DeltaSpin (Self-Adaptive Spin Constraint in VASP)

Zefeng Cai

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2 Docker Image User Guide

For the following commands, replace the values within the angle brackets (<>) with relevant values for your system:

1. Pull the DeltaSpin image from ghcr.io:

```
docker pull ghcr.io/caizefeng/deltaspin:latest
```

2. Start a container using this image. Please note that you should allocate at least 16 GB of shared memory when starting the container from this image:

```
docker run -d --shm-size=<size of shared memory> \
--name=deltaspin ghcr.io/caizefeng/deltaspin:latest
```

3. Enter the container:

```
docker exec -it deltaspin bash
```

4. Navigate to your calculation directory and place the activation code in the same directory as the INCAR, POSCAR, POTCAR, and KPOINTS files:

```
cd /root/DeltaSpin/examples/metal/Fe cp <path to your activation code file> .
```

5. Run vasp_deltaspin:

```
mpirun -np <number of threads> /root/DeltaSpin/bin/vasp_deltaspin
```

3 Getting Started

1. Determine Wigner-Seitz radius (in Å) for each atom type in the system.

Read the RWIGS entry presented in POTCAR and use that value for INCAR. Note that this value should **remain constant** in every single calculation of this system.

(OPTIONAL) Though tedious and not recommended, you can also determine the Wigner-Seitz radius through a Bader-type spin density analysis.

2. Determine the magnetic ground state of the system.

Set up INCAR with the following parameters:

```
RWIGS = <values from the 1st step>
I_CONSTRAINED_M = 2
SCTYPE = 0
LAMBDA = 0
```

and run vasp_deltaspin. When the calculation is finished, the last MW_current in OSZICAR represents the new ground-state RWIGS-defined magnetic moments.

Note that MAGMOM should be close to the ground-state.

You can add the following two tags to increase precision in ANY non-collinear calculation.

```
LASPH = .TRUE.
GGA_COMPAT = .FALSE.
```

3. Constrain the system to the desired magnetic configuration.

Refer to the INCAR template in DeltaSpin/templates directory. Set M_CONSTR to any value you want to constrain the magnetic moments to.

Run vasp_deltaspin.

Note that the ground-state magnetic configuration for reference, throughout all DeltaSpin calculations of this specific material system, is what we obtained in the 2nd step instead of any empirical value or magnetization in OUTCAR ¹. Different definitions of "magnetization" used in VASP will be discussed in another section.

4. (OPTIONAL) Non-Self-Consistent (NSCF) calculation.

Set INCAR-LAMDA to the last value of OSZICAR-lambda from the same SCF calculation that CHGCAR is from.

Also, complete the standard procedure for a typical VASP NSCF calculation, like preparing a CHGCAR and setting INCAR-ICHARG >= 10.

Run vasp_deltaspin.

¹This means that if you set M_CONSTR to the last MW_current in OSZICAR from the 2nd step and apply the constraining, the energy will be exactly the same.

