#### Department of Information Engineering MSc in Computer Engineering (a.y. 2024/2025) University of Pisa

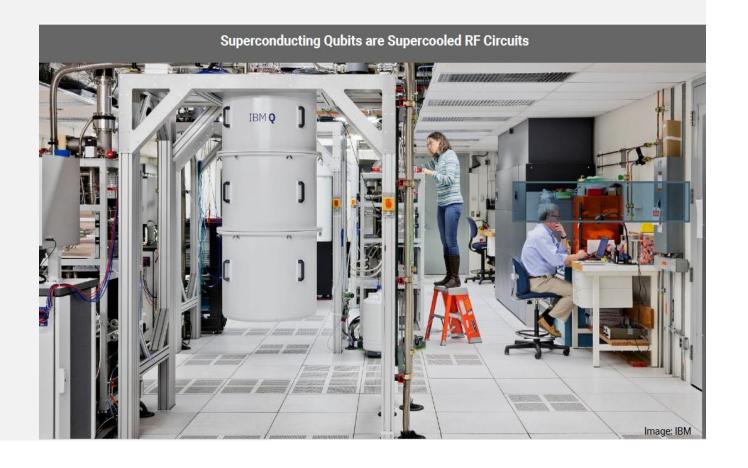
#### Quantum Computing and Quantum Internet

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#### Introduction

- One important thing to remember is that it's impossible to prepare the exact desired pure state with complete accuracy
- Sometimes, the prepared state is pure but doesn't quite match the target state we intended
- At other times, it may be affected by incoherent noise, resulting in a mixed state
- Therefore, it is essential to have a method for quantifying the difference between the actual state and the desired (or target) state
- A convenient tool that tells us how close the real state is to the target state is the **Fidelity**

- The fidelity of states  $\rho$  and  $\sigma$  is defined to be

$$F(\rho,\sigma) = \operatorname{tr}\sqrt{\rho^{1/2}\sigma\rho^{1/2}} \tag{1}$$

- It is certainly not immediately obvious that this is a useful measure of distance between  $\rho$  and  $\sigma$
- It doesn't even look symmetric!
- There is (at least) one important special case where it is possible to give more explicit formulae for the fidelity

- Specifically, let's calculate the fidelity between a pure state  $|\psi\rangle$  (with density matrix  $\sigma=|\psi\rangle\langle\psi|$  ) and an arbitrary state,  $\rho$
- From Equation (1) we see that

$$F(|\psi\rangle, \rho) = \operatorname{tr}\sqrt{\sqrt{\sigma}\rho\sqrt{\sigma}} = \operatorname{tr}\sqrt{\sqrt{|\psi\rangle\langle\psi|\rho\sqrt{|\psi\rangle\langle\psi|}}} \tag{6}$$

- However, from  $\sigma = |\psi\rangle\langle\psi|$  it follows:

$$\sigma^2 = (|\psi\rangle\langle\psi|)^2 = (|\psi\rangle\langle\psi|)(|\psi\rangle\langle\psi|) = |\psi\rangle\langle\psi|\psi\rangle\langle\psi| = |\psi\rangle\langle\psi| = \sigma$$

- Thus

$$\sigma^2 = \sigma \tag{7}$$

- Substituting (7) into (6) leads to the following result

$$F(|\psi\rangle, \rho) = \operatorname{tr}\sqrt{\sqrt{\sigma^2}\rho\sqrt{\sigma^2}} = \operatorname{tr}\sqrt{\sigma\rho\sigma} = \operatorname{tr}\sqrt{|\psi\rangle\langle\psi|\rho|\psi\rangle\langle\psi|}$$
(8)

- Since  $\langle \psi | \rho | \psi \rangle$  is a scalar number and since  $(|\psi\rangle\langle\psi|)^2 = |\psi\rangle\langle\psi|$ , (8) can be rewritten as follows:

$$F(|\psi\rangle,\rho) = \operatorname{tr}\sqrt{\langle\psi|\rho|\psi\rangle|\psi\rangle\langle\psi|} = \operatorname{tr}\sqrt{\langle\psi|\rho|\psi\rangle(|\psi\rangle\langle\psi|)^{2}} = \operatorname{tr}\left(\sqrt{\langle\psi|\rho|\psi\rangle}(|\psi\rangle\langle\psi|)\right)$$
$$= \sqrt{\langle\psi|\rho|\psi\rangle}\operatorname{tr}\left((|\psi\rangle\langle\psi|)\right)$$

