

Fits of 1990/1991 Data with Expostar Including Non Standard Model Extensions

A study of the combined 1990/1991 data sample has been completed with the Expostar fitting program. The analysis is virtually identical to the previously presented work to the 1990 data alone[1], with a few corrections being evaluated more precisely and a more sophisticated treatment of some of the systematic uncertainties being used. As has been previously described, Expostar is a rigorous standard model calculation optimized to fit observable differential and total cross sections. The processes that can be simultaneously calculated, so as to include all the theoretical correlations, are the total hadronic cross section, the differential production cross sections for $\mu^+\mu^-$ and e^+e^- (the latter including t channel exchanges to α^3 and α^4 leading log), the differential τ production cross sections for polarized states, and the heavy quark forward backward asymmetries. These all explicitly include double boson exchange diagrams (boxes) which are particularly important for the asymmetries. This has all been described in detail [2], but just as a reminder, the fit parameters in the standard model are the Z^0 mass, the top quark mass (the higgs mass has a degenerate effect on neutral current processes with the exception of the b quark branching ratio), and the factorized QCD corrections. This last variation is accommodated by adjusting the number of colors, which always multiplies these corrections.

The calculations have been extended to include the effects of extended types of models in the way prescribed by Peskin and Takeuchi [3]. This is done by the addition of constants to linear combinations of the various self energy loop integrals and their derivatives, in such a way as to preserve the underlying gauge theory structure (ward identities etc..). As this structure, the so called S, T and U parameters, was originally arrived at by Peskin and Takeuchi in considering the running star scheme[4], it was straight forward to incorporate it into the Expostar program. In fact it had been discussed in earlier work [5].

The data samples used in these fits are straight forward. The total number of hadronic and LAL events are taken from scanbook. A modified version of the standard common lepton analysis (using a rapidity cut) that divides the common lepton sample into two correlated classes based on the Locci/Cliff Bhabha analysis provides differential distributions of leptons. These samples use the EW group run selection. By analyzing the data in this manner the largest systematic errors in the lepton analyses, τ misidentification, is overcome. Further, the τ efficiency within the solid angle is above 94 percent. All of the lepton efficiencies are evaluated as a function

of angle. The non collinear photon smearing corrections (and the effect of unmeasured ν and π^0) causing errors in the perceived value of $\cos(\theta')$ and rapidity are evaluated at all angles and energies by Monte Carlo (Koralz). The differential τ polarization measurements have been provided to me by J. Conway. The QCD corrections were constrained by the Aleph value of $\alpha_s(M_z) \equiv 0.125 \pm 0.005$, determined from event shape measurements, in some of the fits. The proton antiproton W measurements are taken from the literature[6].

The expected cross section for the LCAL is evaluated inside the fits as the electroweak parameters are varied. The luminosity is a fit parameter allowing for a nonzero contribution to the overall χ^2 from the LCAL. This smooths out the fit and yields a lower overall χ^2 . It also provides a convenient way to include the common luminosity correlation. The effect of the luminosity systematic error is divided into two uncorrelated components and a common component. These are incorporated as constrained fit parameters, rescaling the luminosities. The luminosity treatment was done in conjunction with the group from NBI, many thanks to all concerned. The energy scales are allowed to vary within constraints of 25 and 7 MeV for the 1990 and 1991 data samples. These values are perhaps a bit optimistic. The various nonlinear corrections to the energies have been incorporated.

The table below shows the results of various fits. In all of these the Higgs mass was assumed to be 200 GeV. Changing the Higgs mass to 300 GeV raises the top quark mass by approximately 5 GeV.

N_{had}	$\frac{dN_{\mu,\tau}}{d\cos\theta'}$	$\frac{dN_e}{d\cos\theta'}$	α_s	$\tau_{pol}(\theta')$	M_w	M_z	M_{top}	N_{col}	χ^2/NDF
X	$X _{-0.9}^{+0.9}$	$X _{-0.9}^{+0.9}$	-	-	-	91.190 ± 0.008	206_{-37}^{+31}	2.998 ± 0.013	548/552
-	$X _{-0.9}^{+0.9}$	$X _{-0.9}^{+0.9}$	-	-	-	91.188 ± 0.016	176 ± 58	2.99 ± 0.019	534/537
X	$X _{-0.9}^{+0.9}$	$X _{-0.9}^{+0.9}$	X	-	-	91.190 ± 0.008	197_{-36}^{+27}	3.007 ± 0.004	549/553
X	$X _{-0.9}^{+0.9}$	$X _{-0.9}^{+0.9}$	X	X	-	91.190 ± 0.008	181_{-34}^{+29}	3.007 ± 0.004	554/562
X	$X _{-0.9}^{+0.9}$	$X _{-0.9}^{+0.9}$	X	X	X	91.190 ± 0.007	165_{-30}^{+25}	3.008 ± 0.004	555/563

In the second fit it should be noted that while the leptons make up only 10 percent of the data, the accuracy with which the Z^0 mass can be determined is only a factor of 2 worse than including the hadronic line shape. This is due to the different energy dependences of the interference terms in the Bhabha and μ, τ production cross sections. EXPOSTAR explicitly calculates the full Bhabha matrix element, and the t channel photon exchange interference with the s channel Z^0 exchange decreases with energy, while the s channel photon exchange interference increases. The simultaneous fitting of these processes constrains $S-M_z^2$. This will be an extremely useful technique when the 1992 data is considered, where there will probably not be as much off peak data.

