

# Hamiltonian Monte Carlo



Michael Betancourt  
@betanalpha  
Symplectomorphic, LLC

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Machine Learning Summer School  
London, United Kingdom  
July 23, 2019

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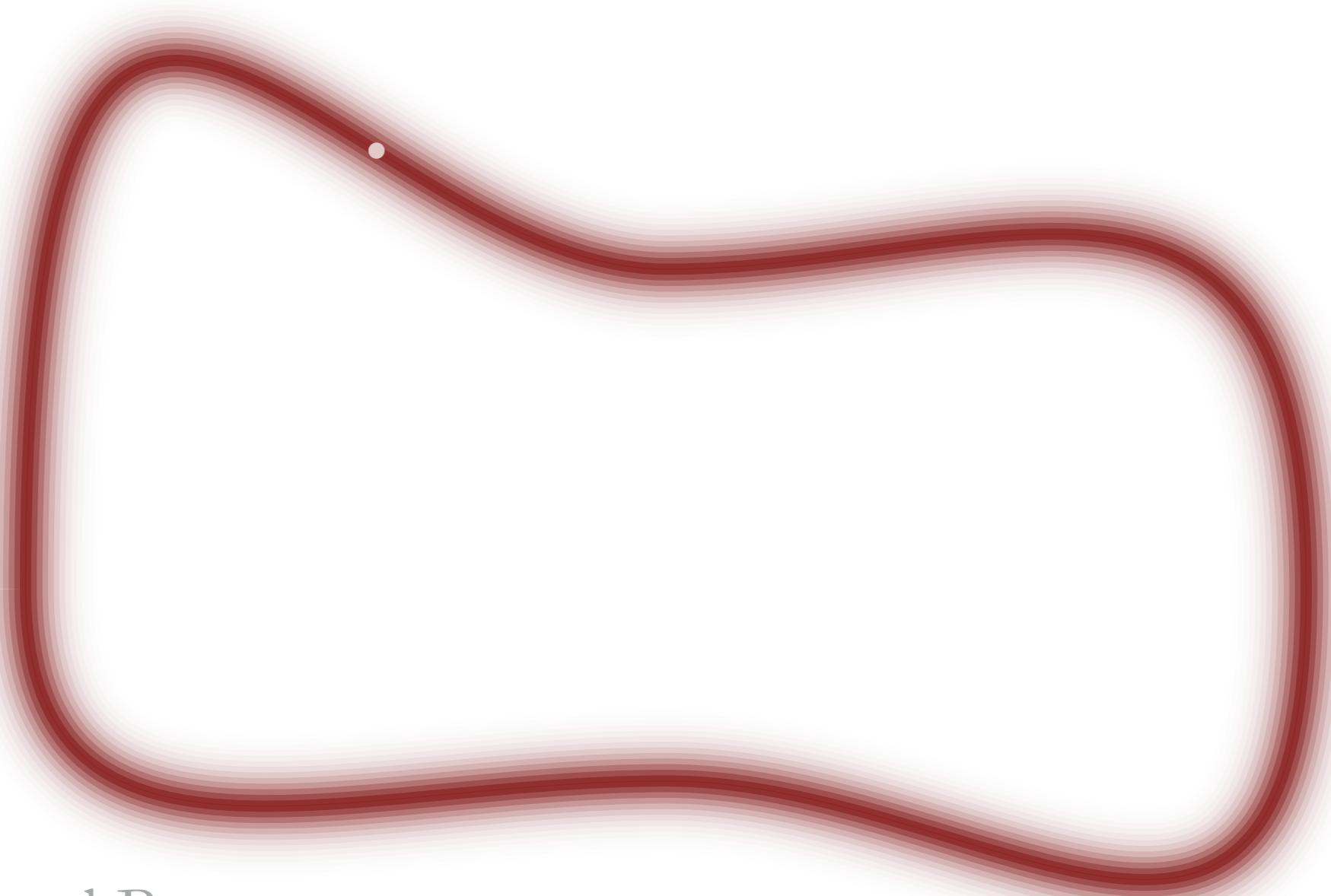


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MCMC needs *coherent* exploration of the typical set.



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Exploration of the typical set is generated by measure-preserving maps over parameter space.

$$\phi_t : Q \rightarrow Q$$

$$(\pi \circ \phi_t^{-1})(A) = \pi(A)$$

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$$\phi_0 = \text{Id}_Q$$

Together these conditions yield *measure-preserving flows* which coherently explore the typical set as desired.

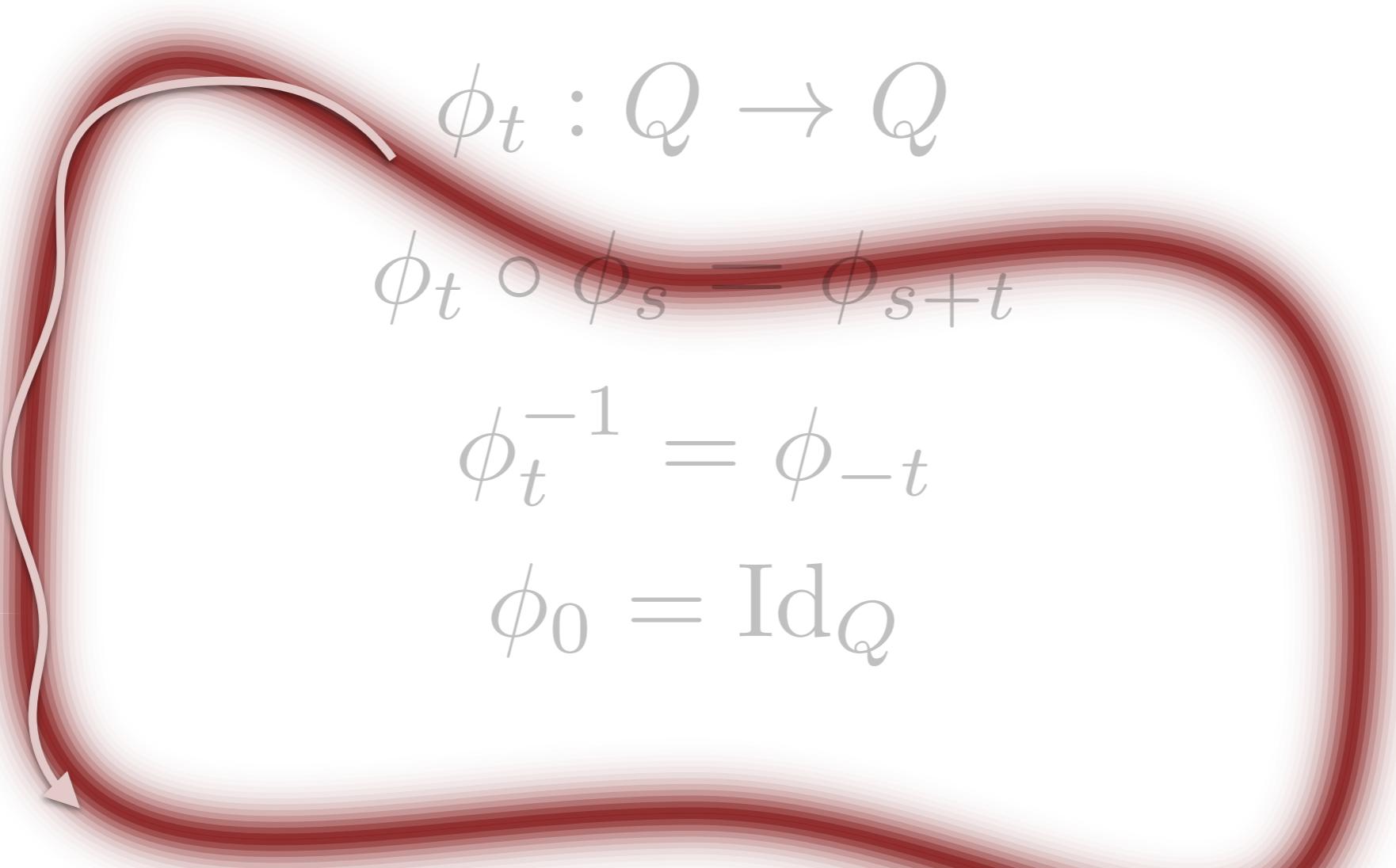
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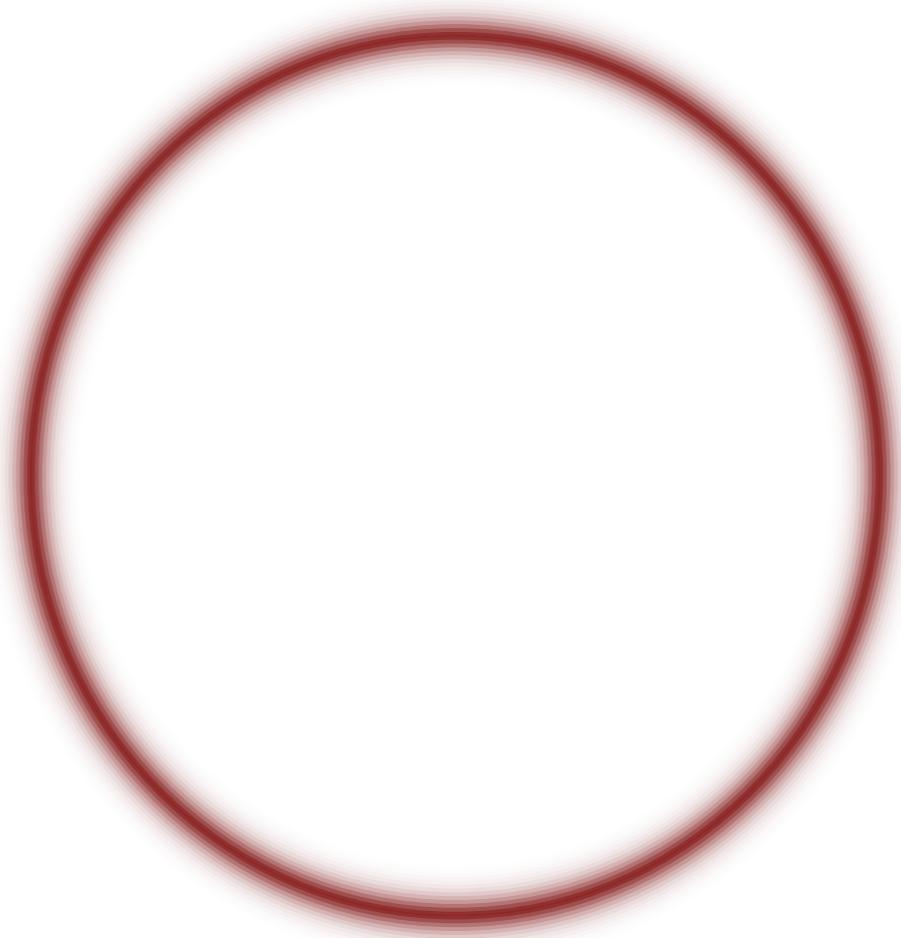
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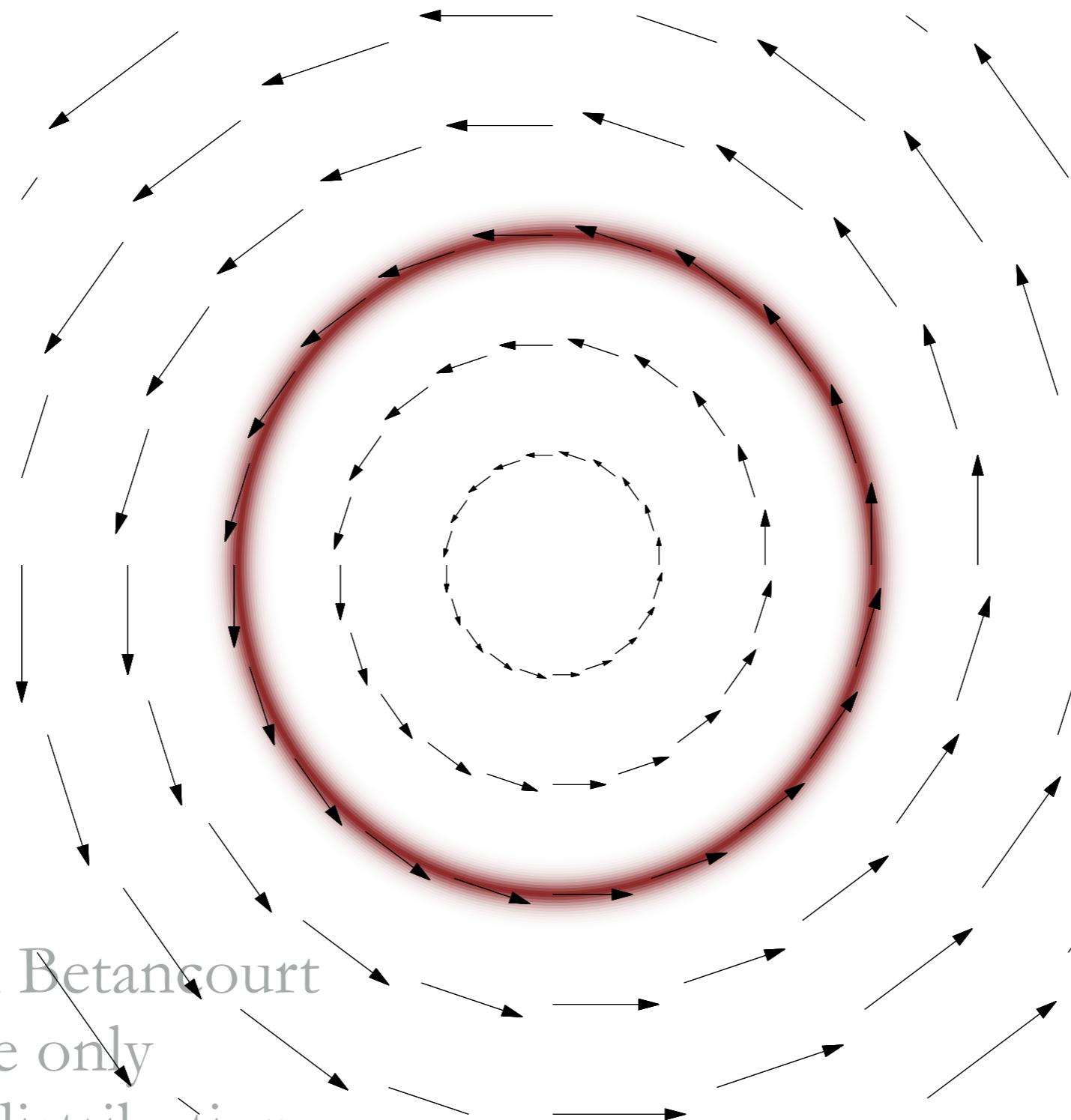
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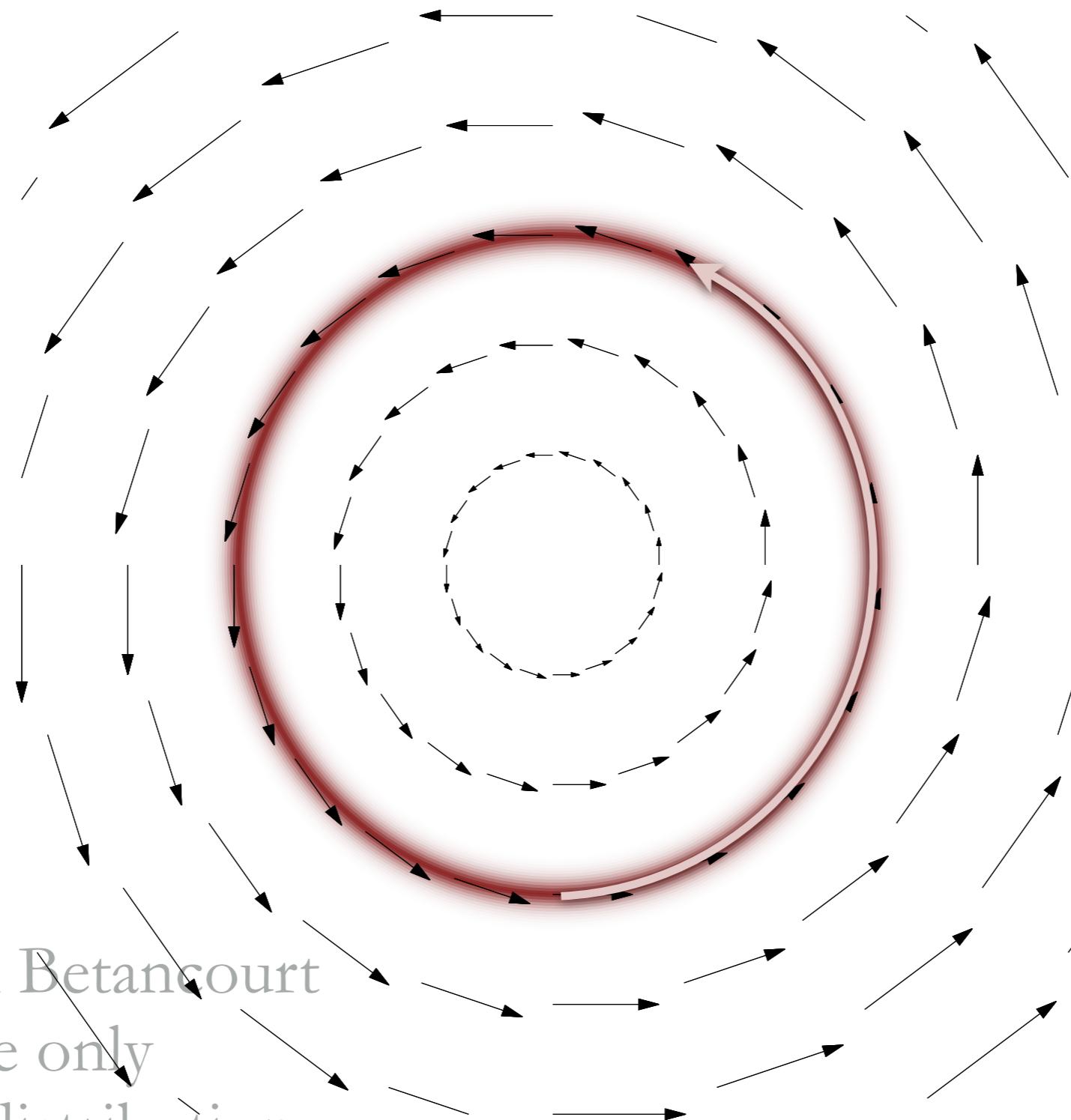
One way to construct the desired exploration is to integrate along a *vector field* aligned with the typical set.



