

PH 108 Formulae

Advait Risbud
Jujhaar Singh
Kartik Gokhale

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MATHEMATICAL PRELIMINARIES

- **Gradients.**

1. Consider a scalar field T in 3-dimensions. If we move in the direction denoted by \vec{dl} , then the change in the value of the scalar field T is denoted by

$$dT = \nabla T \cdot \vec{dl} \quad (1)$$

2. Gradient gives the direction of largest increase in the value of the scalar field.

- **Properties of Grad, Div, and Curl.**

1. $\nabla(f + g) = \nabla f + \nabla g$
2. $\nabla \cdot (\vec{A} + \vec{B}) = \nabla \cdot \vec{A} + \nabla \cdot \vec{B}$
3. $\nabla \times (\vec{A} + \vec{B})$
4. $\nabla(kf) = k(\nabla f)$
5. $\nabla \cdot (k\vec{A}) = k(\nabla \cdot \vec{A})$
6. $\nabla \times (k\vec{A}) = k(\nabla \times \vec{A})$
7. $\nabla(fg) = f(\nabla g) + g(\nabla f)$
8. $\nabla(\vec{A} \cdot \vec{B}) = \vec{A} \times \nabla \times \vec{B} + \vec{B} \times \nabla \times \vec{A} + (\vec{A} \cdot \nabla)\vec{B} + (\vec{B} \cdot \nabla)\vec{A}$
9. $\nabla \cdot (f\vec{A}) = f(\nabla \cdot \vec{A}) + \vec{A} \cdot \nabla f$
10. $\nabla \cdot \vec{A} \times \vec{B} = \vec{B} \cdot (\nabla \times \vec{A}) - \vec{A} \cdot (\nabla \times \vec{B})$
11. $\nabla \times (f\vec{A}) = f(\nabla \times \vec{A}) - \vec{A} \times \nabla f$
12. $\nabla \times (\vec{A} \times \vec{B}) = (\vec{B} \cdot \nabla)\vec{A} - (\vec{A} \cdot \nabla)\vec{B} + \vec{A}(\nabla \cdot \vec{B}) - \vec{B}(\nabla \cdot \vec{A})$

Double derivatives

1. Divergence of the gradient(**laplacian**).

$$\nabla \cdot (\nabla f) = \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} + \frac{\partial^2 f}{\partial z^2}$$

We get the laplacian operator from here as follows,

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \quad (2)$$

2. Curl of a gradient

$$\nabla \times (\nabla f) = 0 \quad (3)$$

3. Gradient of the divergence

$$\nabla(\nabla \cdot \vec{A}) \quad (4)$$

4. Divergence of a curl

$$\nabla \cdot (\nabla \times \vec{A}) = 0 \quad (5)$$

5. Curl of the curl

$$\nabla \times (\nabla \times \vec{A}) = \nabla(\nabla \cdot \vec{A}) - \nabla^2 \vec{A} \quad (6)$$

- **Co-ordinate systems**

1. **Relations between magnitudes**

2-D:

$$x = r \cos \phi \quad (7)$$

$$y = r \sin \phi \quad (8)$$

3-D:

Spherical:

$$x = r \sin \theta \cos \phi \quad (9)$$

$$y = r \sin \theta \sin \phi \quad (10)$$

$$z = r \cos \theta \quad (11)$$

Cylindrical:

$$x = s \cos \phi \quad (12)$$

$$y = s \sin \phi \quad (13)$$

$$z = z \quad (14)$$

2. Unit vectors

In general, the unit vector is the direction of increase in that variable.

2-D:

$$\begin{pmatrix} \hat{r} \\ \hat{\theta} \end{pmatrix} = \begin{pmatrix} \cos \phi & \sin \phi \\ -\sin \phi & \cos \phi \end{pmatrix} \begin{pmatrix} \hat{i} \\ \hat{j} \end{pmatrix} \quad (15)$$

3-D:

Spherical:

$$\begin{pmatrix} \hat{r} \\ \hat{\phi} \\ \hat{\theta} \end{pmatrix} = \begin{pmatrix} \sin \theta \cos \phi & \sin \theta \sin \phi & \cos \theta \\ -\sin \phi & \cos \phi & 0 \\ \cos \theta \cos \phi & \cos \theta \sin \phi & -\sin \theta \end{pmatrix} \begin{pmatrix} \hat{i} \\ \hat{j} \\ \hat{k} \end{pmatrix} \quad (16)$$

Cylindrical:

$$\begin{pmatrix} \hat{s} \\ \hat{\phi} \\ \hat{z} \end{pmatrix} = \begin{pmatrix} \sin \phi & \cos \phi & 0 \\ -\sin \phi & \cos \phi & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} \hat{i} \\ \hat{j} \\ \hat{k} \end{pmatrix} \quad (17)$$

The polar co-ordinate system is orthogonal.