- Since  $\operatorname{tr}(|\psi\rangle\langle\psi|)=1$ , we can conclude that

$$F(|\psi\rangle,\rho) = \sqrt{\langle\psi|\rho|\psi\rangle} \tag{9}$$

#### Pure-State Fidelity

- The fidelity measure has a simple interpretation when both  $\rho = |\phi\rangle\langle\phi|$  and  $\sigma = |\psi\rangle\langle\psi|$  are pure states
- In this specific case, the pure-state fidelity  $F(|\psi\rangle,|\phi\rangle)$  is a measure of how  $|\psi\rangle$  is close to  $|\phi\rangle$

$$F\left(|\psi\rangle,|\phi\rangle\right) = \sqrt{\langle\psi\left|\left(|\phi\rangle\langle\phi\right|\right)|\psi\rangle} = \sqrt{\langle\psi\left|\phi\rangle\langle\phi\right|\psi\rangle} = \sqrt{\langle\psi\left|\phi\rangle\langle\psi\right|\phi\rangle^*} = \sqrt{\left|\langle\psi\left|\phi\rangle\right|^2} = \left|\langle\psi\left|\phi\rangle\right|$$

which is the overlap of the states  $|\psi
angle$  and  $|\phi
angle$ 

#### Pure-State Fidelity

- The **pure-state** fidelity is *symmetric* 

$$F(|\psi\rangle,|\phi\rangle) = |\langle\psi|\phi\rangle|^2 = |\langle\phi|\psi\rangle|^2 = F(|\phi\rangle,|\psi\rangle)$$

and it obeys the following bounds:

$$0 \le F(|\psi\rangle, |\phi\rangle) = |\langle\psi|\phi\rangle|^2 \le 1$$

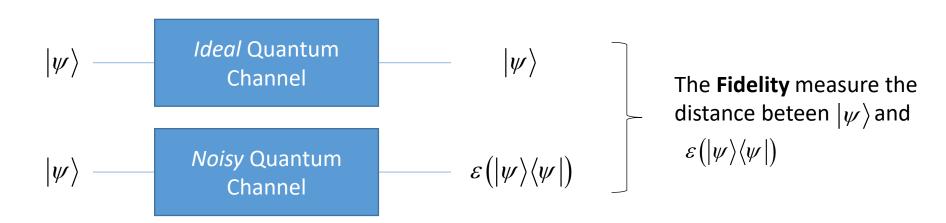
- It is equal to one if and only if the two states are the same, and it is equal to zero if and only if the two states are orthogonal to each other
- The fidelity measure is not a distance measure in the strict mathematical sense because it is equal to one when two states are equal, whereas a distance measure should be equal to zero when two states are equal

- So far, we have primarily considered qubits that evolve uncorrupted in time and space
- In other words, these, ideal, qubits navigate noisless communication channels
- In the real world, qubits and their classical cousins, bits, are regularly subjected to corrupting environmental perturbations

- Suppose a quantum system is in the state  $|\psi\rangle$  and some physical process occurs, changing the quantum system to the state  $\varepsilon(|\psi\rangle\langle\psi|)$
- How well has the physical process preserved the state  $|\psi\rangle$  of the quantum system?
- The fidelity discussed previously can be used now to develop measures of how well a quantum system preserves information

- This type of scenario occurs often in *quantum computation* and *quantum information*
- For example, in the memory of a quantum computer,  $|\psi\rangle$  is the initial state of the memory, and  $\varepsilon$  represents the dynamics that the memory undergoes, including noise processes arising from interaction with the environment

- A second example is provided by a quantum communication channel for transmitting the state  $|\psi\rangle$  from one location to another
- No channel is ever perfect, so the action of the channel is described by a quantum operation  ${\cal E}$



