



Norwegian University of  
Science and Technology

# CUT FINITE ELEMENT METHOD FOR THE CAHN-HILLIARD EQUATION

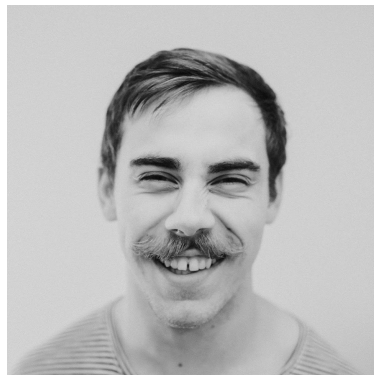
Supervised by André Massing

Isak Hammer

May 4, 2023

# Introducing Myself

- ▶ Isak Hammer, 27 year old, Lofoten
- ▶ Graduate student in Industrial Mathematics
- ▶ Research Focus: Numerical methods for Partial Differential Equations (PDEs).



# Importance and Motivation of the Cahn Hilliard Equation

- ▶ Thermodynamically modelling of a two-component liquid separation<sup>1</sup>.
- ▶ Modelling of so-called lipid rafts in biological membrane dynamics<sup>2</sup>.
- ▶ Droplet dynamics, i.e., coalescence, breakup and movement by coupling with Navier-Stokes<sup>3</sup>.

---

<sup>1</sup>John W Cahn and John E Hilliard. "Free energy of a nonuniform system. III. Nucleation in a two-component incompressible fluid". In: *The Journal of chemical physics* 31.3 (1959), pp. 688–699

<sup>2</sup>Vladimir Yushutin et al. "A computational study of lateral phase separation in biological membranes". In: *International journal for numerical methods in biomedical engineering* 35.3 (2019), e3181

<sup>3</sup>Patrick Zimmermann, Andrew Mawbey, and Tim Zeiner. "Calculation of droplet coalescence in binary liquid–liquid systems: An incompressible Cahn–Hilliard/Navier–Stokes approach using the non-random two-liquid model". In: *Journal of Chemical & Engineering Data* 65.3 (2019), pp. 1083–1094

## The Cahn Hilliard Equation

The general Cahn Hilliard Equation has the form  $u(x, t) : \Omega \times [0, T] \mapsto [-1, 1]$  s.t.

$$\begin{aligned}u_t + \Delta \left( \varepsilon \Delta u - \frac{1}{\varepsilon} f(u) \right) &= 0 \quad \text{in } \Omega \\ \partial_n u &= \partial_n \Delta u = 0 \quad \text{on } \Gamma \\ u &= u_0 \quad \text{on } \Omega\end{aligned}$$

where  $f(s) = F'(s)$  and  $F(s) = \frac{1}{4} (s^2 - 1)^2$  and  $\Omega \subset \mathbf{R}^d, d = 2, 3$ , is a bounded domain.

## Challenges

1. Highly nonlinear and stiff. Often practical applications require  $\varepsilon \ll 1$ .
2. 4th order system.

## Why Finite Element Method (FEM)

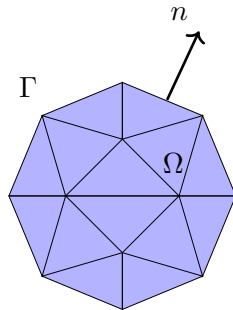
1. **Robust mathematical framework**
2. **Can easily handle complex geometries**
3. **High flexibility of basis functions**
4. **Other:** Supports adaptive refinements, easily adaptable to multi-physics problems ++ .

# The Biharmonic Problem (on a polygon)

Let  $\Omega \approx \Omega_h = \mathcal{T}_h$  be a bounded **polygonal** domain with boundary  $\Gamma$ . Let the biharmonic problem have the form s.t.  
 $u : \Omega \mapsto \mathbb{R}$ ,

$$\begin{aligned}\Delta^2 u + \alpha u &= f(x) && \text{in } \Omega, \\ \partial_n u &= 0 && \text{on } \Gamma, \\ \partial_n \Delta u &= 0 && \text{on } \Gamma.\end{aligned}\tag{1}$$

Here is  $\Delta^2 = \Delta(\Delta)$  the biharmonic operator.



