

The strategy is to bound each term individually using Cauchy-Schwarz (2.2). From this is it easy to see that $|(\text{I})| + |(\text{II})| \lesssim \|v\|_{a_h} \|w\|_{a_h}$. For the terms symmetrical terms (III) and (IV) we apply the inverse inequality (4.10).

$$|(\text{III})| \lesssim \|h^{\frac{1}{2}} \partial_{nn} v\|_{\mathcal{F}_h \cap \Omega} \|h^{-\frac{1}{2}} [\partial_n w]\|_{\mathcal{F}_h \cap \Omega} \lesssim \|v\|_{a_h} \|w\|_{a_h}. \quad (4.29)$$

Here we used that $\|h^{\frac{1}{2}} \partial_{nn} v\|_{\mathcal{F}_h \cap \Omega} \lesssim \|v\|_{A_h}$ thanks to Corollary 4.1. The interior penalty can we easily see that,

$$|(\text{V})| \lesssim \|h^{-\frac{1}{2}} [\partial_n v]\|_{\mathcal{F}_h \cap \Omega} \|h^{-\frac{1}{2}} [\partial_n w]\|_{\mathcal{F}_h \cap \Omega} \lesssim \|v\|_{a_h} \|w\|_{a_h}. \quad (4.30)$$

It remains to handle the symmetry terms (VI) and (VII).

$$|(\text{VI})| \lesssim \|h^{\frac{1}{2}} \partial_{nn} v\|_{\Gamma} \|h^{-\frac{1}{2}} \partial_n w\|_{\Gamma} \lesssim \|v\|_{a_h} \|w\|_{a_h} \quad (4.31)$$

Again, here we used the Corollary 4.1. Finally, using the definition of the norm is it easily to see that,

$$|(\text{VIII})| \lesssim \|\partial_n v\|_{\Gamma} \|\partial_n w\|_{\Gamma} \lesssim \|v\|_{a_h} \|w\|_{a_h}.$$

Hence, we can conclude

$$|a_h(v, w)| \leq \|v\|_{a_h} \|w\|_{a_h} \forall v, w \in V_h. \quad (4.32)$$

Therefore, given that $\|\cdot\|_{a_h} \lesssim \|\cdot\|_{A_h}$, it has been demonstrated that $a_h(\cdot, \cdot)$ is bounded within the $\|\cdot\|_{A_h}$ norm.

Step 2. The goal is to prove (4.25). Let $v \in V_h \oplus V$ and $w \in V_h$. The only difference is that since v is continuous we cannot apply to Corollary 4.1 on the estimates (4.29) and (4.31). However, this is not a problem since $\|h^{\frac{1}{2}} \{\partial_{nn} v\}\|_{\mathcal{F}_h \cap \Omega}$ and $\|h^{\frac{1}{2}} \partial_{nn} v\|_{\Gamma}$ are terms in the norm $\|v\|_{a_h,*}$. Thus, we know that

$$|a_h(v, w)| \leq \|v\|_{a_h,*} \|w\|_{A_h} \quad \forall v, w \in V_h \quad (4.33)$$

with \square

4.4 A priori estimates

A priori error analysis (estimates)

In this section, we will only briefly overview the important assumptions and definitions and then prove that the ghost penalty does affect the convergence rate specifically for the energy norm. However, the C^0 discrete solution is too rough for the standard Lagrange interpolation operator. Hence, this motivates us to introduce a method to interpolate a non-smooth function, the so-called Clément interpolation operator. ↳ this sentence does not make sense.

4.4.1 Energy a priori estimates. If you have only a single substitution, how can safely remove it.

Recall that for $v \in H^1(\mathcal{T}_h)$ these inequalities holds $\forall T \in \mathcal{T}_h$.

$$\begin{aligned} \|v\|_{\partial T} &\lesssim h_T^{-\frac{1}{2}} \|v\|_T + h_T^{\frac{1}{2}} \|\nabla v\|_T, \\ \|v\|_{\Gamma \cap T} &\lesssim h^{-\frac{1}{2}} \|v\|_T + h_T^{\frac{1}{2}} \|\nabla v\|_T. \end{aligned}$$

For proof, see [92, Lemma 4.2]. under, a key idea for what?