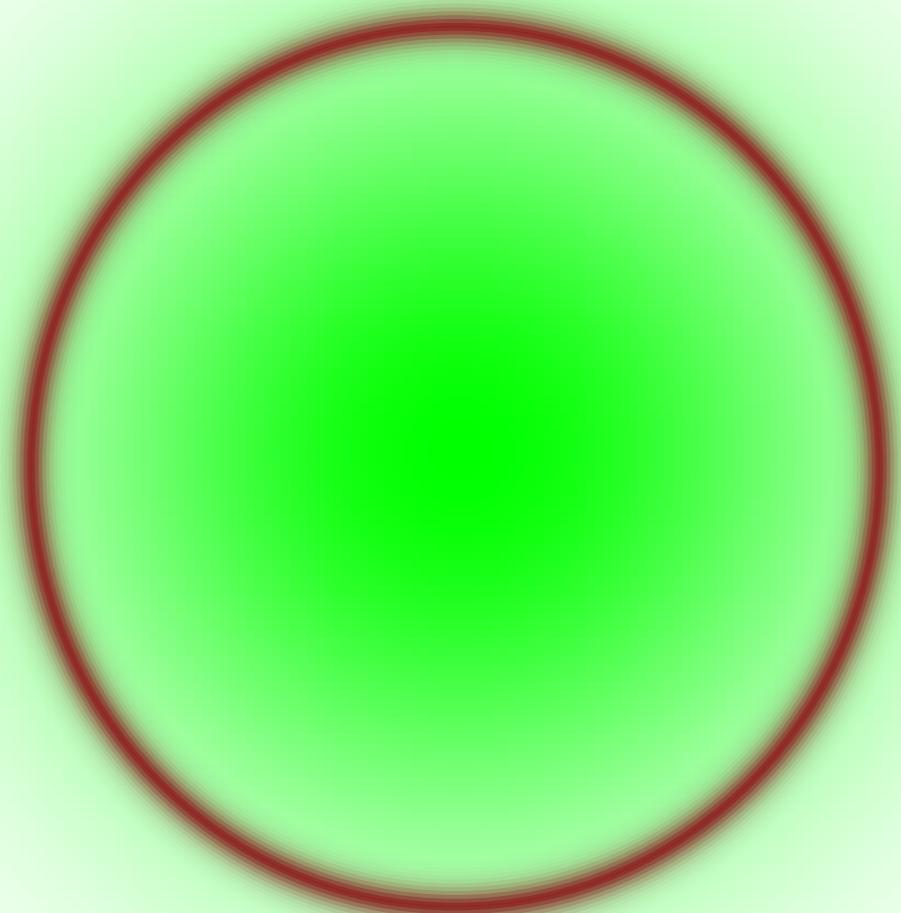
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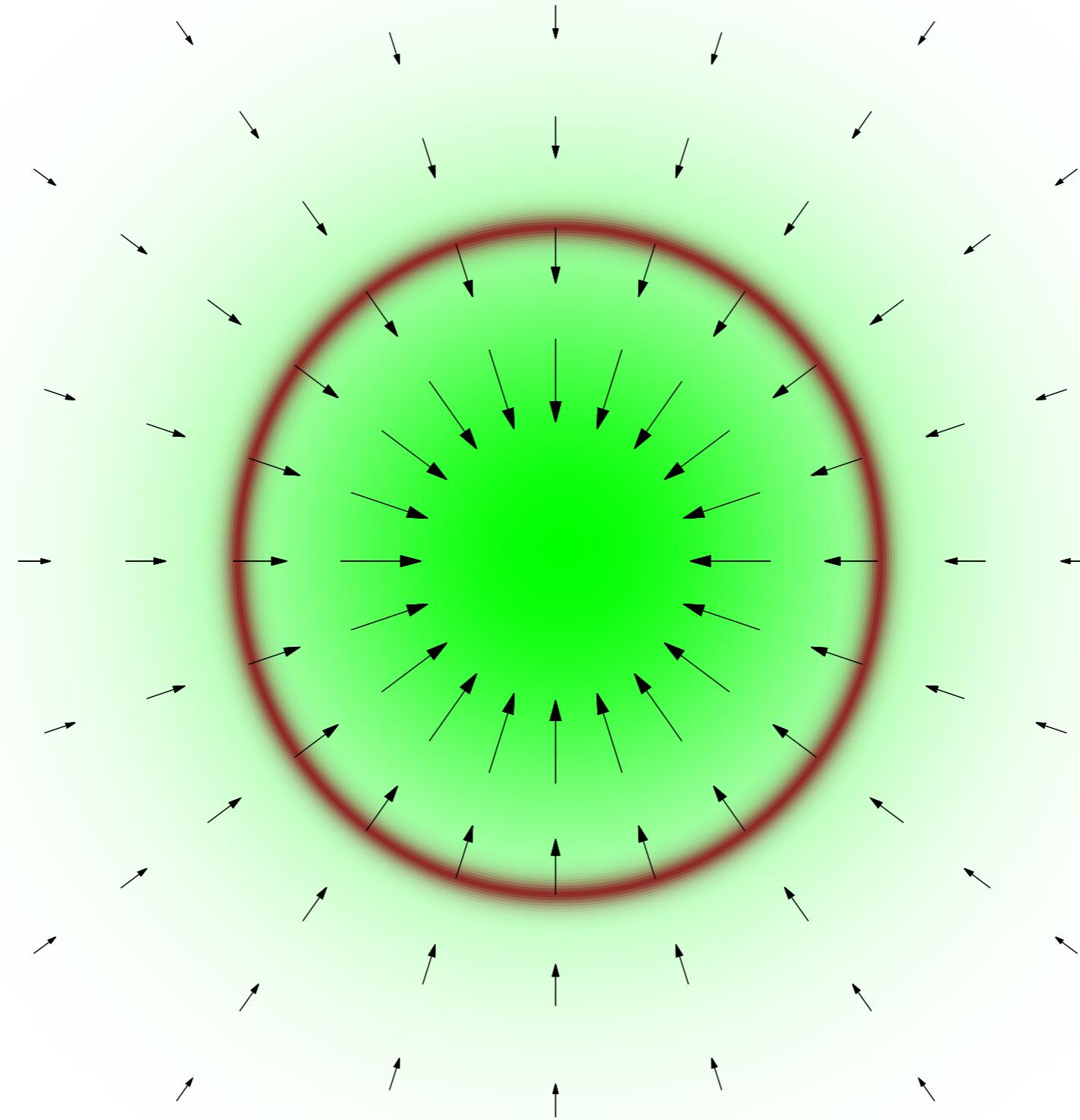


Creating the desired vector field requires transforming available vector fields, such as the gradient.