**Figure:** Illustration of the mesh  $\Omega_h$ , the boundary  $\Gamma$  and the normal vector  $n$ .

# $C^0$ Interior Penalty Method (CIP) for the Biharmonic Problem

The proposed numerical scheme is to find an  $w \in V_h$  .t.

$$a_h(w, v) = l_h(v) = (f, v)_\Omega, \quad \forall v \in V_h.$$

where

$$\begin{aligned} a_h(w, v) = & (\alpha w, v)_\Omega + (\Delta w, \Delta v)_\Omega \\ & + (\llbracket \Delta w \rrbracket, [\partial_n v])_{\mathcal{F}_h} + (\llbracket \Delta v \rrbracket, [\partial_n w])_{\mathcal{F}_h} + \frac{\gamma}{h} ([\partial_n w], [\partial_n v])_{\mathcal{F}_h} \end{aligned}$$

Which is inspired from Brenner2012 <sup>1</sup>

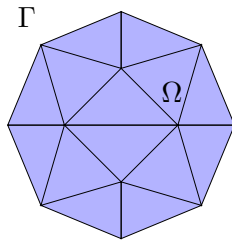
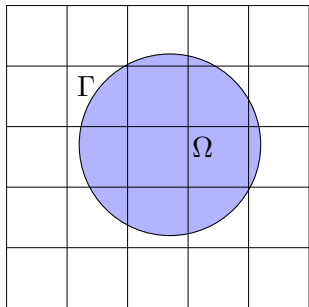
---

<sup>1</sup>Susanne Brenner. *C0 Interior Penalty Methods*. Springer International Publishing, 2012. URL: [https://link.springer.com/content/pdf/10.1007/978-3-642-23914-4\\_2.pdf](https://link.springer.com/content/pdf/10.1007/978-3-642-23914-4_2.pdf)

# Cut Finite Element Method (CutFEM)

## Unfitted mesh vs fitted mesh

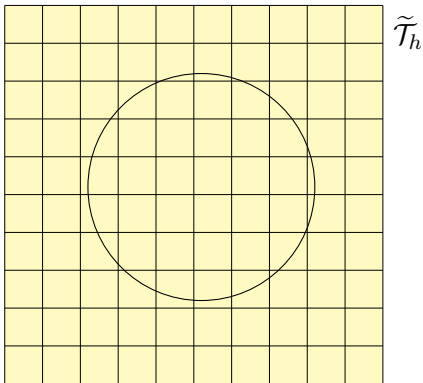
CutFEM is a numerical method for solving partial differential equations (PDEs) using an unfitted mesh.





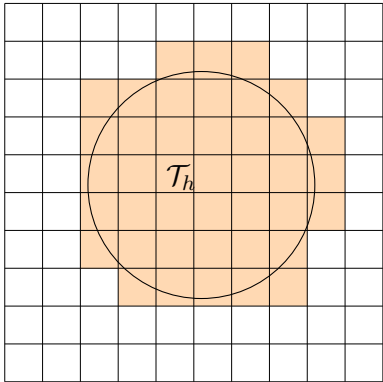
# Cut Finite Element Method

## Background Mesh



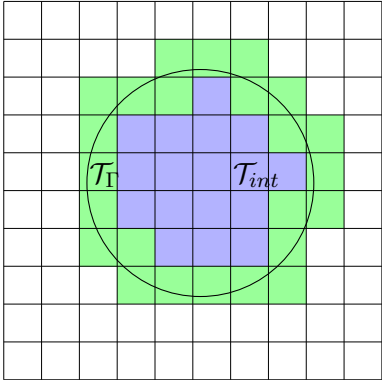
# Cut Finite Element Method

## Active Mesh



# Cut Finite Element Method

## Interior Mesh and Cut Cells



# Cut Finite Element Method

A recent and promising numerical technique for PDEs, has gained significant momentum in the past decade <sup>12</sup>.

- ▶ Complex domains and moving domains efficiently.
- ▶ Utilizing so-called ghost penalties to ensure well-posedness.

---

<sup>1</sup>Erik Burman et al. "CutFEM: discretizing geometry and partial differential equations". In: *International Journal for Numerical Methods in Engineering* 104.7 (2015), pp. 472–501

<sup>2</sup>Ceren Gürkan and André Massing. "A stabilized cut discontinuous Galerkin framework for elliptic boundary value and interface problems". In: *Computer Methods in Applied Mechanics and Engineering* 348 (2019), pp. 466–499

# Cut $C^0$ Interior Penalty Method (CutCIP)

The discretized numerical problem is to solve  $w \in V_h$  such that

$$A(w, v) = a_h(w, v) + g_h(w, v) = l_h(v), \quad \forall v \in V_h.$$

Where the additional bilinear term  $g_h(w, v) : V_h \times V_h \rightarrow \mathbb{R}$  is the so-called **ghost penalty**, which does the numerical regularization to ensure stability on cut cells.

# Cut $C^0$ Interior Penalty Method

My master's thesis is dedicated to demonstrating that the relevant properties remain valid for CutCIP formulation still holds.