- An obvious way of quantifying how well the state  $|\psi\rangle$  is preserved by the channel is to make use of the **Fidelity** measure introduced earlier
- For example, we can compute the fidelity between the starting state  $|\psi\rangle$  and the ending state  $\varepsilon(|\psi\rangle\langle\psi|)$
- For the case of the depolarizing channel, we obtain

$$F(|\psi\rangle,\varepsilon(|\psi\rangle\langle\psi|)) = \sqrt{\langle\psi|\left(p\frac{I}{2} + (1-p)|\psi\rangle\langle\psi|\right)|\psi\rangle} = \sqrt{\langle\psi|p\frac{I}{2}|\psi\rangle + \langle\psi|(1-p)|\psi\rangle\langle\psi|\psi\rangle}$$

$$= \sqrt{p\frac{1}{2}\langle\psi|\psi\rangle + \langle\psi|\psi\rangle\langle\psi|\psi\rangle - p\langle\psi|\psi\rangle\langle\psi|\psi\rangle} = \sqrt{p\frac{1}{2} + 1 - p}$$

$$= \sqrt{1 - \frac{p}{2}}$$

- Thus

$$F(|\psi\rangle,\varepsilon(|\psi\rangle\langle\psi|)) = \sqrt{1-\frac{p}{2}}$$

- This result agrees well with our intuition the higher the probability *p* of depolarizing, the lower the fidelity of the final state with the initial state
- Provided p is very small the **Fidelity** is close to one, and the state  $\varepsilon(|\psi\rangle\langle\psi|)$  is practically indistinguishable from the initial state  $|\psi\rangle$

- Nevertheless, the fidelity  $F(|\psi\rangle,\varepsilon(|\psi\rangle\langle\psi|))$  has some drawbacks which need to be remedied
- In a real quantum memory or quantum communications channel, we don't know in advance what the initial state  $|\psi\rangle$  of the system will be
- However, we can quantify the worst-case behavior of the system by minimizing over all possible initial states

$$F_{\min}\left(\varepsilon\right) = \min_{|\psi\rangle} F\left(|\psi\rangle, \varepsilon\left(|\psi\rangle\langle\psi|\right)\right)$$

- For example, for the *p*-depolarizing channel

$$F_{\min} = \sqrt{1 - \frac{p}{2}}$$

as the fidelity of the channel is the same for all input states  $|\psi
angle$ 

- A more interesting example is the phase damping channel

$$\varepsilon(\rho) = p\rho + (1-p)Z\rho Z$$

### How Well Does a Quantum Channel Preserve Information? $\varepsilon(\rho) = p\rho + (1-p)Z\rho Z$

- For the phase damping channel, the fidelity is given by

$$F(|\psi\rangle,\varepsilon(|\psi\rangle\langle\psi|)) = \sqrt{\langle\psi|(p|\psi\rangle\langle\psi|+(1-p)Z|\psi\rangle\langle\psi|Z)|\psi\rangle}$$

$$= \sqrt{\langle\psi|p|\psi\rangle\langle\psi|\psi\rangle+\langle\psi|(1-p)Z|\psi\rangle\langle\psi|Z|\psi\rangle}$$

$$= \sqrt{p\langle\psi|\psi\rangle\langle\psi|\psi\rangle+(1-p)\langle\psi|Z|\psi\rangle\langle\psi|Z|\psi\rangle}$$

$$= \sqrt{p+(1-p)\langle\psi|Z|\psi\rangle^{2}}$$

# How Well Does a Quantum Channel Preserve Information? $\varepsilon(\rho) = p\rho + (1-p)Z\rho Z$

- The second term under the square root sign is non-negative, and equal to zero when

$$|\psi\rangle = (|0\rangle + |1\rangle)/\sqrt{2} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ 1 \end{bmatrix}$$

- In fact,

$$\langle \psi | Z | \psi \rangle = \frac{1}{2} \begin{bmatrix} 1 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \begin{bmatrix} 1 \\ 1 \end{bmatrix} = \frac{1}{2} \begin{bmatrix} 1 & 1 \end{bmatrix} \begin{bmatrix} 1 \\ -1 \end{bmatrix} = \frac{1}{2} (1-1) = 0$$