3. Distance elements

2-D:

$$d\vec{l} = dr\hat{r} + r d\theta\hat{\theta} \quad (18)$$

3-D:

Spherical:

$$d\vec{l} = dr\hat{r} + r d\theta\hat{\theta} + r \sin\theta d\phi\hat{\phi} \quad (19)$$

Cylindrical:

$$d\vec{l} = ds\hat{s} + s d\phi\hat{\phi} + dz\hat{z} \quad (20)$$

- **Operators in all co-ordinate systems**

Cartesian:

$$\text{Gradient: } \nabla T = \left(\frac{\partial T}{\partial x}\right)\hat{x} + \left(\frac{\partial T}{\partial y}\right)\hat{y} + \left(\frac{\partial T}{\partial z}\right)\hat{z} \quad (21)$$

$$\text{Divergence: } \nabla \cdot v = \frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} + \frac{\partial v_z}{\partial z} \quad (22)$$

$$\text{Curl: } \nabla \times v = \left(\frac{\partial v_z}{\partial y} - \frac{\partial v_y}{\partial z}\right)\hat{x} + \left(\frac{\partial v_x}{\partial z} - \frac{\partial v_z}{\partial x}\right)\hat{y} + \left(\frac{\partial v_y}{\partial x} - \frac{\partial v_x}{\partial y}\right)\hat{z} \quad (23)$$

$$\text{Laplacian: } \nabla^2 v = \frac{\partial^2 v_x}{\partial x^2} + \frac{\partial^2 v_y}{\partial y^2} + \frac{\partial^2 v_z}{\partial z^2} \quad (24)$$

Spherical and Cylindrical:

Gradient:

$$\nabla T = \frac{\partial T}{\partial r} \hat{\mathbf{r}} + \frac{1}{r} \frac{\partial T}{\partial \theta} \hat{\boldsymbol{\theta}} + \frac{1}{r \sin \theta} \frac{\partial T}{\partial \phi} \hat{\boldsymbol{\phi}}. \quad (1.70)$$

Divergence:

$$\nabla \cdot \mathbf{v} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 v_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta v_\theta) + \frac{1}{r \sin \theta} \frac{\partial v_\phi}{\partial \phi}. \quad (1.71)$$

Curl:

$$\begin{aligned} \nabla \times \mathbf{v} = & \frac{1}{r \sin \theta} \left[\frac{\partial}{\partial \theta} (\sin \theta v_\phi) - \frac{\partial v_\theta}{\partial \phi} \right] \hat{\mathbf{r}} + \frac{1}{r} \left[\frac{1}{\sin \theta} \frac{\partial v_r}{\partial \phi} - \frac{\partial}{\partial r} (r v_\phi) \right] \hat{\boldsymbol{\theta}} \\ & + \frac{1}{r} \left[\frac{\partial}{\partial r} (r v_\theta) - \frac{\partial v_r}{\partial \theta} \right] \hat{\boldsymbol{\phi}}. \end{aligned} \quad (1.72)$$

Laplacian:

$$\nabla^2 T = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial T}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial T}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 T}{\partial \phi^2}. \quad (1.73)$$

Gradient:

$$\nabla T = \frac{\partial T}{\partial s} \hat{\mathbf{s}} + \frac{1}{s} \frac{\partial T}{\partial \phi} \hat{\boldsymbol{\phi}} + \frac{\partial T}{\partial z} \hat{\mathbf{z}}. \quad (1.79)$$

Divergence:

$$\nabla \cdot \mathbf{v} = \frac{1}{s} \frac{\partial}{\partial s} (s v_s) + \frac{1}{s} \frac{\partial v_\phi}{\partial \phi} + \frac{\partial v_z}{\partial z}. \quad (1.80)$$

Curl:

$$\nabla \times \mathbf{v} = \left(\frac{1}{s} \frac{\partial v_z}{\partial \phi} - \frac{\partial v_\phi}{\partial z} \right) \hat{\mathbf{s}} + \left(\frac{\partial v_s}{\partial z} - \frac{\partial v_z}{\partial s} \right) \hat{\boldsymbol{\phi}} + \frac{1}{s} \left[\frac{\partial}{\partial s} (s v_\phi) - \frac{\partial v_s}{\partial \phi} \right] \hat{\mathbf{z}}. \quad (1.81)$$

Laplacian:

$$\nabla^2 T = \frac{1}{s} \frac{\partial}{\partial s} \left(s \frac{\partial T}{\partial s} \right) + \frac{1}{s^2} \frac{\partial^2 T}{\partial \phi^2} + \frac{\partial^2 T}{\partial z^2}. \quad (1.82)$$

Theorems of Vector Calculus

1. FTOVC

$$\int_a^b \nabla \vec{v} \cdot d\vec{\mathbf{l}} = \vec{v}(b) - \vec{v}(a) \quad (25)$$

Corollary:

$$\oint_C \nabla \vec{v} \cdot d\vec{\mathbf{l}} = 0 \quad (26)$$

2. Gauss's/Divergence Theorem

$$\oint_S \vec{v} \cdot d\vec{\mathbf{a}} = \int_V (\nabla \cdot \vec{v}) d\tau \quad (27)$$

This states that the flux of a vector field \vec{v} through a closed surface S is equal to the volume integral of the divergence of that field over the volume that S encloses.

3. Stoke's Theorem

$$\oint_{\delta S} \vec{v} \cdot d\vec{\mathbf{l}} = \int_S (\nabla \times \vec{v}) \cdot d\vec{\mathbf{a}} \quad (28)$$

Where δS is the boundary of the surface S .