## Well-posedness

The discrete bilinear form  $a_h$  is wellposed on  $V_h$  if this holds;

$$(Coercivity) \quad A(v, v) \gtrsim \|v\|_A^2 \quad \forall v \in V_h$$

$$(Boundedness) \quad A(v, w) \lesssim \|v\|_A \|w\|_{a_h} \quad \forall v, w \in V_h$$

# Cut $C^0$ Interior Penalty Method Results

## Manufactured solution

In the experiments will we only consider polynomial order  $k = 2$ . We consider the manufactured solution:

$$u_{ex}(\mathbf{x}) = (x_1^2 + x_2^2 - 1)^2 \cos(2\pi x_1) \cos(2\pi x_2)$$

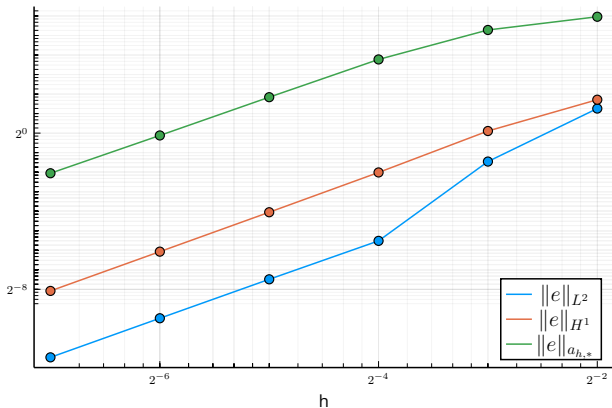
where  $\mathbf{x} = (x_1, x_2)$  and  $\Omega = \{(x_1, x_2) : x_1^2 + x_2^2 \leq 1\}$ . This manufactured solution can be used to test the accuracy of numerical methods for solving the above differential equation.

# Cut $C^0$ Interior penalty method (CutCIP) Results

$n$	$\ e\ _{L^2}$	EOC	$\ e\ _{H^1}$	EOC	$\ e\ _{a_h,*}$	EOC	Cond number	ndofs
4	2.4E+00		3.3E+00		6.2E+01		8.7E+04	8.1E+01
8	3.6E-01	2.72	1.1E+00	1.60	3.9E+01	0.68	5.1E+05	2.4E+02
16	2.2E-02	4.06	2.5E-01	2.12	1.4E+01	1.51	3.7E+06	8.3E+02
32	5.6E-03	1.97	6.0E-02	2.04	3.6E+00	1.93	2.8E+07	3.0E+03
64	1.4E-03	2.00	1.5E-02	2.02	9.2E-01	1.96	2.1E+08	1.1E+04
128	3.5E-04	2.00	3.7E-03	2.01	2.4E-01	1.94	1.7E+09	4.3E+04



# Cut $C^0$ Interior penalty method (CutCIP) Results



**Figure:** The plot presents the  $L^2$  and  $H^1$  error norms and the error in the energy norm ( $\|e\|_{a_h,*}$ ).

# The Cahn Hilliard Equation

## Recall

The problem has the form

$u(x, t) : \Omega \times [0, T] \mapsto [-1, 1]$  s.t.

$$\begin{aligned} u_t + \Delta \left( \varepsilon \Delta u - \frac{1}{\varepsilon} f(u) \right) &= 0 && \text{in } \Omega \\ \partial_n u &= \partial_n \Delta u = 0 && \text{on } \Gamma \\ u &= u_0 && \text{on } \Omega \end{aligned}$$

where  $f(u)$  is a nonlinear function.

## Plan forward

1. We have now a tool to solve the  $\Delta(\Delta u)$  operator
2. Will utilize the time-iteration scheme to solve non-linearity

# The CutCIP Cahn-Hilliard Formulation

Drawing upon the concepts delineated in Feng<sup>1</sup>, the most efficient approach to address the nonlinear term is by employing an implicit-explicit (IMEX) scheme.

## IMEX method on the CutCIP formulation

Let  $u_h^m \in V_h$  for the timesteps  $m = 0, 1, \dots, M$ . Let  $u_h^0 = u_0$  be the initial timestep, then is.

$$(\bar{\partial}_t u_h^m, v_h) + \varepsilon A(u_h^m, v_h) + \frac{1}{\varepsilon} c_h(u_h^{m-1}, v_h) = 0 \quad \forall v_h \in V_h^m.$$

Here is  $c_h(.,.)$  an the nonlinear terms handled in a implicit fashion. The  $\bar{\partial}_t$  operator is simply a finite difference scheme in time-dimension.

---

<sup>1</sup>Xiaobing Feng and Ohannes Karakashian. "Fully discrete dynamic mesh discontinuous Galerkin methods for the Cahn-Hilliard equation of phase transition". In: *Mathematics of computation* 76.259 (2007), pp. 1093–1117

# The CutCIP Cahn-Hilliard Experiments

Implemented using the Gridap FEM framework written in Julia <sup>1</sup>.