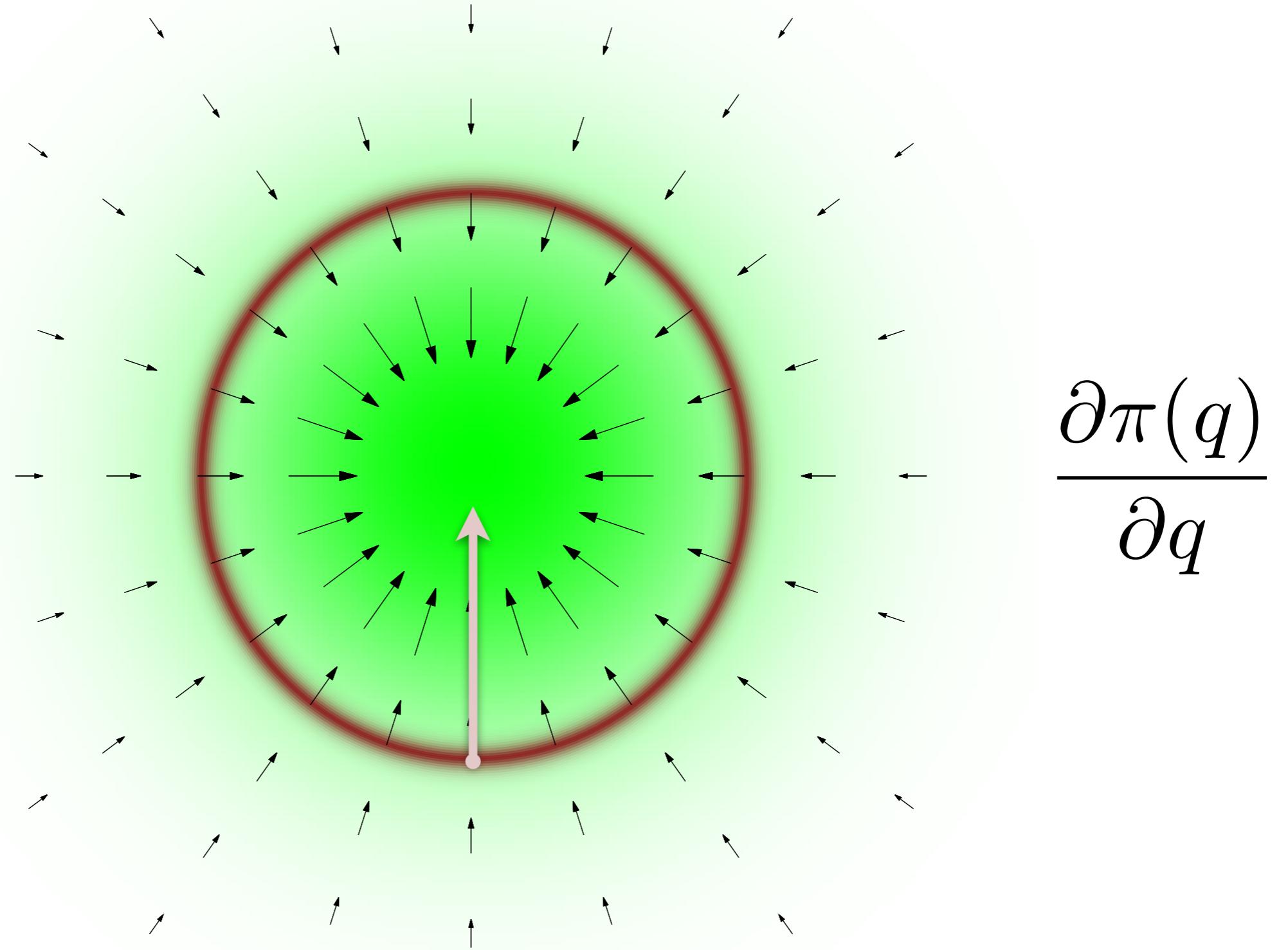
$$\pi(q)$$

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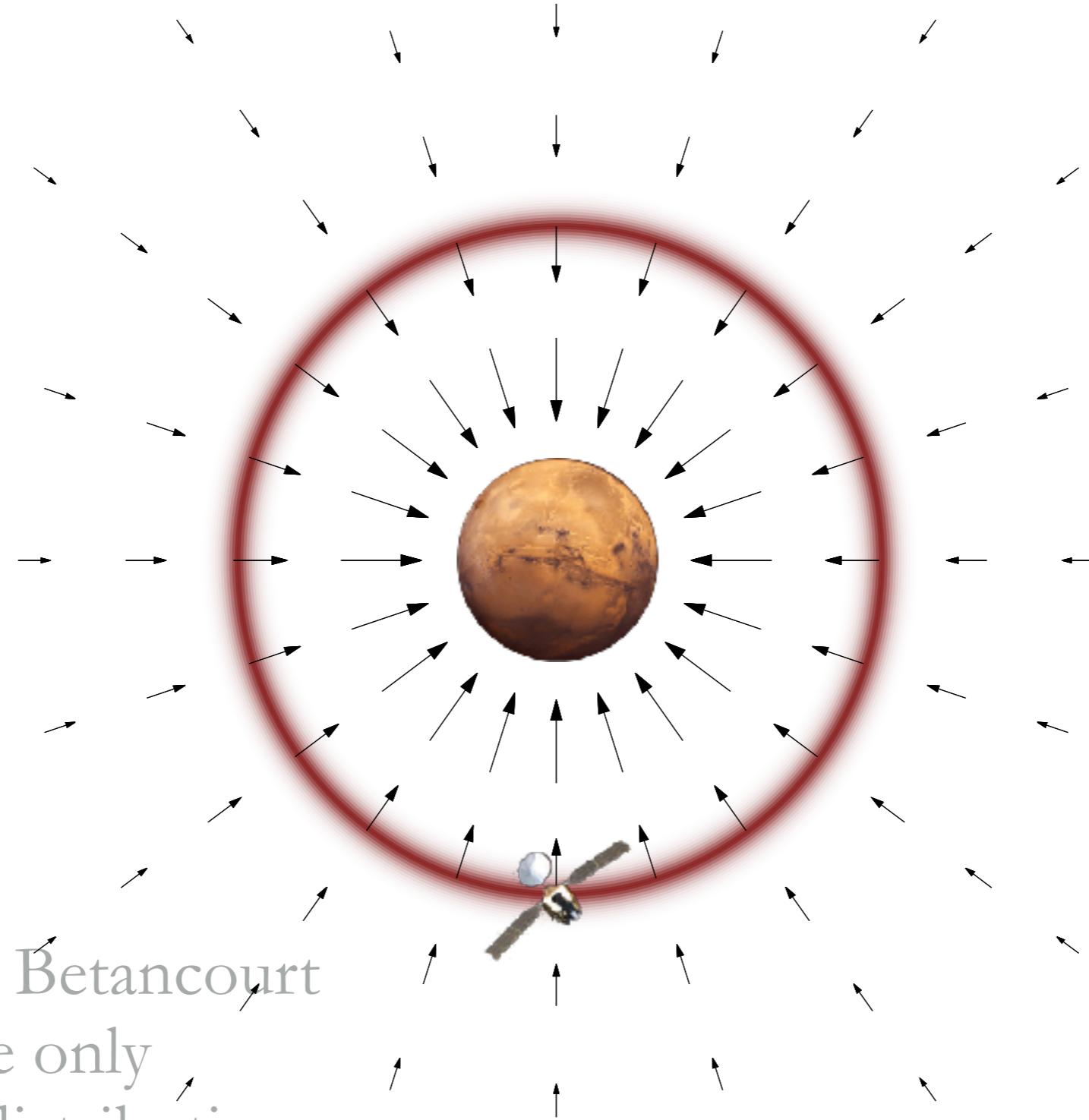


