

# NE 155/255

## Numerical Simulations in Radiation Transport

### Introduction to Monte Carlo

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1/17

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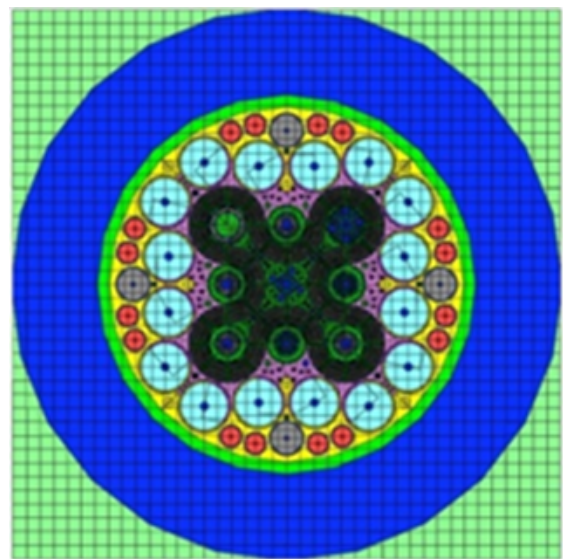
NE 155/255

November 04, 2019

1 / 17

## Learning Objectives

- 1 Define Monte Carlo simulation
- 2 Justify the choice of Monte Carlo for radiation transport
- 3 Understand the mathematical validity of Monte Carlo for radiation
- 4 Understand the major components of Monte Carlo methods transport



**Figure 1:** ATR reactor geometry

Notes derived from Rachel Slaybaugh, Jasmina Vujic, and Paul Wilson.

2/17

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NE 155/255

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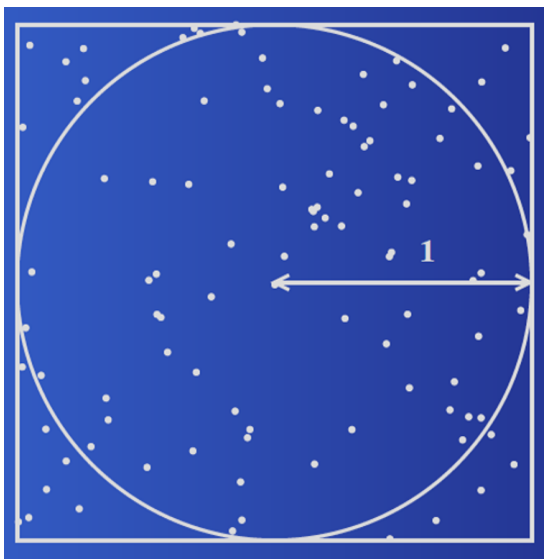
2 / 17

# What are Monte Carlo methods?

- The use of **random processes** to determine a **statistically-expected** solution to a problem
- Random processes can fulfill two roles:
  - Statistical approximation to **mathematical equations**
  - Statistical approximations to **physical processes**
- Construct a random process for a problem
- Carry out a numerical simulation by N-fold sampling from a random number sequence

3/17

## Evaluate $\pi$ by Random Sampling



- Area of square,  $A_s = 4$
- Area of circle,  $A_c = \pi$
- Fraction of random points in circle

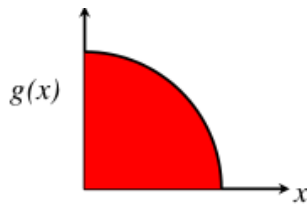
$$p = \frac{A_c}{A_s} = \frac{\pi}{4}$$

- Random points =  $N$
- Random points in circle =  $N_c$ ,  $\therefore$

$$p = \frac{N_c}{N} ; \quad \pi = \frac{4N_c}{N}$$

4/17

## Evaluate $\pi$ by Random Sampling (math)



$$g(x) = \sqrt{1 - x^2} \quad G = \int_0^1 g(x) dx = \frac{\pi}{4}$$

$$G = \int_0^1 g(x) dx = (1 - 0) \overline{g(x)}$$

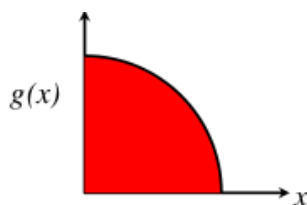
Determine  $\overline{g(x)}$  by random sampling:

for  $k = 1, \dots, N$ , choose  $\hat{x}_k$  randomly on the interval  $(0, 1)$ ,

$$\overline{g(x)} \equiv \frac{1}{N} \sum_{k=1}^N g(\hat{x}_k) = \frac{1}{N} \sqrt{1 - \hat{x}_k^2}$$