- **Dirac Delta Function**

1-D:

$$\delta(x - a) = \begin{cases} 0, & \text{for } x \neq a \\ \infty, & \text{for } x = a \end{cases}$$

Integral of the delta function,

$$\int_b^c \delta(x - a) dx = \begin{cases} 1, & \text{for } b \leq a \leq c \\ 0, & \text{otherwise} \end{cases}.$$

Delta function picks out values as,

$$\int_b^c f(x) \delta(x - a) dx = \begin{cases} f(a), & \text{for } b \leq a \leq c \\ 0, & \text{otherwise} \end{cases}.$$

2 delta functions $D_1(x)$ and $D_2(x)$ are considered equal (for all ordinary $f(x)$) if

$$\int_{-\infty}^{\infty} f(x) D_1(x) dx = \int_{-\infty}^{\infty} f(x) D_2(x) dx \quad (29)$$

Important result:

$$\delta(kx) = \frac{1}{|k|} \delta(x) \quad (30)$$

3-D:

Here \mathbf{r} is the position vector.

$$\delta^3(\mathbf{r}) = \delta(x) \delta(y) \delta(z) \quad (31)$$

Analogous to the 1-D case,

$$\int_{\text{all space}} \delta^3(\mathbf{r}) dx dy dz = 1 \quad (32)$$

$$\int_{\text{all space}} f(\mathbf{r}) \delta^3(\mathbf{r} - \mathbf{a}) dx dy dz = f(\mathbf{a}). \quad (33)$$

Also,

$$\int_V f(\mathbf{r}) \delta^3(\mathbf{r} - \mathbf{a}) dx dy dz = \begin{cases} f(\mathbf{a}), & \text{if } \mathbf{a} \in V \\ 0, & \text{if } \mathbf{a} \notin V \end{cases} \quad (34)$$

Important:

$$\nabla \cdot \left(\frac{\hat{r}}{r^2} \right) = 4\pi\delta^3(\vec{r}) \quad (35)$$

- **Helmholtz' Theorem**

Consider a vector field $\vec{F}(\vec{r})$ which satisfies:

1. $\nabla \cdot \vec{F} = D(\vec{r})$
2. $\nabla \times \vec{F} = \vec{C}(\vec{r})$
3. $\nabla \cdot \vec{C} = 0$
4. Diminishes to zero as $\vec{r} \rightarrow \infty$

Then we can write the field as the sum of divergence-free(solenoidal) and curl-free(irrotational) fields as follows,

$$\vec{F}(\vec{r}) = -\nabla U(\vec{r}) + \nabla \times \vec{W}(\vec{r}) \quad (36)$$

where $U(\vec{r})$ and $\vec{W}(\vec{r})$ are the scalar and vector potentials respectively, and are given by

$$U(\vec{r}) = \frac{1}{4\pi} \int \frac{D(\vec{r}')}{|\vec{r} - \vec{r}'|} d\tau' \quad (37)$$

$$\vec{W}(\vec{r}) = \frac{1}{4\pi} \int \frac{\vec{C}(\vec{r}')}{|\vec{r} - \vec{r}'|} d\tau'. \quad (38)$$

Note: integrals are w.r.t \vec{r}' where \vec{r}' is the location of the source/sink of the field and differential operators are w.r.t \vec{r} where \vec{r} is the location where field/potential is to be calculated.

ELECTROSTATICS

- Maxwell's equations for electrostatic fields:

$$\nabla \cdot \vec{E} = \frac{\rho(\vec{r})}{\epsilon_0} \quad (39)$$

$$\nabla \times \vec{E} = 0 \quad (40)$$

- Coulomb's law:

$$\vec{E}(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r}')}{z^2} \hat{z} d\tau' \quad (41)$$

Where $\vec{z} = \vec{r} - \vec{r}'$ is the separation vector between the source charge(at \vec{r}') to where the field is being calculated(\vec{r}).

- Electrostatic potential:

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r}')}{|\vec{r} - \vec{r}'|} d\tau' \quad (42)$$

Poisson's equation:

$$\nabla^2 V = -\frac{\rho(\vec{r})}{\epsilon_0} \quad (43)$$

In regions of zero charge density we get Laplace's equation:

$$\nabla^2 V = 0 \quad (44)$$

2.1 BOUNDARY CONDITIONS:

– Gauss's Law:

$$\oint \vec{\mathbf{E}} \cdot d\vec{\mathbf{a}} = \frac{Q_{en}}{\epsilon_0} \quad (45)$$

– BCs for fields:

$$E_{above}^{\perp} - E_{below}^{\perp} = \frac{\sigma}{\epsilon_0} \quad (46)$$

$$E_{above}^{\parallel} = E_{below}^{\parallel} \quad (47)$$

$$\vec{\mathbf{E}}_{above} - \vec{\mathbf{E}}_{below} = \frac{\sigma}{\epsilon_0} \hat{\mathbf{n}} \quad (48)$$

– Potential BCs:

$$V_{above} = V_{below} \quad (49)$$

$$E^{\perp} = -\vec{\mathbf{E}} \cdot \hat{\mathbf{n}} \quad (50)$$

$$= -\nabla V \cdot \hat{\mathbf{n}} \quad (51)$$

$$= -\frac{\partial V}{\partial \mathbf{n}} \quad (52)$$

2.2 GENERAL FORMULAE

- Energy density:

Here W=work done in assembling charges in this configuration.

