coincident with  $[\text{O IV}](25.89 \mu\text{m})$  and X-ray emission, indicating that shocks and/or X-rays are responsible for exciting this phase.

*Subject headings:* galaxies: ISM — galaxies:  $\text{H}_2$  — galaxies: individual(M51)

## 1. Introduction

Star formation and galactic evolution are connected via the molecular gas in a galaxy. In the Milky Way, nearly all star formation occurs in molecular clouds (), although not all molecular clouds are actively forming stars. On a more global, galactic scale, star formation may be triggered whenever the molecular gas surface density is enhanced, for example, by a spiral density wave (), by increased pressure or gas density in galactic nuclei (; ), due to hydrodynamic shocks along the leading edge of bars (; ), and in the transition region at the ends of bars (; ). How does this star formation affect the surrounding molecular gas? How is it heated and what is the distribution of the gas temperatures? How does the mass of the warm and hot gas vary from region to region? We address these questions using a spectral map of a radial strip across the spiral galaxy, M51.

M51 (also known as the Whirlpool galaxy and NGC 5194) is a nearby, face-on spiral galaxy that is rich in molecular gas. Its proximity (assumed to be 8.2 Mpc ()), face-on orientation, and grand-design spiral morphology make it the ideal target for many studies of the interstellar medium (ISM) across distinct dynamical, chemical, and physical environments in a galaxy. Studies of the molecular gas within M51 have revealed giant molecular associations (GMAs) along the spiral arms (), a reservoir of molecular gas associated with the nuclear AGN (), and star-formation in molecular clouds triggered by a spiral density wave (). In addition to being well-studied at millimeter and radio wavelengths, M51 has been also studied at X-ray, UV, optical, near-infrared, infrared, and submillimeter wavelengths (; ?; ; ; ). Here we present new mid-infrared spectra obtained with the Infrared Spectrograph (IRS) () on-board the *Spitzer Space Telescope*. Pure rotational emission lines of  $\text{H}_2$  in the mid-infrared are powerful tools for observing the warmer ( $>100$  K) phase of  $\text{H}_2$ .<sup>1</sup> The mid-infrared  $\text{H}_2$  lines are important diagnostics of

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<sup>1</sup>Although we are always exploring the warm phase of  $\text{H}_2$ , in this paper we always refer

the ISM because they allow us to model the  $\text{H}_2$  excitation-temperature, mass (; ), and ortho-to-para ratio (; ). These quantities, in turn, help us constrain the energy injection mechanism (i.e. radiative heating, shocks, turbulence) that heats the warm molecular gas phase in the ISM.

## 2. Observations and Data Reduction

### 2.1. Spectral Data

We mapped a radial strip across M51 using the short-low (SL) and long-low (LL) modules of the *Spitzer* IRS in spectral mapping mode. The radial strips were  $324'' \times 57''$  and  $295'' \times 51''$  in the SL and LL, respectively. Integration times in the SL and LL were 14.6 s. Each slit position was mapped twice and half-slit spacings were used. In total, 1,412 spectra were taken in the SL and 100 were taken in the LL (including background observations). Dedicated off source background observations were taken for the SL observations. Backgrounds for the LL observations were taken from outrigger data collected while the spacecraft was mapping in the adjacent module. The astronomical observation requests (AORs) are available on SST’s Leopard and Spot (Project ID 200138, PI: K. Sheth).

The spectra were assembled from the basic calibration data (BCD) into spectral data cubes for each module using CUBISM (). Background subtraction and bad pixel removal was done within CUBISM. The individual IRS spectra were processed using the S14.0 version of the Spitzer Science Center (SSC) pipeline. In CUBISM, the SL and LL data cubes have  $1.85''$  pixels and  $5.08''$  pixels respectively. This pixel size is half the point spread

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to a warm and hot phase corresponding to temperatures of,  $T = 100 - 300$  K, and  $T = 400 - 1000$  K respectively.

function (PSF) at the red end of the wavelength sampled by the modules. In principle, the PSF should vary with wavelength but since the PSF is undersampled at the blue end of the module, it is approximately constant across a given module. So the approximate resolution of the SL module is  $xx$  arcsecs and of the LL module is  $yy$  arcsecs. \*\* DELETE THAT COLUMN IN TABLE 1\*\*

We created continuum subtracted line-only maps of the  $H_2$  S(0) –  $H_2$  S(5) using a combination of PAHFIT () and our own code. PAHFIT is a spectral fitting routine that decomposes IRS low resolution spectra with the main advantage being that it allows one to recover the full line flux of any blended features. Several  $H_2$  lines are blended: the  $H_2$  S(1) with the  $17.0\ \mu\text{m}$  PAH complex,  $H_2$  S(2) with the [Ne II]( $12.8\ \mu\text{m}$ ), and the  $H_2$  S(5) with the [Ar II]( $6.9\ \mu\text{m}$ ). We first concatenated SL1 and SL2, and LL1 and LL2 data cubes into two cubes, one for SL and one for LL. We smoothed each map in the cubes by a  $3 \times 3$  pixel box, conserving the flux, to increase the signal-to-noise ratio of the spectra. Then, for each pixel, we extracted a spectrum and ran PAHFIT to decompose it. We saved the location of the pixel along with the PAHFIT output and reconstructed de-blended features and line maps for all of the mid-infrared features. In this paper we focus only on the  $H_2$  line maps.

## 2.2. Ancillary Data: CO (J=1–0), $H\alpha$ , and X-ray Observations

The BIMA (Berkely Illinois Maryland Array) CO ( $J = 1 - 0$ ) map was acquired as a part of the BIMA Survey of Nearby Galaxies (SONG) (; ). At the distance of M51, the SONG beam ( $5''.8 \times 5''.1$ ) subtends  $220\ \text{pc} \times 190\ \text{pc}$ . The  $H\alpha$  + [N II] image of M51 was observed at Kitt Peak as part of *Spitzer Infrared Nearby Galaxies Survey* (SINGS). The native pixel scale for that image is  $0''.3$ . Additional information about the  $H\alpha$  image can be found in the SINGS Fourth Data Release (DR4). X-ray emission from M51 was observed by the Advanced CCD Imaging Spectrometer (ACIS) on the *Chandra X-Ray Observatory* on

20 June 2000. The total integration time was 14,865 seconds. The resolution of the image is  $\sim 1''$ . Further details of the observations are presented in ().

### 3. Results

#### 3.1. Morphology of H<sub>2</sub> Emission

\*\*\* MUST PUT A CROSS FOR THE NUCLEAR POSITION. ALSO CONVERT ALL JPGS TO \*\*\* EPS FOR SUBMISSION. KEEP THE JPGS THOUGH FOR TALKS / ETC. ALSO RENAME \*\*\* ALL FIGS to FIG1, FIG2 BECAUSE THIS IS THE FORMAT REQUESTED BY APJ.

We have detected and mapped H<sub>2</sub> emission from the six lowest pure rotational H<sub>2</sub> lines (Figure ??). The maps reveal remarkable differences in the distribution of the H<sub>2</sub> emission in M51. H<sub>2</sub> S(0) emission is strongest in the northwest spiral arm peaking at an intensity of  $3.66 \times 10^{-18}$  W/m<sup>2</sup> and decreases by a factor of 2 in the nuclear region. In contrast, the H<sub>2</sub> S(1) emission peaks in the nucleus of the galaxy at an intensity of  $1.03 \times 10^{-17}$  W/m<sup>2</sup> and has an extension of equal intensity towards the northwest spiral arm. In the spiral arm itself, the H<sub>2</sub> S(0) peak is offset from the H<sub>2</sub> S(1) emission by  $10''.2$  ( $\sim 380$  pc). These offsets cannot be explained by the difference in the resolution between the SL and LL map \*\* CHECK THIS BY SMOOTHING THE SL MAP TO LL RESOLUTION AND VERIFYING THAT THE TWO PEAKS ARE STILL DISTINCT. \*\* We find H<sub>2</sub> S(0) and H<sub>2</sub> S(1) emission as far as 5 – 6 kpc from the nucleus of the galaxy. In the outer spiral arm, the H<sub>2</sub> S(0) intensity is a factor of 2 times lower than in the inner northwest spiral arm and the H<sub>2</sub> S(1) intensity is a factor of 5 times lower than in the nucleus.

\*\* I SEE THAT YOU HAVE SEVERAL VERSIONS OF FIGURES AND I WOULD USE THE \*\* ONES WITH THE NUCLEAR CROSS AND LABEL ALL THE H2Sx FIGS

WITH THE \*\* SAME LABELS

The H<sub>2</sub> S(2) – H<sub>2</sub> S(5) maps show different molecular gas morphology within M51 through each H<sub>2</sub> line. The strongest H<sub>2</sub> S(2) emission is from the nucleus at  $2.21 \times 10^{-18}$  W/m<sup>2</sup>. We also see bright H<sub>2</sub> S(2) emission from the inner northwest spiral arm at half the intensity of the nuclear peak. The H<sub>2</sub> S(3) peak at the nucleus is  $1.35 \times 10^{-17}$  W/m<sup>2</sup>, a factor of  $\sim 6$  greater than the H<sub>2</sub> S(2) nuclear peak. There is also a linear bar-like structure in H<sub>2</sub> S(3) emission across the nucleus of the galaxy at a PA $\sim$ -10°. The emission peaks in the H<sub>2</sub> S(2) and H<sub>2</sub> S(3) maps are not spatially coincident. For instance, in the inner spiral arm there is a H<sub>2</sub> S(2) peak coincident with the CO peak whereas the H<sub>2</sub> S(3) peaks further down the spiral arm. Offsets like these suggest that there may be variations in the excitation temperature from region to region within a galaxy, and even within a spiral arm.

The H<sub>2</sub> S(4) and H<sub>2</sub> S(5) lines are brightest at the nucleus with intensities of 3.05 and  $8.04 \times 10^{-18}$  W/m<sup>2</sup> respectively. The H<sub>2</sub> S(4) line shows emission in the nucleus and in the spiral arm to the west. In the spiral arm to the west of the nucleus, the H<sub>2</sub> S(4) intensity is  $2.11 \times 10^{-18}$  W/m<sup>2</sup>. This is notable because the spiral arm to the southwest of the nucleus is very bright in CO and studies have revealed very high molecular gas column densities in the southwest inner spiral arm (). H<sub>2</sub> S(5) emission is asymmetric in the nucleus and mimics the morphology of the H<sub>2</sub> S(3) line with extended emission to the north of the nucleus. The differences in the morphology of H<sub>2</sub> emission are indicative of changes in the H<sub>2</sub> excitation-temperature across the galaxy, which we discuss in the next section.

### 3.1.1. *Note on Uncertainties in the H<sub>2</sub> Line Intensity Maps*

\*\* THIS SUBSECTION IS UNNECESSARY - WE DON'T HAVE A GOOD HANDLE ON \*\* IRS UNCERTAINTIES ANYWAYS AND THERE ARE SYSTEMATICS

THAT ARE NOT \*\* REALLY TAKEN INTO ACCOUNT HERE SO I RECOMMEND  
DELETING IT ENTIRELY \*\*

The uncertainty in the H<sub>2</sub> line intensity varies as a function of line intensity across each map. The uncertainty in the H<sub>2</sub> S(0) line intensity ranges from 20 – 40 % across the nuclear region and spiral arms. The uncertainty in the spiral arm intensity at greater distances from the center of the strip is 80 – 90 %. The H<sub>2</sub> S(1) line map has lower uncertainties with the uncertainty being  $\sim 10$  % in the nuclear region and 20 – 30 % in the spiral arms. The H<sub>2</sub> S(2) and H<sub>2</sub> S(3) lines show similar uncertainties to each other. In the nuclear region their uncertainty is 10 – 30 % and in the spiral arms their uncertainty is 50 – 80 %. The H<sub>2</sub> S(4) and H<sub>2</sub> S(5) lines only show emission from the nuclear region; their uncertainties in the line intensity are  $\sim 35$  % and 25 %, respectively.