The comparison of the first and third fits illustrates the correlation between the

QCD correction and the deduced top mass. When the QCD correction is constrained by the ALEPH value of α_s , it is increased. To compensate for the increase in the effective width, $\sin\theta_w$ must increase, requiring a smaller top quark mass. The inclusion of the W mass measurements pulls the top quark mass by almost one σ . The value of the top quark mass is consistently 10 GeV higher than the official result presented at Dallas.

The consistently higher Z^0 masses found compared to the official results (3 MeV) is probably due to the fitting of the luminosity, and the omission on my part of the light pair production from the structure functions. The correlation coefficients between the mass and the luminosities varies between 10% and 25%, and so probably account for the rest. The second factor would lower the mass by 1.3 MeV[7].

A series of fits, like the second to last fit in the above table, were made to study the sensitivity to the higgs mass in a standard model constrained fit. The higgs mass was varied from 50 to 1000 GeV. All of the fits had 562 degrees of freedom. As can be seen from the table below, the ALEPH data alone in fact seems to prefer a higher higgs mass, but with weak statistical significance. The χ^2 is virtually constant between 100 and 500 GeV. This is consistent with the simulated study described in the EXPOSTAR paper[2]. The $b\bar{b}$ data does have some sensitivity, as the top quark dependence of the loops is virtually canceled by the vertex corrections, leaving only higgs mass dependence. It should be noted that all perturbative calculations make little sense for higgs masses above 600 GeV as the higgs sector becomes strongly interacting and perturbative calculations cannot converge.

M_{higgs}	50GeV	100GeV	200GeV	300GeV	500GeV	1000GeV
χ^2	554.28	553.96	553.73	553.71	553.49	553.35

Including the W mass measurements does not change the conclusion as the results were as follows:

M_{higgs}	50GeV	100GeV	200GeV	500GeV
χ^2	555.69	555.38	555.10	555.00

Thus the 1991/1992 data sample would indicate a higher top quark mass and Higgs mass than had been extracted from the 1990 data alone. As will be shown this can easily be understood in terms of the fits to the extension parameters, S and T.

As usual, an additional fit was performed to see if the data preferred a form different than that predicted by the Standard Model. All of the Aleph measurements discussed were used. The fit variables were M_Z , N_ν , N_c , and the vector and axial lepton coupling multipliers. These last two parameters rescale the standard model vector and axial couplings for leptons. The results are shown below assuming a top quark mass of 181 GeV and a Higgs mass of 200 GeV.

M_Z	N_ν	N_c	k_v	k_a	$\chi^2/N.D.F.$
91.191 ± 0.008	3.017 ± 0.05	3.008 ± 0.005	$.930 \pm 0.095$	1.0028 ± 0.0025	552/560

These results show that the Standard Model with 3 light neutrinos not only describes the data well, but is the *preferred* solution.

After the standard model parameters have been optimized I considered the ability of these data to constrain the S and T parameters. Thus, these are the values of these constants *beyond* the contributions due to the standard model. The method of analysis is similar to the fit just discussed, which yielded the number of neutrinos and the lepton multipliers. The value of the top mass is frozen at the preferred values of 181 GeV and 165 GeV (when the proton antiproton data is used) and a simultaneous fit to the Z^0 mass, the number of colors, and S and T is performed. The variation of the QCD correction should not be overlooked as it is correlated to the extension parameters. Minos was also used to evaluate the full error on S and T and these errors were the same as the parabolic errors. The contour plot of the correlated extension parameters is then evaluated within Minuit, but of course the other correlated parameters are not varied.

The significance of this is that the theoretical impact on the parameter S can be evaluated for many types of models (technicolor models for example) without knowledge of the masses of the new particles of those models[5],[3],[8]. As a historical note, the actual fit parameters used are like those discussed in the earliest references[5] and most of the later references[8]. The first two are identically S (called del3tec) and T (called rhotec) and the third, called del1tec, is actually S+U. These names are rather obvious as they are addition to the running coupling correction functions Δ_{ρ^*} , Δ_{3^*} and Δ_{1^*} , which form the core of the * scheme. The first two functions control the q^2 evolution of ρ^* and G_F^* . The third can be evaluated from the other two if custodial SU(2) symmetry is correct. This means that the Higgs sector has a global SU(2) symmetry, ie the reason that the tree level ρ parameter is 1. The third function controls the running of the product of $\rho^*G_F^*$, thus is clearly not independent in the MSM.

These parameters measure only *deviations* from the MSM. In such terms the contribution to S from a one generation SU(4) technicolor model would be around 1.8, and a double SU(4) technicolor model would be at S of 0.5. These values can be seen on the reproduction of the graph from the second Peskin/Takeuchi paper[3]. The errors of this analysis are considerably smaller than those presented by Peskin and Takeuchi. However the central value has shifted to the upper right hand quadrant, thus explaining the preferred high Higgs mass.