$$W = \int \frac{1}{2} \epsilon E^2 d\tau \quad (53)$$

Where u_E =energy density of the field,

$$u_E = \frac{1}{2} \epsilon E^2 \quad (54)$$

- Field at the surface of a conductor:

$$\vec{\mathbf{E}} = \frac{\sigma}{\epsilon_0} \hat{\mathbf{n}} \quad (55)$$

- Electrostatic pressure:

$$P = \frac{\sigma^2}{2\epsilon_0} \quad (56)$$

- Capacitance

–

$$Q = CV$$

- Capacitance of a spherical conductor of radius R :

$$C = 4\pi\epsilon_0 R$$

POTENTIALS

3.1 SOLUTIONS TO LAPLACE EQUATION

1. In 1-D:

$$\frac{d^2V(x)}{dx^2} = 0 \quad (57)$$

$$V(x) = mx + c \quad (58)$$

$$\implies V(x) = \frac{1}{2}[V(x+a) + V(x-a)] \quad (59)$$

The last equation allows us to conclude that there can be no maximum or minimum within the region of interest i.e. the domain of x .

2. In 2-D:

$$\frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} = 0 \quad (60)$$

$$V(x, y) = \frac{1}{2\pi R} \oint_C V(x, y) dl \quad (61)$$

Again there can be no maxima or minima in the domain of interest, only at the boundary.

3. In 3-D:

Let the domain under consideration be a region of volume V enclosed by a surface of area S .

$$\frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = 0 \quad (62)$$

$$V(x, y, z) = \frac{1}{4\pi R^2} \oint_S V(x, y, z) da \quad (S \text{ is a sphere here}) \quad (63)$$

No maxima or minima can exist in this volume. Earnshaw's theorem follows from this fact as no charge configuration can give a potential distribution with a point

of stable equilibrium hence a charged particle cannot be trapped by other fixed charges.

3.2 UNIQUENESS THEOREMS

Theorems

1. First uniqueness theorem: The solution to LE in some volume V is uniquely determined if the potential is specified on the boundary surface.
2. Second uniqueness theorem: In a volume V surrounded by conductors and containing a specified charge density ρ , the electric field is uniquely determined if the total charge on each conductor is given.

The domain definition means that the conductors are not a part of the domain and the charge density is specified everywhere else (see diagram in Griffiths).

3.3 METHOD OF IMAGES

- Image charge replicates potential distribution due to conducting surface.
- Image charge cannot be in the region where original charge is present.
- Image of image charge must be taken if multiple surfaces are present.
- Combination of image charge and original charge gives potential distribution in region due to equivalent conducting plane and charge.
- Energy of infinite conducting sheet and charge ' q '.

$$U = \frac{1}{2} \left(- \frac{kq^2}{2d} \right) \quad (64)$$

3.4 SEPARATIONS OF VARIABLES (CARTESIAN)

The Laplacian in Cartesian co-ordinates is

$$\frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = 0 \quad (65)$$

We take

$$V(x, y, z) = X(x)Y(y)Z(z) \quad \text{hence (65) becomes} \quad (66)$$

$$\frac{1}{X} \frac{\partial^2 X}{\partial x^2} + \frac{1}{Y} \frac{\partial^2 Y}{\partial y^2} + \frac{1}{Z} \frac{\partial^2 Z}{\partial z^2} = 0 \quad (67)$$

Since all three terms are in independent variables, each of them must be a constant such that their sum always equals zero. If variation is in all three directions i.e. none of the 3 terms are zero then the positive one can be taken as $k^2 + l^2$ and in case of 2-D variation the positive one can be taken as k^2 .

Which derivative should be taken positive and which negative is best seen through examples in Griffiths or tutorials. Some mathematical ideas however are important.

3.4.1 Common solutions of DEs

i

$$\frac{1}{X} \frac{d^2 X}{dx^2} = k^2 \quad (68)$$

$$\implies X(x) = Ae^{kx} + Be^{-kx} \quad (69)$$

ii

$$\frac{1}{X} \frac{d^2 X}{dx^2} = -k^2 \quad (70)$$

$$X(x) = A \sin(kx) + B \cos(kx) \quad (71)$$

iii Hyperbolic trig functions:

$$\text{a } \sinh(x) = \frac{e^x - e^{-x}}{2}$$

$$\text{b } \cosh(x) = \frac{e^x + e^{-x}}{2}$$

3.4.2 Fourier's trick

As far as sinusoidal functions are concerned the trick involves finding the coefficients of a Fourier series using orthogonality of the sine functions. If

$$V(y) = \sum_{n=1}^{\infty} C_n \sin\left(\frac{n\pi y}{a}\right) \quad \text{we get,} \quad (72)$$

$$\int_0^a V(y) \sin\left(\frac{n'\pi y}{a}\right) dy = \sum_{n=1}^{\infty} C_n \int_0^a \sin\left(\frac{n\pi y}{a}\right) \sin\left(\frac{n'\pi y}{a}\right) dy \quad (73)$$

From the orthogonality of the sine function we have

$$\int_0^a \sin\left(\frac{n'\pi y}{a}\right) \sin\left(\frac{n\pi y}{a}\right) dy = \begin{cases} 0 & \text{if } n' \neq n \\ \frac{a}{2} & \text{if } n' = n \end{cases} \quad (74)$$

So basically if we want the n^{th} coefficient of the sine series then we should just multiply by the n^{th} sine function to both sides and integrate within the limits of the variable. Here I have taken the boundaries to be at 0 and a because that is the principal range for $\cos\left(\frac{n\pi y}{a}\right)$ (the choice of $V(y)$ and sine functions is all inspired by the topic at hand).