4 INCAR Template

```
SYSTEM = NiO
#GGA = PE
 #Electronic minimization
#Electronic minimi:
PREC = Accurate
of FFT grid.
ENGUT = 600
EDIFF = 1E-9
LORBIT = 11
LREAL = .FALSE.
ISTART = 0
NPAR = 4
#ICHARG = 1
VORYBUN = 1
                                                   # Accurate is always preferred since non-collinear magnetization is very sensitive to the completeness
 #ICHARG = 1
VOSKOWN = 1
LWAVE = .FALSE.
LCHARG = .FALSE.
NELMIN = 10
 #IBRION = 2 #OPT
 #EDIFFG = -1E-3 #OPT
#ISIF = 3 #OPT
#NSW = 500 #OPT
 #POTIM = 0.2 #OPT
 #DOS
ISMEAR = -5
SIGMA = 0.2
#NEDOS = 3000 #DOS
 # The DEFINITION of atomic spins which are constrained (M_CONSTR in INCAR, MW in constrmag.F/OSZICAR):
 # \int_{\Omega_{I}} \vec{m}(\mathbf{r}) W_{I}(\mathbf{r}) d\mathbf{r}
! IMPORTANT!
 #Non-collinear
 ISYM = -1
RWIGS = 1.286 0.820
GGA_COMPAT = .FALSE.
                                           # Wigner-Seitz radius (in angstrom) for each atom type in the system.
# Restore the full lattice symmetry for gradient corrected functionals. Recommended.
# Non-spherical contribution to the gradient of the density. Recommended.
# Does not matter in non-collinear calculation.
 LASPH = .TRUE.
ISPIN = 2
 LNONCOLLINEAR = .TRUE.
LSORBIT = .TRUE.
!IMPORTANT!
NSC = 100

NSCMIN = 2

SCDIFF = 1E-8

SCCONVB_GRAD = 1.9 -1
                                                  # Break condition for local gradients of spins w.r.t. Lagrangian coefficients (local field) -1: off
           for the element
IDECAY_GRAD = 2
SCDECAY_GRAD = 0.9
SCGRADB = 0.1
NGRAD = 2
NGRAD_STEP = 200 400
                                                   # Decay policy of gradient break condition 0: no decay 1: exponential decay 2: step down # Exponential base of decayed gradient break condition (IDECAY_GRAD = 1) # Lower bound of exponentially decayed gradient break condition # Number of "steps" the gradient break condition curve includes (IDECAY_GRAD = 2) # Indices in the outer loop where the gradient break condition discontinuously decay # SCCONVB_GRAD for corresponding "steps"
NGRAD_VALUE = 1.7 -1 \
1.5 -1
                                                   # Exponential decay of convergence criterion -1: no decay >0,<1: exponential base
# Lower bound of criterion decay, not gonna work if SCDECAY = -1
# Break condition for Lagrangian coefficients difference between two inner steps -1: this condition
SCDECAY = -1
SCDIFFB = 1E-7
SCCONVB = -1
is off
INISC = 0.01
SCCUT = 3
NELMSCI = 1
NELMSCT = 0
TALGRSC = 1
                                                  # Initial trial step size
                                                   # Initial trial step size
# Restriction of step size
# Number of electronic steps BEFORE SASC(L)
# Number of INTERMEDIATE normal electronic steps
# Inner optimization algorithm 1: CG(F-R) 2: CG(P-R) 3: CG(H-S) 4: CG(D-Y)
# Inner diagonalization algorithm 1: sub-space rotation 2: Blocked-Davidson 3: sub-space rotation
 IALGOSC = 1
 IALGOSC_DIAG = 1
with B-D for the last step

LCUTSC_TRIAL = .TRUE.  # Whether trial step size is updated according to last-step optimal value  TRUE: update FALSE: do not update

LDESC = .FALSE.  # Debug mode  TRUE: on FALSE: off
#SASC(Q)
```

```
LAMBDDA_Q = 10 10 10 10 10 10 10 # Initial Lagrangian coefficients for SASC(Q)

CONSTRI_Q = 11 1 1 1 1  # Whether the component is constrained or not 0: not constrained 1: constrained 
NSC_Q = 500  # Maximum number of steps for SASC(Q) iteration

SCDIFF_Q = 3E-4  # Convergence criterion of iteration (iteration iteration)

EDIFF_Q = 1E-6  # Global break condition (EDIFF) for SASC(Q)

INISC_Q = 0.1  # Initial trial step size

#Orbital moments
#LORBMOM = .TRUE.

#Mixer

AMIX = 0.2

EMIX = 0.0001

AMIX_MAG = 0.8

EMIX_MAG = 0.8

EMIX_MAG = 0.0001

#LSDA+U

LDAU = .TRUE.

LDAUTYPE = 2

LDAUL = 2 -1

LDAUU = 5.3 0.0

LDAUJ = 0.0 0.0

LDAUJPRINT = 2

LMAXMIX = 4
```

5 Selected Multi-step Use Cases

5.1 Calculate the effective field of the original ground-state configuration under Zeeman field

1. Relax the system to reach the ground-state magnetic configuration without an external field.

```
Bfield = 0 0 0
I_CONSTRAINED_M = 2
SCTYPE = 0
LAMBDA = 0
```

2. Apply a finite external field B. Constrain the size and direct of the magnetic moments using DeltaSpin.

```
Bfield = <B_x> <B_y> <B_z>
I_CONSTRAINED_M = 2
SCTYPE = 1
LAMBDA = <number of atoms>*0
M_CONSTR = <the last "M_current" from step 1>
```

5.2 Calculate the effective field of a direction-constrained magnetic configuration

1. Constrain the direction of the magnetic moments using the original VASP algorithm that is compatible with DeltaSpin.

```
I_CONSTRAINED_M = 1
SCTYPE = 0
LAMBDA = <some positive real number>
```

2. Constrain the size and direct of the magnetic moments using DeltaSpin.

```
I_CONSTRAINED_M = 2
SCTYPE = 1
LAMBDA = <number of atoms>*0
M_CONSTR = <the last "M_current" from step 1>
```

5.3 Calculate the electronic structure of a fully-constrained configuration

1. Self-consistently constrain the size and direct of the magnetic moments using DeltaSpin.

```
ISTART = 0
ICHARG = 2
LAMBDA = <number of atoms>*0
LWAVE = .TRUE.
LCHARG = .TRUE.
```

 $2. \ \,$ Change to a desired k-point grid/path. Run Delta Spin non-self-consistently.

```
ISTART = 1
ICHARG = 11
LAMBDA = <the last "lambda" from step 1>
```

6 Strategies

When the energy difference (dE) keeps fluctuating around a value above the normal EDIFF (approximately $1E+02 \text{ eV} \Leftrightarrow 1E-02 \text{ eV}$), consider the following steps:

- 1. TEST FOR A SUITABLE GRADIENT CRITERION, i.e. SCCONVB_GRAD. Modifying this value might have a significant impact on convergence. Some values of SCCONVB_GRAD may result in a complete corruption of the diagonalization process. Note that the same element in different systems may have different optimal gradient criteria. For example, the optimal SCCONVB_GRAD for chromium (Cr) is $0.5~(\mu_B^2/\text{eV})$ in a monolayer CrI₃, but it is $2.0~\text{in CrB}_2$. Normally, a simple GRID SEARCH is sufficient to find a decent value for SCCONVB_GRAD.
- 2. SLOW DOWN THE MIXING. Decrease the values of AMIX and AMIX_MAG, and usually, it is necessary to increase NELM at the same time. To avoid out-of-memory issues caused by the increasing number of electronic steps, you may also need to allocate more memory when specifying the computing resources.