A key idea is to utilize between the relationship between the physical space Ω and the active mesh \mathcal{T}_h . Assume that Ω has a boundary Γ in C^1 , then does it exist an bounded extension operator,

$$(.)^e : W^{m,q}(\Omega) \rightarrow W^{m,q}(\Omega^e) \quad \rightarrow \text{what is } \Omega^e? \text{ Just make it } \mathbb{R}^d.$$

for all $v \in W^{m,q}(\Omega)$ where $0 < m \leq \infty$ and $1 \leq q \leq \infty$ which satisfies

$$\begin{aligned} v^e|_{\Omega} &= v, \\ \|v^e\|_{m,q,\Omega^e} &\lesssim \|v\|_{m,q,\Omega}. \end{aligned}$$

This extension theorem primarily utilizes [93, Theorem 9.7] by Brezis, with valuable context provided by [94, p.181, p. 185]. word choice. Just refer to [93], [94] for proofs.

① Rewrite the introducing outline of this section.

Outline should include:

- Describe goal: for proposed method, we want to derive a prior error estimates with respect to both the energy and L^2 -norm
- Briefly outline how to reach goal
 - Construction of a suitable (quasi-) interpolation operator.
here we use the Clement quasi-interpolation operator, which in contrast to the std. Lagrange nodal interpolation operator is also defined for less regularity functions $w \in L^2$.
 - In combination with discrete coercivity this allows you to derive an a priori error estimate in the energy norm
 - Finally you use a std. duality argument (Lubich-Nitsch) to derive the L^2 error estimate.

Why? We don't need that...
We can simply stick to $\mathfrak{L}^e = \mathfrak{L}^d$.

Lipschitz boundary is enough.

Remark. Be aware that this theorem requires that we have a sufficiently smooth domain, hence, emphasizing the assumption of a sufficiently smooth boundary. Also keep in mind that the original theorem states the mapping $(\cdot)^e : W^{m,q}(\Omega) \rightarrow W^{m,q}(\mathbb{R}^d)$, but we restrict ourselves to a slightly bigger extension $\Omega^e \supset \Omega$. This is useful because we are now able to extend the function to the active set, while preserving the Sobolev regularity.

TODO: Not sure if this argument holds. At least from my point of view is this generalization is not trivial.

Now construct the extended Sobolev space Ω_h^e that $\Omega_h^e = T_h \subset \Omega^e$. We define an unfitted Clément interpolator $C_h^e : H^m(\Omega_h^e) \rightarrow V_h$ s.t. $C_h^e v := C_h v^e$. We can immediately observe that the interpolation satisfies the global error estimates, that is,

→ down line →
she is supposed to be at

$$\|v - C_h^e v\|_{m,2,\mathcal{T}_h} \lesssim h^{l-m} \sum_{T \in \mathcal{T}_h} \|v\|_{l,2,\omega(T)}, \quad m \leq l \leq k+1$$

had to swap in and out

$$\|v - C_h^e v\|_{m,2,\mathcal{F}_h} \lesssim h^{l-m-\frac{1}{2}} \sum_{T \in \mathcal{T}_h} \|v\|_{l,2,\omega(F)}, \quad m + \frac{1}{2} \leq l \leq k+1$$

set collection of triangles.
so you have to write

$$\|v - C_h^e v\|_{m,2,\Gamma} \lesssim h^{l-m-\frac{1}{2}} \sum_{T \in \mathcal{T}_h} \|v\|_{l,2,\omega(T)}, \quad m + \frac{1}{2} \leq l \leq k+1$$

$\mathcal{T}_h = \bigcup_{T \in \mathcal{T}_h} T$.