3.5 SEPARATION OF VARIABLES (SPHERICAL POLAR)

The Laplacian in spherical co-ordinates is:

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) + \frac{1}{r^2 \sin(\theta)} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial V}{\partial \theta} \right) + \frac{1}{r^2 \sin^2(\theta)} \frac{\partial^2 V}{\partial \phi^2} = 0 \quad (75)$$

while considering problems with azimuthal symmetry the equation reduces to:

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) + \frac{1}{r^2 \sin(\theta)} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial V}{\partial \theta} \right) = 0 \quad (76)$$

Once again we substitute V as $V(r, \theta) = R(r)P(\theta)$ and divide the Laplacian by V to get:

$$\frac{1}{R} \frac{\partial}{\partial r} \left(r^2 \frac{\partial R}{\partial r} \right) + \frac{1}{P \sin(\theta)} \frac{\partial}{\partial \theta} \left(\sin(\theta) \frac{\partial P}{\partial \theta} \right) = 0 \quad (77)$$

3.5.1 Solution to radial part

$$R(r) = A_l r^l + \frac{B_l}{r^{l+1}} \quad (78)$$

3.5.2 Solution to polar part

$$P(\theta) = P_l(\cos \theta) \quad (79)$$

Here $P_l(\cos \theta)$ is the l^{th} Legendre Polynomial. These can be found using Rodrigue's Formula:

$$P_l(x) = \frac{1}{2^l l!} \frac{d^l (x^2 - 1)^l}{dx^l} \quad \text{where } x = \cos \theta \quad (80)$$

Some of them are:

$$P_0(x) = 1$$

$$P_1(x) = x$$

$$P_2(x) = \frac{3x^2 - 1}{2}$$

$$P_3(x) = \frac{5x^3 - 3x}{2}$$

Legendre polynomials are also orthogonal and [Fourier's trick](#) is applicable here as well.

$$\int_{-1}^1 P_l(x) P_{l'}(x) dx = \int_0^\pi P_l(\cos \theta) P_{l'}(\cos \theta) \sin \theta d\theta \quad (81)$$

$$= \begin{cases} 0 & \text{if } l' \neq l \\ \frac{2}{2l+1} & \text{if } l' = l \end{cases} \quad (82)$$

Note that when integrating with respect to θ you must multiply by $\sin \theta$ as well. So the solution to Laplace's equation in the case for azimuthal symmetry is as follows:

$$V(r, \theta) = \sum_{l=0}^{\infty} \left(A_l r^l + \frac{B_l}{r^{l+1}} \right) P_l(\cos \theta) \quad (83)$$

In problems where surface charge density is asked at a bounding surface, the following result proves useful:

$$\left. \frac{\partial V_{out}}{\partial \mathbf{n}} - \frac{\partial V_{in}}{\partial \mathbf{n}} \right|_{boundary} = -\frac{\sigma}{\epsilon_0} \quad (84)$$

In the above equation V_{in} and V_{out} are the functions in form of series expansions or otherwise for the potential inside and outside a boundary.

3.6 MULTIPOLE EXPANSION

Using the multipole expansion we can find the potential and by extension the field due to complex charge distributions by breaking them into simpler assemblies of charge. Potential at a point at a distance r from a charge distribution with density $\rho(\vec{r}')$ is given by the multipole expansion as follow:

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \sum_{n=0}^{\infty} \frac{1}{r^{(n+1)}} \int (r')^n P_n(\cos(\theta)) \rho(\vec{r}') d\tau' \quad (85)$$

Note: the integral is with respect to the primed co-ordinates and the θ is the angle between \vec{r} and \vec{r}' so it is also a variable w.r.t the integral.

The $n=0$ term is the potential due to the monopole which basically considering a point charge equal to the total charge of the configuration placed at origin and its potential calculated using Coulomb's law at \vec{r} .

$$V_{mon}(\vec{r}) = \frac{1}{4\pi\epsilon_0 r} \int \rho d\tau \quad (86)$$

If the total charge is zero then the monopole **contribution** will be zero however there may still be a net non-zero potential at \vec{r} due to the remaining multipole terms. $\int \rho d\tau$ is called the **monopole moment**. The potential due to the dipole is the second term and is as follows:

$$V_{dip}(\vec{r}) = \frac{1}{4\pi\epsilon_0 r^2} \int r' P_1(\cos \theta) \rho(\vec{r}') d\tau' \quad , \text{ which can be modified as} \quad (87)$$

$$= \frac{1}{4\pi\epsilon_0 r^2} \hat{r} \cdot \int \vec{r}' \rho(\vec{r}') d\tau' \quad (88)$$

The dipole moment is thus given by

$$\vec{p} = \int \vec{r}' \rho(\vec{r}') d\tau' \quad (89)$$

Since dipole moment is a vector, it follows the laws of vector addition and so multiple moments can be added vectorially to get a resultant. The discussion on physical/ideal dipoles as well as shift of origin is quite important and may be referred from section 4.2 and 4.3 of Griffiths.

