

# Quantum Computing

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## 1 Linear Algebra

### 1.3 The Pauli matrices

$$\sigma_0 = I$$

$$\sigma_1 = \sigma_x = X = \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}$$

$$\sigma_2 = \sigma_y = Y = \begin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix}$$

$$\sigma_3 = \sigma_z = Z = \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

### 1.6 Adjoints and Hermitian operators

**Definition 1.1** (Hermitian). An operator  $A$  is Hermitian if  $A = A^\dagger$ .

**Theorem 1.2.** Two eigenvectors of a Hermitian operator with different eigenvalues are orthogonal.

**Definition 1.3.** A matrix is normal if  $AA^\dagger = A^\dagger A$ .

**Theorem 1.4.** A normal matrix is Hermitian iff it has real eigenvalues.

**Definition 1.5** (Unitary). A matrix  $U$  is unitary if  $U^\dagger U = I$ .

**Definition 1.6** (Positive and Positive definite). A positive operator  $A$  is defined to be an operator such that for any vector  $|v\rangle$ ,  $\langle v|A|v\rangle$  is a real, non-negative number. If  $\langle v|A|v\rangle > 0$  then  $A$  is positive definite.

## 1.7 Tensor products

## 1.8 Operator functions

## 1.9 The commutator and anti-commutator

**Definition 1.7** (Commutator). The *commutator* between two operators  $A$  and  $B$  is defined to be

$$[A, B] = AB - BA.$$

If  $[A, B] = 0$  then we say that  $A$  commutes with  $B$ .

**Definition 1.8** (Anti-commutator). The *anti-commutator* between two operators  $A$  and  $B$  is defined to be

$$\{A, B\} = AB + BA.$$

If  $\{A, B\} = 0$  then we say that  $A$  anti-commutes with  $B$ .

**Definition 1.9** (Simultaneously Diagonalizable). Hermitian operators  $A$  and  $B$  are said to be *simultaneously diagonalizable* if there exist some orthonormal set of vectors  $\langle i |$  such that  $A = \sum a_i |i\rangle \langle i|$  and  $B = \sum b_i |i\rangle \langle i|$ .

**Theorem 1.10** (Simultaneous diagonalization theorem). Suppose  $A$  and  $B$  are Hermitian operators. Then  $[A, B] = 0$  iff  $A$  and  $B$  are simultaneously diagonalizable.

Here are some facts

$$[X, Y] = 2iZ; [Y, Z] = 2iX; [Z, X] = 2iY$$

$$[A, B]^\dagger = [B^\dagger, A^\dagger]$$

**Theorem 1.11.** Suppose  $A$  and  $B$  are Hermitian. Then  $i[A, B]$  is also Hermitian.

## 1.10 The polar and singular value decompositions

**Theorem 1.12** (Polar decompositions). Let  $A$  be a linear operator on a vector space  $V$ . Then there exists unitary  $U$  and positive operators  $J$  and  $K$  such that

$$A = UJ = KU,$$

where the unique positive operators  $J$  and  $K$  satisfying these equations are  $J = \sqrt{A^\dagger A}$  and  $K = \sqrt{AA^\dagger}$ . Moreover, if  $A$  is invertible then  $U$  is unique.

**Theorem 1.13** (Singular value decomposition). Let  $A$  be a square matrix. Then there exist unitary matrices  $U$  and  $V$ , and a diagonal matrix  $D$  with non-negative entries such that

$$A = UDV.$$

The diagonal elements of  $D$  are known as the *singular values* of  $A$ .

## 2 The postulates of quantum mechanics

### 2.1 State Space

**Postulate 2.1.** Associated to any isolated physical system is a complex vector space with inner product (that is, a Hilbert space) known as the *state space* of the system. The system is completely described by its *state vector*, which is a unit vector in the system's state space.

### 2.2 Evolution

**Postulate 2.2.** The evolution of a *closed* quantum system is described by a *unitary transformation*. That is, the state  $|\phi\rangle$  of the system at time  $t_1$  is related to the state  $|\phi'\rangle$  of the system at time  $t_2$  by a unitary operator  $U$  which depends only on the times  $t_1$  and  $t_2$ ,

$$|\phi'\rangle = U |\phi\rangle.$$

The  $X$  Pauli matrix is often referred to as the not gate or the *bit flip* gate. It will send  $|0\rangle$  to  $|1\rangle$  and  $|1\rangle$  to  $|0\rangle$ .

The  $Z$  Pauli matrix is called the *phase flip* gate. It leaves  $|0\rangle$  invariant but sends  $|1\rangle$  to  $-|1\rangle$ .

**Definition 2.1** (Hadamard Gate). An interesting unitary operator is the *Hadamard gate*. This had the matrix representation

$$H = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix}$$

**Postulate 2.2** (Revised). The time evolution of the state of a closed quantum system is described by the *Schrödinger equation*,

$$i\hbar \frac{d|\phi\rangle}{dt} = H |\phi\rangle.$$

In this equation,  $\hbar$  is a physical constant known as *Planck's constant* whose value must be experimentally determined. The exact value is not important to us. In practice, it is common to absorb the factor  $\hbar$  into  $H$ , effectively setting  $\hbar = 1$ .  $H$  is a fixed Hermitian operator known as the *Hamiltonian* of the closed system.

### 2.3 Quantum measurement

We introduce Postulate 3 in order to describe the effects that measurement have on a system.

**Postulate 2.3.** Quantum measurements are described by a collection  $\{M_m\}$  of *measurement operators*. These are operators acting on the state space of the system being measured. The index  $m$  refers to the measurement outcomes that may occur in the experiment. If the state of the quantum system is  $|\phi\rangle$  immediately before the measurement then the probability that result  $m$  occurs is given by

$$p(m) = \langle \phi | M_m^\dagger M_m | \phi \rangle,$$

and the state of the system after the measurement is

$$\frac{M_m |\phi\rangle}{\sqrt{\langle \phi | M_m^\dagger M_m | \phi \rangle}}.$$

The measurement operators satisfy the *completeness equation*,

$$\sum_m M_m^\dagger M_m = I.$$

The completeness equation expresses the fact that probabilities sum to one:

$$I = \sum p(m) = \sum \langle \phi | M_m^\dagger M_m | \phi \rangle.$$