$$\frac{\partial \pi(q)}{\partial q}$$

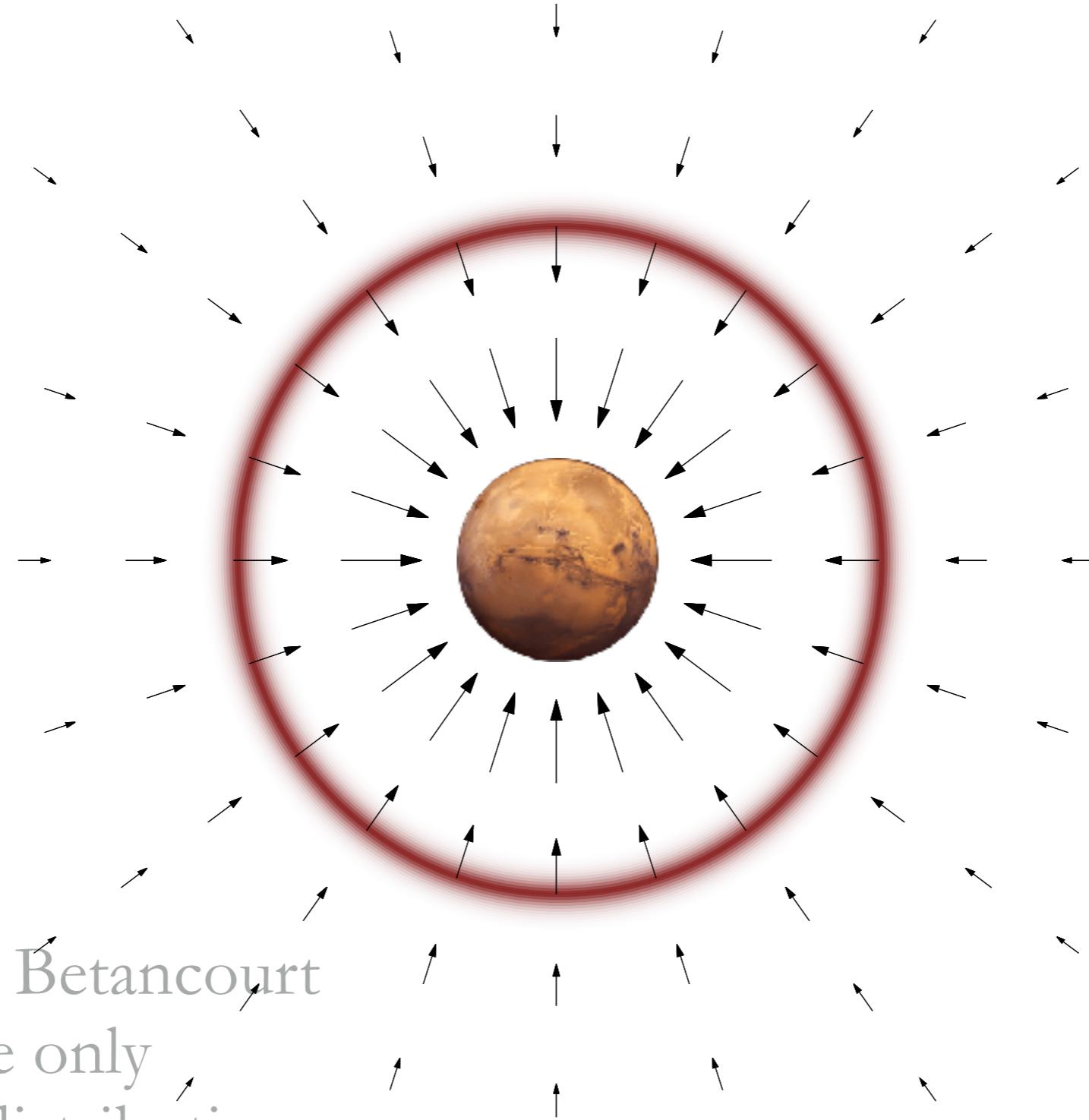
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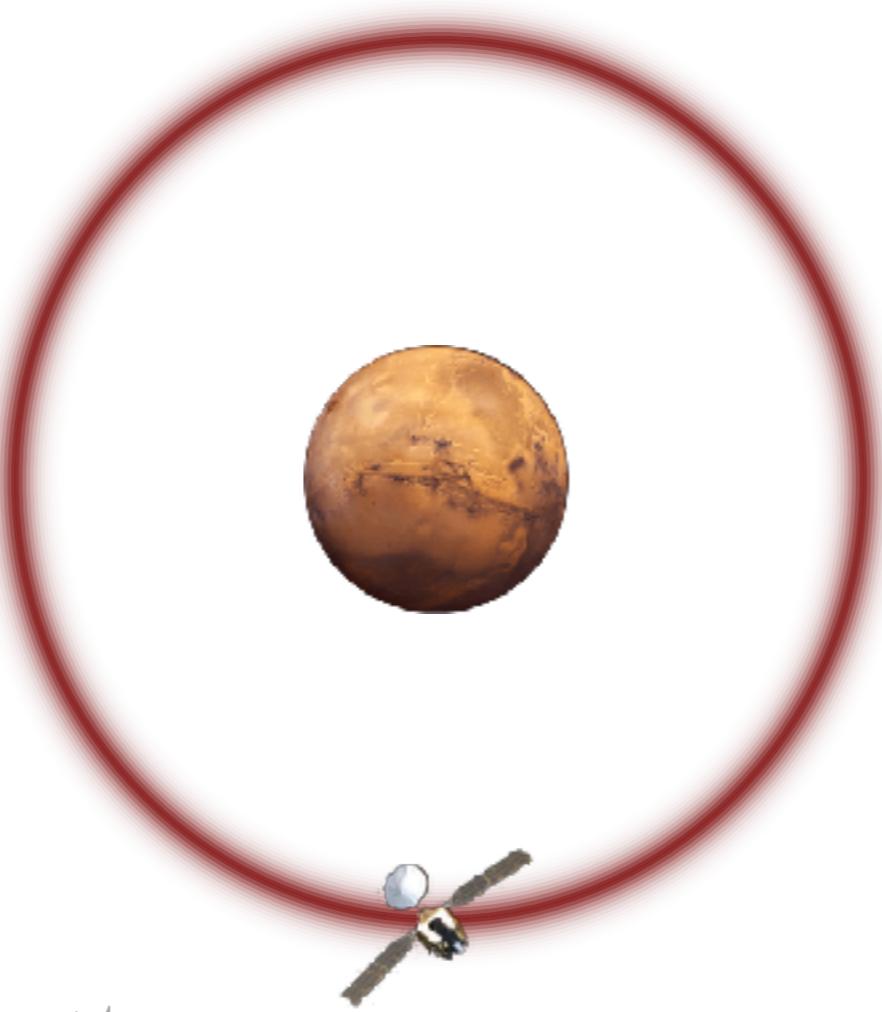
Differential geometry informs this transformation, although a physical analogy can be more intuitive.



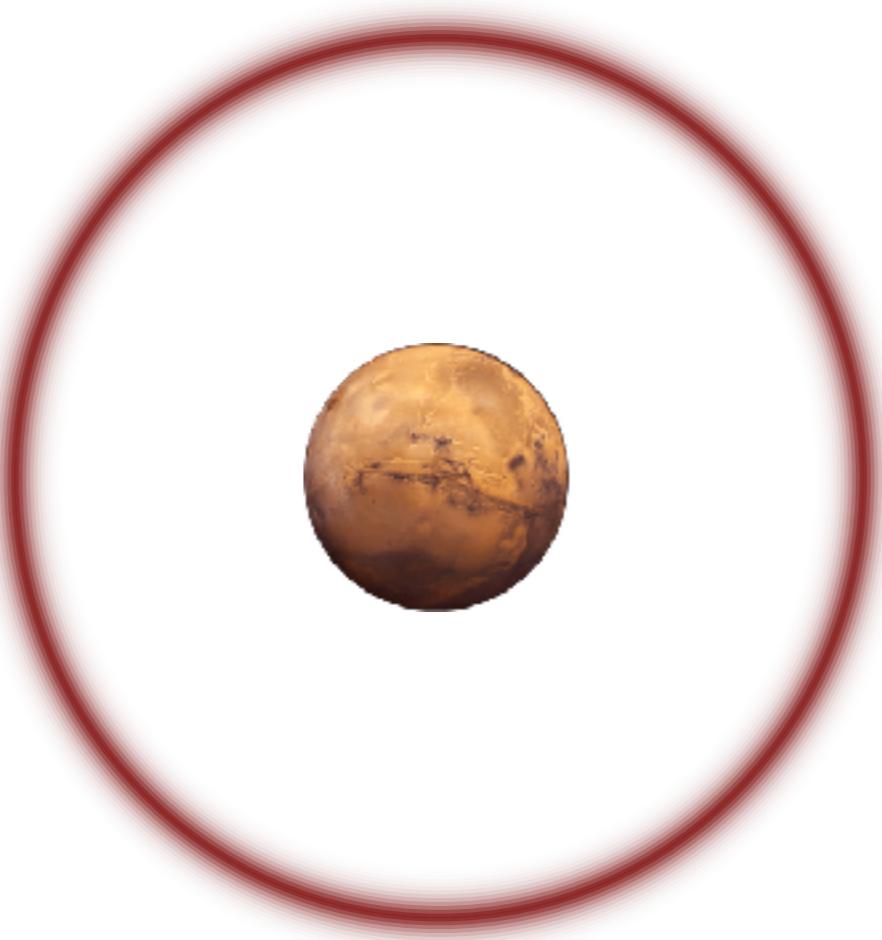
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Too little and we still crash into the planet.



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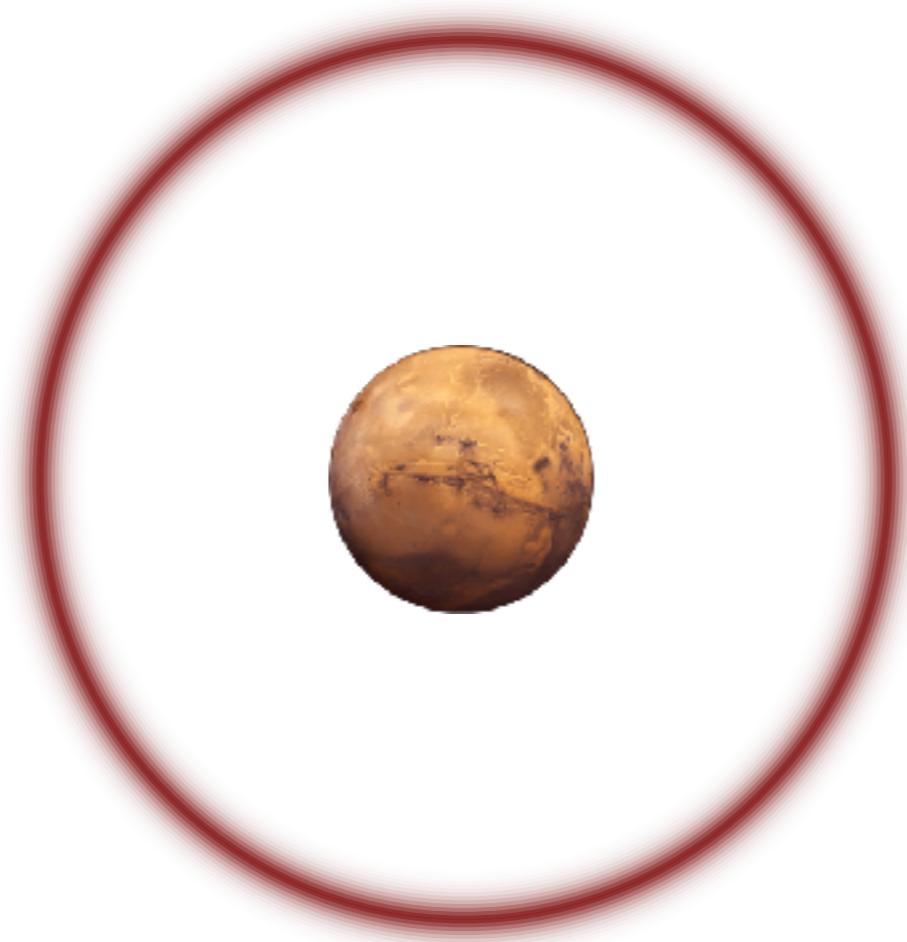


Too much and we fly off to infinity.



