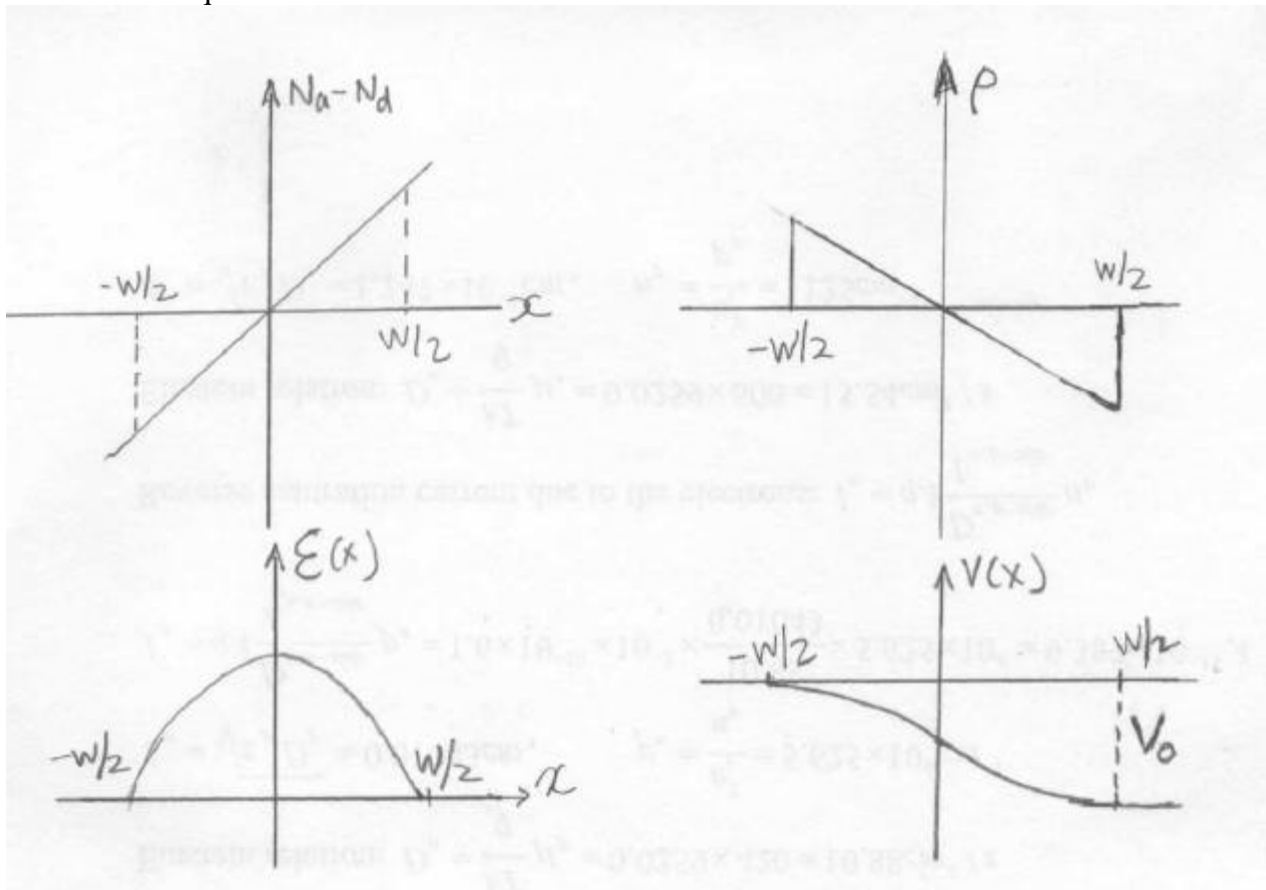


Due: Monday, March 17, 2003

1. When a prolonged diffusion or a high-energy implantation is conducted to form a p/n junction. The doping profile near the junction is usually graded, and the step-junction approach is no longer suitable to find the relationship between the width of the depletion region and the contact potential. However, the underlying principle used to establish equations 5-13 to 5-23 remains intact, and they can still be used to determine similar equations for the graded junction. Assume that the doping profile varies as  $N_a - N_d = Gx$  where  $G$  is  $10^{20}/\text{cm}^4$  in a linear junction.
- (a) Find and plot the electric field,  $e(x)$ , for  $-\infty < x < +\infty$ .
- (b) Determine the relationship between the width of the depletion region and contact potential for the junction at equilibrium.

