#### Glassy Dynamics in Physics & Beyond (Landscape & Glassy Dynamics)

•Dynamical Mean-Field Theory of Glassy dynamics Dynamical mean-field equation via the dynamic cavity method

•How & Why dynamics slow down in rough high-dimensional landscape? Analysis of the Langevin dynamics for the p-spin spherical model

•Rough-landscape with a signal ("crystal hunting")
Dynamics for the spiked-tensor model and generalization

# Rough Landscape and Generalized Spin-Glasses

$$\frac{\sum_{\langle i,j\rangle} J_{ij} s_i s_j}{\langle i,j\rangle} \xrightarrow{\text{More complex}} \sum_{\langle i_1,\cdots,i_p\rangle} J_{i_1,\cdots,i_p} s_{i_1} s_{i_2} \cdots s_{i_p}$$

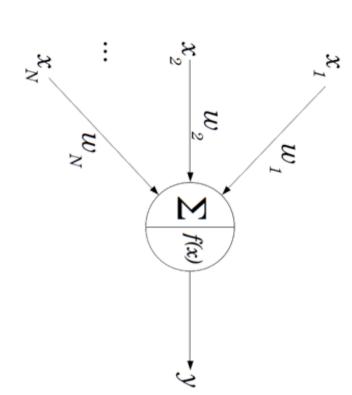
$$\frac{J_{i_1,\cdots,i_p}}{J_{i_1,\cdots,i_p}^2} = \frac{p!}{2N^{p-1}}$$

$$\frac{ds_i(t)}{dt} = -\frac{\partial V}{\partial s_i} + \frac{1}{(p-1)!} \sum_{i_2 \neq \cdots i_p} J_{i,i_2,\cdots,i_p} s_{i_2} \cdots s_{i_p} + \xi_i(t)$$

DMFT:

$$\frac{ds(t)}{dt} = -\frac{\partial V}{\partial s} + \frac{p(p-1)}{2} \int_0^t R(t, t') C^{p-2}(t, t') s(t') dt' + \xi(t)$$
$$\langle \xi(t) \xi(t') \rangle = 2T\delta(t - t') + \frac{p}{2} C^{p-1}(t, t')$$

#### Perceptron



$$E(\lbrace w_i \rbrace) = \sum_{\mu=1}^{M} f(x_{\mu} \cdot w)$$
$$w^2 = N \qquad M = \alpha N$$

- DMFT for Langevin dynamics
  Agoritsas, Biroli, Urbani, Zamponi 2018
  SGD for online learning
- SGD for online learningMace, Coolen '97
- •DMFT for Glauber dynamics Horner '92

### Non-Conservative Dynamics (No Landscape )

sheared glasses Mean-field models of recurrent neural networks, large-interacting eco-systems and

$$\frac{dh_i}{dt} = -h_i + \sum_{l \neq i}^{N} J_{il} S_l = -h_i + \sum_{l \neq i}^{N} J_{il} \tanh(gh_l)$$

$$=rac{1}{N}$$
  $\overline{J_{ji}J_{ji}}=0$ 

DMFT: 
$$\frac{dh}{dt} = -h + \xi(t) \qquad \langle \xi(t)\xi(t')\rangle = \langle \tanh(gh(t)) \tanh(gh(t'))\rangle$$

(Crisanti, Sompolinsky '87,...; Kadmon, Sompolinsky PRX 2015) Transition to chaos by increasing g

#### P-spin spherical model

$$E = -\sum_{\langle i_1, \cdots, i_p \rangle} J_{i_1, \cdots, i_p} S_{i_1} S_{i_2} \cdots S_{i_p}$$

$$\sum_{i=1}^{N} s_i^2 = N$$

$$\frac{ds_i(t)}{dt} = -\frac{\partial E}{\partial s_i} + \xi_i(t)$$

#### P-spin spherical model

$$E = -\sum_{\langle i_1, \dots, i_p \rangle} J_{i_1, \dots, i_p} s_{i_1} s_{i_2} \dots s_{i_p} \qquad \sum_{i=1}^N s_i^2 = N$$

$$\frac{ds_i(t)}{dt} = -\frac{\partial E}{\partial s_i} + \xi_i(t) - \lambda(t) s_i \qquad \lambda(t) = -p \frac{E(t)}{N} + T$$

#### P-spin spherical model

$$E=-\sum_{\langle i_1,\cdots,i_p \rangle} J_{i_1,\cdots,i_p} s_{i_1} s_{i_2} \cdots s_{i_p} \sum_{i=1}^N s_i^2 = 1$$

$$\sum_{i=1}^{N} s_i^2 = N$$

$$\frac{ds_i(t)}{dt} = -\frac{\partial E}{\partial s_i} + \xi_i(t) - \lambda(t)s_i$$

$$\lambda(t) = -p \frac{E(t)}{N} + T$$

$$\frac{ds_i(t)}{dt} = -\lambda(t)s_i + \frac{1}{N} - \sum_{i=1}^{N} \frac{1}{s_i} \frac{1}$$

$$(t) = -p\frac{E(t)}{N} + T$$

$$J = -\lambda(t)s_i + \frac{1}{(p-1)!} \sum_{i_2 \neq \dots i_p} J_{i,i_2,\dots,i_p} s_{i_2} \dots s_{i_p} + \xi_i(t)$$