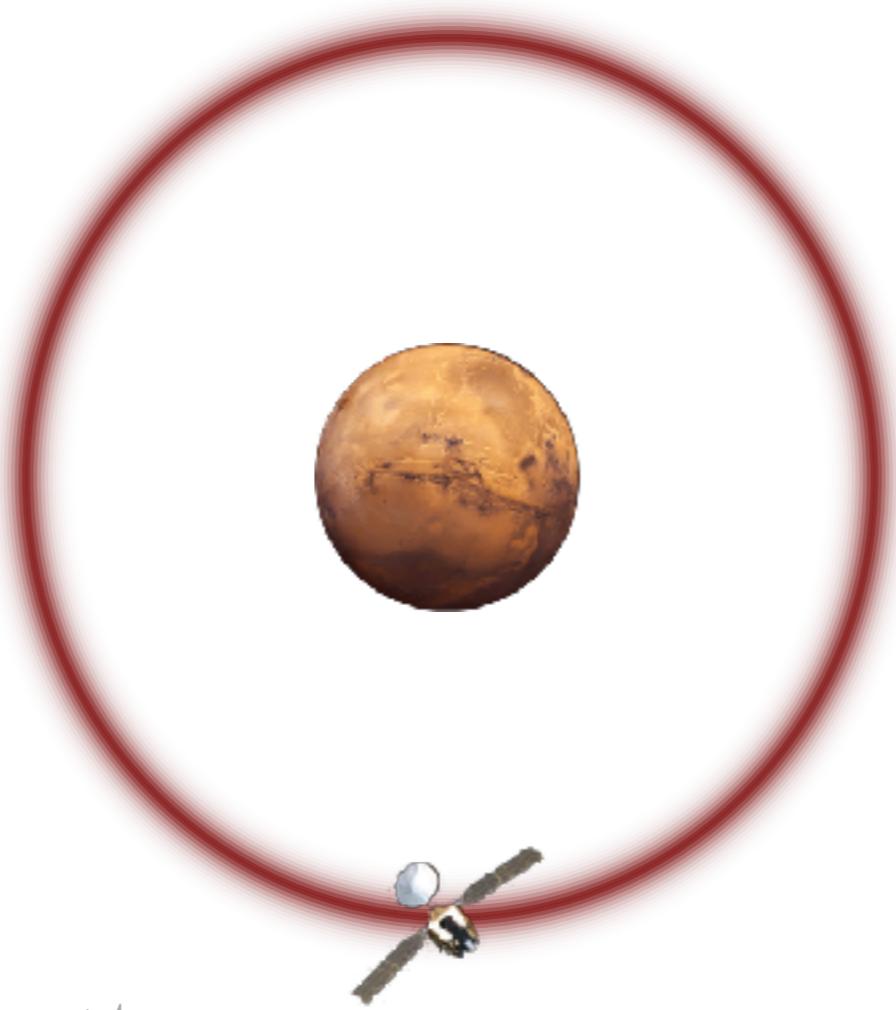
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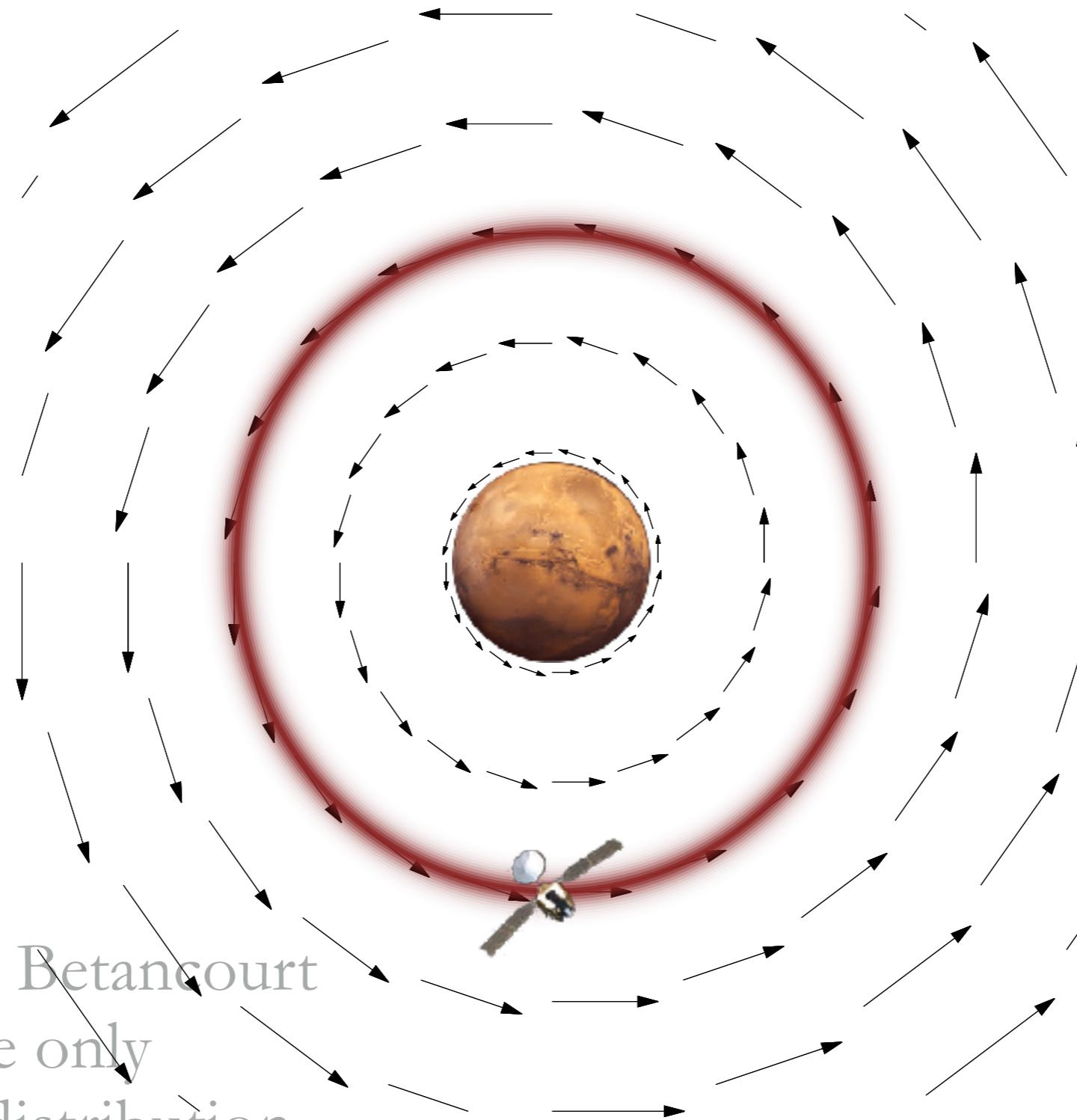


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Just enough, however, aligns the gradients with the typical set and yields the desired orbital trajectory.



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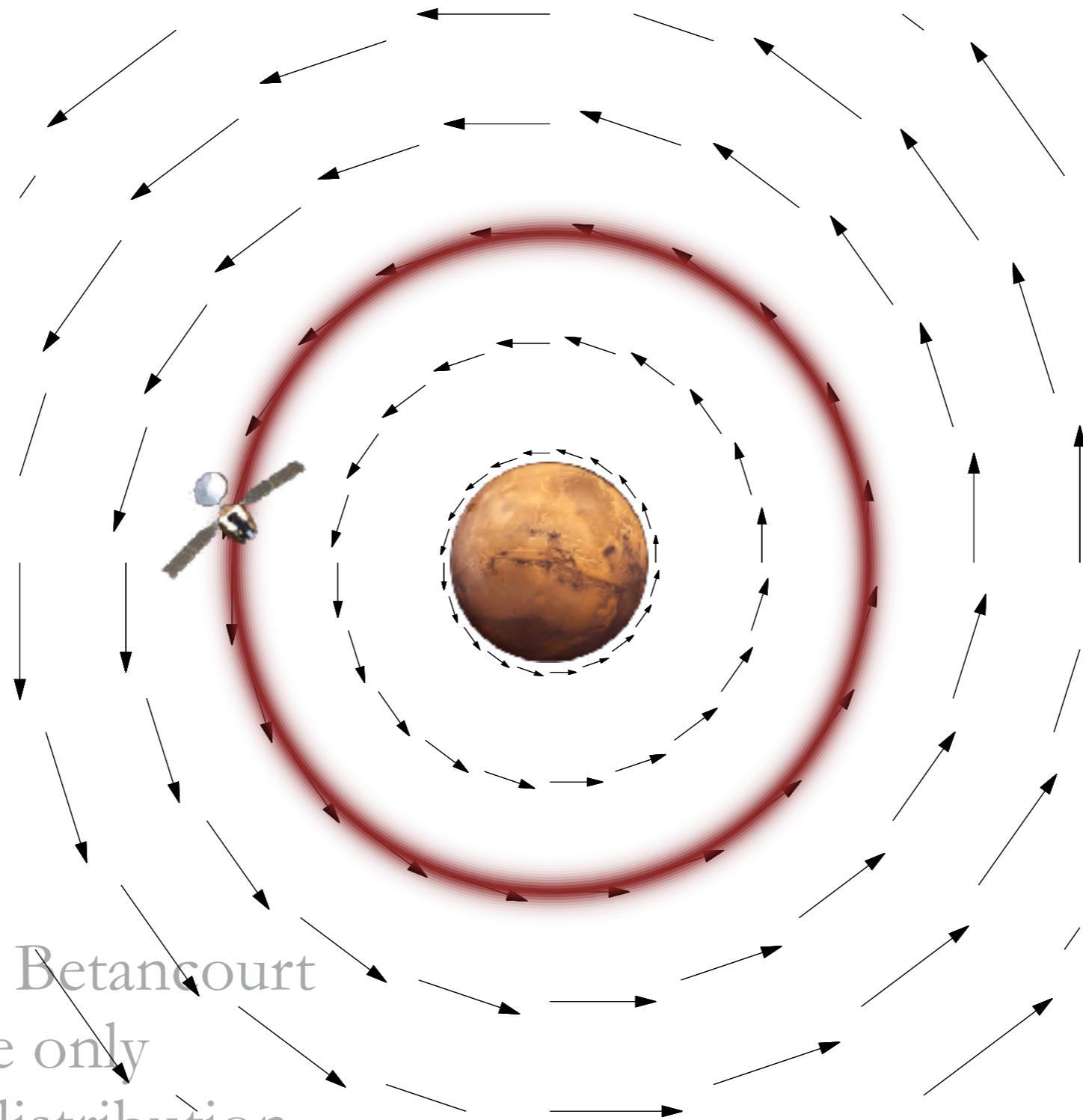


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Hamiltonian Monte Carlo is the formal procedure for adding momentum to generate coherent exploration.

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$$q \rightarrow (q, p)$$

$$\pi(q) \rightarrow \pi(q, p) = \pi(p \mid q) \pi(q)$$

Lifting the target distribution gives a *Hamiltonian* that decomposes into a *potential energy* and a *kinetic energy*.

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*Any* choice of kinetic energy generates coherent exploration through the expanded phase space.

$$\frac{dq}{dt} = \frac{\partial K}{\partial p}$$

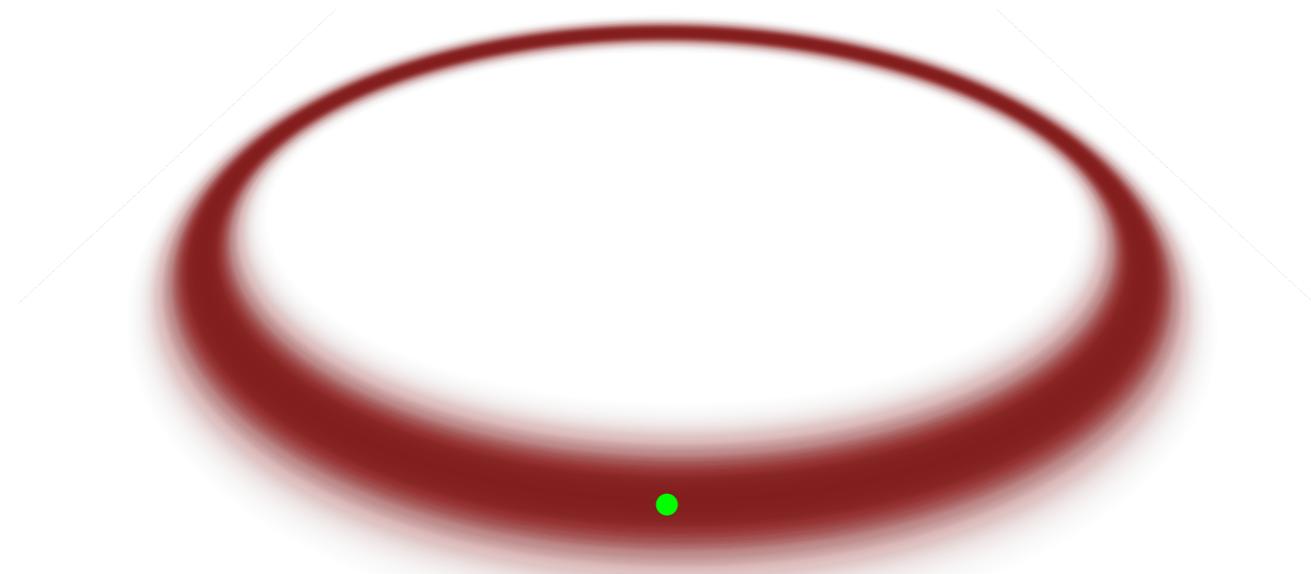
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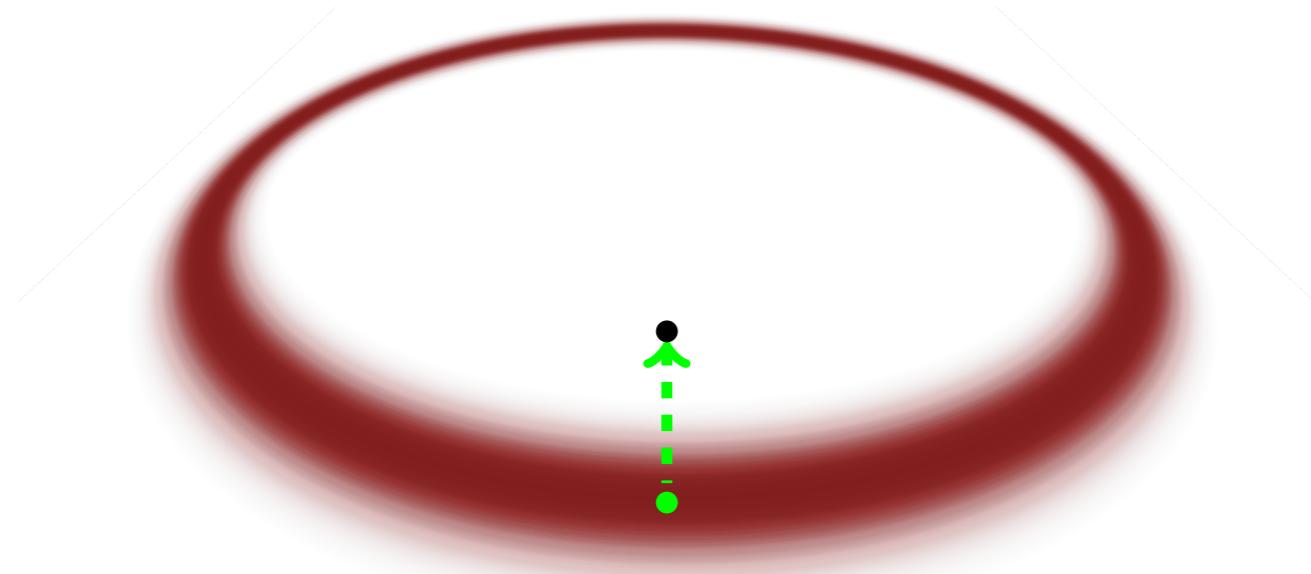
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We then construct a Markov transition by lifting into, exploring, and projecting from the expanded space.

