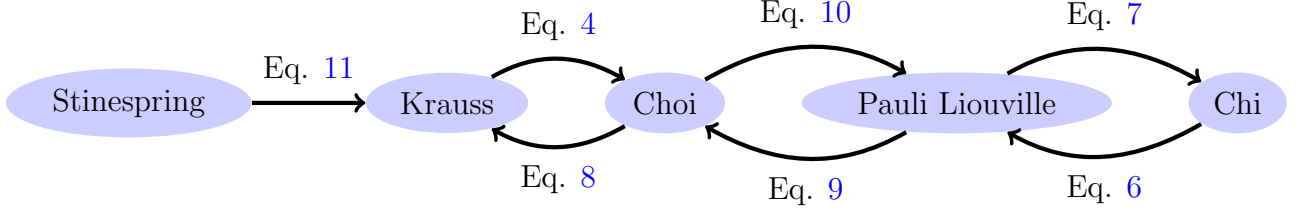


# 1 Converting between channel representations



1. Krauss operators  $\rightarrow$  Pauli Liouville matrix.

Input:  $\{K_i\}_{k=1}^r$ , where each of  $K_i$  are  $2 \times 2$  matrices. Output:  $\Gamma$ , a  $4 \times 4$  real matrix.

The entries of  $\Gamma$  are given by the following expression.

$$\Gamma_{i,j} = \text{Tr}(\mathcal{E}(P_i)P_j) \quad (1)$$

where  $\mathcal{E}$  is the CPTP map described by the input Krauss operators, as in (Eq. ??) and  $P_i$  is one of the Pauli matrices  $\{I, X, Y, Z\}$ . So, the process matrix can be expressed in terms of the Krauss operators as

$$\Gamma_{i,j} = \sum_{k=1}^r \text{Tr}(K_k P_i K_k^\dagger P_j). \quad (2)$$

2. Krauss operators  $\rightarrow$  Choi Matrix

Input:  $\{K_i\}_{k=1}^r$ , where each of  $K_i$  are  $2 \times 2$  matrices. Output:  $J$ , a  $4 \times 4$  complex matrix.

The Choi matrix of a CPTP map  $\mathcal{E}$  (specified by the input Krauss operators) is the result of applying  $\mathcal{E}$  on the first qubit of a bell state, i.e,

$$J = \frac{1}{4} \sum_{i,j=0}^1 \mathcal{E}(|i\rangle\langle j|) \otimes |i\rangle\langle j| \quad (3)$$

$$= \frac{1}{4} \sum_{i,j=0}^1 \sum_{k=1}^r K_k |i\rangle\langle j| K_k^\dagger \otimes |i\rangle\langle j| \quad (4)$$

3.  $\chi$  matrix  $\rightarrow$  Pauli Liouville Matrix

Input:  $\chi$ , a  $4 \times 4$  complex matrix. Output:  $\Gamma$ , a  $4 \times 4$  real matrix.

The action of a CPTP map  $\mathcal{E}$  on an input state  $\rho$  is specified using its  $\chi$  matrix as

$$\mathcal{E}(\rho) = \sum_{l,m} \chi_{l,m} P_l \rho P_m. \quad (5)$$

Using the definition of the process matrix  $\Gamma$ , corresponding to the CPTP map  $\mathcal{E}$ , in (Eq. 1), we have

$$\Gamma_{i,j} = \sum_{l,m} \text{Tr}(P_l P_i P_m P_j) \chi_{l,m}.$$

Defining a  $16 \times 16$  matrix  $\Omega$ , such that  $\Omega_{4 \times i + j, 4 \times l + m} = \text{Tr}(P_l P_i P_m P_j)$ , we have

$$|\Gamma\rangle\rangle = \Omega|\chi\rangle\rangle, \quad (6)$$

where  $|\Gamma\rangle\rangle$  and  $|\chi\rangle\rangle$  are the (column) vectorized forms of  $\Gamma$  and  $\chi$  respectively.

4. Pauli Liouville Matrix  $\rightarrow \chi$  matrix.

We can easily obtain  $|\chi\rangle\rangle$  by inverting the relation in Eq. 6, giving

$$|\chi\rangle\rangle = \Omega^{-1}|\Gamma\rangle\rangle. \quad (7)$$

5. Choi Matrix  $\rightarrow$  Krauss operators.

Input:  $J$ , a  $4 \times 4$  complex matrix. Output:  $\{K_i\}_{i=1}^r$ , where each of  $K_i$  are  $2 \times 2$  matrices.

Eq. 4 can be expressed as

$$J = \sum_{k=1}^r |K_k\rangle\rangle\langle\langle K_k|, \quad (8)$$

where  $|K_k\rangle\rangle$  is the (column) vectorized form of the Krauss operator  $\frac{1}{2}K_k$ . Performing a singular value decomposition of  $J$ , we find  $J = \sum_k \lambda_k |\lambda_k\rangle\rangle\langle\langle \lambda_k|$ . Hence the Krauss operator  $K_k$  can be constructed from un-vectorizing  $2|\lambda_k\rangle\rangle$ .

6. Pauli Liouville Matrix  $\rightarrow$  Choi Matrix

Input:  $\Gamma$ , a  $4 \times 4$  real matrix. Output:  $J$ , a  $4 \times 4$  complex matrix.