The evaluation of these parameters is unfortunately not entirely unique. In fact one must always subtract off the values evaluated at some reference point. In Peskin and Takeuchis' paper their reference point is defined by a Z mass of 91.174 GeV, a

top mass of 150 GeV and a Higgs mass of 1 TeV. Even with this subtraction, the renormalization scheme dependence remains as I tested by evaluating the standard model variations with both the old * scheme (just loops), and the new one (**), which uses also universal boxes and vertices in defining the running couplings. While the errors are large compared to this difference, clearly a scheme could be chosen to produce really any value one likes, the same problem that exists with widths, $\sin\theta_{w,*,-}$, etc. The standard model predicts cross sections, given the masses. But proceeding undaunted, the fits using the ALEPH data with and without the W mass measurement yields:

ALEPH	M_w	M_Z	N_c	S	T	$\chi^2/N.D.F.$
X	-	91.191 ± 0.007	3.007 ± 0.004	$0.794 \pm .723$	$.531 \pm 0.663$	553/561
X	X	91.191 ± 0.006	3.007 ± 0.004	$0.737 \pm .603$	$.687 \pm 0.633$	554/562

The contour plots from the fits (attached) are in the upper right hand quadrant. (Note that the contour plots have the axes flipped with respect to Peskins') If one looks at the attached graphs copied from reference[3], the reason for the high top and higgs masses are immediately clear. The standard model is an extremely narrow vertical band around $S = 0$. Thus an upper right hand quadrant value will require a high top and Higgs mass. A fit, in general, will yield some value of S, pulling the Higgs mass to its lowest or highest value, with no ability to differentiate in between, as the MSM range is far too narrow to allow a real minimum to exist. If the analysis is done in terms of the usual fit parameters of scheme dependent widths and so on, what you are doing is looking at the intersection in the ST plane of nearly parallel diagonal bands. This is illustrated by the reproduced graphs also taken from referece[3]. The contour plots of T (called rhotec) versus S (del3tec) for the first two extension fits are attached. Thus it is clear that with the inclusion of the 1991 data sample the situation in the ST plane and hence any understanding about the higgs sector has changed substantially.

Fitting to all three of the extension parameters simultaneously using all of the ALEPH data and the $p\bar{p}$ data yields the following after untangling S and U.

M_Z	N_c	S	T	U	$\chi^2/N.D.F.$
$91.191 \pm .007$	$3.007 \pm .004$	$.765 \pm .744$	$.698 \pm .659$	$-.68 \pm 1.03$	554/561

The value of U indicates that there is no indication of a custodial SU(2) symmetry breaking.

Just to check, two fits to just the 1990 data were also made, after optimizing the top quark mass to their preferred values of 130 GeV and 137 GeV. For these fits the τ polarization was not used as the values I have are the sum of the 1990 and 1991 data.

ALEPH	M_w	M_Z	N_c	S	T	$\chi^2/N.D.F.$
X	-	91.176 ± 0.010	3.008 ± 0.005	$-.56 \pm 1.1$	-1.13 ± 1.59	258/255
X	X	91.176 ± 0.010	3.009 ± 0.005	$-.36 \pm .88$	$-.33 \pm 0.98$	259/256

The value in the lower left quadrant explains the conclusions of the low higgs mass and the lower top mass than the current data indicates. While it is unclear as to whether these parameters really have any meaning being RS dependent, they do make the top quark and Higgs mass dependence of the data easy to visualize. I suspect that a value in the upper right quadrant may have a bearing on the likelihood of observing supersymmetry at LEP 200. High mass supersymmetry models with no SU(2) breaking exactly duplicate the standard model. It would therefore seem reasonable that a model with low masses would deviate from the standard model in the S T plane by the maximal amount allowed in MSSM. This suspicion was confirmed by a conversation with J. Ellis. I will try to investigate this in the near future, by including true MSSM corrections to loops, vertices and maybe even boxes. This is required if the scheme dependence is to be reasonably contained.

References

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D44 1591 (1991). G. Altarelli and R. Barbieri, Phys. Lett. **B253** 161 (1991). G. Altarelli, R. Barbieri and S. Jadach, Nucl. Phys. **B269** 3 (1992).