Solutions:

- (a), Please refer to section 5.6.4 and Figure 5-39 (page 220) to obtain a general understanding of the device. We are going to use depletion approximation in the junction region and quasi-neutral approximation outside junction region. The example in section 5.6.4 has  $N_d - N_a = Gx$ , but we have  $N_a - N_d = Gx$ . Therefore, in our case, the  $\mathbf{r}$ ,  $E_x$ , and  $V$  will be the mirror image of those in figure 5-39 about the  $x$ -axis. After obtaining a general understanding of the device, we can solve the problem through steps similar to equations 5-13 to 5-23:



Since it's a symmetrical device,  $x_{n0} = x_{p0} = W/2$ . For  $x$  outside depletion region ( $|x| > W/2$ ), electrical field is 0. For  $x$  within depletion region ( $|x| < W/2$ ), start with Poisson equation,

$$\frac{dE(x)}{dx} = \frac{q}{e}(p - n + N_d^+ - N_a^-) \xrightarrow{\text{Depletion approximation}} \frac{dE(x)}{dx} = \frac{q}{e}(N_d^+ - N_a^-) = -\frac{q}{e}Gx$$

The general solution to above differential equation:  $E_x = -\frac{qG}{2e}x^2 + \text{Const}$

Applying boundary condition:  $x = \frac{W}{2}$ ,  $E_x = 0$ , we get  $\text{Const} = \frac{qG}{8e}W^2$ .

Therefore,  $E_x = -\frac{qG}{2e}x^2 + \frac{qG}{8e}W^2$ , for  $|x| < \frac{W}{2}$

But we still need to solve for W: following the method used in section 5.2.3 (page 166), we integrate  $E_x$  over depletion region to get  $V_0$ :

$$\begin{aligned} V_0 &= -\int_{-W/2}^{W/2} E_x dx = -\int_{-W/2}^{W/2} \left( -\frac{qG}{2e}x^2 + \frac{qG}{8e}W^2 \right) dx \\ &= -\left[ -\frac{qG}{6e}x^3 + \frac{qG}{8e}W^2x \right]_{-W/2}^{W/2} \quad (*) \\ &= -\frac{qGW^3}{12e} = -1.277 \times 10^{12} W^3 \text{ (volts)} \end{aligned}$$

Noticing  $V_0$  is also the Fermi-level difference across the depletion region:

$$V_0 = -0.0259 \ln \frac{p_p n_n}{n_i^2} = -0.0259 \ln \frac{(GW/2)^2}{n_i^2} = -0.0259 \ln (1.111 \times 10^{19} W^2) \quad (**)$$

Combining (\*) and (\*\*),  $V_0 = -0.0259 \ln (1.111 \times 10^{19} W^2) = -1.277 \times 10^{12} W^3$

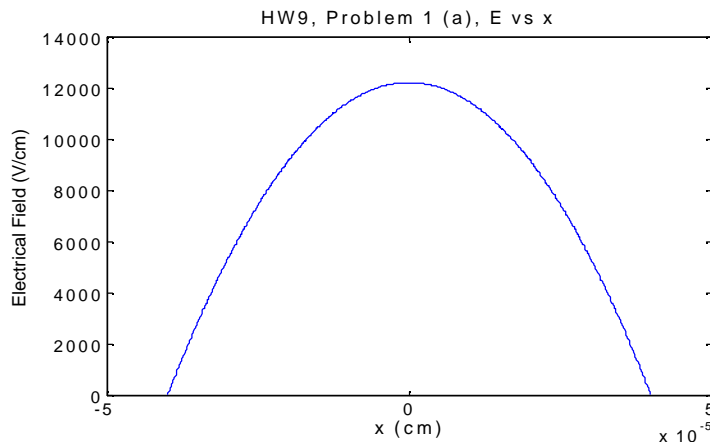
Use a graphic calculator or a computer program, the reasonable answer is:

$$W = 7.97 \times 10^{-5} \text{ cm} = 0.797 \text{ } \mu\text{m}; V_0 = 0.646 \text{ V}.$$