3.6.1 Field due to a perfect dipole

Net field due to a dipole \vec{p} in spherical co-ordinates is given by

$$E(\vec{r}) = \frac{p}{4\pi\epsilon_0 r^3} \left[(2 \cos \theta) \hat{r} + (\sin \theta) \hat{\phi} \right] \quad (90)$$

and in co-ordinate independent form is

$$E(\vec{r}) = \frac{1}{4\pi\epsilon_0 r^3} \left[3(\vec{p} \cdot \hat{r})\hat{r} - \vec{p} \right] \quad (91)$$

MAGNETOSTATICS

4.1 BIOT-SAVART LAW

For line currents:

$$d\vec{B} = \frac{\mu_0 I}{4\pi} \frac{d\vec{l} \times \hat{z}}{r^2} \quad (92)$$

In case of current per unit width $K(r')$ the law looks like:

$$d\vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} \frac{K(\vec{r}') \times \hat{z}}{r^2} da \quad (93)$$

In case if 3-D currents:

$$d\vec{B}(\vec{r}) = \frac{\mu_0}{4\pi} \frac{\vec{J}(\vec{r}') \times \hat{z}}{r^2} d\tau' \quad (94)$$

The source of the electrostatic field is a charge $\rho d\tau$, similarly the source for the magnetic field is a current element $\vec{J} d\tau$. For finite sources \vec{B} falls off as $1/r^2$ and for infinite sources as $1/r$. This leads us to the result that

$$\nabla \cdot \vec{B} = 0 \quad (95)$$

This means that there are no magnetic monopoles/charges which leads us to the result that all magnetic field lines should close on themselves, so we arrive at the result that

$$\nabla \times \vec{B} = \mu_0 \vec{J}(\vec{r}) \quad (96)$$

In all the above results \vec{r}' represents the location of magnetic sources. \vec{r} represents the location where field is to be calculated. r represents $|\vec{r} - \vec{r}'|$. In electrostatics, the charges were at rest similarly in magnetostatics the current is steady. This gives us the continuity equation

$$\frac{d\rho}{dt} + \nabla \cdot \vec{J} = 0 \quad (97)$$

This implies that charge cannot teleport from one place to another. It has to follow a continuous path between the two points.

4.2 MAGNETIC VECTOR POTENTIAL

From Helmholtz's theorem we arrive at the fact that Magnetic fields should have a corresponding vector potential. The following identity holds between $\vec{\mathbf{A}}$ and $\vec{\mathbf{B}}$:

$$\boxed{\vec{\mathbf{B}}(\vec{\mathbf{r}}) = \nabla \times \vec{\mathbf{A}}(\vec{\mathbf{r}})} \quad (98)$$

Hence,

$$\vec{\mathbf{B}} = \nabla \times (\nabla \times \vec{\mathbf{A}}) \quad (99)$$

Using the gauge transformation we can get a simplified form of (99) by making the assumption that $\nabla \cdot \mathbf{A} = 0$,

$$\nabla^2 \vec{\mathbf{A}} = -\mu_0 \vec{\mathbf{J}}(\vec{\mathbf{r}})$$

which is the vector form of Laplace's equation the solution of which we already know to be of the form (42). Only difference is that we replace ϵ_0 by $\frac{1}{\mu_0}$ and ρ by $\vec{\mathbf{J}}$. Note that it is a vector integral so one must integrate component wise.

4.3 BOUNDARY CONDITIONS

- Magnetic Field BCs:

$$\nabla \cdot \mathbf{B}(\vec{\mathbf{r}}) = 0 \quad (100)$$

$$\nabla \times \mathbf{B}(\vec{\mathbf{r}}) = \mu_0 \vec{\mathbf{J}}(r) \quad (101)$$

If we consider a general 2-D boundary, the following conditions apply

$$B_{above}^{\perp} = B_{below}^{\perp} \quad (102)$$

$$B_{above}^{\parallel} - B_{below}^{\parallel} = \mu_0 K \quad (103)$$

$$\vec{\mathbf{B}}_{above} - \vec{\mathbf{B}}_{below} = \mu_0 \vec{\mathbf{K}} \times \hat{\mathbf{n}} \quad (104)$$

- Potential BCs: If we invoke the gauge transformation and get a divergence less vector potential then the following conditions hold

$$\vec{A}_{above} = \vec{A}_{below} \quad (105)$$

the derivative however is not continuous,

$$\frac{\partial \vec{A}_{above}}{\partial \mathbf{n}} - \frac{\partial \vec{A}_{below}}{\partial \mathbf{n}} = -\mu_0 \vec{K} \quad (106)$$

DIELECTRICS

5.1 ELECTRIC FIELDS IN MATTER

Matter, which essentially comprises atoms, contains charges in the forms of electrons and protons. These react to an external electric field and internally displace inside the material to produce an opposing dipole moment. This is polarisation.

Polarization is the net dipole moment per unit volume. It is a vector in 3-D space. So is the electric field. The quantity connecting the electric field to the polarization is a second-rank tensor called the polarizability tensor and is a property of the object.

$$\begin{bmatrix} p_x \\ p_y \\ p_z \end{bmatrix} = \alpha \begin{bmatrix} E_x \\ E_y \\ E_z \end{bmatrix} \quad (107)$$

It can be shown using energy conservation and symmetry that the polarizability tensor for materials is symmetric. Hence, using the spectral theorem, it is diagonalisable. However, in this course, we will only study substances whose polarizability matrix is already diagonal and all diagonal elements are equal. These are called linear dielectrics and the polarisation is parallel to the net electric field.