### 3.2. Mapping H<sub>2</sub> Excitation-Temperature and Mass across M51

#### 3.2.1. Modeling H<sub>2</sub> Excitation-Temperature and Mass

The pure rotational lines of molecular hydrogen provide a powerful probe of the conditions of the ISM by placing constraints on the energy injection that excites H<sub>2</sub> (). Following ; we modeled the H<sub>2</sub> excitation and mass across M51.

First, we smoothed the H<sub>2</sub> S(1) – H<sub>2</sub> S(5) maps were smoothed to the resolution of the H<sub>2</sub> S(0) map. The maps were then interpolated to the same spatial grid. Excitation diagrams across the strip were derived from the Boltzman equation using the formulation of Rigopoulou et al. (2002),

$$N_i/N = (g(i)/Z(T_{\text{ex}}))\exp(-T_i/T_{\text{ex}}) \quad (1)$$

where  $g(i)$  is the statistical weight of state  $i$ ,  $Z(T_{\text{ex}})$  is the partition function,  $T_i$  is the



energy level of a given state, and  $T_{\text{ex}}$  is the excitation temperature.  $N$  and  $N_i$  are the total column density and the column density of a given state  $i$  and  $N_i$  is determined directly from the measured extinction-corrected flux by

$$N_i = 4\pi \times \text{flux}(i) / (\Omega A(i) h \nu(i)) \quad (2)$$

where  $A(i)$  is the Einstein  $A$ -coefficient,  $\nu(i)$  is the frequency of state  $i$ ,  $\Omega$  is the solid angle of the beam, and  $h$  is Planck’s constant. Table 1 lists the values for the wavelength, rotational state, Einstein  $A$ -coefficient, energy, and statistical weight of the pure rotational levels of  $\text{H}_2$ .

\*\* PUT THE NAME OF THE REGION ON THE PLOTS. ALSO PUT THE NUMBERS OF \*\* THE TRANSITIONS ON ONE OF THEM. \*\* ALSO HOW DO I KNOW FROM THESE PLOTS THAT  $\text{OPR} = 3$ ? YOU NEED TO SAY \*\* HERE THAT IF  $\text{OPR} = 3$  YOU EXPECT THE FOLLOWING IN THE BOLTZMANN \*\* DIAGRAM AND IF NOT THEN YOU EXPECT SUCH AND SUCH...

\*\* OVERALL THIS PARA IS VERY CONFUSING. AT ONE POINT  $\text{OPR}$  IS 3 BUT \*\* SOMEHOW THIS ALSO SHOWS A RANGE OF TEMPERATURES? VERY CONFUSED. \*\* THIS NEXT PARAGRAPH REALLY NEEDS TO BE REWRITTEN.

In Figure ?? we present excitation diagrams from three different regions across the M51 strip. The excitation diagrams exhibit an ortho-to-para ratio ( $\text{OPR}$ ) of 3 in the nuclear region. This is in agreement with the value of the  $\text{OPR}$  of the nuclear region determined by SINGS (). Outside the nucleus, the lower  $J$  ( $\text{H}_2 \text{ S}(0) - \text{H}_2 \text{ S}(3)$ ) levels exhibit an  $\text{OPR}$  of 3. The  $\text{H}_2 \text{ S}(4)$  measurement shows significant scatter in the excitation diagrams outside of the nuclear region of M51. This would indicate that the  $\text{OPR}$  is less than 3; however, due to the low signal-to-noise ratio of the  $\text{H}_2 \text{ S}(4)$  map, we do not believe that the  $\text{OPR}$  determined

from the  $\text{H}_2$  S(4) intensity reflects the OPR of the warm  $\text{H}_2$ . The excitation-diagrams also exhibit a change in slope as a function of rotational level. This indicates that a continuous distribution of  $\text{H}_2$  temperatures is being sampled within the beam.

\*\* REWRITE PARA ABOVE

\*\* DO YOU ASSUME TEMPERATURES FIRST AND DRAW THE LINE OR GET A \*\* TEMPERATURE FROM THE LINES THEMSELVES? THAT WILL DETERMINE HOW THE \*\* NEXT PARA SHOULD START. HERE I AM ASSUMING YOU GET A BEST FIT TO \*\* THE HIGH J LINES AND CALL THAT YOUR TEMPERATURE. THEN SUBTRACT \*\* THAT CONTRIBUTION TO LOW J LINES AND THEN FIT A TEMP TO THOSE? \*\* CORRECT?

The simplest model beyond a single temperature model is a two temperature model, which we call a warm and hot phase from here on. To determine the hot phase temperature we do a least squares fit AT EVERY PIXEL? to the  $\text{H}_2$  S(2) –  $\text{H}_2$  S(5) lines. We subtract the contribution of this hot phase temperature from the lower J lines and do a least squares fit to the  $\text{H}_2$  S(0) –  $\text{H}_2$  S(2) lines. \*\* HOW DO YOU THEN GET THE MASS? \*\*

### 3.2.2. Warm and Hot $\text{H}_2$ Mass Distributions

In Figure ?? we present the warm (*left*) and hot (*right*)  $\text{H}_2$  mass distributions across the M51 strip<sup>2</sup> \*\* NO MAKE THE SIZES THE SAME AND DELETE THIS SENTENCE – Also note that the sizes of the plots of the warm and hot  $\text{H}_2$  mass phases are different. \*\*

\*\* ITS BECOMING PAINFUL TO WRITE EVERYWHERE.. ARE YOU SURE

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<sup>2</sup>Note that the non-rectangular shape of the strip is due to the offset between the SL and LL strips.

YOU \*\* NEED IT? WHY IS  $H_2$  NOT SUFFICIENT - -SAME GOES FOR EVERY TIME  
\*\* YOU USE MATHRM ?? AND PLEASE REDEFINE COMMANDS SO YOU CAN  
USE A \*\* SHORTHAND NOTATION FOR COMMONLY USED LABELS., E.G. \*\* \*\*  
\*\* AT THE BEGINNING OF THE DOCUMENT.

The highest gas surface density for the warm  $H_2$  phase is in the inner northwest spiral arm at 11 \*\* CHANGE ALL  $PC^2 TOPC^{-2} AND REMOVED DIVISIONS$   $SM_{\odot} pc^{-2}$  warm  $H_2$ . The gas surface density in the outer northwest and southeast spiral arms is maximum at the center of the spiral arms at  $3.5 M_{\odot}/pc^2$  and  $1.0 M_{\odot}/pc^2$  respectively. \*\* PEAKS CANNOT BE BIMODAL BY DEFINITION - I REWORDED THAT SENTENCE!  
The hot phase surface density is highest in the nucleus and interior to the inner spiral arm at  $0.24 M_{\odot}/pc^2$ . The gas surface density of the hot phase in the spiral arms is 1/3–1/5th that of the nuclear region.

### 3.2.3. Warm and Hot $H_2$ Excitation-Temperature Distributions

\*\* DO YOU MEAN MASS OR SURFACE DENSITY – PLEASE GET THE NOMENCLATURE \*\* CORRECT BASED ON THE UNITS YOU ARE USING. IF PER  $PC^2$  IT SHOULD BE \*\* SURFACE DENSITY.

\*\* THIS SECTION NEEDS TO BE REWRITTEN. WRITE IT AS FOLLOWS. FIG XX \*\* AND YY SHOW IN GREYSCALE THE DISTRIBUTION OF TEMPERATURES \*\* DETERMINED FROM THE HIGH AND LOW TRANSITIONS AS DESCRIBED ABOVE. \*\* OVERLAID ON IT ARE CONTOURS OF THE  $H_2$  GAS SURFACE DENSITY. IN BOTH \*\* CASES WE FIND THAT THE TEMPERATURE AND GAS SURFACE DENSITY ARE \*\* INVERSELY CORRELATED. – THAT’S ALL YOU YOU NEED TO WRITE. YOU \*\* NEED TO REALLY UNDERSTAND WHY THE

INTERARM REGIONS HAVE SUCH HIGH \*\* TEMPS ACCORDING TO THIS FIGURE. IT REALLY CANNOT BE REAL BECAUSE \*\* OUR S/N IS LOWEST IN THESE REGIONS.

In Figure ?? we compare the warm  $\text{H}_2$  mass distribution to the excitation-temperature. The gas is warmest in the nuclear region of M51 with the temperature ranging from 175 – 190 K in the center of the galaxy. The gas is cooler within the spiral arms of M51 than in the inter-arm regions. The temperature in this spiral arm decreases from 175 K at the outer surface layers to 154 K in the most dense region of the arm indicating that the warm  $\text{H}_2$  excitation-temperature is cooler at higher  $\text{H}_2$  densities. In the outer northwest and southeast spiral arms (at 5 – 6 kpc from the center of the galaxy each) molecular the gas excitation-temperature ranges from 150 K at the center to about 175 K near the edge of the spiral arms; again indicating that high density regions are cooler than low density regions.

The excitation-temperature distribution of the hot  $\text{H}_2$  is shown in Figure ?? and reveals a different picture than the excitation-temperature distribution of the warm  $\text{H}_2$ . The excitation-temperature in the nuclear region ranges from 587 to 612 K. The excitation-temperature in the northwest inner spiral arm is the coolest at 538 K. The hottest excitation-temperatures are observed in the inter-arm regions where there is the least amount of  $\text{H}_2$ .

\*\* I DON'T REALLY BELIEVE THIS –<sub>i</sub> DO YOU? YOU MAY BE MISLED BY THE \*\* LOW S/N IN THESE MAPS IN THE INTERARM REGION ESP. FOR THE HIGHEST J \*\* LINES. In contrast, in the southeast inter-arm region, the excitation temperature is in excess of 900 K.

## 4. Discussion

### 4.1. Comparison to Previous Studies of Warm and Hot H<sub>2</sub> in Galaxies

Previous studies have used aperture-averages over entire galactic nuclei to derive the physical conditions of the molecular gas (; ; ). In M51, find that within the central 330 arcsec<sup>2</sup> (ADD PHYSICAL SCALE IN XX PC<sup>2</sup>), the warm H<sub>2</sub> phase has a maximum temperature of 180 K and a total mass of  $M_{\text{warm}} = 1.46 \times 10^6 M_{\odot}$ . In the central 412 arcsec<sup>2</sup> of M51, we find a corresponding maximum of 186 K and a total mass of  $1.63 \times 10^6 M_{\odot}$ , consistent with the SINGS results. The SINGS team also measured the excitation-temperature of the hot phase (though they do not measure the mass in the hot phase) and find a hot H<sub>2</sub> excitation-temperature of 521 K. Over the same region, we find a maximum excitation-temperature of 584 K.