- Thus, for the phase damping channel the minimum fidelity is

$$F_{\min} = \sqrt{p}$$

- You might wonder why we minimized over pure states in the definition of  $F_{min}$
- After all, might not the quantum system of interest start in a **mixed state**  $\rho$ ?
- For example, a quantum memory might be entangled with the rest of the quantum computer, and therefore would start out in a mixed state
- Fortunately, the **joint concavity of the fidelity** (not proved in the lecture) can be used to show that allowing mixed states does not change  $F_{min}$
- Suppose that  $\rho = \sum_i \lambda_i |i\rangle\langle i|$  is the initial state of the quantum system, it can be proved that

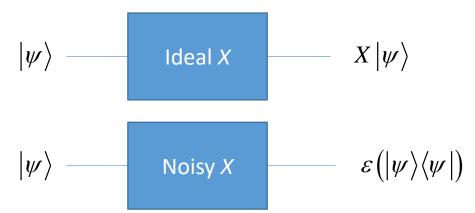
$$F \rho, \varepsilon \rho \geq F_{\min}$$

- Of course, we are interested not only in protecting quantum states as they are transmitted through a quantum communications channel, but also as they dynamically undergo computation
- Suppose, for example, that as part of a quantum computation we attempt to implement a quantum gate described by the unitary operator *U*
- As described in a previous lecture, any such attempt will inevitably encounter some (hopefully not too severe) noise, so the correct description of the gate is using a quantum operation  $\varepsilon$

- A natural measure of how successful our gate has been is the gate fidelity,

$$F \ U,\varepsilon \equiv \min_{|\psi\rangle} F \ U|\psi\rangle,\varepsilon \ |\psi\rangle\langle\psi|$$

- Suppose, for example, that we try to implement an X gate on a single qubit, but instead implement the noisy operation  $\varepsilon \rho = 1 - p X \rho X + p Z \rho Z$ , for some small noise parameter p



- Then the gate fidelity for this operation is given by

$$F X, \varepsilon \equiv \min_{|\psi\rangle} \sqrt{\langle \psi | X} \quad 1 - p \ X \rho X + p Z \rho Z \ X | \psi\rangle$$

$$= \min_{|\psi\rangle} \sqrt{\langle \psi | X} \quad 1 - p \ X | \psi\rangle \langle \psi | X + p Z | \psi\rangle \langle \psi | Z \ X | \psi\rangle$$

$$= \min_{|\psi\rangle} \sqrt{\langle \psi | X} \quad 1 - p \ X | \psi\rangle \langle \psi | X X | \psi\rangle + p Z | \psi\rangle \langle \psi | Z X | \psi\rangle$$

$$= \min_{|\psi\rangle} \sqrt{1 - p \ \langle \psi | X X | \psi\rangle \langle \psi | X X | \psi\rangle + p \langle \psi | X Z | \psi\rangle \langle \psi | Z X | \psi\rangle}$$

- Since  $X^2 = I \rightarrow \langle \psi | XX | \psi \rangle = 1$ 

- Keeping also into consideration that ZX = -XZ the gate fidelity for this operation is given by

$$F X, \varepsilon = \min_{|\psi\rangle} \sqrt{1 - p \langle \psi | XX | \psi \rangle \langle \psi | XX | \psi \rangle + p \langle \psi | XZ | \psi \rangle \langle \psi | ZX | \psi \rangle}$$
$$= \min_{|\psi\rangle} \sqrt{1 - p - p \langle \psi | XZ | \psi \rangle^{2}}$$

- Since XZ = -iY

$$F X, \varepsilon = \min_{|\psi\rangle} \sqrt{1 - p \langle \psi | XX | \psi \rangle \langle \psi | XX | \psi \rangle + p \langle \psi | XZ | \psi \rangle \langle \psi | ZX | \psi \rangle}$$
$$= \min_{|\psi\rangle} \sqrt{1 - p \langle \psi | Y | \psi \rangle^{2}} = \sqrt{1 - p}$$