7 Compatibility

Using SCTYPE=0, *DeltaSpin* is now compatible with almost all functionalities from the original VASP, including

- 1. Constrained local moments approach (direction, I_CONSTRAINED_M = 1)
- 2. Constrained local moments approach (size and direction, I_CONSTRAINED_M = 2)
- 3. Spin spiral (LSPIRAL=.TRUE.)

8 Discussion on Magnetization in VASP

OUTCAR magnetization (magnetization (x,y,z) in OUTCAR), is the integral of spin density in an area $\Omega_{\mathbf{R}}^{\mathrm{AE}}$, where AE partial wave has integer occupancy, which can be slightly different from the augmentation sphere $\Omega_{\mathbf{R}}$.

$$\mathbf{M}_{\text{OUTCAR}} = \int_{\Omega_{\mathbf{p}}^{\text{AE}}} \rho^{1}(r) \tag{1}$$

$$\rho^{1}(r) = \sum_{n,(i,j)} f_{n} \langle \tilde{\Psi}_{n} | \tilde{p}_{i} \rangle \langle \phi_{i} | r \rangle \sigma \langle r | \phi_{j} \rangle \langle \tilde{p}_{j} | \tilde{\Psi}_{n} \rangle$$
 (2)

where

$$\int_{\Omega_{\mathbf{R}}^{\mathbf{AE}}} \langle \phi_i | r \rangle \langle r | \phi_j \rangle = \delta_{ij} \tag{3}$$

Because

$$\rho(r) = \tilde{\rho}(r) + \rho^{1}(r) - \tilde{\rho}^{1}(r)$$

$$= \sum_{n} f_{n} \langle \tilde{\Psi}_{n} | r \rangle \sigma \langle r | \tilde{\Psi}_{n} \rangle + \rho^{1}(r) - \sum_{n, (i, j)} f_{n} \langle \tilde{\Psi}_{n} | \tilde{p}_{i} \rangle \langle \tilde{\phi}_{i} | r \rangle \langle r | \tilde{\phi}_{j} \rangle \sigma \langle \tilde{p}_{j} | \tilde{\Psi}_{n} \rangle$$

$$(4)$$

and

$$\sum_{i} |\tilde{\phi}_{i}\rangle\langle \tilde{p}_{i}| = 1 \text{ within } \Omega_{\mathbf{R}}$$
 (5)

If we assume $\Omega_{\mathbf{R}}^{AE} \approx \Omega_{\mathbf{R}}$, we have

$$\mathbf{M}_{\text{OUTCAR}} = \int_{\Omega_{\mathbf{R}}^{\text{AE}}} \rho^{1}(r) \approx \int_{\Omega_{\mathbf{R}}} \rho(r)$$
 (6)

 ${\tt OSZICAR}\ magnetization\ ({\tt M_current}\ in\ {\tt OSZICAR}),$ is the integral of spin density in a user-defined real-space sphere.

$$\mathbf{M}_{\text{OSZICAR}} = \int_{\Omega_{\mathbf{R}}'} \rho(r) \tag{7}$$

Weighted OSZICAR magnetization (MW_current in OSZICAR), is the weighted integral of spin density in a user-defined real-space sphere. Prevalent weight functions are Bessel functions.

$$\mathbf{M}_{\text{OSZICAR}}^{\text{W}} = \int_{\Omega_{\mathbf{R}}'} W(r) \, \rho(r) \tag{8}$$

If we define $\Omega'_{\mathbf{R}}$ using WS radii (RWIGS) in POTCAR, that is $\Omega'_{\mathbf{R}} \approx \Omega_{\mathbf{R}}$, "OUTCAR magnetization" and "OSZICAR magnetization" will be closer (closer than partial core radius RPCOR and outmost cutoff radius RCORE), but still have a notable difference. That's because these two have totally different origins.

Most importantly, both of them are based on the muffin-tin approximation and the summation of all atomic magnetic moments won't be equal to total magnetization. It's because there always exist some gaps in between those spheres.

The only well-defined quantity in DFT is "total magnetization", the all-space integral of charge density, which is mag in OSZICAR.

To get access to a new definition of atomic magnetization, the summation of which is exactly the all-space integral, one could apply Bader analysis on spin density.

9 Disclaimer

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