\*\* BUT YOUR REAL STRENGTH IS THE SPATIAL RESOLUTION AND YOU SHOULD \*\* REALLY REWRITE THE TOP PARAGRAPH TO REFLECT THAT. SO SAY THAT WE \*\* HAVE THE ABILITY TO DO SPATIALLY RESOLVED STUDIES BUT IT WOULD BE \*\* GOOD TO COMPARE OUR NUMBERS TO THE AVERAGES FOUND BY PREVIOUS \*\* STUDIES. THEN SAY THAT THE SPATIALLY RESOLVED NUMBERS MATCH THE \*\* AVERAGED STUDIES SO MUCH SO THAT THE AVERAGED STUDIES MUST BE \*\* DOMINATED BY THE BRIGHTEST REGIONS.. OR IT COULD BE THAT YOU ARE \*\* AVERAGING THE NUMBERS YOU QUOTE ABOVE – IS THAT THE CASE? I AM \*\* CONFUSED A BIT HERE.

\*\* I DON'T SEE THE POINT OF THE PARA YOU HAD HERE BELOW. I WOULD DELETE IT. In a survey of Seyfert galaxies, Rigopoulou et al. (2002) find that the warm-to-hot phase mass fraction varies between  $5 \times 10^3 - 1 \times 10^6$ . This is significantly higher than the warm-to-hot mass fraction (Figure ??) than we measure within the central

412 arcsec<sup>2</sup> of M51, 14.8. To measure the hot H<sub>2</sub> mass, we use the H<sub>2</sub> S(2) – H<sub>2</sub> S(5) lines, whereas Rigopoulou et al. (2002) used the H<sub>2</sub> S(5) – H<sub>2</sub> S(7) lines. Thus, the hot H<sub>2</sub> phase that they measured is preferentially offset to higher H<sub>2</sub> excitation-temperatures and lower H<sub>2</sub> masses.

#### 4.2. The Warm-to-Hot H<sub>2</sub> Mass Ratio

\*\* NOW WE GO BACK TO DESCRIBING WARM VS. HOT? I THOUGHT WE WERE IN \*\* THE DISCUSSION SECRION HERE. THIS IS REPETITIOUS. DIDN'T YOU \*\* ALREADY DO THIS BEFORE? JUST REDUCE THE FOLLOWING PARA TO ONE \*\* LINE. SAY THAT THE WARM TO HOT RATIO VARIES (PERHAPS YOU CAN SHOW \*\* THE DIVISION OF THE TWO MAPS RATHER THAN OVERLAYING THEM WHICH DOES \*\* NOT CONVEY QUANTITATIVE INFO. YOU JUST NEED TO SAY THIS IS THE \*\* RATIO SEE...

In Figure ??, we compare the warm H<sub>2</sub> mass distribution to the hot H<sub>2</sub> mass distribution. The warm H<sub>2</sub> mass distribution peaks at in the northwest inner spiral arm and the hot H<sub>2</sub> mass distribution peaks at in the nucleus and in the region interior to the northwest spiral arm. It is evident from the figure that the warm-to-hot H<sub>2</sub> mass ratio is not constant across the galaxy but is lowest in the nucleus of the galaxy (at 12) and increases to 170 and 136 in the southeast and northwest spiral arms, respectively. The differences between the warm and hot H<sub>2</sub> mass distributions and the variations in the morphology of M51 as a function of rotational energy level suggest different origins to the warm and hot H<sub>2</sub> phases and that the primary excitation mechanisms of the warm and hot H<sub>2</sub> phases differ.

\*\* THEN IN THIS PARA JUST SAY - THE MOST OBVIOUS HEATING

MECHANISM FOR \*\* THE MOLECULAR GAS IS STAR FORMATION OR NUCLEAR ACTIVITY. THE FORMER \*\* CAN BE TRACED USING HALPHA AND HTE LATTER USING XRAYs AND OIV LINE. \*\* IT WOULD ALSO BE INTERESTING TO SEE HOW THE WARM GAS WAS \*\* DISTRIBUTED RELATIVE TO THE COLD GAS - SO WE ALSO COMPARE OUR DATA \*\* TO CO DATA FROM SONG.

In order to understand the  $H_2$  excitation and where each excitation mechanism is dominant, we made comparisons of the  $H_2$  line intensity and mass distributions to diagnostics of three excitation mechanisms: UV photons emitted in dense photon dominated regions (PDRs) and H II regions, shocks, and X-rays. We use CO ( $J = 1 - 0$ ) emission to determine the location of the  $H_2$  relative to locations dense PDRs,  $H\alpha$  imagery to identify H II regions (), [O IV](25.98  $\mu m$ ) line emission as a diagnostic of shocks (), and the 0.5 – 10.0 keV X-ray band to distinguish X-ray dominated regions (XDRs). The following section compares the  $H_2$  mass distributions to CO ( $J = 1 - 0$ ),  $H\alpha$ , [O IV], and X-ray emission.

#### 4.3. Distinguishing the $H_2$ Excitation Mechanisms

\*\* AS A GUIDE TO DECREASING THE VERBOSITY OF THE ARTICLE, TRY AND \*\* REDUCE THE NEXT THREE SUBSECTIONS TO THREE PARAGRAPHS OF NO MORE \*\* THAN 10 SENTENCES (AND NO CHEATING WITH RUN-ON SENTENCES). THIS \*\* WILL MAKE THIS PAPER MUCH MORE COMPACT AND TRACTABLE. I WILL TAKE \*\* OUT THE SALIENT SENTENCES FROM THE FIRST SUBSECTION HERE AND PUT \*\* THEM INSIDE \*\*\* - THESE OUGHT TO BE THE 10 SENTENCES OR LESS YOU \*\* FOCUS ON.

#### 4.3.1. $H_2$ Excitation by UV Photons from PDRs: Comparison of $H_2$ to CO Emission

In star-forming regions,  $H_2$  exists within the PDR and deep into the molecular cloud. \*\* Owing to a low dissociation energy of 4.5 eV,  $H_2$  formation generally does not occur within a PDR until the radiation field becomes sufficiently weak. \*\* Understanding  $H_2$  in PDRs and implementing  $H_2$  into photoionization codes has become the emphasis of many theoretical models. Recent advances in the CLOUDY photoionization code have included  $H_2$  and modeled the structure of star-forming regions while treating the H II region and PDR as one continuous cloud (; ). Kaufman et al. (2006) used Starburst99/Mappings to model  $H_2$  pure rotational line emission from PDRs to probe the conditions of dense PDRs in star-forming regions. \*\* They show that within galaxies, where the telescope beam size is generally kiloparsecs across,  $H_2$  emission could serve to probe the surface layers of dense molecular clouds. \*\*

\*\* HERE ARE YOU SAYING THAT THERE IS NO  $H_2$  WHERE THERE IS CO?? THAT IS NOT RIGHT. PERHAPS YOU MEAN TO SAY CO TRACES THE COLD GAS WHERE NO MIDIR  $H_2$  EMISSION IS SEEN – BUT THERE IS  $H_2$  THERE? RIGHT? \*\* CO is found deep within a molecular cloud where the temperatures are too cold to excite  $H_2$  emission. In these regions, CO is collisionally excited by the more abundant  $H_2$  molecule. CO is connected to  $H_2$  in star-forming regions because at the surface layers of the molecular clouds,  $H_2$  is excited and CO emits dipole rotational lines due to heating by the ionizing radiation of newborn massive stars ().

In Figure ??, we compare the warm (*left*) and hot  $H_2$  (*right*) mass distributions to CO ( $J = 1 - 0$ ) emission. The brightest CO emission is seen in the spiral arms where the greatest amounts of warm  $H_2$  mass are also found. \*\* The most striking result is that in the inner spiral arms, we see that the CO is offset toward the nucleus from the warm  $H_2$  mass. The offset between the peaks in CO and warm  $H_2$  mass is  $7''.2$  in the northwest spiral



arms and  $5''$  in the southeast spiral arms. We believe that these offsets are real with one possible explanation being that the  $H_2$  is tracing the regions of active star-formation within the giant molecular associations. \*\* REDUCE ABOVE SENTENCES INSIDE \*\* I IN THIS PARA TO 2 SENTENCES

We compare CO to the hot  $H_2$  mass distribution (Figure ??, *right*) and we see that the hot  $H_2$  mass is most abundant in the nuclear region, interior to the CO bright spiral arms. In northwest spiral arm, the hot  $H_2$  mass is offset from the CO by  $2''.6$ . We believe that the offsets between the CO and the hot  $H_2$  mass are real, however, they are likely due to the dominance of other excitation mechanisms (shocks or X-rays) in the nuclear region of M51.

\*\*\* THIS PARA REALLY CONFUSED ME. WHAT IS BELIEVABLE AND WHAT IS NOT. \*\*\* WHAT IS THE CO DEPTH IS NOT SUFFICIENT. LET'S TALK ABOUT THIS \*\*\* OVER THE PHONE OR ICHAT

In Figure ??, we compare the CO intensity to the  $H_2$  S(0) –  $H_2$  S(3) line intensity maps. We see that the  $H_2$  S(0) contours trace the CO spiral arms with one notable difference being that we detect  $H_2$  S(0) emission far ( $5 - 6$  kpc) from the nucleus of M51, where there is no CO detected. The  $H_2$  S(1) –  $H_2$  S(3) contours also trace the CO spiral arms; however, within the spiral arms, the  $H_2$  is not necessarily aligned with the CO. For example, in the comparison of  $H_2$  S(3) to CO, we see that in the inner spiral arms, the  $H_2$  contours trace the CO; however, within the northwest arm, we see that there is strong  $H_2$  emission offset to the west by  $9''.3$  from the bright CO emission in the same spiral arm. The offsets between the peaks in  $H_2$  and CO emission are not systematic in any direction.

\*\* The CO intensity in the nucleus of the galaxy is much fainter than in the bright CO spiral arms. \*\* JUST SAYS CO - $\dot{L}$  COLD GAS AND NO COLD GAS IN NUCLEUS. \*\* This is interesting because  $H_2$  S(1) –  $H_2$  S(3) emission is brightest in the nucleus of M51 where  $H_2$  emission from the higher  $J$  lines is likely due to more energetic processes such as

shock excitation and X-rays (which we discuss in §4.3.3 and §4.3.4).

*4.3.2. H<sub>2</sub> Excitation by UV Photons from H II Regions: Comparison of H<sub>2</sub> to H $\alpha$  Emission*

\*\* SENTENCES LIKE THE FIRST FOUR BELOW ARE UNNECESSARY IN A PAPER! \*\* OKAY IN A THESIS BUT STILL MAKES A PAPER WAAAY TO VERBOSE. JUST \*\* STICK TO THE RELEVANT, IMPORTANT STUFF THAT IS ADDING TO THE \*\* SCIENCE. ONLY GIVE EXTRA INFO IF ABSOLUTELY NEEDED TO SET UP AN \*\* ARGUMENT.

\*\* AGAIN REDUCE THIS SUBSECTION TO 10 SENTENCES AND 1 PARA ONLY. \*\* BOTTOM LINE IS YOU ARE SAYING WARM H<sub>2</sub> NOT CORRELATED WITH HALPHA \*\* RIGHT? WHAT IS THE HALPHA WAS EXTINGUED? WHAT ABOUT PASCHEN ALPHA \*\* - IS THAT A BETTER CORRELATION?

H II regions are sites of recent massive star formation. They illuminate bright nebulae in distant galaxies and outline the spiral arms. H II regions emit prodigious amounts of UV radiation at energies  $\geq 13.6$  eV capable of exciting and ionizing H<sub>2</sub>. H $\alpha$  imagery is often used to identify and map H II regions in star-forming galaxies. Scoville et al. (2001) used H $\alpha$  and Pa $\alpha$  imagery to identify and characterize over 1,350 H II regions in M51. \*\*

In Figure ??, we compare the warm (*left*) and hot (*right*) H<sub>2</sub> mass distributions to H $\alpha$  emission. In general, the warm and hot H<sub>2</sub> concentrations are not cospatial with the H $\alpha$  emission regions in the spiral arms with the one exception being that the warm H<sub>2</sub> mass in the inner spiral arms appears to trace the H $\alpha$  emission. The warm H<sub>2</sub> mass contours show that local peaks in H<sub>2</sub> mass are found within the dust lanes. An example of this is in the northwest spiral arms where we see the H<sub>2</sub> mass offset from the H $\alpha$  spiral arms with local

peaks being found in the dust lanes.