All the formulations are summarised below

5.2 FORMULATION

- Atomic Polarizability:

Electron cloud of an atom responds to an external electric field

$$\vec{p} = (3\epsilon_0 \cdot V) \vec{E}_{ext} \quad (108)$$

where atomic polarizability(α) = $3\epsilon_0$

- Bound Charges:

$$\rho_b = -\nabla \cdot \vec{P} \quad (109)$$

$$\sigma_b = \vec{P} \cdot \hat{n} \quad (110)$$

- Displacement Vector:

$$\vec{D} = \epsilon_0 \vec{E} + \vec{P} \quad (111)$$

$$\rho_f = \nabla \cdot \vec{D} \quad (112)$$

- Boundary Conditions:

$$D_{above}^{\perp} - D_{below}^{\perp} = \sigma_f \quad (113)$$

$$D_{above}^{\parallel} - D_{below}^{\parallel} = P_{above}^{\parallel} - P_{below}^{\parallel} \quad (114)$$

- Linear Dielectric

$$\vec{P} = \epsilon_0 \chi \vec{E} \quad (115)$$

$$\vec{D} = \epsilon_0 (1 + \chi) \vec{E} = \epsilon \vec{E} \quad (116)$$

$$(117)$$

- Energy in a Dielectric:

$$W = \frac{1}{2} \iiint \vec{D} \cdot \vec{E} d\tau \quad (118)$$

5.3 GAUSS'S LAW IN DIELECTRICS

$$\nabla \cdot \vec{\mathbf{D}} = \rho_f \quad (119)$$

The free charge here is the one we control, if some charge is present initially then that is the free charge. This free charge polarises the material. However, we cannot conclude a Coulomb's law for the displacement vector. We can safely take $\vec{\mathbf{D}} = \frac{1}{4\pi} \int \rho_f / |\vec{\mathbf{r}} - \vec{\mathbf{r}}'| d\tau$ when **some symmetry is present** (such that $\nabla \times \vec{\mathbf{D}} = 0$).

ELECTRODYNAMICS AND MAXWELL'S LAWS

6.1 OHM'S LAW

When a constant electric field is applied, the charges in a conductor start moving. One would expect that the velocity and thus the current should grow with time, since the electric field is accelerating the charge carriers. But in reality, we see that the current density is proportional to the electric field applied. This is because the charge carriers collide with each other. Thus their average velocity is a constant and proportional to the applied electric field. This gives rise to Ohm's Law.

$$\vec{J} = \sigma(\vec{E} + \vec{v} \times \vec{B}) \quad (120)$$

The velocity of the charge carriers is often small and thus we ignore the second term.

$$\vec{J} = \sigma \vec{E} \quad (121)$$

An alternate, and more familiar, form of this equation is given in (122)

$$V = IR \quad (122)$$

For steady currents and uniform conductivity,

$$\nabla \cdot \vec{E} = \frac{1}{\sigma} \nabla \cdot \vec{J} = 0 \quad (123)$$

This means that any unbalanced charge resides on the surface. Hence we can apply the methods we learnt earlier in this case as well.

The conductance can be given using Drude's model, where n is the number of molecules per unit volume, f is the number of free electrons per molecule and q is the charge of an electron.

$$\sigma = \frac{nfq^2\tau}{2m} E \quad (124)$$

Since the work done per unit charge is V and the charge flowing per unit time is I , we can write the power as given in eqrefeq:joule. This is called Joule Heating Law.

$$P = VI = I^2R \quad (125)$$

6.2 ELECTRO-MAGNETIC INDUCTION

6.2.1 Electromotive Force

In a circuit, the power source applies a force inside it, which is responsible for the potential difference across the circuit. This is also responsible for what we call EMF(ϵ). We can represent this using the following equations.

$$\vec{f} = \vec{f}_{source} + E \quad (126)$$

$$\epsilon = \oint \vec{f} \cdot d\vec{l} = \oint \vec{f}_{source} \cdot d\vec{l} \quad (127)$$

Taking a and b to be the terminals of the power source, we can further write the following expression,

$$\epsilon = V = \int_a^b \vec{E} \cdot d\vec{l} = \int_a^b \vec{f}_{source} \cdot d\vec{l} = \oint \vec{f}_{source} \cdot d\vec{l} \quad (128)$$

6.2.2 Induction

We define magnetic flux as,

$$\phi \equiv \int \vec{B} \cdot d\vec{a} \quad (129)$$

A change in the magnetic flux gives rise to an induced electric field. This electric field generates an EMF, which in turn causes flow of current. The following equations govern this.

- The Flux Rule:

This is the master equation in a sense, relating magnetic flux to induced EMF

$$\epsilon = -\frac{d\phi}{dt} \quad (130)$$

- Faraday's Law: Differential form

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad (131)$$

- Faraday's Law: Integral form

$$\oint \vec{E} \cdot d\vec{l} = - \int \frac{\partial \vec{B}}{\partial t} \cdot d\vec{a} \quad (132)$$

6.2.3 Induced Electric Field

If you look closely at the equation (131) along with the property of the induced electric field that $\nabla \cdot \vec{E} = 0$, we can find parallels with the equations we have for magnetic field. Thus, we say that the change in flux serves as a source for the induced electric field the way that \vec{J} is a source for the magnetic field.