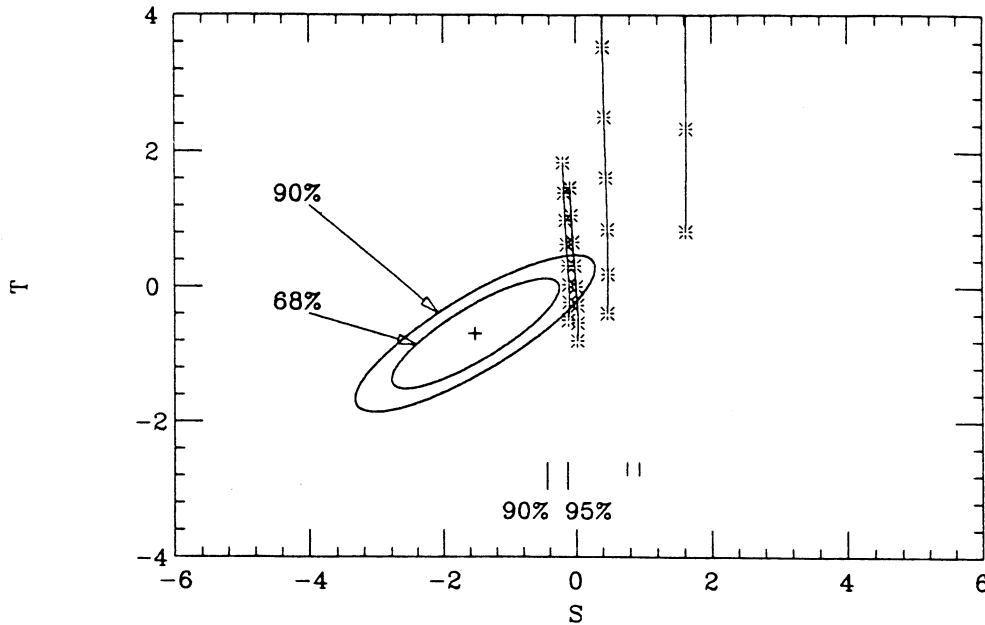


FIG. 15. Contours of the likelihood function of S and T corresponding to 68% and 90% probability, computed from the measurements listed in Table I. The long notches at the bottom of the figure correspond to the 90% and 95% confidence upper limits on S in an unbiased analysis. The shorter notches show the locations of these limits if one imposes *a priori* that $S > 0$.

$m_H = 1$ TeV. For different values of m_t and m_H , the position of the likelihood contours on the S - T plane will be different. However, the shapes and sizes of these contours would be the same. This gives a convenient way to plot the influence of the reference m_t and m_H on the S - T analysis: We simply hold the position of the likelihood

contours fixed and plot the relative position of the origin with respect to these contours. As we vary m_t , this relative position then sweeps out a contour in the S - T plane which roughly follows the displacements (4.4) but gives a more accurate accounting for small values of m_t . In Figs. 15 and 16, we have plotted the contours corre-