In Figure ??, we compare the  $\text{H}_2$  S(0) –  $\text{H}_2$  S(3) line intensity maps to  $\text{H}\alpha$  emission. Comparison of the  $\text{H}_2$  S(0) map to  $\text{H}\alpha$  reveals that the strongest  $\text{H}_2$  emission in the northwest and southeast inner spiral arms is coincident with  $\text{H}\alpha$  emission; however, the other  $\text{H}_2$  S(0) spiral arms show the strongest emission in the dust lanes, offset from the  $\text{H}\alpha$  spiral arms. The largest offsets are seen in the southwest spiral arm where the  $\text{H}_2$  S(0) emission is offset from the  $\text{H}\alpha$  spiral arm by  $\sim 15''$  (560 pc).  $\text{H}_2$  S(1) emission appears to follow the dust lanes and the  $\text{H}_2$  S(1) intensity subsides into the  $\text{H}\alpha$  spiral arms.  $\text{H}_2$  S(2) and  $\text{H}_2$  S(3) emission is also found in the dust lanes; however, there are instances (such as in the southeast spiral arm) where the  $\text{H}_2$  emission appear to be found straddling the dust lane and  $\text{H}\alpha$  spiral arm.

These results are in contradiction to a previous study of rovibrational  $\text{H}_2$  line emission in Seyfert galaxies by Quillen et al. (1999) who used *HST* NICMOS to map the  $\text{H}_2(1-0)\text{S}(1)$  ( $2.121\ \mu\text{m}$ ) and  $\text{H}_2(1-0)\text{S}(3)$  ( $1.957\ \mu\text{m}$ ) lines in 10 Seyfert galaxies. They found that the  $\text{H}_2$  emission is generally coincident with  $\text{H}\alpha$  emission and near the dust lanes (though offset from them). One explanation for this could be that the rovibrational lines trace a much hotter (1000 – 5000 K) molecular gas which can be found associated with H II regions whereas the cooler  $\text{H}_2$  traced by the lowest pure rotational transitions is found in dense PDRs and molecular clouds associated with star forming regions.

#### 4.3.3. $\text{H}_2$ Excitation by Shocks: Comparison of $\text{H}_2$ to $[\text{O IV}](25.89\ \mu\text{m})$ Emission

\*\* THIS SUBSECTION IS ACTUALLY WELL WRITTEN AND GETS TO THE POINT \*\* IMMEDIATELY - QUITE IN CONTRAST TO THE PREVIOUS TWO. SO TRY TO \*\* FOLLOW THIS MODEL. BUT IT IS STILL TOO LONG AND VERBOSE

SO CUT IT \*\* DOWN TO ONE PARA IF POSSIBLE.

The [O IV](25.89  $\mu\text{m}$ ) line can be excited in shocks (), the stellar winds of massive Wolf-Rayet stars (), or by an active galactic nucleus (AGN)(). Though the [O IV] line is blended with the [Fe II](25.99  $\mu\text{m}$ ) line in *Spitzer* IRS low resolution spectra, PAHFIT can deblend the two lines and in mapping the H<sub>2</sub> S(0) and H<sub>2</sub> S(1) line in the LL data cubes, we also mapped the [O IV] line.

In Figure ??, we compare the [O IV] intensity to the warm (*left*) and hot (*right*) H<sub>2</sub> distributions. The [O IV] emission is brightest in the nuclear region at  $8.75 \times 10^{-18} \text{ W/m}^2$  and the peak is coincident with the nuclear peak in the mass of the hot H<sub>2</sub>. [O IV] emission subsides from the nucleus to the inner spiral arm by 50 %. We resolve weaker [O IV] emission within the warm and hot H<sub>2</sub> spiral arms. The [O IV] intensity in the spiral arms is a factor of  $\sim 6$  lower in the spiral arms than the peak intensity found in the nucleus.

The [O IV] emission within the nuclear region of M51 is likely due to the weak Seyfert II nucleus () and is possibly associated with shocked gas from the outflows of the AGN. The peak of the [O IV] emission coincides with the nuclear peak in hot H<sub>2</sub> mass indicating that the hot H<sub>2</sub> phase in the nuclear region of the galaxy is AGN or shock heated. With  $2.43 \times 10^5 M_{\odot}$  of hot H<sub>2</sub> in the central 0.58 kpc<sup>2</sup>, it is unlikely that the hot H<sub>2</sub> is fueling the central AGN, but is excited by the AGN or shocks produced by it. In the nuclear region we observe a factor of 12 times greater warm H<sub>2</sub> mass. The warm H<sub>2</sub> mass is much greater within the spiral arms than within the nucleus and the warm-to-hot mass ratio is lowest in the nuclear region where the [O IV] intensity is greatest. In the nuclear region, shocks appear to be a more efficient means to excite the hot H<sub>2</sub> phase than the warm H<sub>2</sub> phase.

#### 4.3.4. $H_2$ Excitation by X-rays: Comparison of $H_2$ to X-ray Emission

M51 has been extensively studied in X-rays by ASCA (), Newton XMM (), and the Chandra X-ray Observatory (). These observations of M51 have revealed more than 80 X-ray sources. Candidates for these X-ray sources include neutron stars, black holes, supernova remnants (SN 1994I), and a low-luminosity AGN (; ). X-ray emission from the nuclear region of M51 has been studied by Terashima and Wilson (2001). They observe X-ray emission from the nucleus, the extranuclear cloud (XNC, to the south of the nucleus), and the northern loop. A radio jet has been observed connecting the nucleus of M51 to the XNC in 6 cm imagery (). The jet emanates from the south of the elongated nucleus and is shock heating ISM.

In Figure ??, we compare the smoothed 0.5 - 10 keV band X-ray image to the warm (*left*) and hot (*right*)  $H_2$  mass distributions. The 0.5 - 10 keV band has been smoothed to the resolution of the warm and hot  $H_2$  mass distributions and the nucleus, XNC, and northern loop are indistinguishable in the smoothed image. X-ray emission is most intense from the nucleus and decreases into the northwest spiral arm that contains the greatest  $H_2$  mass. There appears to be very little connection between the 0.5 - 10 keV X-ray band and the warm  $H_2$  mass distribution.

The peak in X-ray emission is coincident with the hot  $H_2$  mass peak. The most intense 0.5 - 10 keV X-ray emission originates from the nucleus and is oriented north-to-south, similar to the [O IV](25.89  $\mu$ m) emission. The peak in X-ray emission is located within the peak in hot  $H_2$  mass suggesting that X-rays play an important role in exciting the hot  $H_2$  phase. While there is a correlation between X-ray emission and the hot  $H_2$  phase,  $H_2$  excitation by X-rays cannot be distinguished from  $H_2$  excitation by shocks.

We further investigate X-ray excited  $H_2$  emission in M51 by comparing the 0.5 – 10.0 keV X-ray band to the  $H_2$  S(2) –  $H_2$  S(5) line intensity maps (Figure ??). The  $H_2$  S(2) –

H<sub>2</sub> S(5) line intensities all peak at the X-ray source within the nucleus. The X-ray source fits within the H<sub>2</sub> contours because the X-ray image resolution is slightly higher than the resolution of each of the H<sub>2</sub> maps.

The morphology of the nuclear H<sub>2</sub> emission appears to be correlated with the X-ray source. The H<sub>2</sub> S(2) intensity peaks around the X-ray nucleus and the intensity decreases to the south by 70 % in the southern XNC. The bar structure seen in the H<sub>2</sub> S(3) emission is aligned with X-ray emission from the nucleus, XNC, and northern loop. The H<sub>2</sub> S(3) intensity peaks between the X-ray nucleus and the XNC. To the north of the nucleus, the contours follow the X-ray loop. To the south of the nucleus (into the XNC), H<sub>2</sub> S(3) intensity decreases by 80 %. The H<sub>2</sub> S(4) and H<sub>2</sub> S(5) line intensities are highest in the nucleus which coincides with strong X-ray emission from the nucleus and XNC.

\*\* SAME COMMENT ON SECTION ABOVE.. TOO VERBOSE. 1 PARAGRAPH PLEASE.

## 5. Conclusions

\*\* WE CAN WORK ON THIS LATER. AGAIN NEED TO MAKE IT CONCISE.

We have spectrally mapped a strip across M51 using the *Spitzer* IRS low resolution modules. We used the spatially resolved spectra to map H<sub>2</sub> S(0) – H<sub>2</sub> S(5) line intensities across the strip. We find:

1. The morphology of H<sub>2</sub> emission in M51 varies with H<sub>2</sub> rotational level. H<sub>2</sub> S(0) emission is strongest in the spiral arms of the galaxy while the higher  $J$  transitions show the strongest emission towards the nucleus. The H<sub>2</sub> S(1) intensity is strongest in the nuclear region and in the northwest spiral arms, however, the peak in H<sub>2</sub> S(0) intensity in the

northwest spiral arm is offset from the peak in  $\text{H}_2$  S(0) intensity by  $10''.2$ . The  $\text{H}_2$  S(2) and  $\text{H}_2$  S(3) maps show  $\text{H}_2$  emission in the nucleus, spiral arms, and inter-arm regions of M51 and bar structure aligned north-to-south is apparent in  $\text{H}_2$  S(3) emission.  $\text{H}_2$  S(4) and  $\text{H}_2$  S(5) emission is resolved in the nuclear region of M51.

2. The different morphologies of  $\text{H}_2$  emission in M51 indicate significant spatial variations in  $\text{H}_2$  excitation-temperature and mass. Excitation diagrams reveal that the  $\text{H}_2$  exists in a continuous distribution of temperatures across the galaxy. Using the low  $J$  lines to trace the warm ( $T = 100 - 300$  K)  $\text{H}_2$ , we find that the warm  $\text{H}_2$  excitation-temperature is highest in the nuclear region at 192 K and the warm  $\text{H}_2$  mass peaks in the northwest inner spiral arm at a mass density of  $11 \text{ M}_\odot/\text{pc}^2$ . Using the higher  $J$  lines to trace the hot ( $T = 400 - 1000$  K)  $\text{H}_2$ , we find that the hot  $\text{H}_2$  excitation-temperature is lowest in the inner spiral arms ( $500 - 550$  K) and increases to  $\sim 600$  K in the nucleus, where the largest hot  $\text{H}_2$  mass densities are found to be  $0.24 \text{ M}_\odot/\text{pc}^2$ .

3. The warm and the hot  $\text{H}_2$  mass distributions are not cospatial and the warm-to-hot mass ratio varies across M51. The hot mass distribution shows two peaks, one in the nucleus of M51 and one located interior to the northwest inner spiral arm of M51. The warm mass distribution peaks in the northwest spiral arm and is offset from the hot mass peak by  $11''.4$ . The warm-to-hot mass ratio varies across the galaxy with the ratio being  $\sim 15$  in the nucleus and increasing to  $> 100$  in the spiral arms. Variations in the warm-to-hot  $\text{H}_2$  mass ratio and differences in the morphology of the  $\text{H}_2$  emission across M51 indicate that the primary excitation mechanism differs for the warm and hot  $\text{H}_2$  mass phases as a function of location within the galaxy.