Plug W into the  $E_x$  expression above, we get:

$$\begin{aligned} E_x &= -\frac{qG}{2e}x^2 + \frac{qG}{8e}W^2 = -7.66 \times 10^{12} x^2 + 1.22 \times 10^4 \text{ (V/cm)} & \text{for } |x| < 3.99 \times 10^{-5} \text{ cm} \\ E_x &= 0 & \text{for } |x| > 3.99 \times 10^{-5} \text{ cm} \end{aligned}$$

The plot is shown below:



(b), As shown in part (a),

$$V_0 = -\frac{qGW^3}{12e} = -1.277 \times 10^{12} W^3 (\text{volts})$$

2. An abrupt Si p-n junction has the following properties at 300 K:

p-side

n-side

$$A = 10^{-4} \text{cm}^2$$

$$N_a = 2 \times 10^{17} / \text{cm}^3$$

$$N_d = 1 \times 10^{16} / \text{cm}^3$$

$$\tau_n = 0.1 \mu\text{s}$$

$$\tau_p = 10 \mu\text{s}$$

$$\mu_p = 180 \text{cm}^2/\text{V-s}$$

$$\mu_n = 1080 \text{cm}^2/\text{V-s}$$

$$\mu_n = 600 \text{cm}^2/\text{V-s}$$

$$\mu_p = 400 \text{cm}^2/\text{V-s}$$

- (a) Determine the contact potential  $V_0$  of the junction, and the depletion widths at equilibrium, under a forward bias of  $V_0/2$  and under a reverse bias of  $4V_0$ .
- (b) Draw the band diagrams qualitatively at equilibrium and under forward and reverse bias showing the varying depletion widths, Fermi level and the quasi-Fermi levels.
- (c) Calculate the reverse saturation current due to holes, due to electrons and the total reverse saturation current.
- (d) Calculate the total injected minority carrier current for  $V = V_0/3$  and  $V_0/2$ .
- (e) Poor heat dissipation often leads to a rise of the device temperature. Assuming  $\mu$ 's and  $\tau$ 's do not change with temperature, repeat part (c) for a temperature of  $T = 350 \text{K}$  at which the intrinsic carrier concentration  $n_i = 4 \times 10^{11} / \text{cm}^3$ .

Solutions:

$$(a), V_0 = 0.0259 \ln \frac{N_a N_d}{n_i^2} = 0.772 \text{V}$$

$$W_0 = \sqrt{\frac{2eV_0}{q} \left( \frac{1}{N_a} + \frac{1}{N_d} \right)} = \sqrt{\frac{2 \times 11.8 \times 8.85 \times 10^{-14} \times 0.772}{1.6 \times 10^{-19}} \left( \frac{1}{2 \times 10^{17}} + \frac{1}{1 \times 10^{16}} \right)} = 0.33 \text{mm}$$