5/17

## Evaluate $\pi$ by Random Sampling (physics)



$$g(x) = \sqrt{1 - x^2} \quad G = \int_0^1 g(x) dx = \frac{\pi}{4}$$

$G$  = area under curve,  
= fraction of unit square under curve

for  $k = 1, \dots, N$ , chose  $\hat{x}_k, \hat{y}_k$  randomly on the interval  $[0, 1]$ ,

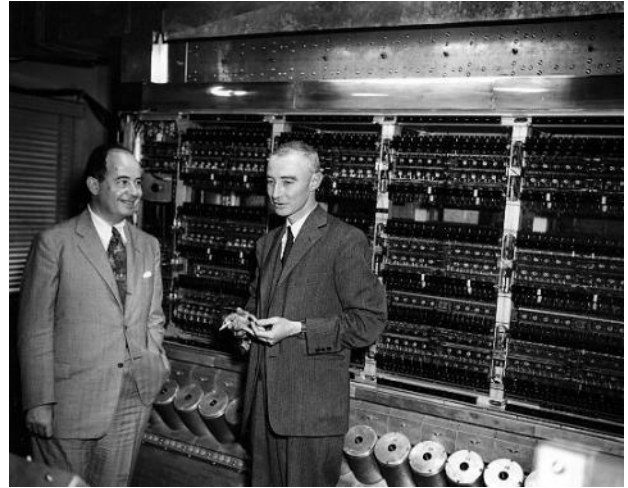
$m_N = \#$  of times in  $N$  trials that  $\hat{x}_k^2 + \hat{y}_k^2 \leq 1$ ,

$$G = \frac{m_N}{N}$$

6/17

# Manhattan Project

- The first human-engineered nuclear detonation, the Trinity Test in New Mexico.
- Active: 1942–1945
- Branch: U.S. Army Corps of Engineers
- Monte Carlo Pioneers:
  - Enrico Fermi,
  - Stanislaw Ulam,
  - John von Neumann,
  - Robert Richtmeyer,
  - Nicholas Metropolis



**Figure 2:** Oppenheimer, von Neumann, MANIAC

Nicholas Metropolis, S. Ulam. "The Monte Carlo Method," *Journal of the American Statistical Association*, **44**, No. 247, 335-341 (Sep. 1949).

7/17

## General Purpose MC Codes

- **MCNP**: developed at LANL, distributed via RSICC, <http://rsicc.ornl.gov>
- **Geant4**: developed by a large collaboration in the HEP community, <http://geant4.web.cern.ch/geant4/>
- **OpenMC**: originally developed at MIT, <https://docs.openmc.org/en/stable/>
- **Serpent**: Developed by Dr. Jaakko Leppanen, VTT, Finland, <http://montecarlo.vtt.fi/>
- **Shift**: developed at ORNL, distributed via RSICC, <http://rsicc.ornl.gov>
- **Mercury**: developed at LLNL, <https://wci.llnl.gov/simulation/computer-codes/mercury>

8/17

## Why/When Monte Carlo?

- Applications that are mathematically equivalent to *integration over many dimensions*
  - Analytic integration may be impossible
  - Deterministic numerical integration may be slow and/or require error prone approximations
- However, statistically accurate results can require **significant computer time**
- Fortunately, Monte Carlo and parallel computing go well together
- We also have Variance Reduction methods

9/17

## What is MC Radiation Transport?

Simulate many independent particles in a system

- Treat each physical process as a *probabilistic process*
- *Randomly sample* each process using an independent stream of random numbers
- Follow each particle from birth until it no longer matters
- Accumulate the contributions of each particle to find the statistically-expected mean behavior and variance

10/17

# Mathematical Validity

- Consider particles with a phase space describing position,  $\vec{r}$ , and velocity,  $\vec{v}$
- A neutral particle can be transmitted from one position to another at a constant velocity

$$T(\vec{r}' \rightarrow \vec{r}, \vec{v})$$

- A particle can undergo a collision at a single position that changes its velocity

$$C(\vec{r}, \vec{v}' \rightarrow \vec{v})$$

11/17

## Contributions After 0 Collisions

- Consider a particle born from a source described by

$$Q(\vec{r}', \vec{v}')$$

- This particle will contribute to the flux at  $(\vec{r}, \vec{v})$  before any collisions

$$\psi_0(\vec{r}, \vec{v}) = \int_{\vec{r}'} Q(\vec{r}', \vec{v}') T(\vec{r}' \rightarrow \vec{r}, \vec{v}) d\vec{r}'$$