4. CO emission is offset from the warm H<sub>2</sub> mass in the inner spiral arms of M51.

These apparent offsets are real and are possibly associated with the regions of active star formation within the molecular clouds. In the spiral arms, the H<sub>2</sub> S(0) – H<sub>2</sub> S(3) contours trace the CO; however, within the spiral arms, the peaks in H<sub>2</sub> can be offset from the peaks in CO intensity. In the nucleus, the H<sub>2</sub> S(1) – H<sub>2</sub> S(3) lines are brightest and the CO intensity is a factor of  $\sim 2.5$  weaker than in the spiral arms suggesting that H<sub>2</sub> emission from the higher  $J$  lines is excited by shocks of X-rays.

5. Comparing the distributions of H<sub>2</sub> to H $\alpha$  reveals that the warm and hot H<sub>2</sub> mass is found in the dust lanes rather than at or around the H $\alpha$  emission regions with the one exception being that the warm H<sub>2</sub> mass in the inner spiral arms is coincident with bright H $\alpha$  emission. This is in contradiction to previous studies of spatially resolved rovibrational H<sub>2</sub> line emission in Seyferts that found the H<sub>2</sub> emission to be coincident with the H $\alpha$  emission.

6. The peaks in [O IV](25.89  $\mu\text{m}$ ) intensity and X-ray intensity are both coincident with the peak in hot H<sub>2</sub> mass in the nucleus of M51 suggesting that the hot H<sub>2</sub> in the nucleus is primarily excited by the AGN, shocks (possibly associated with the AGN), or X-rays associated with the AGN. The spatial distributions of the [O IV] emission and X-ray surface brightness are very similar, but a primary excitation mechanism (shocks or X-rays) of the hot H<sub>2</sub> mass phase cannot be distinguished. Further comparison of the H<sub>2</sub> S(2) – H<sub>2</sub> S(5) intensity maps to the X-ray surface brightness reveal that the nuclear H<sub>2</sub> emission is associated with X-rays and the bar structure apparent in the H<sub>2</sub> S(3) map is aligned with



the nucleus, XNC, and northern loop.

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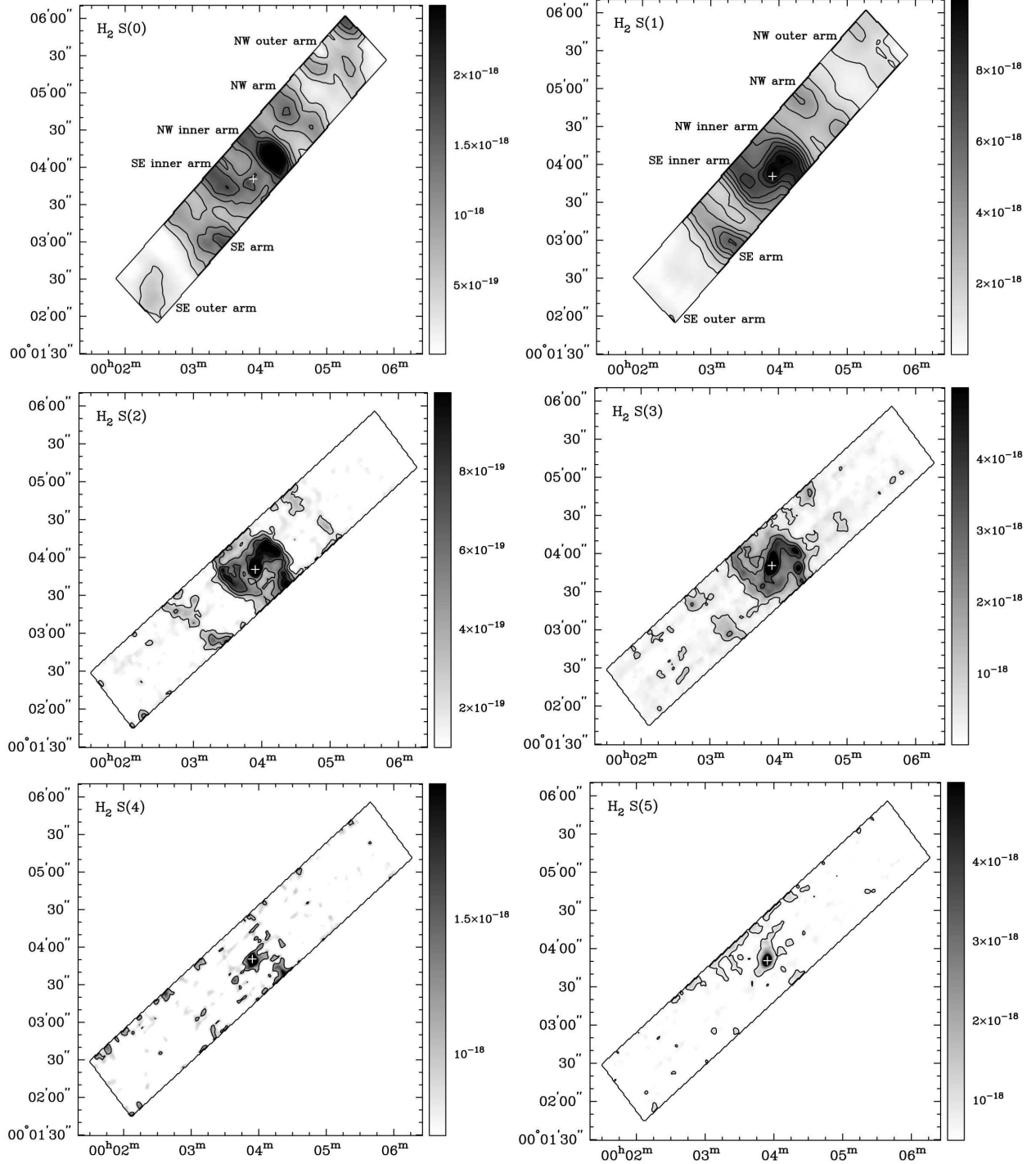


Fig. 1.— Maps of the H<sub>2</sub> S(0) (*top left*), H<sub>2</sub> S(1) (*top right*), H<sub>2</sub> S(2) (*middle left*), H<sub>2</sub> S(3) (*middle right*), H<sub>2</sub> S(4) (*bottom left*), and H<sub>2</sub> S(5) (*bottom right*) intensity across the SL and LL strips that we mapped with the Spitzer IRS. The H<sub>2</sub> S(0) and H<sub>2</sub> S(1) maps are created from the LL data cubes. The H<sub>2</sub> S(2), H<sub>2</sub> S(3), H<sub>2</sub> S(4), and H<sub>2</sub> S(5) maps are created from the SL data cube. The grey-scale is in units of W/m<sup>2</sup>. Contour levels are at 2.9

× 10<sup>-18</sup>, 2.2 × 10<sup>-18</sup>, 1.8 × 10<sup>-18</sup>, 1.5 × 10<sup>-18</sup>, 1.1 × 10<sup>-18</sup>, 7.3 × 10<sup>-19</sup>, and 3.7 × 10<sup>-19</sup>

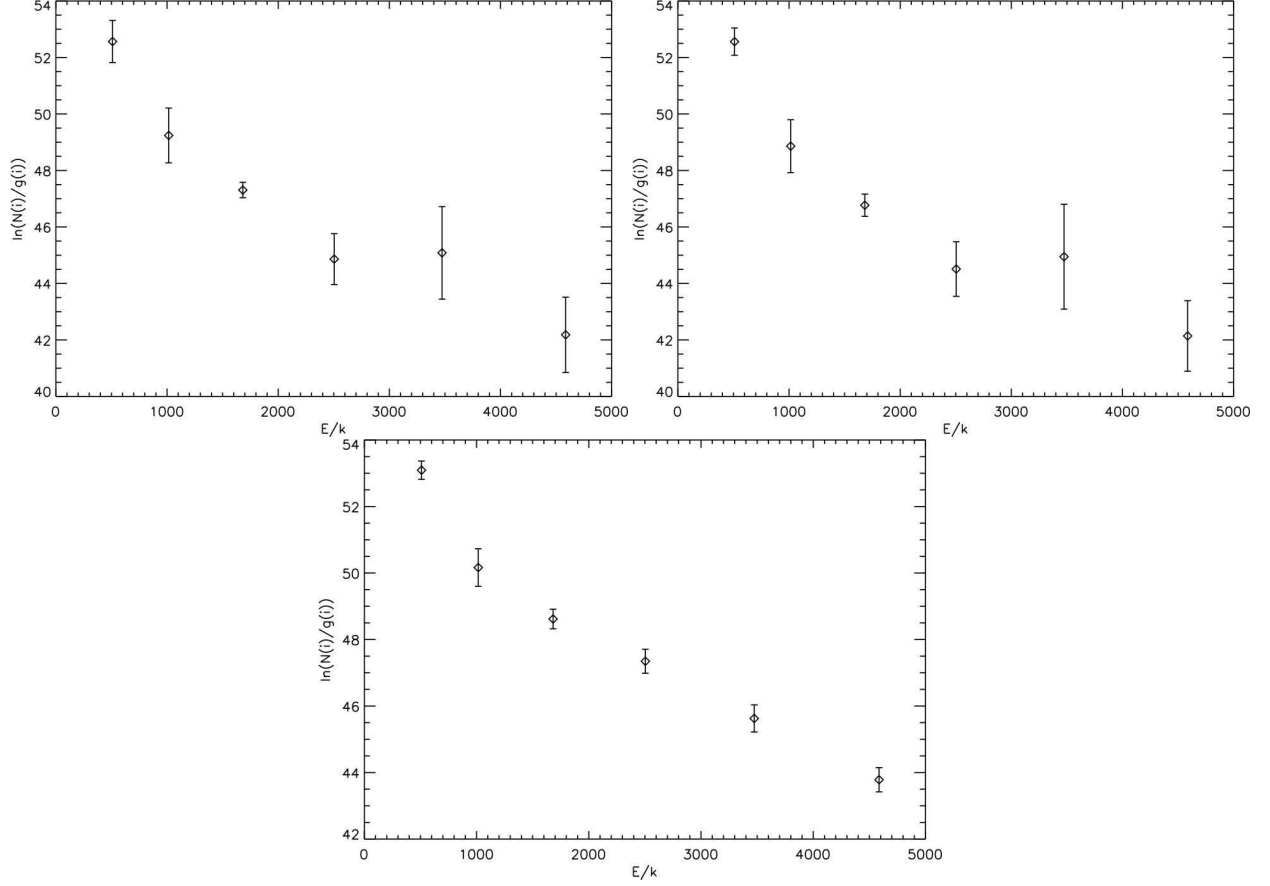


Fig. 2.— Excitation diagrams taken from 3 different regions along the M51 strip. The top two excitation diagrams are taken from regions within the southeast and northwest spiral arms that are  $10''.2$  in diameter ( $1.13 \times 10^5 \text{ pc}^2$ ) and centered at (RA, Dec) of (202.45, 47.21) and (202.49, 47.18), respectively. The excitation diagram at the bottom is taken from from the nuclear region. The aperture is  $10''.2$  in diameter ( $1.13 \times 10^5 \text{ pc}^2$ ), centered at (RA, Dec) of (202.47, 47.19).