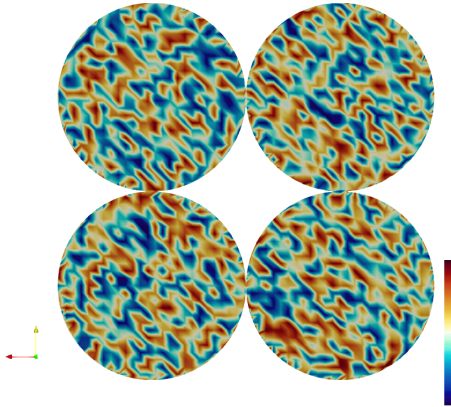
## Simulation parameters

- ▶ Physical domain  $\Omega$  is a 4 discs of radius  $R = 1$  with distance  $d = 0.999$ , i.e. they are touching!
- ▶ Initial data is  $u_0 = \text{random}(-1, 1)$  in physical domain  $\Omega$ .

---

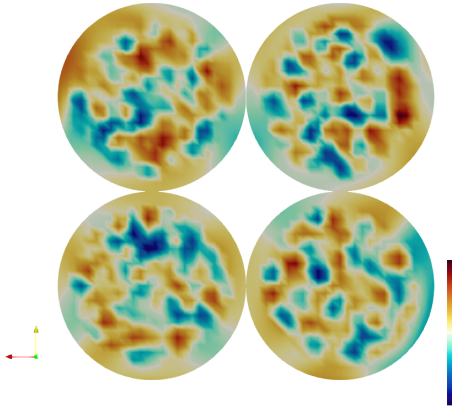
<sup>1</sup>Santiago Badia and Francesc Verdugo. "Gridap: An extensible finite element toolbox in julia". In: *Journal of Open Source Software* 5.52 (2020), p. 2520

# The CutCIP Cahn-Hilliard Experiments



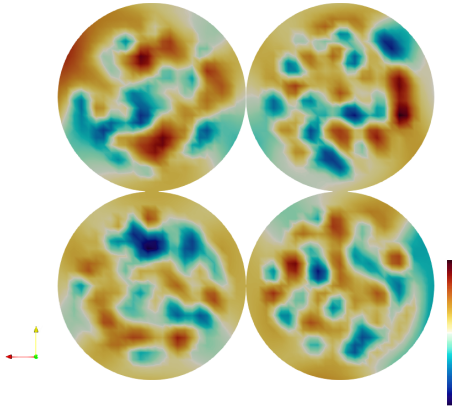
**Figure:** Iteration 0

# The CutCIP Cahn-Hilliard Experiments



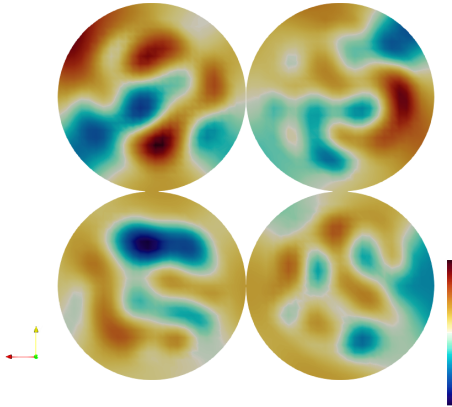
**Figure:** Iteration 1

# The CutCIP Cahn-Hilliard Experiments



**Figure:** Iteration 10

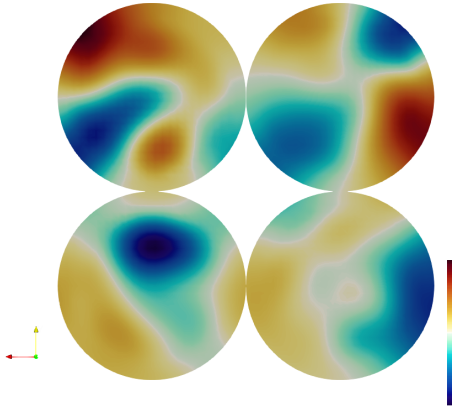
# The CutCIP Cahn-Hilliard Experiments



**Figure:** Iteration 50

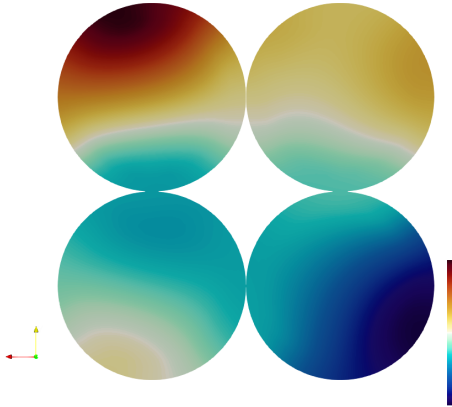


# The CutCIP Cahn-Hilliard Experiments



**Figure:** Iteration 200

# The CutCIP Cahn-Hilliard Experiments



**Figure:** Iteration 1000

# Further work

1. Adaptive time steps.
2. Further numerical validation.
3. Extend the method to handle moving domains.

# Questions?