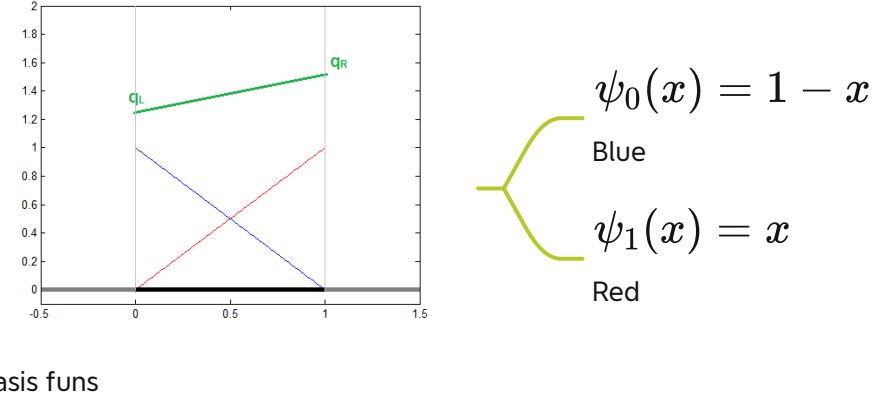


Module 2: A Simple 1D DG Solver

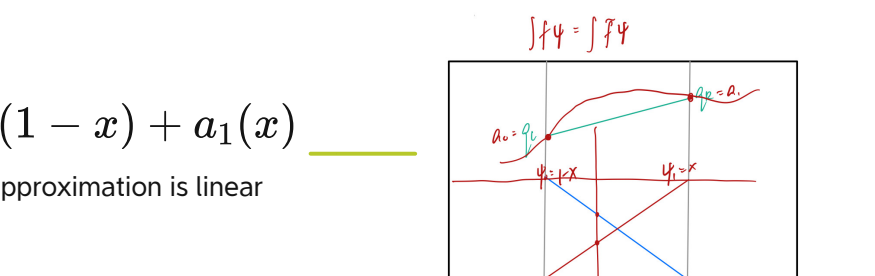
Module 2: A Simple 1D DG Solver

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Linear Solution Approximation: Bases and Resultant Approximation



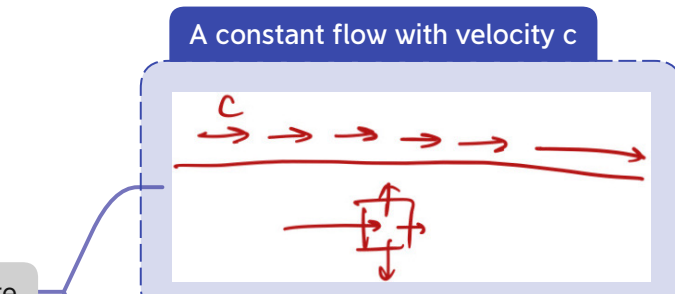
$\psi_0(x) = 1 - x$
Blue
 $\psi_1(x) = x$
Red



$\bar{q} = a_0(1 - x) + a_1(x)$
Resultant approximation is linear

Upwind Flux

We must have... everywhere



$c > 1$
Let's assume that

$c\bar{q}^-$
the numerical flux is then... at all boundaries

Note that there are

two distinct numerical fluxes
 K for element
 $cq^-(x_L)$ on the left it is
 $cq^-(x_R)$ on the right is
 $q^-(x_L) = q^-(x_R)$ where not necessarily

Stiffness Matrix

$$\int_k \frac{\partial q}{\partial t} \phi_j dx + [f(q)\phi_j]_{x_L}^{x_R} - \int_k f(q) \frac{d\phi}{dx} dx = 0$$

Similar to mass matrix...
Substitute our approximation & bases into the 3rd term

The integral for all weighting bases
Combine to form K Stiffness Matrix
we can analytically integrate for this simple case

Similar to the mass matrix...
$$\sum_{i=0}^M ca_i(t) \int_k \psi_i(x) \phi_j'(x) dx = cKa$$

we have a vector of our element approximation coefficients s.t.

Let's substitute in our expressions
for approximation & the 1st weighting basis & analytically integrate
$$\int_k \psi_i(x) \phi_j'(x) dx = \int_0^1 \psi_i(x) \phi_j'(x) dx = \int_0^1 \begin{bmatrix} (1-x) \cdot (-1) & x \cdot (-1) \end{bmatrix} = \begin{bmatrix} -\frac{1}{2} & -\frac{1}{2} \end{bmatrix}$$

Again, the 2'2 matrix

Semi-discrete system

We can now express the semi-discrete system...
$$Ma' - \hat{f} - cKa = 0$$

compactly in matrix/vector form

For each time we have a starting with... from ICs

We'd like to calculate the time rate of change of our coefficients
to use in our time discretization
$$a' = M^{-1}[cKa - \hat{f}]$$

so solve for a'

Investigation- h-Convergence

What do we expect... for h-convergence
How coarse can our elements be and... still get a "good enough" answer
What effects does the shape/smoothness of the solution have on convergence

Investigation- Stability(CFL)

As was seen... too large of a time step lead to issue why?
What happens if a parcel of quantity moves fast enough to skip elements between timesteps?
Methods
Explicit methods $CFL = \frac{c\Delta t}{\Delta x} \leq C_{max} = 1$ need to ensure
Implicit methods relax this restriction to permit higher values

$M = 1$
We must have two basis functions because

We shall choose two ramp functions, let's see what they look like...