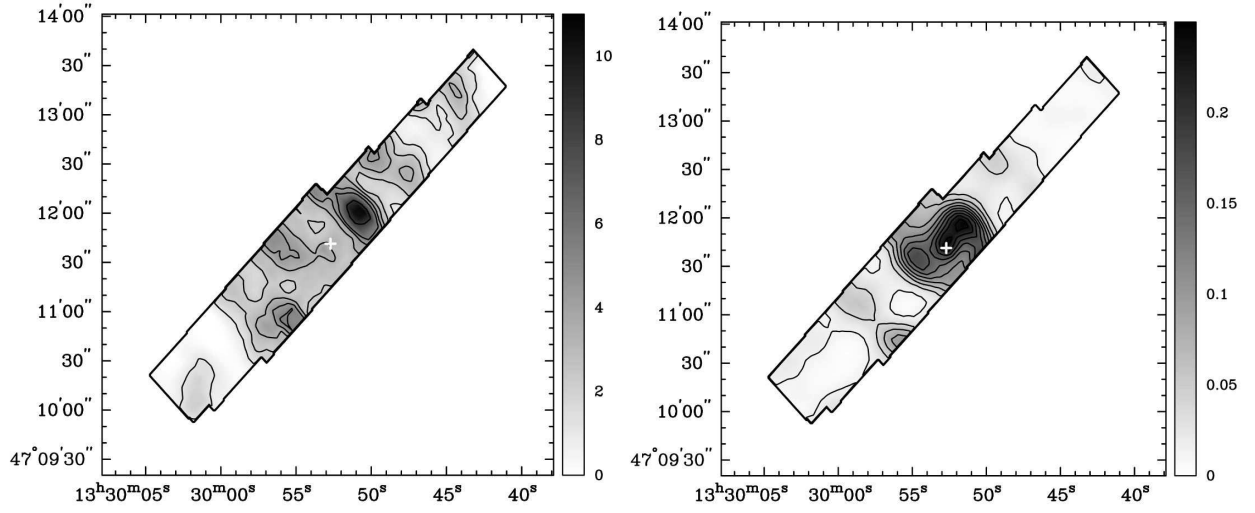


Fig. 3.— Shown are the warm ( $T = 100 - 300$  K)  $H_2$  (*left*) and hot ( $T = 400 - 1000$  K)  $H_2$  (*right*) mass distributions. The mass distributions are in units of  $M_{\odot}/pc^2$ . Contours are overplotted for clarity. The warm  $H_2$  mass contour levels are at 8.85, 5.55, 4.43, 3.32, 2.21, and 1.10  $M_{\odot}/pc^2$ . The hot  $H_2$  contour levels are at 10 % of 0.25  $M_{\odot}/pc^2$ . The hot  $H_2$  mass distribution is derived from the fit to the  $H_2$  S(2) –  $H_2$  S(5) lines and the warm  $H_2$  mass distribution is derived from the fit to the  $H_2$  S(0) –  $H_2$  S(2) lines, corrected for the contribution of the hot  $H_2$  mass phase.

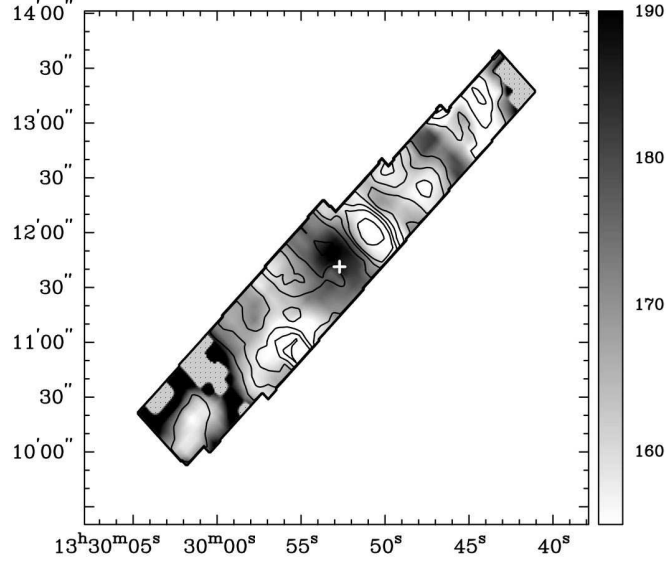


Fig. 4.— The warm ( $T = 100 - 300$  K)  $H_2$  mass distribution compared to the warm  $H_2$  excitation-temperature. The warm  $H_2$  excitation-temperature and mass distributions are derived from the fit to the excitation diagrams across the strip for the  $H_2$  S(0) –  $H_2$  S(2) lines, corrected for the contribution of the hot ( $T = 400 - 1000$  K)  $H_2$  phase. Mass density contour levels are at 8.85, 5.55, 4.43, 3.32, 2.21, and 1.10  $M_\odot/\text{pc}^2$  (same as in Figure ??). The grey-scale represents the excitation-temperature distribution (in units of Kelvin). The non-rectangular shape to the map is due to the slight offset of the *Spitzer* IRS SL strip relative to the LL strip.



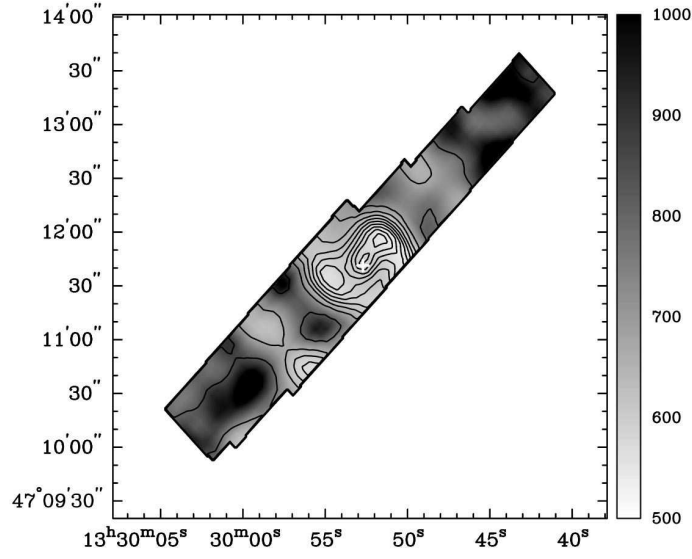


Fig. 5.— The hot ( $T = 400 - 1000$  K)  $H_2$  mass distribution compared to the hot  $H_2$  excitation-temperature. The hot  $H_2$  excitation-temperature and mass distributions are derived from the fit to the excitation diagrams across the strip for the  $H_2$  S(2) –  $H_2$  S(5) lines. Mass density contour levels are at 10% of  $0.25 M_\odot/\text{pc}^2$  (same as in Figure ??). The grey-scale represents the excitation-temperature distribution (in units of Kelvin).

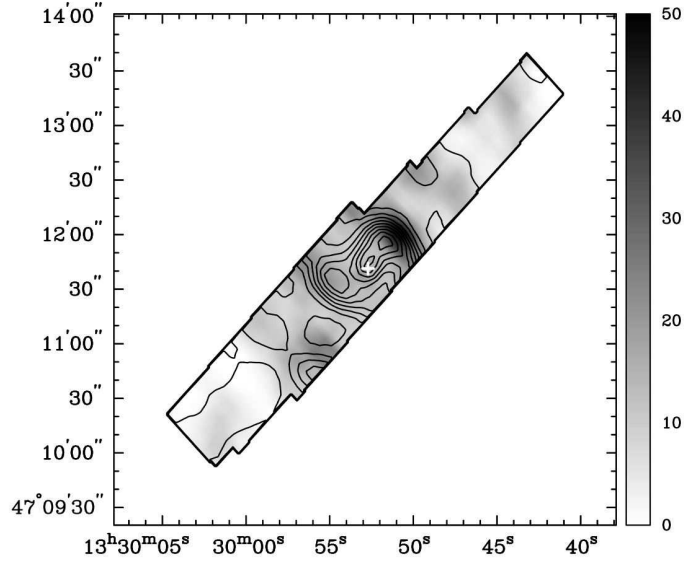


Fig. 6.— The warm ( $T = 100 - 300$  K)  $H_2$  mass (in grey-scale) compared to the hot ( $T = 400 - 1000$  K)  $H_2$  mass (in contours). Contours levels for the hot  $H_2$  mass distribution are at 10% of the maximum mass density  $0.25 M_\odot/\text{pc}^2$ . The grey-scale is in units of  $M_\odot/\text{pc}^2$ .

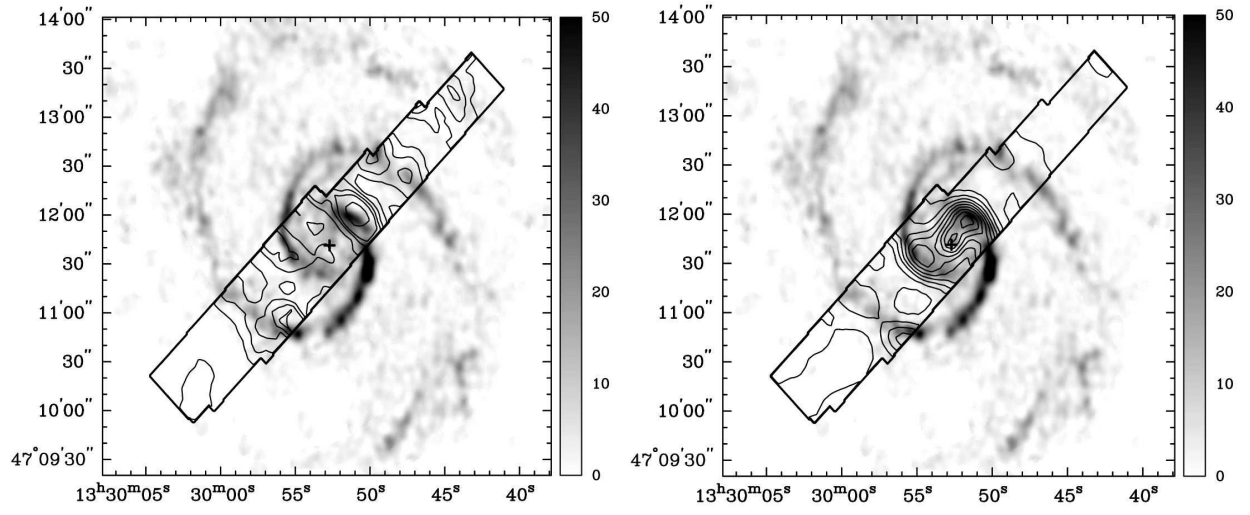


Fig. 7.— *Left*: Comparison of CO intensity (in grey-scale) to the warm ( $T = 100 - 300$  K)  $H_2$  mass (in contours). The CO intensity is in units of  $\text{Jy beam s}^{-1}$ . The warm  $H_2$  mass contours are the same as in Figures ?? and ?. *Right*: Comparison of CO intensity (in grey-scale) to the hot ( $T = 400 - 1000$  K)  $H_2$  mass (in contours). The CO intensity is in units of  $\text{Jy beam s}^{-1}$ . The hot  $H_2$  mass contours are the same as in Figures ?? and ?.