12/17

## Contributions After 1 Collision

- The uncollided particles,  $\psi_0(\vec{r}', \vec{v}')$ , could undergo 1 **collision** and then be **transmitted** to the point  $(\vec{r}, \vec{v})$

$$\psi_1(\vec{r}, \vec{v}) = \underbrace{\int_{\vec{r}'} \left[ \underbrace{\int_{\vec{v}'} \psi_0(\vec{r}', \vec{v}') C(\vec{r}', \vec{v}' \rightarrow \vec{v}) d\vec{v}'}_{\text{collision}} \right]}_{\text{transmission}} T(\vec{r}' \rightarrow \vec{r}, \vec{v}) d\vec{r}'$$

13/17

## Contributions After $k$ Collisions

- Particles that have undergone  $k$  collisions,  $\psi_k(\vec{r}', \vec{v}')$ , could undergo another **collision** and then be **transmitted** to the point  $(\vec{r}, \vec{v})$

$$\psi_{k+1}(\vec{r}, \vec{v}) = \underbrace{\int_{\vec{r}'} \left[ \underbrace{\int_{\vec{v}'} \psi_k(\vec{r}', \vec{v}') C(\vec{r}', \vec{v}' \rightarrow \vec{v}) d\vec{v}'}_{\text{collision}} \right]}_{\text{transmission}} T(\vec{r}' \rightarrow \vec{r}, \vec{v}) d\vec{r}'$$

14/17

## Combine Collision and Transmission Kernels

$$\vec{p} = (\vec{r}, \vec{v}) \quad \text{and}$$

$$R(\vec{p}' \rightarrow \vec{p}) \equiv C(\vec{r}', \vec{v}' \rightarrow \vec{v}) T(\vec{r}' \rightarrow \vec{r}, \vec{v})$$

$$\psi_{k+1}(\vec{r}, \vec{v}) = \int_{\vec{p}_k} \psi_k(\vec{p}_k) R(\vec{p}_k \rightarrow \vec{p}_{k+1}) d\vec{p}_k$$

$$\psi_{k+1}(\vec{r}, \vec{v}) = \int_{\vec{p}_k} \left[ \int_{\vec{p}_{k-1}} \psi_{k-1}(\vec{p}_{k-1}) R(\vec{p}_{k-1} \rightarrow \vec{p}_k) d\vec{p}_{k-1} \right] R(\vec{p}_k \rightarrow \vec{p}_{k+1}) d\vec{p}_k$$

... and so on ...

$$\psi_{k+1}(\vec{r}, \vec{v}) = \int_{\vec{p}_k} \int_{\vec{p}_{k-1}} \cdots \int_{\vec{p}_0} \psi_0(\vec{p}_0) R(\vec{p}_0 \rightarrow \vec{p}_1) d\vec{p}_0 \cdots$$

$$\psi_{k-1}(\vec{p}_{k-1}) R(\vec{p}_{k-1} \rightarrow \vec{p}_k) d\vec{p}_{k-1} R(\vec{p}_k \rightarrow \vec{p}_{k+1}) d\vec{p}_k$$

15/17

## Sum Over All Collisions

$$\psi(\vec{p}) = \sum_{k=0}^{\infty} \psi_k(\vec{p})$$

Arriving at the *integral form* of the transport equation

$$\psi(\vec{r}, \vec{v}) = \int_{\vec{r}'} \left[ \int_{\vec{v}'} \psi(\vec{r}', \vec{v}') C(\vec{r}', \vec{v}' \rightarrow \vec{v}) d\vec{v}' \right] T(\vec{r}' \rightarrow \vec{r}, \vec{v}) d\vec{r}'$$

16/17



# Mathematical Validity

$$\Psi_k(\vec{p}) = \int \int \cdots \int \Psi_0(\vec{p}_0) R(\vec{p}_0 \rightarrow \vec{p}_1) R(\vec{p}_1 \rightarrow \vec{p}_2) \cdots R(\vec{p}_{k-1} \rightarrow \vec{p}_k) d\vec{p}_0 d\vec{p}_1 \cdots d\vec{p}_{k-1}$$

- Integration over many variables
- Generate a “history”  
(sequence of states  $\vec{p}_0, \vec{p}_1, \dots, \vec{p}_k$ )
  - Randomly sample from source:  $\Psi_0(\vec{p}_0)$
  - Randomly sample for each of  $k$  transitions:  $R(\vec{p}_{k-1} \rightarrow \vec{p}_k)$
- Average for result  $A$  by averaging of  $M$  histories

$$\langle A \rangle = \int A(\vec{p}) \Psi(\vec{p}) d\vec{p} = \frac{1}{M} \sum_{m=1}^M \left[ \sum_{k=1}^{\infty} A(\vec{p}_{k,m}) \Psi(\vec{p}_{k,m}) \right]$$