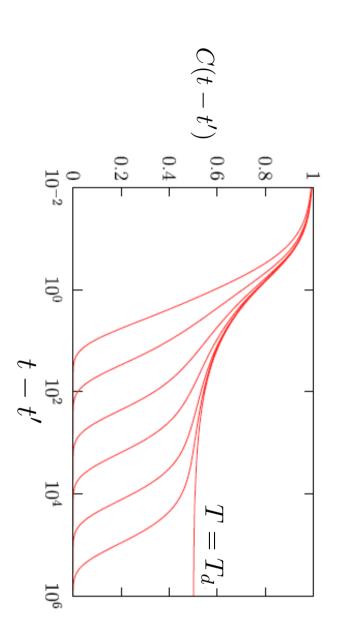
#### P-spin spherical model

$$\frac{ds_i(t)}{dt} = -\lambda(t)s_i + \frac{1}{(p-1)!} \sum_{i_2 \neq \dots i_p} J_{i,i_2,\dots,i_p} s_{i_2} \dots s_{i_p} + \xi_i(t)$$

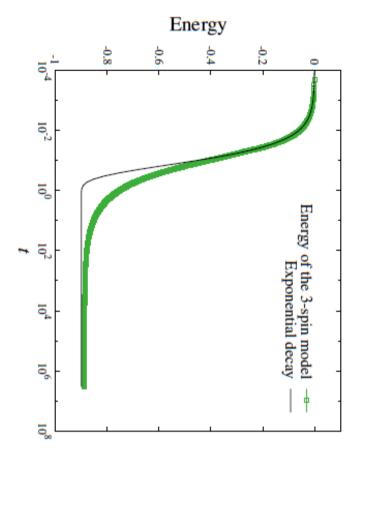
#### DMFT:

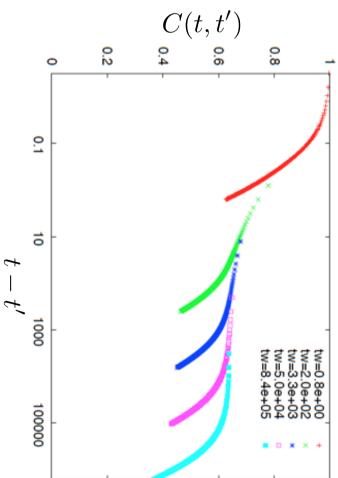
$$\frac{ds(t)}{dt} = -\lambda(t)s(t) + \frac{p(p-1)}{2} \int_0^t R(t,t')C^{p-2}(t,t')s(t')dt' + \xi(t)$$
$$\langle \xi(t)\xi(t') \rangle = 2T\delta(t-t') + \frac{p}{2}C^{p-1}(t,t')$$

P-spin spherical model: EQUILIBRIUM

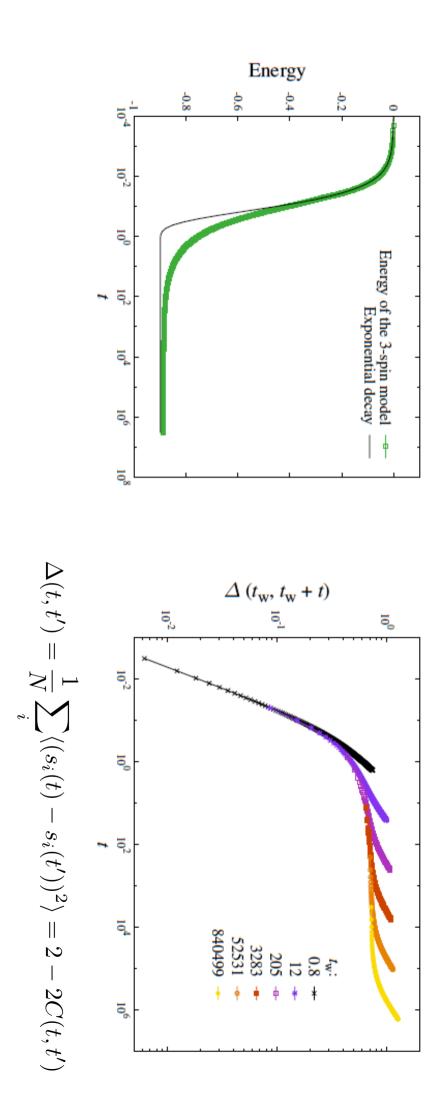


P-spin spherical model: AGING

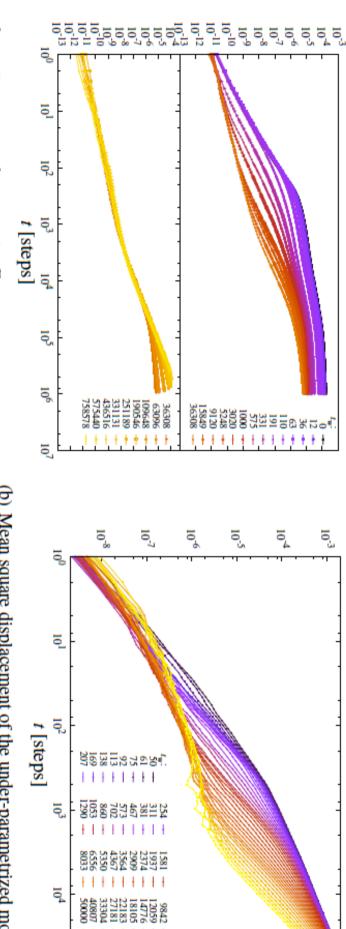




P-spin spherical model: AGING



# Correlation-Functions in DNNs



(b) Fully Connected on MNIST, B = 128,  $\alpha = 0.01$ .

(b) Mean square displacement of the under-parametrized model.

Over-parametrized

Under-parametrized

M. Baity-Jesi, L. Sagun, M. Geiger, S. Spigler, G. Ben Arous, C. Cammarota, Y. LeCun, M. Wyart, G. Biroli (2018) arXiv1803.06969 & ICML 2018

Explanation in terms of changes in the landscape: L. Sagun More on the transition: M. Wyart