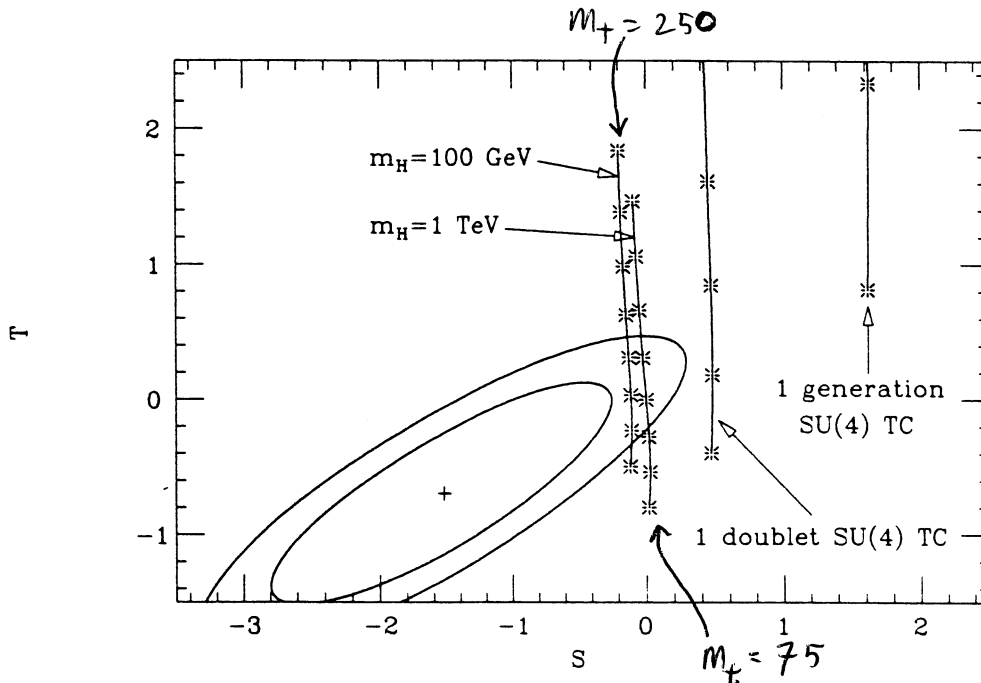


FIG. 16. Enlarged version of Fig. 15, showing the comparison of the region preferred by the fit with the predictions of the minimal standard model and two technicolor models. The values of S and T for the minimal standard model are computed as described in the text, for Higgs-boson masses for 100 GeV and 1 TeV, as a function of the top-quark mass. The stars denote values of m_t from 75 to 250 GeV in 25 GeV steps. The values of S in technicolor models are the values for $N_{TC} = 4$ from Sec. VII. The values of T due to technicolor are computed from (8.2), as an indication of the possible size of this effect. Again, the stars denote values of m_t increasing from 75 GeV in steps of 25 GeV.


```
FCN=      552.6972      FROM MINOS      STATUS=SUCCESSFUL      427 CALLS      1367 TOTAL
EDM=    0.33E-05      STRATEGY=1      ERROR MATRIX UNCERTAINTY=    0.6%
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EXT	PARAMETER		PARABOLIC	MINOS ERRORS	
NO.	NAME	VALUE	ERROR	NEGATIVE	POSITIVE
1	ZMASS	91.191	0.62367E-02		
3	NUM COL	3.0071	0.45265E-02		
10	RHOTE ^C	0.53267	0.66095		
11	DEL3TEC	0.79487	0.74882	-0.75105	0.75720
13	HADNORM	0.99941	0.28672E-02		
14	ABSNORM	0.99994	0.21848E-02		
15	LEPNORM	1.0003	0.25206E-02		
16	BHASYS	1.0001	0.10671E-03		
18	ESCALE90	1.0001	0.93401E-04		
19	ESCALE91	0.99999	0.46064E-04		
20	ABSNORM90	1.0058	0.32687E-02		
21	ABSNORM91	0.99860	0.18633E-02		
22	LUMIN_1	480.26	4.1805		
23	LUMIN_2	517.74	4.2222		
24	LUMIN_3	446.94	3.4457		
25	LUMIN_4	3609.6	11.718		
26	LUMIN_5	551.38	3.6143		
27	LUMIN_6	590.03	4.3295		
28	LUMIN_7	633.07	4.8826		
29	LUMIN_8	677.14	4.9017		
30	LUMIN_9	800.98	5.3031		
31	LUMIN_10	755.14	4.7809		
32	LUMIN_11	4620.5	13.558		
33	LUMIN_12	697.08	4.0378		
34	LUMIN_13	677.63	4.5276		
35	LUMIN_14	769.03	5.1490		
36	LUMIN_15	2952.8	9.6639		

* * * * *

Y-AXIS: PARAMETER 11: DEL3TEC
X=0

2.233	7	6	55	44*33	222
2.113	77	66	55	44	33 222 22
1.993	7	66	55:44	33	222 2
1.873	66	55	44	33* 22	2
1.753	6	5	4	3 *22	22
1.634	66	55	44	33 22	2 3
1.514	6	55	44:33	22 111	2 3
1.394	55	44	33	22*111	1 22 33
1.274	5	44	33	22 11	1 2 33
1.154	55	4	3:22	11 11	22 3 4
1.034	5	44	33:2	11*	1 22 33 4
0.9147	44	33	22	1 *	11 2 33 44
0.7949	*4*33*22*11**11*22*3*44*				
0.6751	44	3	2:11	* 1	22 33 4 5
0.5552	4	33	22:1	*11	2 33 44 5
0.4354	3	22	11	11	2233 44 55
0.3156	33	2	1	11	22 3 44 55
0.1958	33	22	1	111*22	3344 55 6
0.7600E-01	3	2	111	22	33 4 55 66
-0.4381E-01	3--2---+---22-33-4455-66-				
-0.1636	22	:	22*33	44	5 66 7
-0.2834	2	:	222	33 44	5566 77
-0.4032	2	22	33	44	55 6 77
-0.5231	222222	33*44	55	6677	8
-0.6429	22	:	33	44	55 6677 88

I	I	I
-0.7892		1.855