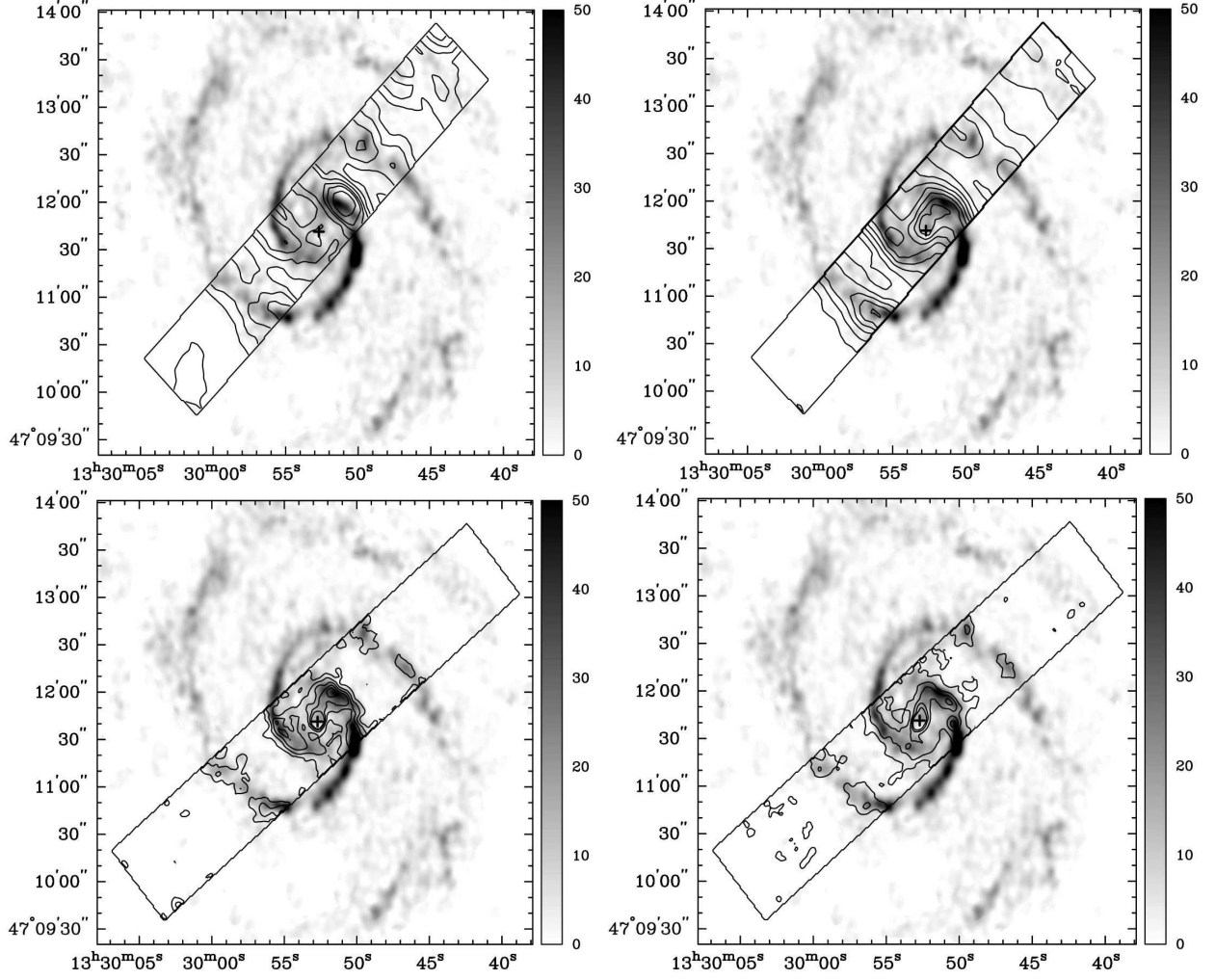


Fig. 8.— Comparison of the CO emission to the  $\text{H}_2$  S(0) (*top left*),  $\text{H}_2$  S(1) (*top right*),  $\text{H}_2$  S(2) (*bottom left*), and  $\text{H}_2$  S(3) (*bottom right*) emission. The CO emission maps are in units of  $\text{Jy beam s}^{-1}$ . Contour levels for  $\text{H}_2$  S(0),  $\text{H}_2$  S(1),  $\text{H}_2$  S(2), and  $\text{H}_2$  S(3) are the same as in Figure ??.

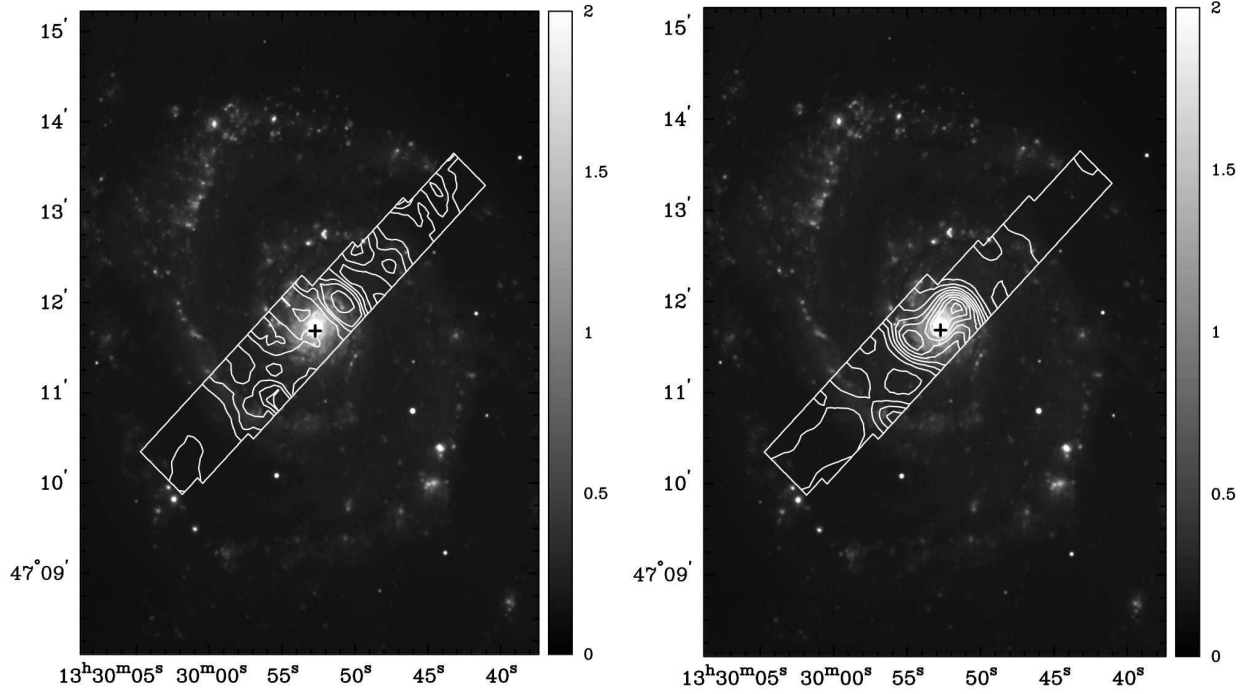


Fig. 9.— *Left*: Comparison of  $H\alpha$  (in grey-scale) to the warm ( $T = 100 - 300$  K)  $H_2$  mass (in contours). The  $H\alpha$  image is in units of counts/sec. The warm  $H_2$  mass contours are the same as in Figures ?? and ?. *Right*: Comparison of  $H\alpha$  (in grey-scale) to the hot ( $T = 400 - 1000$  K)  $H_2$  mass (in contours). The  $H\alpha$  image is in units of counts/sec. The hot  $H_2$  mass contours are the same as in Figures ?? and ?.

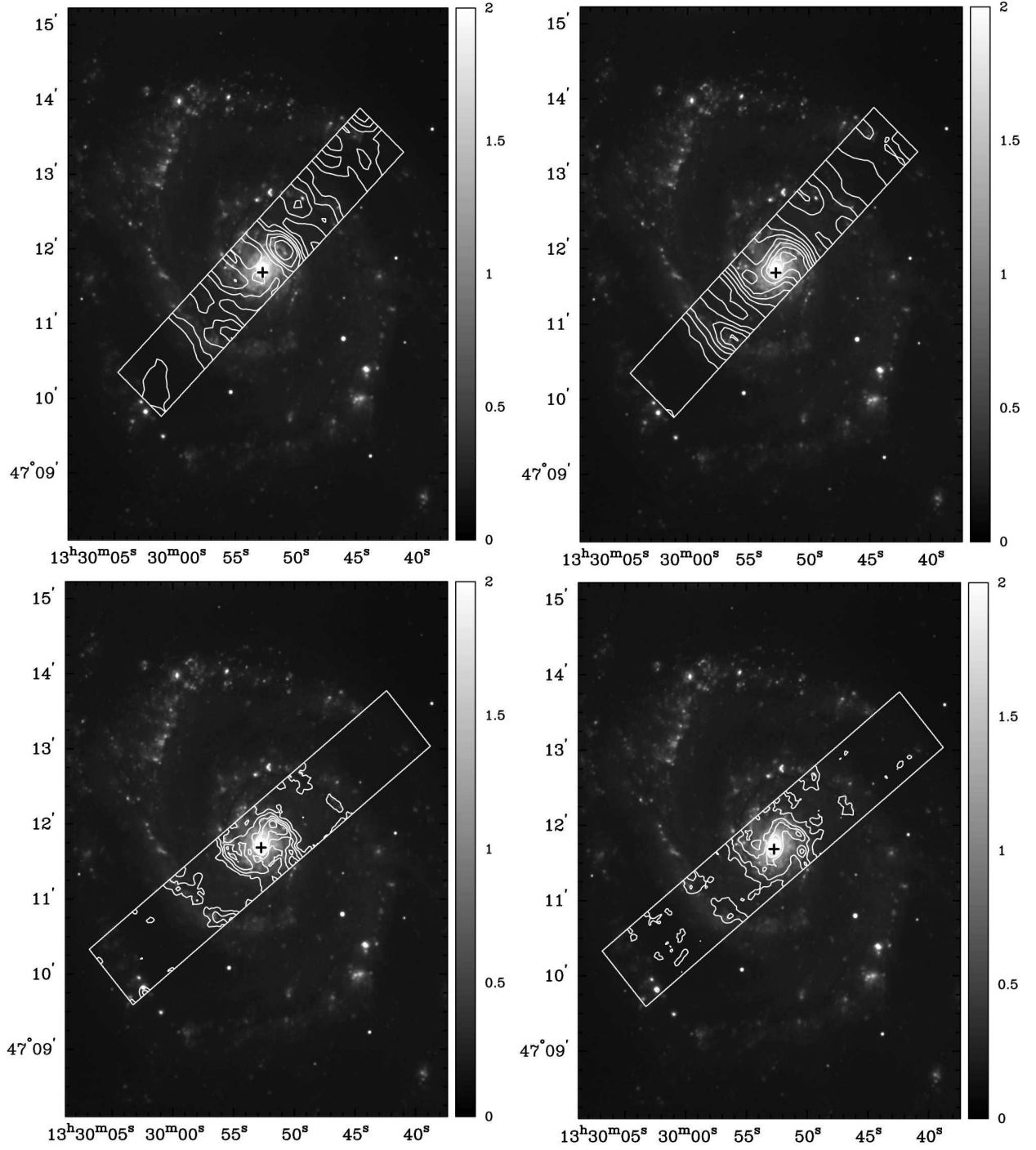


Fig. 10.— Comparison of  $H\alpha$  emission to the  $H_2$  S(0) (*top left*),  $H_2$  S(1) (*top right*),  $H_2$  S(2) (*bottom left*), and  $H_2$  S(3) (*bottom right*) emission. The  $H\alpha$  image is in units of counts/s. Contour levels for  $H_2$  S(0),  $H_2$  S(1),  $H_2$  S(2), and  $H_2$  S(3) are the same as in Figure ??.