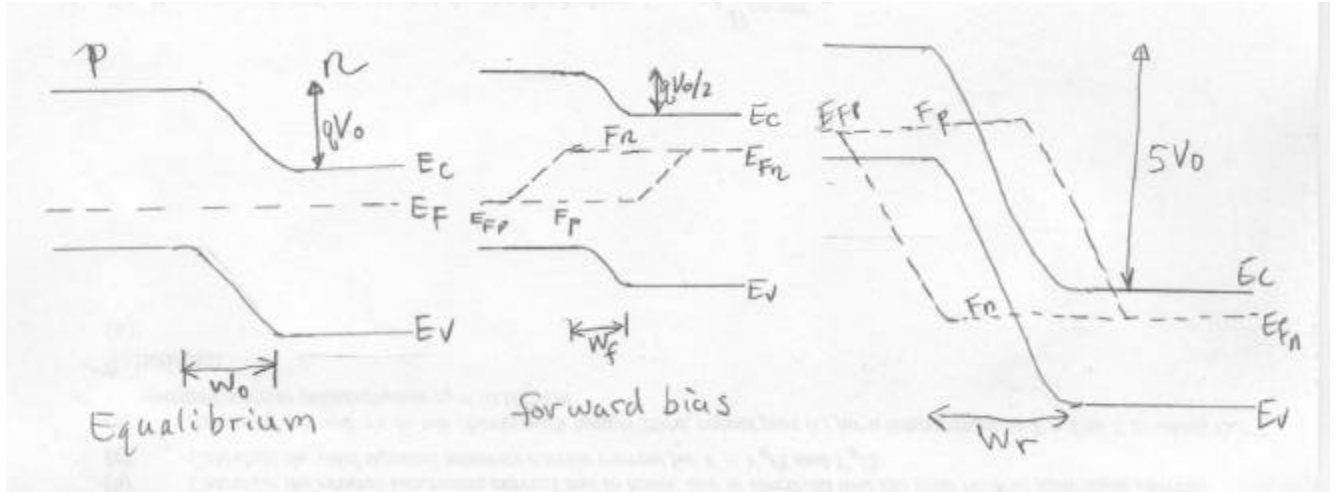
Under forward bias:

$$W_f = \sqrt{\frac{2e(V_0 - V_f)}{q} \left( \frac{1}{N_a} + \frac{1}{N_d} \right)} = \sqrt{\frac{2 \times 11.8 \times 8.85 \times 10^{-14} \times 0.386}{1.6 \times 10^{-19}} \left( \frac{1}{2 \times 10^{17}} + \frac{1}{1 \times 10^{16}} \right)} = 0.23 \text{mm}$$

Under reverse bias:

$$W_r = \sqrt{\frac{2e(V_0 + V_r)}{q} \left( \frac{1}{N_a} + \frac{1}{N_d} \right)} = \sqrt{\frac{2 \times 11.8 \times 8.85 \times 10^{-14} \times 5 \times 0.772}{1.6 \times 10^{-19}} \left( \frac{1}{2 \times 10^{17}} + \frac{1}{1 \times 10^{16}} \right)} = 0.73 \text{mm}$$

(b),



(c),  $D_p(n\text{-side}) = 0.0259 \times 400 = 10.36 \text{ cm}^2 / \text{s}$      $D_n(p\text{-side}) = 0.0259 \times 600 = 15.54 \text{ cm}^2 / \text{s}$

$$L_p(n\text{-side}) = \sqrt{D_p \tau_p} = \sqrt{10.36 \times 10 \times 10^{-6}} = 1.018 \times 10^{-4} \text{ cm}$$

$$L_n(p\text{-side}) = \sqrt{D_n \tau_n} = \sqrt{15.54 \times 0.1 \times 10^{-6}} = 1.25 \times 10^{-3} \text{ cm}$$

$$I_p = qA \frac{D_p}{L_p} p_n = qA \frac{D_p}{L_p} \frac{n_i^2}{n_n} = 3.66 \times 10^{-16} \text{ A}$$

$$I_n = qA \frac{D_n}{L_n} n_p = qA \frac{D_n}{L_n} \frac{n_i^2}{p_p} = 2.24 \times 10^{-16} \text{ A}$$

$$I_{\text{total}} = I_n + I_p = 5.90 \times 10^{-16} \text{ A}$$

(d), For  $V = V_0/3$ ,  $I = I_0(e^{qV/kT} - 1) = I_0(e^{qV_0/3kT} - 1) = 1.22 \times 10^{-11} \text{ A}$

For  $V = V_0/2$ ,  $I = I_0(e^{qV/kT} - 1) = I_0(e^{qV_0/2kT} - 1) = 1.75 \times 10^{-9} \text{ A}$