```

0.7892      0.5327      1.000
X-AXIS: PARAMETER 10: RHOTEC      ONE COLUMN= 0.1058
FUNCTION VALUES: F(I)= 552.7      + 1.000      *I**2

```

T

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3	NUM COL	3.0069	0.44417E-02	
10	RHOTEC	0.68887	0.64421	
11	DEL3TEC	0.73751	0.61382	-0.61519 0.61550
13	HADNORM	0.99919	0.28948E-02	
14	ABSNORM	0.99987	0.21871E-02	
15	LEPNORM	1.0005	0.24835E-02	
16	BHASYS	1.0001	0.10604E-03	
18	ESCALE90	1.0001	0.94632E-04	
19	ESCALE91	0.99999	0.45812E-04	
20	ABSNORM90	1.0057	0.32977E-02	
21	ABSNORM91	0.99856	0.18655E-02	
22	LUMIN_1	480.38	4.1795	
23	LUMIN_2	517.88	4.2192	
24	LUMIN_3	447.06	3.4568	
25	LUMIN_4	3610.6	11.735	
26	LUMIN_5	551.54	3.6628	
27	LUMIN_6	590.19	4.3540	
28	LUMIN_7	633.24	4.8870	
29	LUMIN_8	677.30	4.8875	
30	LUMIN_9	801.17	5.2643	
31	LUMIN_10	755.32	4.7361	
32	LUMIN_11	4621.8	13.550	
33	LUMIN_12	697.28	4.0716	
34	LUMIN_13	677.81	4.5638	
35	LUMIN_14	769.23	5.1791	
36	LUMIN_15	2953.6	9.6506	

Y-AXIS: PARAMETER 11: DEL3TEC
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1.916 8 77 66 55 44 33
1.818 77 66 55 44*33 22222
1.720 77 66:55 44 33 22 2
1.621 7 66 55 44 33 222 2
1.523 7 6 55 44 33*22 2 3
1.425 66 5:44 33 22 2 3
1.327 66 55:4 33 22 11 22 33
1.229 6 55 4433 22* 1111 2 3
1.130 55 44 3 22 111 11 22 33
1.032 5 44:33 2 11 1 2 33 4
0.9339 55 4 33 22 1* 1 22 3 44
0.8357 5 4433 22 11* 11 2 33 4
0.7375 *44*3*22*11***11*2233*445
0.6393 44 33:2 11 *11 22 3 44 5
0.5411 4 3 22 1 *1 22 3344 55
0.4429 4 33 2 1 1122 33 4 55
0.3447 33 22 11 111 2 33 4455 6
0.2465 3 2: 1111 *2233 44 5 66
0.1482 33 22: 11 22 3 44 5566
0.5003E-01 3 2 : 22 3344 55 6 7
-0.4818E-01 3--2+---222-3344-55-6677
-0.1464 2 : 22 33 4 55 6677
-0.2446 2 :222 33 4455 66 7 8
-0.3428 2222 33*4455 66 7788
-0.4410 : 333 44 5 66 77 8

I I I
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0.6889

X-AXIS: PARAMETER 10: RHOTEC ONE COLUMN= 0.1031
FUNCTION VALUES: F(I)= 553.8 + 1.000 *I**2

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10	RHOTE	0.69815	0.65917	-0.15747E-03	0.55786E-03
11	DEL3TEC	0.76544	0.74484	0.12068E-03	0.12243E-02
12	DEL1TEC	0.84648E-01	1.2523	-0.44378E-03	-0.70367E-03
13	HADNORM	0.99919	0.28975E-02	0.19072E-05	-0.69295
14	ABSNORM	0.99988	0.21964E-02	0.38131E-05	0.37386
15	LEPNORM	1.0005	0.25740E-02	-0.21985E-05	-0.37567
16	BHASYS	1.0001	0.10729E-03	0.19018E-06	7.9805
18	ESCALE90	1.0001	0.96972E-04	-0.10887E-07	-3.5433
19	ESCALE91	0.99999	0.43624E-04	0.27876E-07	-12.496
20	ABSNORM90	1.0057	0.32840E-02	-0.86728E-06	0.11187
21	ABSNORM91	0.99856	0.19000E-02	-0.14972E-05	0.63768
22	LUMIN_1	480.38	4.1702	-0.84270E-06	-0.50382
23	LUMIN_2	517.88	4.2353	0.27544E-07	-0.31022
24	LUMIN_3	447.06	3.5394	0.89126E-06	-0.25153
25	LUMIN_4	3610.6	11.760	-0.43506E-06	0.23763
26	LUMIN_5	551.54	3.7634	-0.99224E-06	0.63439
27	LUMIN_6	590.20	4.3362	0.46320E-07	1.0630
28	LUMIN_7	633.25	4.8890	-0.32916E-06	0.31729
29	LUMIN_8	677.30	4.8892	-0.86280E-06	-0.71209
30	LUMIN_9	801.17	5.3066	0.56685E-06	-0.13552
31	LUMIN_10	755.33	4.7406	0.10762E-05	-0.69349E-01
32	LUMIN_11	4621.9	13.627	-0.92173E-06	-0.66211
33	LUMIN_12	697.29	3.9185	-0.65274E-07	0.92053
34	LUMIN_13	677.82	4.4408	0.49014E-06	1.0003
35	LUMIN_14	769.23	5.1301	0.12728E-06	0.49071
36	LUMIN_15	2953.6	9.6993	-0.66938E-06	-0.64889

Y-AXIS: PARAMETER 11: DEL3TEC
X=0

