

# 第四章 态和力学量的表象

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1. 态的表象.

$$\Psi(x,t) = \frac{1}{(2\pi\hbar)^{1/2}} \int C(p,t) e^{\frac{i}{\hbar} p x} dp \quad \text{坐标表象}$$

$$C(p,t) = \frac{1}{(2\pi\hbar)^{1/2}} \int \Psi(x,t) e^{-\frac{i}{\hbar} p x} dx \quad \text{动量表象}$$

若力学量  $Q$  的本征态是  $u_1(x), u_2(x), \dots, u_n(x), \dots$ , 本征值  $Q_1, Q_2, \dots$

$$\Psi(x,t) = \sum_n a_n(t) u_n(x)$$

$$a_n(t) = \int \Psi(x,t) u_n^*(x) dx$$

$$\Psi = \begin{pmatrix} a_1(t) \\ a_2(t) \\ \vdots \\ a_n(t) \\ \vdots \end{pmatrix} \quad |a_n|^2 \text{ 是在 } \Psi(x,t) \text{ 态中测量力学量 } Q \text{ 所得结果为 } Q_n \text{ 的概率}$$

2. 算符的矩阵表示

$$\begin{aligned} \text{若 } \hat{F} \Psi(x,t) &= \Phi(x,t) \\ \downarrow & \quad \quad \quad \searrow \\ \Psi(x,t) = \sum_m a_m(t) u_m(x) & \quad \quad \quad \Phi(x,t) = \sum_m b_m(t) u_m(x) \end{aligned}$$

$$\text{则矩阵元: } F_{nm} = \int u_n^*(x) \hat{F} u_m(x) dx = \langle u_n | \hat{F} | u_m \rangle$$

Hermite 算符的矩阵是 Hermite 矩阵:

$$F = F^\dagger \quad \rightarrow \quad F_{mn}^\dagger = F_{nm}^*$$

算符在自身表象中是一个对角矩阵, 对角线上的元素是其本征值.

4. 么正变换 (一个表象到另一个表象的变换)

$$\text{么正矩阵: } S^\dagger = S^{-1}$$

么正变换 ( $A$  表象  $\rightarrow B$  表象):

$$\begin{cases} \text{算符} & F' = S^\dagger F S \\ & \downarrow \quad \quad \quad \searrow \\ & \hat{F} \text{ 在 } B \text{ 表象中的矩阵} \quad \quad \quad \hat{F} \text{ 在 } A \text{ 表象中的矩阵} \\ \text{态列} & b = S^\dagger a \\ & \downarrow \quad \quad \quad \searrow \\ b = \begin{pmatrix} b_1(t) \\ b_2(t) \\ \vdots \\ b_n(t) \\ \vdots \end{pmatrix} & \quad \quad \quad a = \begin{pmatrix} a_1(t) \\ a_2(t) \\ \vdots \\ a_n(t) \\ \vdots \end{pmatrix} \end{cases} \quad \begin{cases} \Psi(x,t) = \sum_n a_n(t) \psi_n(x) \\ \Psi(x,t) = \sum_n b_n(t) \varphi_n(x) \end{cases}$$

么正变换不改变算符的本征值和矩阵的迹.

对角化么正矩阵的求法:

$$\text{例: 在 } A \text{ 表象中, } F = \begin{pmatrix} 0 & e^{i\theta} \\ e^{-i\theta} & 0 \end{pmatrix}$$

① 本征值方程:

$$\begin{pmatrix} 0 & e^{i\theta} \\ e^{-i\theta} & 0 \end{pmatrix} \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} = \lambda \begin{pmatrix} a_1 \\ a_2 \end{pmatrix}$$

② 久期方程:

$$\begin{vmatrix} -\lambda & e^{i\theta} \\ e^{-i\theta} & -\lambda \end{vmatrix} = \lambda^2 - 1 = 0$$

$$\lambda = 1 \text{ 或 } -1$$

③ 求本征函数.

$$(1) \lambda = 1 \text{ 代入得 } a_1 = e^{i\theta} a_2$$

$$\psi_1 = a_2 \begin{pmatrix} e^{i\theta} \\ 1 \end{pmatrix}$$

$$\psi_1^\dagger \psi_1 = 1$$

$$a_2 = \frac{1}{\sqrt{2}}$$

$$\psi_1 = \frac{1}{\sqrt{2}} \begin{pmatrix} e^{i\theta} \\ 1 \end{pmatrix}$$

$$(2) \lambda = -1 \text{ 代入}$$

$$\psi_2 = \frac{1}{\sqrt{2}} \begin{pmatrix} e^{i\theta} \\ -1 \end{pmatrix}$$

$$\textcircled{4} S = \begin{pmatrix} \frac{1}{\sqrt{2}} e^{i\theta} & \frac{1}{\sqrt{2}} e^{i\theta} \\ \frac{1}{\sqrt{2}} & -\frac{1}{\sqrt{2}} \end{pmatrix}$$

$$F' = S^\dagger F S = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

5. Dirac 符号

$$\langle \psi | \varphi \rangle = \int \psi^*(x) \varphi(x) dx$$

$$\text{若 } \hat{F} |n\rangle = f_n |n\rangle, \langle n | n' \rangle = \delta_{nn'}$$

$$\text{则 } |\psi\rangle = \sum_n a_n |n\rangle$$

$$a_n = \langle n | \psi \rangle \quad (|\psi\rangle \text{ 在基矢 } |n\rangle \text{ 上的投影})$$

$$|n\rangle \langle n| \text{ 为投影算符: } |n\rangle \langle n | \psi \rangle = \langle n | \psi \rangle |n\rangle = a_n |n\rangle$$

$$\sum_n |n\rangle \langle n| = 1$$

6. 占有态表象.

$$\hat{a} = \left( \frac{m\omega}{2\hbar} \right)^{1/2} \left( \hat{x} + \frac{i}{m\omega} \hat{p} \right)$$

$$\hat{a}^\dagger = \left( \frac{m\omega}{2\hbar} \right)^{1/2} \left( \hat{x} - \frac{i}{m\omega} \hat{p} \right)$$

$$[\hat{a}, \hat{a}^\dagger] = 1$$

在线性谐振子问题中.

$$\begin{cases} \hat{a} |n\rangle = \sqrt{n} |n-1\rangle \\ \hat{a}^\dagger |n\rangle = \sqrt{n+1} |n+1\rangle \end{cases}$$

$$|n\rangle = \frac{1}{\sqrt{n!}} (\hat{a}^\dagger)^n |0\rangle$$

$$\hat{N} = \hat{a}^\dagger \hat{a}$$

$$\hat{N} |n\rangle = \hat{a}^\dagger \hat{a} |n\rangle = n |n\rangle$$