Combining the expression for the bell state in the Pauli basis,

$$\frac{1}{4}(|00\rangle + |11\rangle)(\langle 00| + \langle 11|) = \frac{1}{4} \sum_{i=1}^4 P_i \otimes P_i^T,$$

where  $P_i \in \{I, X, Y, Z\}$  is a Pauli matrix, with Eq. 3, we find

$$J = \frac{1}{4} \sum_i \mathcal{E}(P_i) \otimes P_i^T.$$

Applying the definition of the Pauli Liouville representation of  $\mathcal{E}$  in Eq. 1, we find

$$J = \sum_{i,j} \Gamma_{i,j} P_j \otimes P_i^T. \quad (9)$$

7. Choi Matrix  $\rightarrow$  Pauli Liouville Matrix.

Input:  $J$ , a  $4 \times 4$  complex matrix. Output:  $\Gamma$ , a  $4 \times 4$  real matrix.

Multiplying on either sides of Eq. 9 by  $P_a \otimes P_b^T$ , we find

$$J \cdot (P_a \otimes P_b^T) = \sum_{i,j} \Gamma_{i,j} P_j P_a \otimes P_i^T P_b^T.$$

Taking the trace of the above equation yields

$$\Gamma_{b,a} = \text{Tr}(J \cdot (P_a \otimes P_b^T)). \quad (10)$$

8. Strinespring dilation  $\rightarrow$  Krauss operators.

The action of any CPTP map  $\mathcal{E}$  on an input state  $\rho$  can be derived from unitary dynamics  $U$  on a larger Hilbert space, that comprises of the system  $\rho$  and its environment, which is initially in some state  $|0\rangle_E\langle 0|_E$ . Under the unitary evolution described by  $U$ , we find

$$\rho \otimes |0\rangle_E\langle 0|_E \rightarrow U [\rho \otimes |0\rangle_E\langle 0|_E] U^\dagger.$$

Since the state of the environment is essentially unknown, we must discard it resulting in a partial trace over the environment basis, denoted by  $\{|e_k\rangle_E\}$ , resulting in

$$\begin{aligned} \rho \otimes |0\rangle_E\langle 0|_E &\rightarrow \text{Tr}_E (U [\rho \otimes |0\rangle_E\langle 0|_E] U^\dagger) \\ &= \rho_{ab} \sum_{k=1}^r \langle e_k | U | e_0 \rangle \rho \langle e_0 | U^\dagger | e_k \rangle. \end{aligned} \quad (11)$$

The above expression is the Krauss representation of the channel where the Krauss operators are  $\{\langle e_k | U | e_0 \rangle\}_{k=1}^r$ .

```

Pavithrans-MacBook-Pro:chflow pavithran$ ./chFlow.sh
>> chan rand 0.2
Note: the current channel is in the "krauss" representation.
>> chrep choi
>> chprint
Choi representation
[[ 0.453+0.j    0.045-0.082j -0.063-0.093j  0.352+0.034j]
 [ 0.045+0.082j  0.098+0.j    0.025-0.031j -0.036+0.082j]
 [-0.063+0.093j  0.025+0.031j  0.047+0.j   -0.045+0.082j]
 [ 0.352-0.034j -0.036-0.082j -0.045-0.082j  0.402+0.j   ]]
XXXXXXXX
>> chrep krauss
>> chprint
Krauss representation
E_1
[[ 0.916-0.j    -0.119+0.21j ]
 [ 0.030+0.214j  0.843-0.096j]]
E_2
[[ 0.220+0.067j -0.038+0.044j]
 [ 0.357+0.j    -0.263+0.045j]]
E_3
[[ -0.066+0.105j  0.176-0.j   ]
 [ 0.124-0.077j  0.106-0.047j]]
E_4
[[ 0.004-0.001j  0.009-0.j   ]
 [-0.004+0.004j -0.003+0.003j]]
XXXXXXXX
>> chrep chi
>> chprint
Chi representation
[[ 0.78+0.j    -0.0497+0.17j -0.00539-0.0584j  0.0256+0.0344j]
 [-0.0497-0.17j  0.0969+0.0j -0.0305-0.0256j  0.0316-0.00539j]
 [-0.00539+0.0584j -0.0305+0.0256j  0.0474+0.0j -0.00575+0.0497j]
 [ 0.0256-0.0344j  0.0316+0.00539j -0.00575-0.0497j  0.0754+0.0j]]
XXXXXXXX
>> chrep process
>> chprint
Pauli Liouville representation
[[ 1.0+0.0j -8.99e-18+0.0j -3.3e-16+0.0j -1.21e-18+0.0j]
 [-0.199+0.0j  0.754+0.0j -0.13+0.0j -0.0536+0.0j]
 [-0.0216+0.0j  0.00768+0.0j  0.655+0.0j -0.351+0.0j]
 [ 0.103+0.0j  0.18+0.0j  0.328+0.0j  0.711+0.0j]]
XXXXXXXX
>> chrep choi
>> chprint
Choi representation
[[ 0.453+0.j    -0.063-0.093j  0.045-0.082j  0.352+0.034j]
 [-0.063+0.093j  0.047+0.j    0.025+0.031j -0.045+0.082j]
 [ 0.045+0.082j  0.025-0.031j  0.098+0.j   -0.036+0.082j]
 [ 0.352-0.034j -0.045-0.082j -0.036-0.082j  0.402+0.j   ]]
XXXXXXXX
>>

```

Figure 1: The `chrep` command in `chflow` can be used to change channel representations. The usage of the command is `chrep <representation>`, where `<representation>` should be replaced by one of `{krauss, process, choi, chi, stine}`, which results in the current channel being stored in the specified representation.