```

2.196      6666  555 : 4444
2.076      66  5555 4444
1.957      555  4444 33333333
1.838      555  444 : 3333
1.719      55  444  3333
1.600      444  333: 222222
1.480      444  3333 :22222 222
1.361      44  33  2222 2
1.242      333  222 : 2
1.123      333  222 1111111 22
1.004      33  22  111 1 2
0.8846     3  222 111 : 11 22
0.7654     **22**11*****11**22**
0.6463     22  11 : 111 222 3
0.5271     2  1  111 22 333
0.4079     22  1111111 222 33
0.2887     2 : 222 333 4
0.1696     22  2222 333 444
0.5039E-01 -22222222222+-333--444--
-0.6879E-01      3333 444 5
-0.1880      3333 444 555
-0.3071     3  3333333 : 444 5555
-0.4263     3333 4444 555 66
-0.5455      44444: 555 666
-0.6647     44444444 5555 6666 7
          I          I          I
-2.420      0.8465E-01      2.589

```

X-AXIS: PARAMETER 12: DEL1TEC ONE COLUMN= 0.2004
FUNCTION VALUES: F(I)= 553.8 + 1.000 *I**2

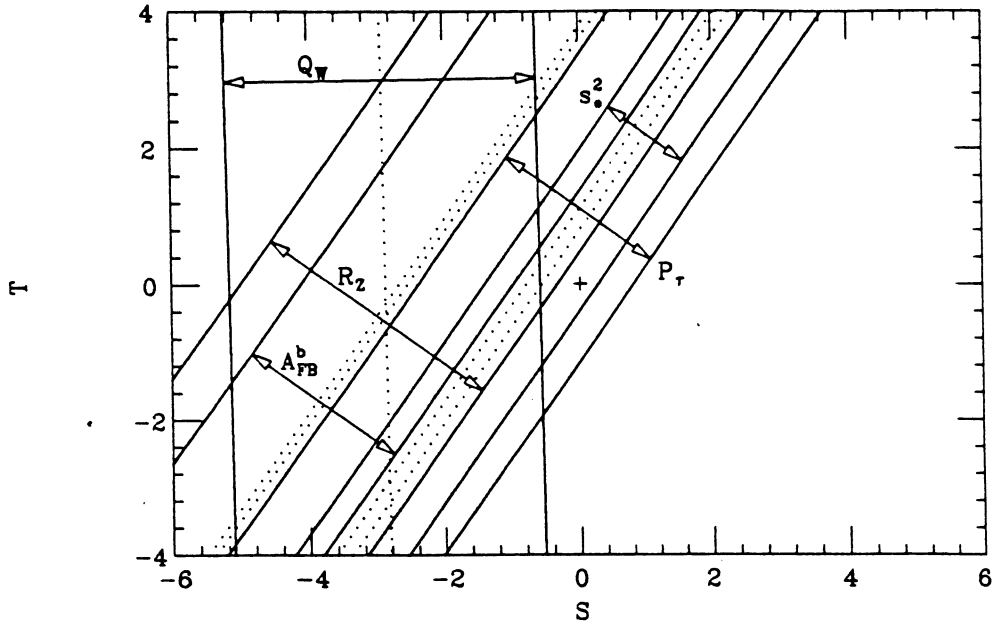


FIG. 14. Bands in the S - T plane allowed, within 1σ errors, by the last five measurements listed in Table I. These observables belong to the second and third classes discussed in the text.

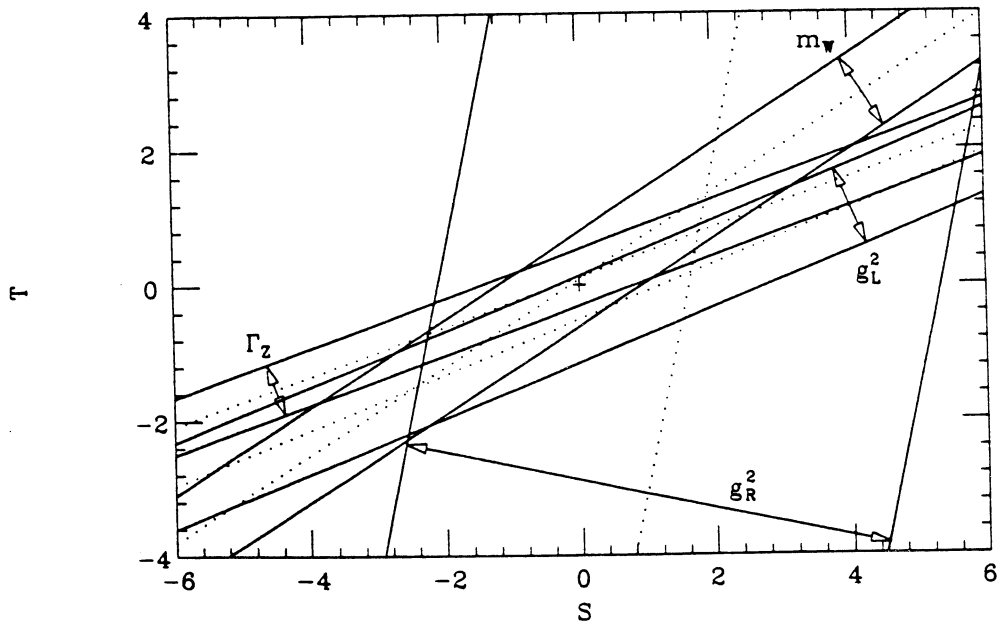


FIG. 13. Bands in the S - T plane allowed, within the 1σ errors, by the first four measurements listed in Table I. These observables belong to the first class discussed in the text.

1990 aleph only

FCN= 257.8940 FROM MINOS STATUS=SUCCESSFUL 437 CALLS 1061 TOTAL
EDM= 0.24E-05 STRATEGY= 1 ERROR MATRIX ACCURATE

EXT NO.	PARAMETER NAME	VALUE	PARABOLIC ERROR	MINOS ERRORS
				NEGATIVE POSITIVE
1	ZMASS	91.176	0.99328E-02	
3	NUM COL	3.0084	0.50153E-02	
10	RHOTE ^C	-0.55530	1.1046	
11	DEL3TEC	-1.1266	1.5859	-1.5257 1.6756
13	HADNORM	1.0021	0.40503E-02	
14	ABSNORM	0.99997	0.24151E-02	
15	LEPNORM	0.99936	0.28229E-02	
16	BHASYS	1.0000	0.11519E-03	
20	ABSNORM90	1.0020	0.40099E-02	
22	LUMIN ₁	481.50	4.3545	
23	LUMIN ₂	519.26	4.5420	
24	LUMIN ₃	447.73	3.8724	
25	LUMIN ₄	3611.4	12.255	
26	LUMIN ₅	553.69	4.1635	
27	LUMIN ₆	592.67	4.7913	
28	LUMIN ₇	636.22	5.1908	

Y-AXIS: PARAMETER 11: DEL3TEC
X=0

S

1.918	77 66 5 44*33: 222 2
1.665	7 6 55 44 33 :22 2
1.411	7 66 55 44 33 22 22
1.157	6 5 4 3* 22 2
0.9034	66 55 44 33*22: 2
0.6496	66 55 44 33 22 : 22 3
0.3959	6 55 44 33 22 111 2 3