V_h
We will choose in this simple example to be the space of all polynomials of order 1 and less

We can think of the approximate solution
...belonging to some finite vector space
This vector space must have a set of basis vectors ψ s.t.
$$\bar{q}(x) = \sum_{i=0}^M a_i \psi_i(x)$$

$q(x)$
The exact sol'n is not usually obtainable
The best we can do is an approximate sol'n
 $\bar{q}(x)$

Linear Solution Approximation: Definition

In order to solve for the approximation
Our weak form sol'n
$$\int \frac{\partial q}{\partial t} \phi_j + \int \frac{\partial f}{\partial x} \phi_j = 0, \forall j$$

must be satisfied for each weighting function

This means our weighting functions are ramp functions as well
 $\phi_0(x) = \psi_0(x)$
 $\phi_1(x) = \psi_1(x)$

We will discuss this in greater detail
but as in a FEM...
..DG chooses the space of weighting function to be the same
 $\phi \in V_h$ as the approximation space

We have now defined our approximation space
But what about the weighting function in the weak form ϕ ?

Test Function Choice (Galerkin)

If the time derivative...
of each of the approximation coefficients
$$\sum_{i=0}^M \frac{da_i(t)}{dt} \int_k \psi_i(x) \phi_j(x) dx = Ma'$$

is a vector a' then

Let's substitute in our expressions
for approximation & the 1st weighting basis & analytically integrate
$$\int_k \psi_0(x) \phi_0(x) dx = (1-x)(1-x) \Big|_0^1 = \frac{1}{3}$$

$$\int_k \psi_1(x) \phi_0(x) dx = x(1-x) \Big|_0^1 = \frac{1}{6}$$

$$\int_k \psi_0(x) \phi_1(x) dx = (1-x)x \Big|_0^1 = \frac{1}{6}$$

$$\int_k \psi_1(x) \phi_1(x) dx = x^2 \Big|_0^1 = \frac{1}{3}$$

$$M_{j,i} = \int \begin{bmatrix} \psi_0 \phi_0 & \psi_1 \phi_0 \\ \psi_0 \phi_1 & \psi_1 \phi_1 \end{bmatrix} = \begin{bmatrix} \frac{1}{3} & \frac{1}{6} \\ \frac{1}{6} & \frac{1}{3} \end{bmatrix}$$

Do this for each weighting basis to get 2'2 matrix:

We now have explicit formula...
$$\int_k \frac{\partial \bar{q}}{\partial t} \phi_j dx = \sum_{i=0}^M \frac{da_i(t)}{dt} \int_k \psi_i(x) \phi_j(x) dx$$

for $\psi_i(q)$ and ϕ_j , substituting

From convention for dynamic systems
the left integral in the weak form is called the:
$$\int_k \frac{\partial \bar{q}}{\partial t} \phi_j dx + [f(\bar{q})\phi_j]_{x_L}^{x_R} - \int_k f(\bar{q}) \frac{d\phi}{dx} dx = 0, \forall j \leq N$$

Mass Matrix

So what do we get?
$$\hat{f} = \begin{bmatrix} -cq_{k-1}(x_R) \\ cq_k(x_R) \end{bmatrix} = c \begin{bmatrix} -q_{k-1}(x_R) \\ q_k(x_R) \end{bmatrix}$$

numerical flux

ϕ_j
...will be zero sometimes depending on j

Just need to calculate now the...
$$[\hat{f}(q)\phi]_{x_L}^{x_R} = cq_k(x_R)\phi_j(x_R) - cq_{k-1}(x_R)\phi_j(x_L)$$

numerical flux vector

Putting it all Together

We must... choose Δt carefully Can you predict why
We may want to... periodically save for plotting
We can overwrite... our previous timestep's coefficients with the new ones to save memory

Let's use a Forward Euler approach!
$$a_{t+1} = a_t + a'_t \Delta t$$

"Very simple" explicit ODE solver

Now we need to complete the method by discretizing in time, the semi-discrete system

Time Discretization: Forward Euler

How do... h convergence t convergence relate?
Do we expect periodic solutions to stay steady over time
How big can we make Δt

Investigation- t-Convergence