6.2.4 Lenz's Law

This law is used to find the direction of induced current. It says that the magnetic flux through an object has a "tendency" to remain constant. Thus the induced current is such that the induced magnetic field opposes the change in flux.

6.2.5 Mutual Inductance

Say we place two circuits side by side and let current flow in them. Clearly they will generate magnetic fields and this magnetic field must produce some flux in the circuits. This gives rise to mutual inductance(M), which is defined as

$$\phi_2 = M_{21}I_1 \quad (133)$$

$$\phi_1 = M_{12}I_2 \quad (134)$$

But $M_{21} = M_{12}$, thus we can remove the subscripts. This mutual flux also produces induced EMF if I_1 or I_2 are

6.2.6 Self Inductance

The current in the circuit itself too gives a rise to a magnetic field and that field causes a flux inside the loop of the circuit as well. Now we define self inductance(L) as

$$\phi = LI \quad (135)$$

6.3 ENERGY STORED IN MAGNETIC FIELD

- In terms of inductance and current:

$$W = \frac{1}{2}LI^2 \quad (136)$$

- In terms of magnetic field:

$$W = \frac{1}{2\mu_0} \left[\int_V B^2 d\tau - \oint_S \nabla \cdot (\vec{A} \times \vec{B}) d\tau \right] \quad (137)$$

$$= \frac{1}{2\mu_0} \left[\int_V B^2 d\tau - \oint_S (\vec{A} \times \vec{B}) \cdot d\vec{a} \right] \quad (138)$$

$$= \frac{1}{2\mu_0} \int_{all \ space} B^2 d\tau \quad (139)$$

6.4 MAXWELL'S EQUATIONS

$$\nabla \cdot \vec{E} = \frac{1}{\epsilon_0} \rho \quad (140)$$

$$\nabla \cdot \vec{B} = 0 \quad (141)$$

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad (142)$$

$$\nabla \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t} \quad (143)$$

6.4.1 Displacement current

The displacement current density is given by the following equation,

$$\vec{J}_d = \epsilon_0 \frac{\partial \vec{E}}{\partial t} \quad (144)$$

6.4.2 *Maxwell's Equations in free space*

In free space, we can reduce Maxwell's Equations to,

$$\nabla^2 \vec{E} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2} \quad (145)$$

$$\nabla^2 \vec{B} = \mu_0 \epsilon_0 \frac{\partial^2 \vec{B}}{\partial t^2} \quad (146)$$

$$(147)$$

The solutions of these equations are of the form $f(x - ct)$ where,

$$c = \frac{1}{\sqrt{\mu_0 \epsilon_0}} \quad (148)$$

EM WAVES

The solutions to the 3 dimensional EM-wave equation are

$$\vec{E}(r, t) = \vec{E}_0 \cos(\vec{k} \cdot \vec{r} - \omega t) \quad (149)$$

$$\vec{B}(r, t) = \vec{B}_0 \cos(\vec{k} \cdot \vec{r} - \omega t) \quad (150)$$

Here \vec{k} is called the *wave vector* and represents the direction of propagation and its magnitude denotes the wave number of the wave.

$$|\vec{k}| = \frac{2\pi}{\lambda} \quad (151)$$

By imposing $\nabla \cdot \vec{E} = 0$ & $\nabla \cdot \vec{B} = 0$ we get that \vec{E} & $\vec{B} \perp \vec{k}$. This means that the undulations of the electric and magnetic field are perpendicular to the direction of propagation of the wave. In a vacuum, E & B are in phase.

$$\vec{B}_0 = \frac{1}{c}(\hat{k} \times \vec{E}_0) \quad (152)$$

Here \vec{E} , \vec{B} and \vec{k} form a right handed co-ordinate system.

7.1 ENERGY IN EM WAVES

The energy density of the electric and magnetic fields is equal due to (152), which means that half the energy of an EM wave is carried by E and half by B . The energy density of an electromagnetic field is given by

$$\boxed{u_{EM} = u_E + u_B} \quad (153)$$

Poynting's Theorem

The rate of energy transfer (per unit volume) from a region of space equals the rate of work done on a charge distribution plus the energy flux leaving that region.

$$-\frac{d}{dt} \left(\int u_{EM} d\tau \right) = \int_V (\vec{J} \cdot \vec{E}) d\tau + \oint \vec{S} \cdot d\vec{a} \quad (154)$$

Here u_{EM} is the energy density of the electromagnetic field and $U = \int u d\tau$. \vec{S} is the Poynting vector. (154) can be written in differential form as.

$$-\frac{d}{dt}(u_{EM}) = \vec{J} \cdot \vec{E} + \nabla \cdot \vec{S} \quad (155)$$

Poynting Vector is given by:

$$\vec{S} = \frac{1}{\mu_0} (\vec{E} \times \vec{B}) \quad (156)$$

Its magnitude represents the intensity (energy per unit cross-sectional area per unit time) of the wave and its direction is the direction of propagation of the wave.

$$\vec{S} = \frac{1}{c\mu_0} E_0^2 \cos^2(\vec{k} \cdot \vec{r} - \omega t) \hat{k} \quad (157)$$

Frequencies for EM waves are generally very large so we cannot measure quantities for a specific time. Instead we consider the time averaged quantities, so the time averaged Poynting vector is,

$$\langle \vec{S} \rangle = \left(\frac{E_0^2}{2c\mu_0} \right) \hat{k} \quad (158)$$