(e),  $D_p(n\text{-side}) = 0.0302 \times 400 = 12.07 \text{ cm}^2 / \text{s}$      $D_n(p\text{-side}) = 0.0302 \times 600 = 18.1 \text{ cm}^2 / \text{s}$

$$L_p(n\text{-side}) = \sqrt{D_p \tau_p} = \sqrt{12.07 \times 10 \times 10^{-6}} = 1.098 \times 10^{-4} \text{ cm}$$

$$L_n(p\text{-side}) = \sqrt{D_n \tau_n} = \sqrt{18.1 \times 0.1 \times 10^{-6}} = 1.345 \times 10^{-3} \text{ cm}$$

$$I_p = qA \frac{D_p}{L_p} p_n = qA \frac{D_p}{L_p} \frac{n_i^2(350\text{K})}{n_n} = 2.81 \times 10^{-13} \text{ A}$$

$$I_n = qA \frac{D_n}{L_n} n_p = qA \frac{D_n}{L_n} \frac{n_i^2}{p_p} = 1.72 \times 10^{-13} \text{ A}$$

$$I_{\text{total}} = I_n + I_p = 4.53 \times 10^{-13} \text{ A}$$

3. Assume that an abrupt Si p-n junction with area  $10^{-4} \text{ cm}^2$  has  $N_A = 5 \times 10^{16} / \text{cm}^3$  on the p-side and  $N_D = 2 \times 10^{17} / \text{cm}^3$  on the n side.

The diode has a forward bias 0.5 volts.

- (a) Using mobility values from Fig. 3-23 (or better yet from the information sheet used for HR Exam) and assuming that  $\tau_n = \tau_p = 1 \mu s$ , plot  $I_p$  and  $I_n$  versus distance on a diagram such as Fig. 5-17, including both sides of the junction. Neglect recombination within the space charge region, W.
- (b) Plot  $\delta n(x_p)$  and  $\delta p(x_n)$ .
- (c) Determine the separation of quasi-Fermi levels at the position  $5 \mu m$  into the quasi-neutral region in both sides of the junction.

(a), Using Eqn (5-21) to (5-23), we obtain:

$$V_0 = 0.814V$$

From graph,  $m_{p,n-side} = 180 cm^2 / Vs$ , therefore,  $D_p = 0.0259 m_p = 4.662 cm^2 / s$

$$m_{n,p-side} = 810 cm^2 / Vs, \text{ therefore, } D_n = 0.0259 m_n = 20.98 cm^2 / s$$

$$L_n = \sqrt{t_n D_n} = 4.58 \times 10^{-3} cm, \quad n_p = \frac{n_i^2}{p_p} = 4500 cm^{-3}$$

$$L_p = \sqrt{t_p D_p} = 2.16 \times 10^{-3} cm, \quad p_n = \frac{n_i^2}{p_p} = 1125 cm^{-3}$$

$$\Delta n_p = n_p (e^{qV/kT} - 1) = 4500 \times (e^{0.5/0.0259} - 1) = 1.09 \times 10^{12} cm^{-3}$$

$$\Delta p_n = p_n (e^{qV/kT} - 1) = 1125 \times (e^{0.5/0.0259} - 1) = 2.72 \times 10^{11} cm^{-3}$$

In p-region,

$$I_n(x_p) = qA \frac{D_n}{L_n} \Delta n_p e^{-x_p/L_p} = 1.6 \times 10^{-19} \times 10^{-4} \frac{20.98}{4.58 \times 10^{-3}} 1.09 \times 10^{12} \times e^{-x_p(m)/45.8}$$

$$= 7.98 \times 10^{-8} \times e^{-x_n(m)/45.8} A$$

In n-region:

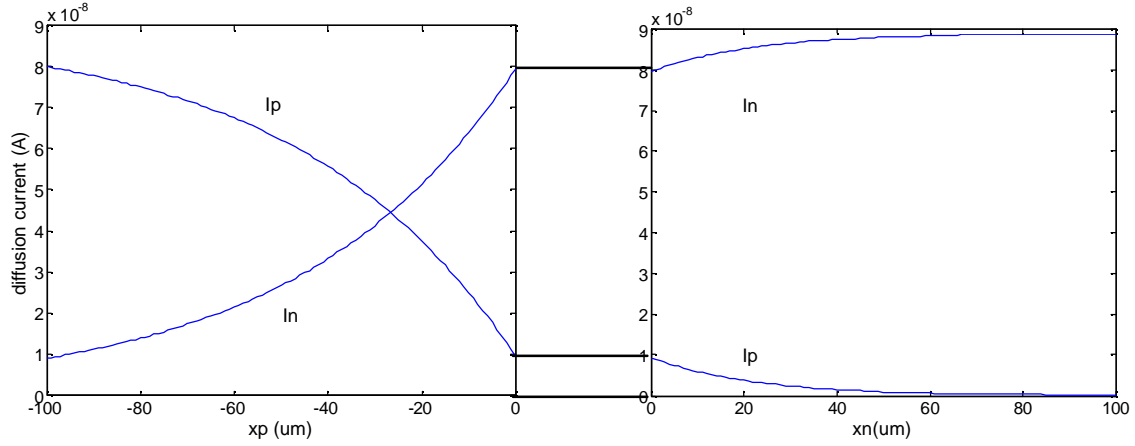
$$I_p(x_n) = qA \frac{D_p}{L_p} \Delta p_n e^{-x_n/L_p} = 9.37 \times 10^{-9} \times e^{-x_n(m)/21.6} A$$

$$I_{total} = I_n + I_p = 8.89 \times 10^{-8} A$$

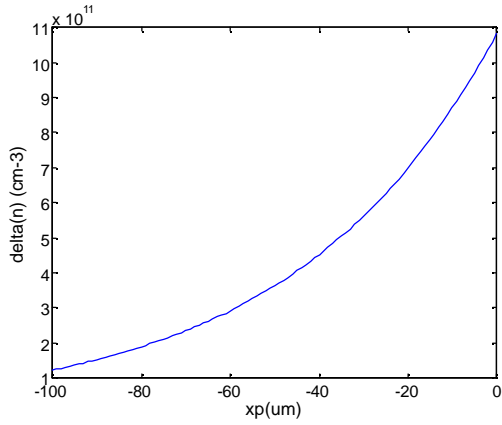
$$I_p(x_n) = I_{total} - I_{n,p-side}$$

$$I_n(x_p) = I - I_p$$

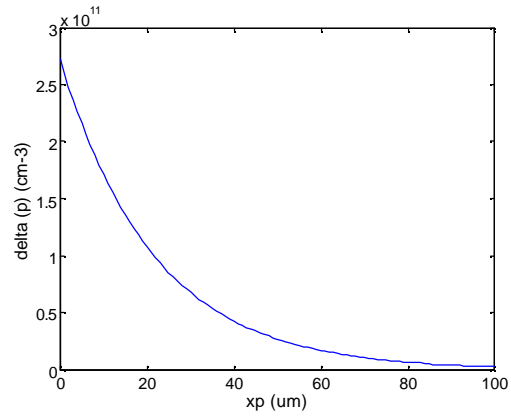
Based on the above calculation, we can draw a graph similar to Figure 5-17



(b),  $d_n(x_p) = 1.09 \times 10^{12} e^{-x_p(\text{mm})/45.8} \text{cm}^{-3}$



$d_p(x_n) = 2.72 \times 10^{11} e^{-x_n(\text{mm})/21} \text{cm}^{-3}$



(c), At  $5\mu\text{m}$  into p-side,  $d_n(x_p = 5\text{mm}) = 1.08 \times 10^{12} e^{-5/45.8} \text{cm}^{-3} = 9.68 \times 10^{11} \text{cm}^{-3}$

There is negligible change in majority carrier concentration.

$$F_n - F_p = 0.0259 \ln \frac{np}{n_i^2} = 0.497 \text{eV}$$

At  $5\mu\text{m}$  into n-side,  $d_p(x_n = 5\text{mm}) = 2.72 \times 10^{11} e^{-5/27.88} \text{cm}^{-3} = 2.15 \times 10^{11} \text{cm}^{-3}$

$$F_n - F_p = 0.0259 \ln \frac{np}{n_i^2} = 0.494 \text{eV}$$