0.1421	6 5 4 33 22* 11 1 22 33
-0.1116	-55-44-3-22-111+-1--2-33-
-0.3654	55 44 33 2 11 :11 22 3 4
-0.6191	5 44 33 22 1* :1 22 3344
-0.8729	5 4 33 22 1100 11 2 33 4
-1.127	*44*3*22*1100011*2233*445
-1.380	44 33 2 11 0011:2233 4455
-1.634	4 33 22 1 11 22 3 4455
-1.888	4 3 22 11 11 22 3344 5 6
-2.142	33 2 1 11*22:3344 5566
-2.395	33 22 1111 22 3344 55667
-2.649	3 2 11 22 3344 556677
-2.903	3 22 2233344 556677
-3.157	2 22233 44 556677 8
-3.410	2 22 33*445556677 88
-3.664	2 222 33 4455 6677 889
-3.918	2222 33344455:6677 8899
-4.172	33 44 55 667788899A

I I I
-2.764 -0.5553 1.654

X-AXIS: PARAMETER 10: RHOTE^C ONE COLUMN= 0.1767
FUNCTION VALUES: F(I)= 257.9 + 1.000 *I**2

T

FCN= 259.1442 FROM MINOS STATUS=SUCCESSFUL 126 CALLS 753 TOTAL
EDM= 0.29E-05 STRATEGY= 1 ERROR MATRIX ACCURATE

EXT NO.	PARAMETER NAME	VALUE	PARABOLIC ERROR	MINOS ERRORS
				NEGATIVE POSITIVE
1	ZMASS	91.176	0.99310E-02	
3	NUM COL	3.0086	0.50060E-02	
10	RHOTE ^C	-0.32516	0.98245	

11	DEL3TEC	-0.36085	0.88391	-0.88396	0.88404
13	HADNORM	1.0020	0.40408E-02		
14	ABSNORM	1.0002	0.23955E-02		
15	LEPNORM	0.99916	0.28068E-02		
16	BHASYS	1.0000	0.11516E-03		
20	ABSNORM90	1.0019	0.40011E-02		
22	LUMIN_1	481.55	4.3505		
23	LUMIN_2	519.31	4.5370		
24	LUMIN_3	447.78	3.8658		
25	LUMIN_4	3612.1	12.049		
26	LUMIN_5	553.85	4.1460		
27	LUMIN_6	592.84	4.7738		
28	LUMIN_7	636.39	5.1741		

Y-AXIS: PARAMETER 11: DEL3TEC

X=0

S

1.336	88 7 66 55 44 :33
1.195	8 77 6 55 44*333 2222
1.053	77 66 5 44 33: 222 2
0.9120	77 66 5544 33 :22 2
0.7706	7 66 55 4 33* 22 2 3
0.6291	66 55 44 3 222 2 3
0.4877	6 55 44 3322 : 22 33
0.3463	6655 44 33 2* 1111 2 33
0.2048	6 5 44 33 22*11 11 22 3 4
0.6342E-01	-55-4-33-22-11+-1--2-33-4
-0.7800E-01	55 4433 22 11 :11 22 3 44
-0.2194	5 44 3 22 11* :1 22 3344
-0.3609	544*33*2*11***11*2*33*4*5
-0.5023	4 33 22 1 *11 2233 4455
-0.6437	44 3 22 11 11:22 3 4455
-0.7851	4 33 2 1 11 22 3344 5 6
-0.9266	33 22 11 11*22 3344 5566
-1.068	3 2 1111 22:33 4 5566
-1.209	33 22 22 33 4455 6 7
-1.351	3 2 22*33 4455 6677
-1.492	3 2 22 33:4455 6677
-1.634	2 22 33 44 5 66 7 8
-1.775	2 222 33*44 5566 7788
-1.917	2222 33 44:5566 7788
-2.058	33 44 5566 7788 9

I	I	I
-2.290		1.640

-0.3252

X-AXIS: PARAMETER 10: RHOTEC

ONE COLUMN= 0.1572

FUNCTION VALUES: F(I)= 259.1 + 1.000 *I**2

T