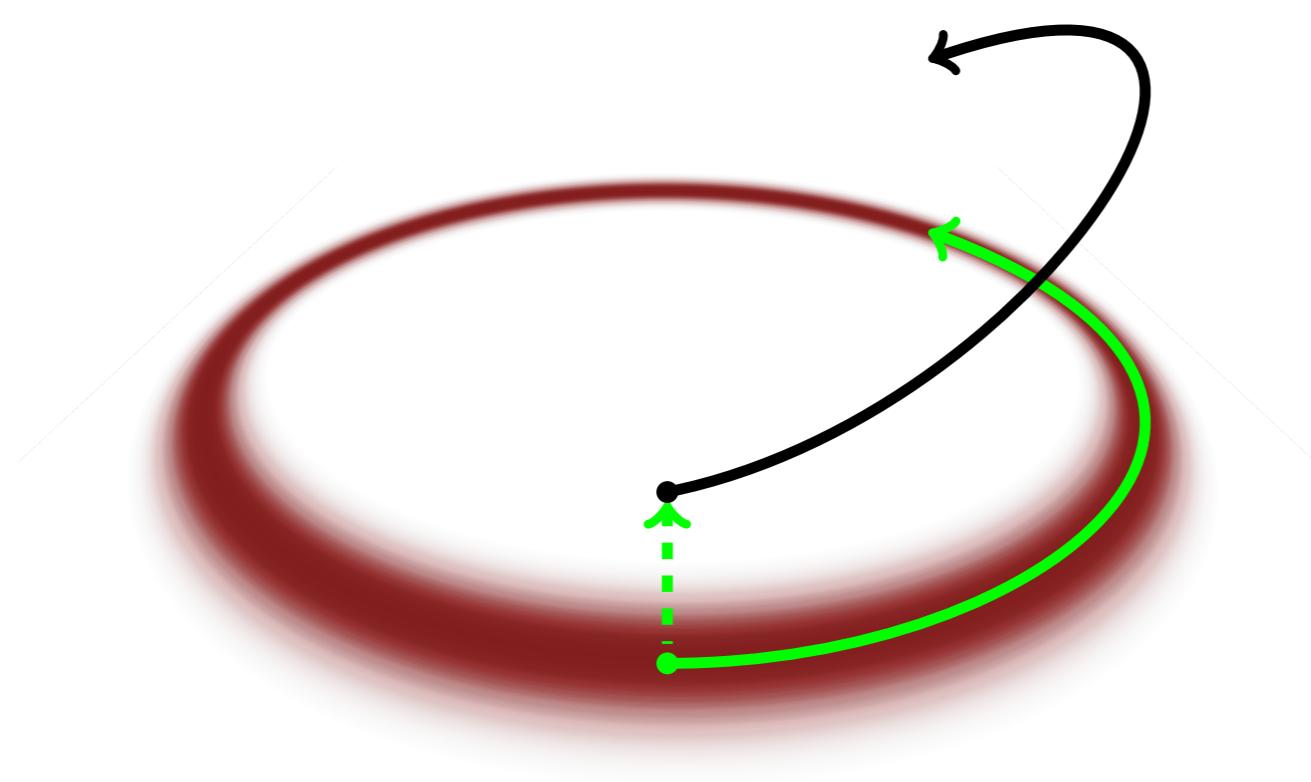


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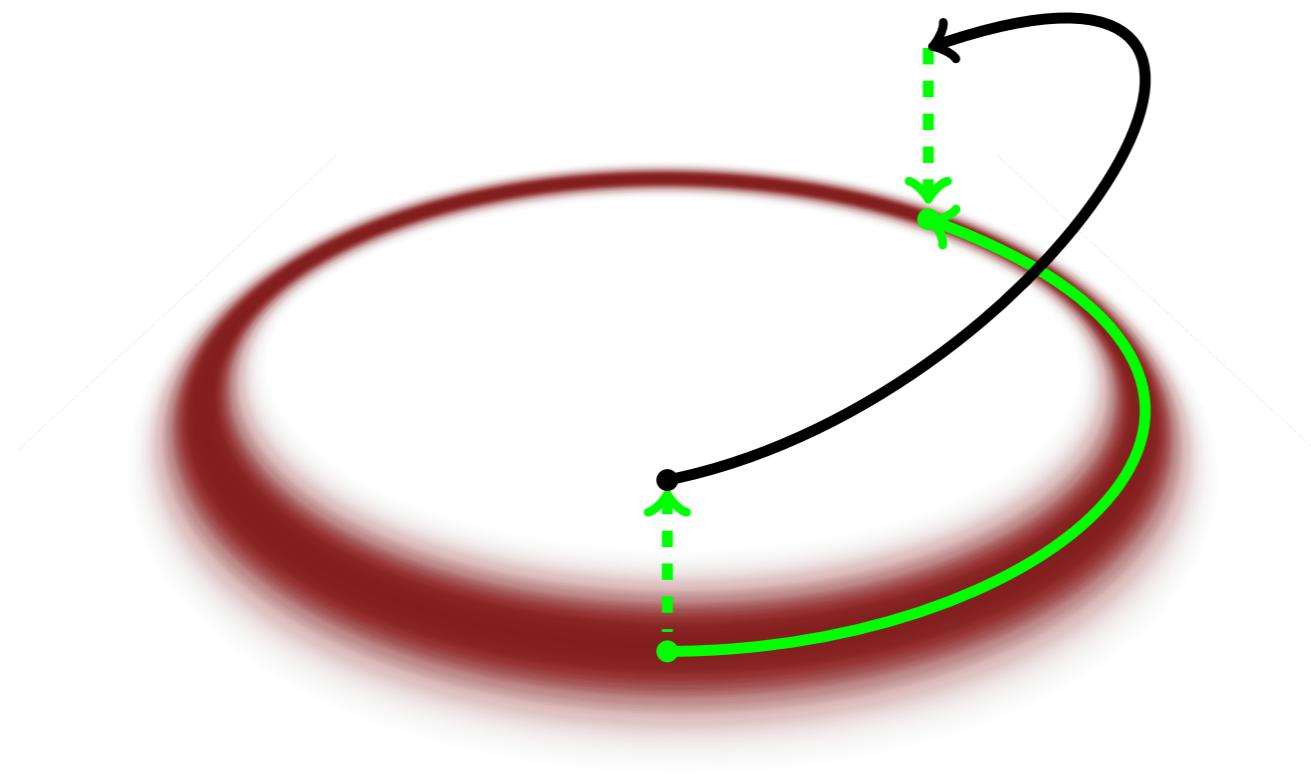


$$p \sim \pi(p \mid q)$$

$$t \sim U[0, T]$$

$$(q', p') \rightarrow \phi_t(q, p)$$

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In practice we have to approximate the exact Hamiltonian trajectories with numerical trajectories.

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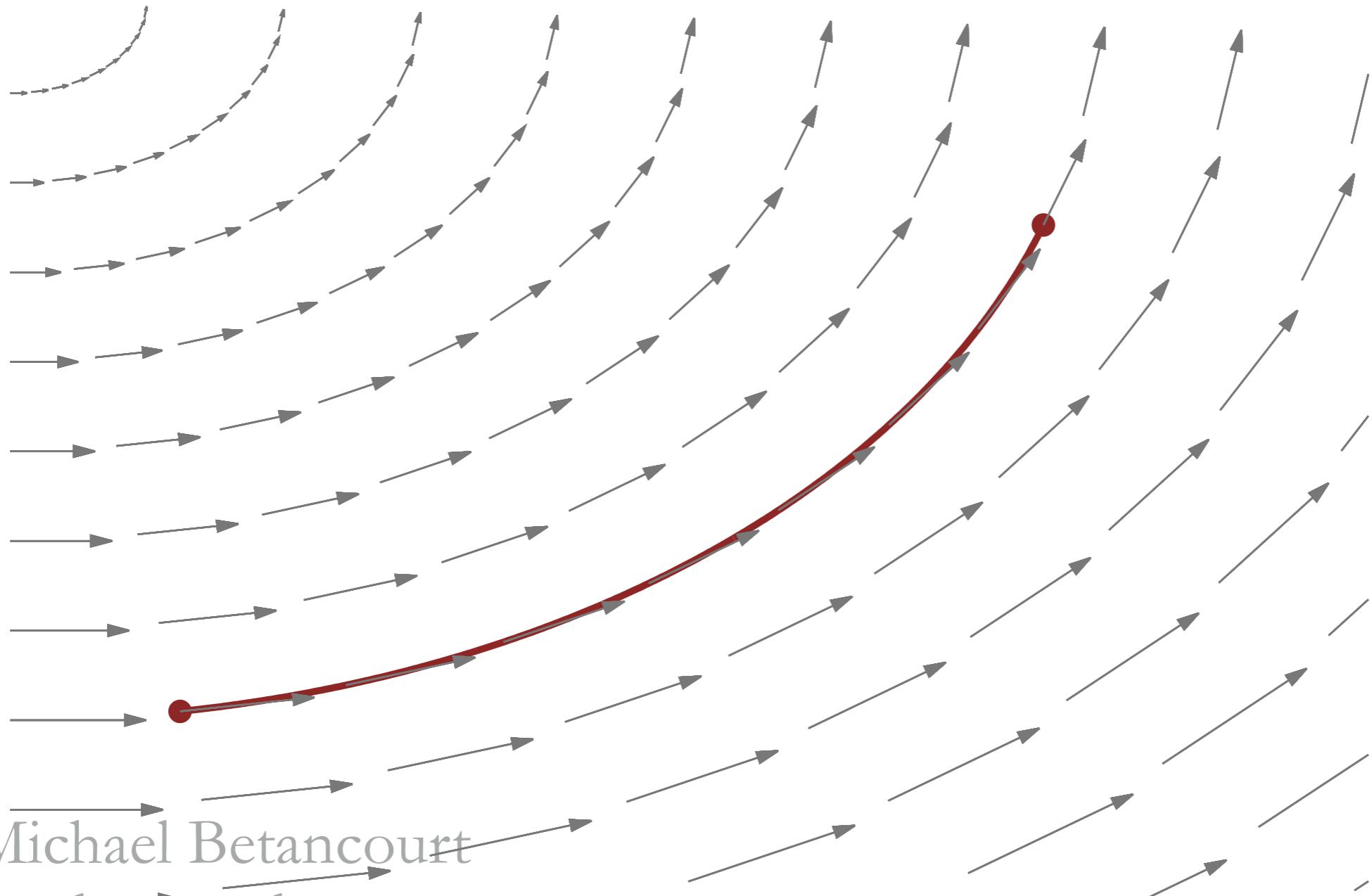
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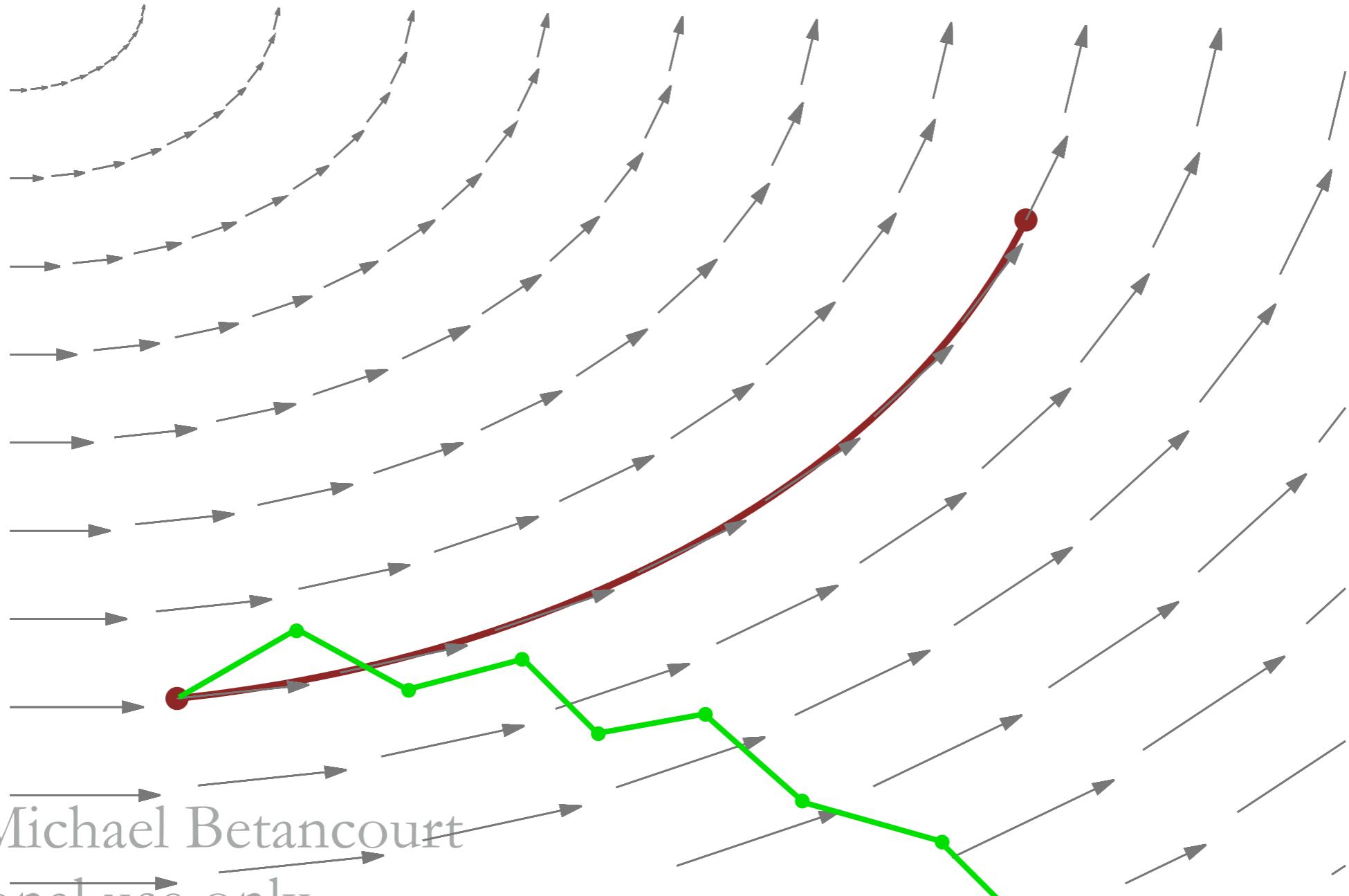
$$q \rightarrow q + \epsilon \frac{\partial K}{\partial p}$$