Don't forget extension symbol, otherwise v on T_h does not make sense. Or you state that it is implicitly understood that $(\cdot)^e$ is applied whenever v is considered and argue that outside its original domain, $C\|v\|_{s,T_h}$. Hence

$$\|v - C_h^e v\|_{m,2,\mathcal{T}_h} \lesssim h^{l-m} \|v\|_{l,2,\mathcal{T}_h}, \quad m \leq l \leq k+1$$

$$\|v - C_b^e v\|_{m, 2, \mathcal{F}_l} \lesssim h^{l-m-\frac{1}{2}} \|v\|_{l, 2, \mathcal{T}_l}, \quad m + \frac{1}{2} \leq l \leq k+1$$

$$\|v - C_h^e v\|_{m,2,\Gamma} \lesssim h^{l-m-\frac{1}{2}} \|v\|_{l,2,\mathcal{T}_h}, \quad m + \frac{1}{2} \leq l \leq k+1$$

Maybe hard to argue (4.36) to hold on Γ , but may be related to some generalization of (4.11) and (4.9). Anyhow, (4.35) and (4.36) was never used in the proof of Lemma 4.4 since we used inverse estimates and ended up with (4.34) instead on all of them.

Naturally can we see this is the tools we need to construct an estimate for the energy norm.

Lemma 4.4. Let $u \in H^s(\Omega)$ for $s \geq 3$ be a exact solution to (??). Then we have

$$\|u - C_h u\|_{c_h, *} \lesssim h^{s-2} \|u\|_{H^s(\Omega)}$$

Proof. By definition is

$$\begin{aligned} \|u - C_h^e u\|_{a_h, *}^2 &= \|\alpha|^{\frac{1}{2}}(u - C_h^e u)\|_{\mathcal{T}_h \cap \Omega}^2 + \|D^2(u - C_h^e u)\|_{\mathcal{T}_h \cap \Omega}^2 \\ &\quad + \gamma \|h^{-\frac{1}{2}} [\partial_n(u - C_h^e u)]\|_{\mathcal{F}_h \cap \Omega}^2 + \gamma \|h^{-\frac{1}{2}} \partial_n(u - C_h^e u)\|_{\Gamma}^2 \\ &\quad + \|h^{\frac{1}{2}} \{\partial_{nn}(u - C_h^e u)\}\|_{\mathcal{F}_h \cap \Omega}^2 + \|h^{\frac{1}{2}} \partial_{nn}(u - C_h^e u)\|_{\Gamma}^2. \end{aligned}$$

The strategy is to bound each term individually. Starting with the first two terms we get

$$\begin{aligned} \|\alpha^{\frac{1}{2}}(u - C_h^e u)\|_{\mathcal{T}_h \cap \Omega}^2 &\lesssim \|(u - C_h^e u)\|_{0,2,\mathcal{T}_h}^2 \\ &\lesssim h^{2(s-0)} \|u\|_{s,\mathcal{T}_h}^2 \lesssim h^{2(s-2)} \|u\|_{s,\mathcal{T}_h}^2 \\ \|D^2(u - C_h^e u)\|_{\mathcal{T}_h \cap \Omega}^2 &\lesssim \|u - C_h^e u\|_{2,\mathcal{T}_h}^2 = \|u - C_h^e u\|_{2,2,\mathcal{T}_h}^2 \\ &\lesssim h^{2(s-2)} \|u\|_{s,\mathcal{T}_h}^2. \end{aligned}$$

① You notation / index use in element operator is quite messed up (jj)

Also, for later use, already start using approximation index.

Fix element estimates both here and in Section 2.

For $\psi \in H^k(\Omega_h)$ and $C_w: L^2(\Omega_h) \rightarrow P_c^k(\Omega_h)$

and $r = \min\{s, k+1\}$ we have that

- $\| \psi - C_w \psi \|_{L^2} \lesssim \rho_{\omega_T}^{r-l} |\psi|_{r, w(T)} \quad 0 \leq l \leq r$
- $\| \psi - C_w \psi \|_{L^2}^2 \lesssim \rho_{\omega_T}^{r-l-\frac{1}{2}} |\psi|_{r, w(T)} \quad 0 \leq l \leq r - \frac{1}{2}$.
- Also adapt / check remaining section / lemma 4.4 Thm, 4.7.