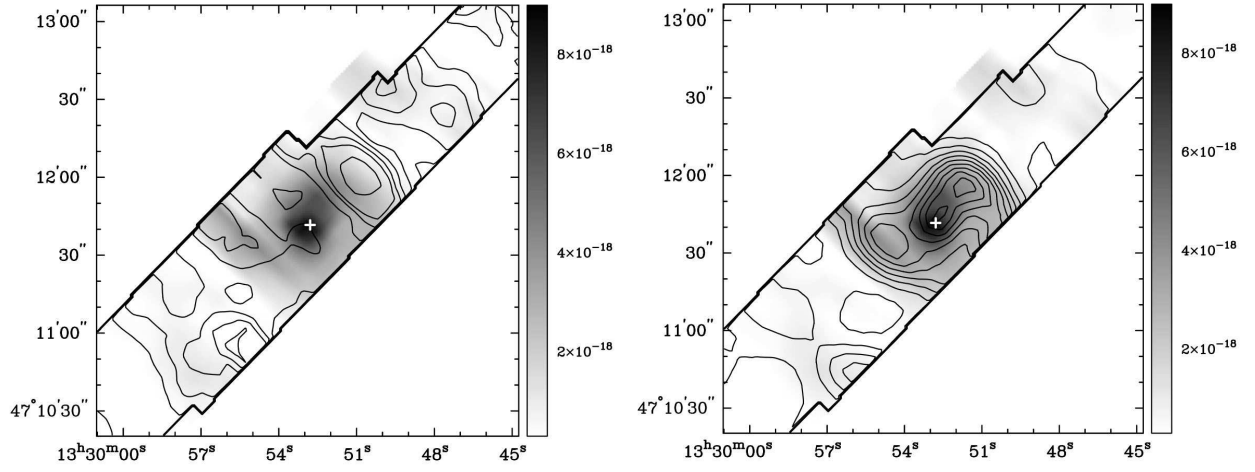


Fig. 11.— *Left*: Comparison of the [O IV](25.89  $\mu\text{m}$ ) emission (in grey-scale) to the warm ( $T = 100\text{ K} - 300\text{ K}$ )  $\text{H}_2$  mass distribution (in contours). Hot  $\text{H}_2$  mass contours are the same as in Figures ?? and ?. *Right*: Comparison of the [O IV](25.89  $\mu\text{m}$ ) emission (in grey-scale) to the hot ( $T = 400 - 1000\text{ K}$ )  $\text{H}_2$  mass distribution (in contours). Hot  $\text{H}_2$  mass contours are the same as in Figures ?? and ?. The [O IV](25.89  $\mu\text{m}$ ) emission is in units of  $\text{W}/\text{m}^2$ .

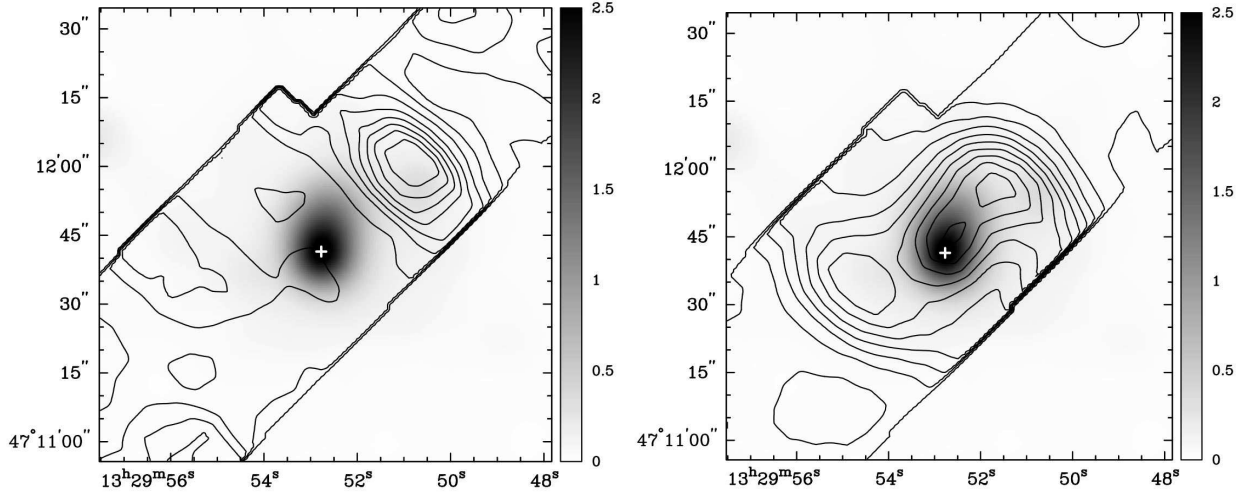


Fig. 12.— *Left*: Comparison of the smoothed 0.5 – 10 keV X-ray emission band (in grey-scale) to the warm ( $T = 100 - 300$  K)  $\text{H}_2$  mass distribution (in contours). The X-ray image has been smoothed to the same resolution as the warm  $\text{H}_2$  mass map. X-ray emission is in units of counts.  $\text{H}_2$  mass contours are the same as in Figures ?? and ?. *Right*: Comparison of the smoothed 0.5 – 10 keV X-ray emission band (in grey-scale) to the hot ( $T = 400 - 1000$  K)  $\text{H}_2$  mass distribution (in contours). The  $\text{H}_2$  mass distribution contours are the same as in Figures ?? and ?.



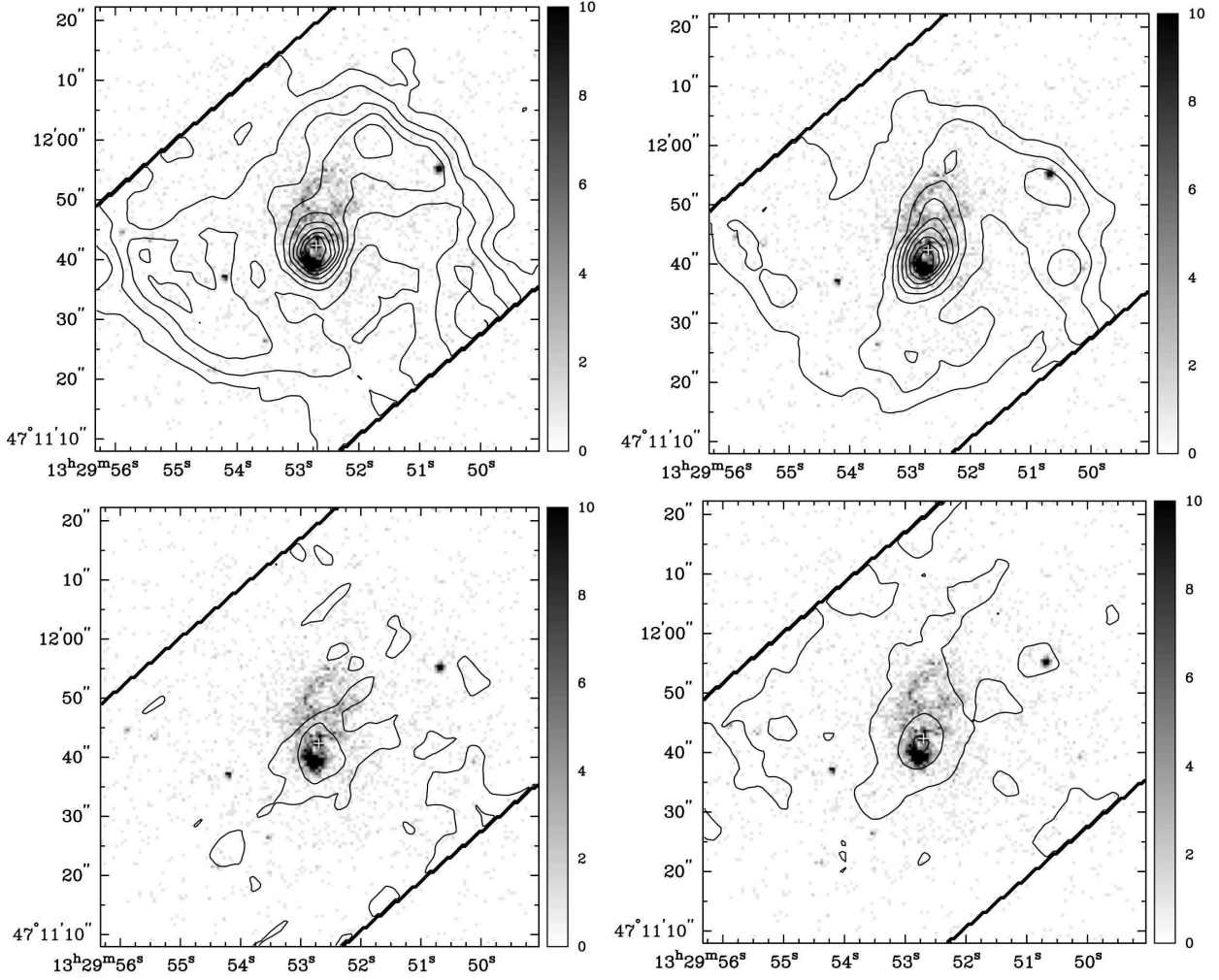


Fig. 13.— Comparison of the 0.5 – 10 keV X-ray emission band (in grey-scale) to the H<sub>2</sub> S(2) (*top left*), H<sub>2</sub> S(3) (*top right*), H<sub>2</sub> S(4) (*bottom left*), and H<sub>2</sub> S(5) (*bottom right*) emission in the nuclear region of M51. X-ray emission is in units of counts. The H<sub>2</sub> S(2) and H<sub>2</sub> S(3) emission contours are at 10 % of their peak values ( $2.20 \times 10^{-18}$  and  $1.35 \times 10^{-17}$  W/m<sup>2</sup>, respectively). The H<sub>2</sub> S(4) contours are at  $2.0 \times 10^{-18}$  and  $1.0 \times 10^{-18}$  W/m<sup>2</sup> and the H<sub>2</sub> S(5) contours are at  $7.3 \times 10^{-18}$ ,  $4.0 \times 10^{-18}$ , and  $8.0 \times 10^{-19}$  W/m<sup>2</sup>.

Table 1. H<sub>2</sub> Parameters

| Transition               | Wavelength ( $\mu\text{m}$ ) | Rotational State (J) | Energy (E/k) | A ( $\text{s}^{-1}$ )  | Statistical Weight (g) |
|--------------------------|------------------------------|----------------------|--------------|------------------------|------------------------|
| H <sub>2</sub> (0-0)S(0) | 28.22                        | 2                    | 510          | $2.94 \times 10^{-11}$ | 5                      |
| H <sub>2</sub> (0-0)S(1) | 17.04                        | 3                    | 1015         | $4.76 \times 10^{-10}$ | 21                     |
| H <sub>2</sub> (0-0)S(2) | 12.28                        | 4                    | 1682         | $2.76 \times 10^{-9}$  | 9                      |
| H <sub>2</sub> (0-0)S(3) | 9.66                         | 5                    | 2504         | $9.84 \times 10^{-9}$  | 33                     |
| H <sub>2</sub> (0-0)S(4) | 8.03                         | 6                    | 3474         | $2.64 \times 10^{-8}$  | 13                     |
| H <sub>2</sub> (0-0)S(5) | 6.91                         | 7                    | 4586         | $5.88 \times 10^{-8}$  | 45                     |

Note. — The statistical weight (g) is  $(2J+1)(2I+1)$  where  $I$  equals 1 for odd J transitions (ortho transitions) and  $I$  equals 0 for even J transitions (para transitions).