$$p \rightarrow p - \epsilon \left( \frac{\partial K}{\partial q} + \frac{\partial V}{\partial q} \right)$$

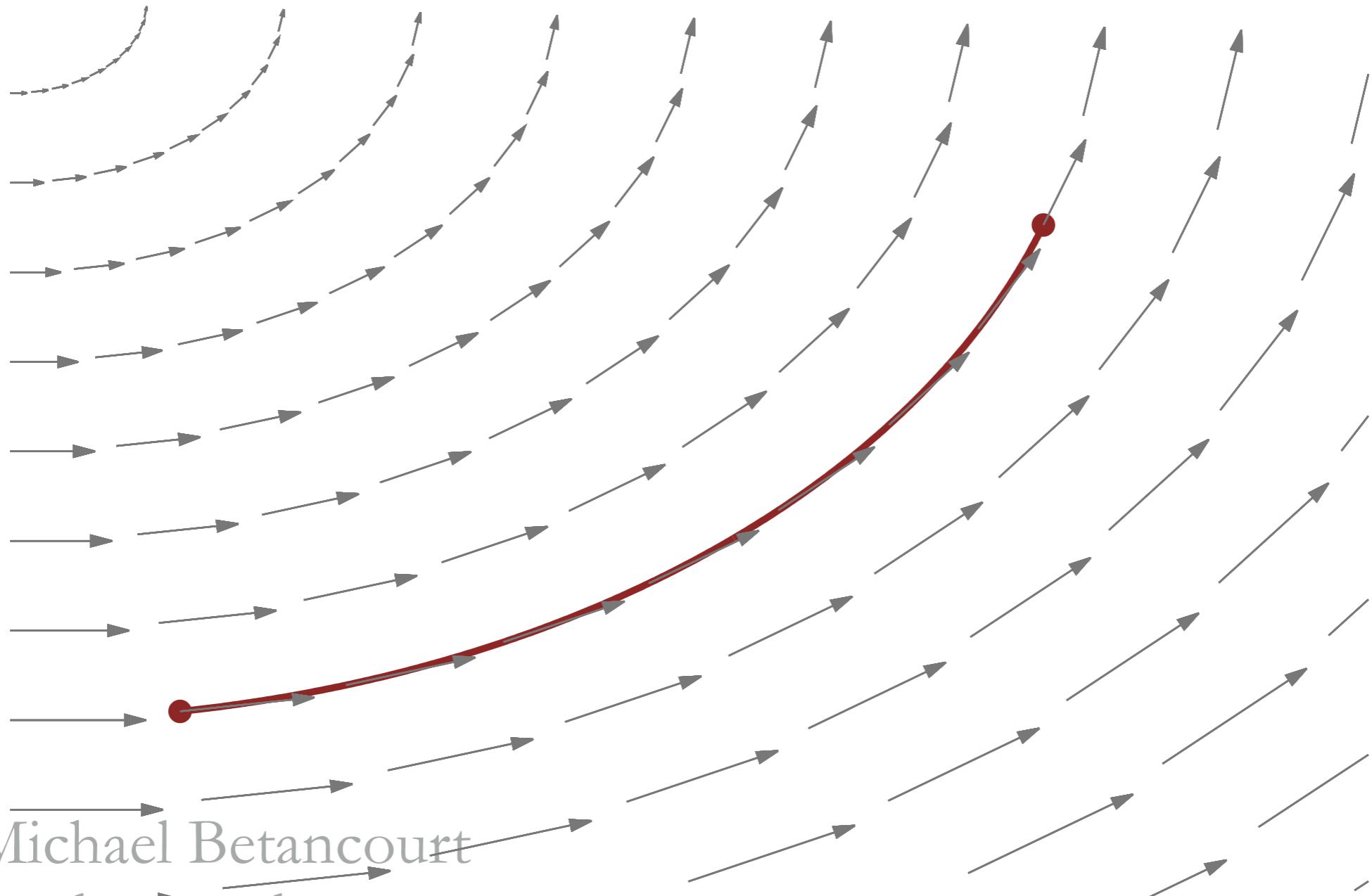
Naive numerical integrators drift away from the exact trajectory, limiting the time for which we can explore.



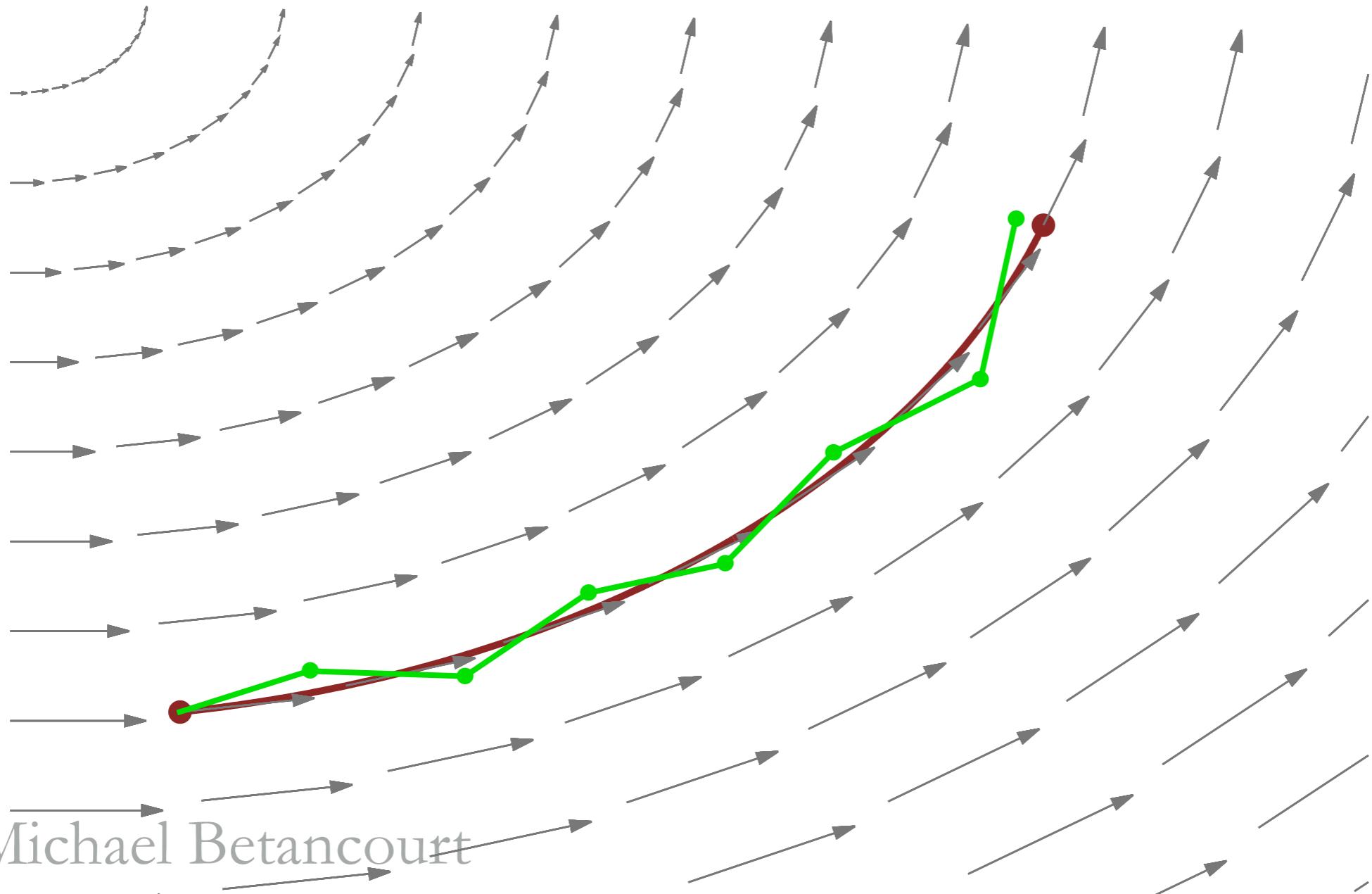
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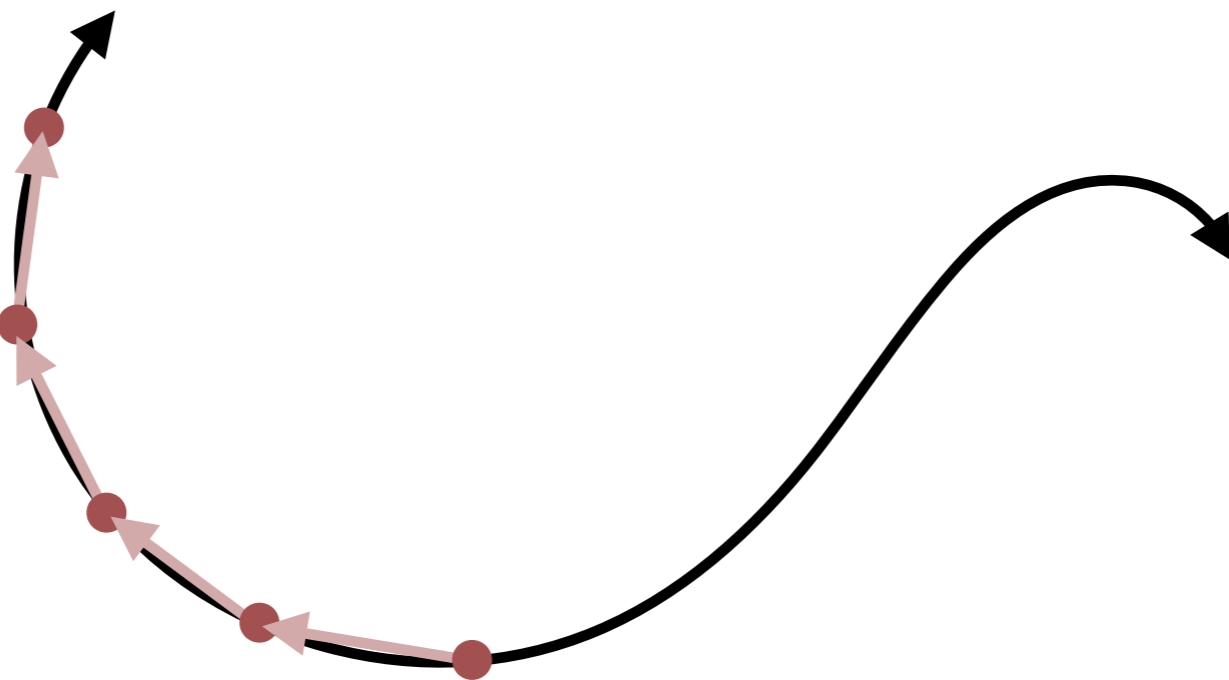
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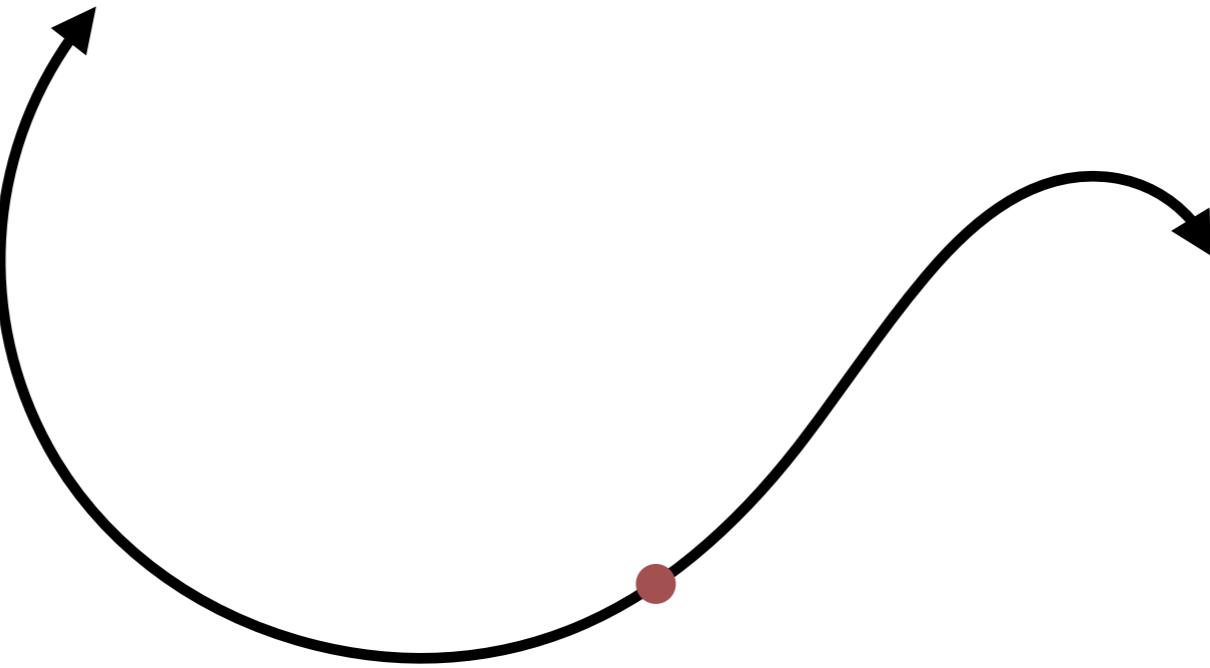
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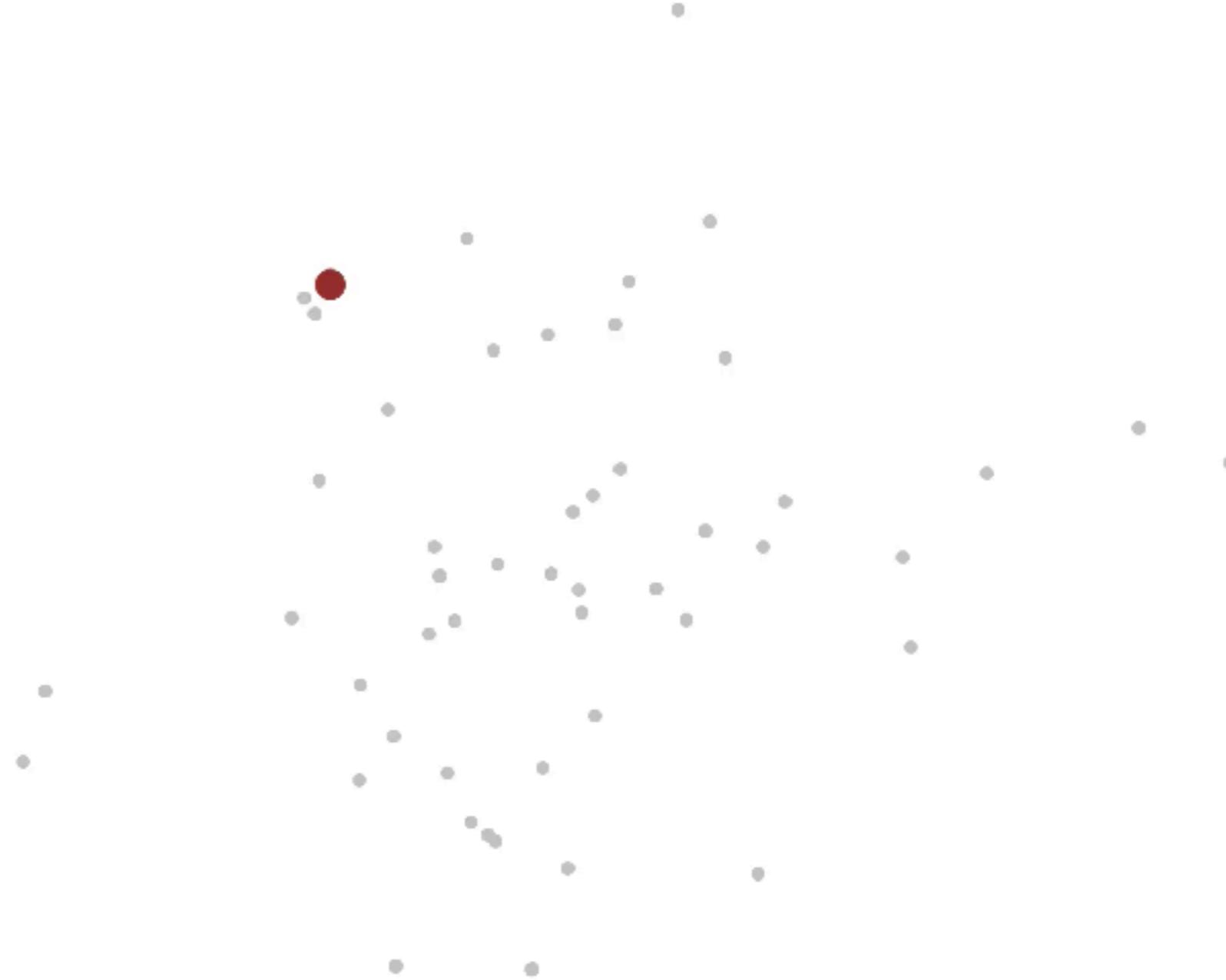
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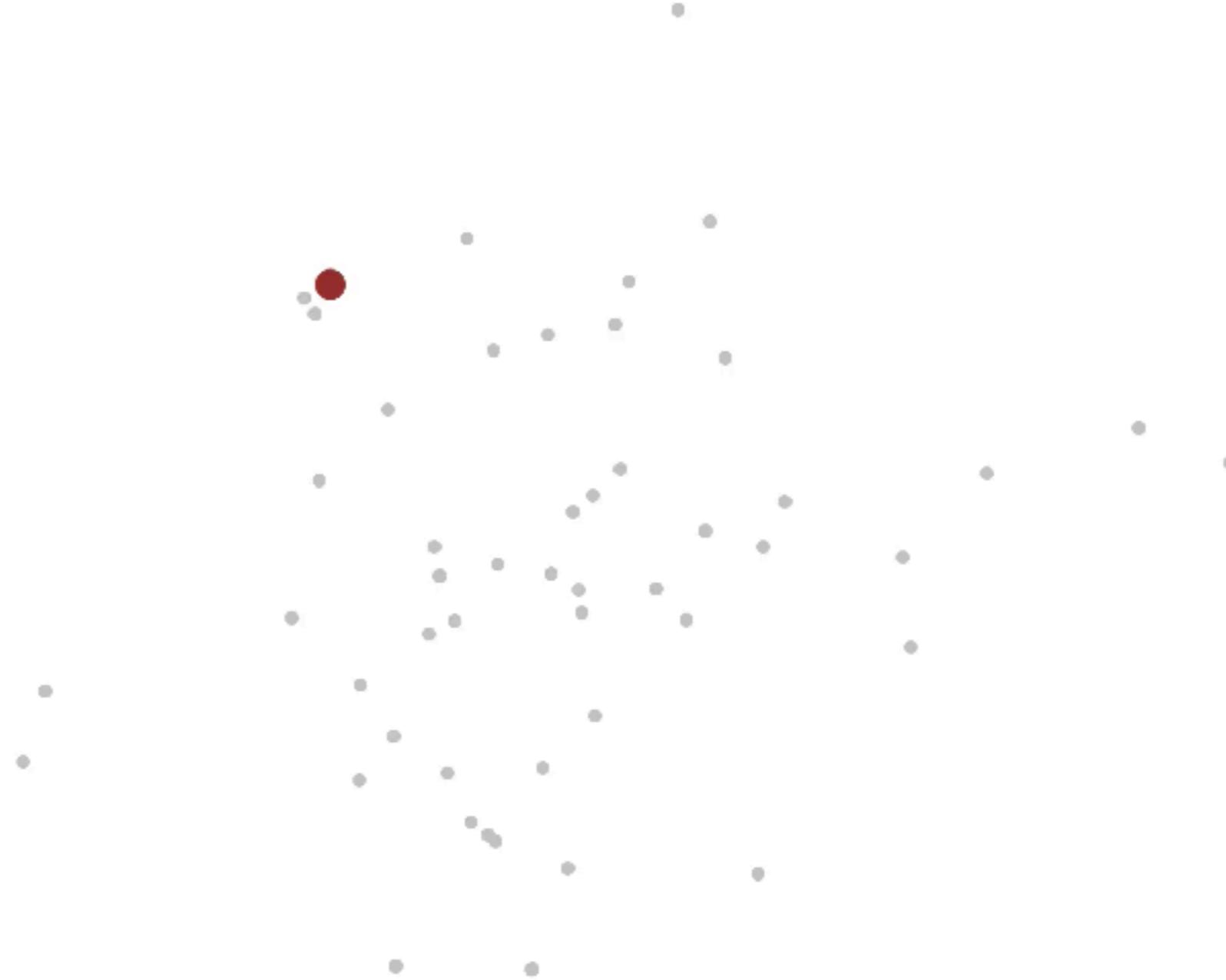
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