# Numerical Methods in Engineering and Applied Science

Lecture 8. Multistep methods.

The model problem that we consider here is the following **initial value problem**: For a function  $\mathbf{f}:[0,T]\times\mathbb{R}^N\to\mathbb{R}^N$  and  $\mathbf{u}_0\in\mathbb{R}^N$  find a differentiable function  $\mathbf{u}(t)$  such that

$$\begin{cases} \frac{\mathrm{d}\boldsymbol{u}}{\mathrm{d}t} = \boldsymbol{f}(t, \boldsymbol{u}(t)), \\ \boldsymbol{u}(0) = \boldsymbol{u}_0. \end{cases}$$
 (1)

Most commonly it describes the **time evolution** of some quantity. The numerical methods for it are called **time stepping** or **time marching** methods. These methods evaluate the numerical solution  $u^{n+1}$  at a time instant  $t_{n+1}$  using the information from previous time instants  $t_n$ ,  $t_{n-1}$  etc. We consider three groups of such methods:

- Taylor series methods;
- Multistage (Runge-Kutta) methods;
- Multistep methods.

## Multistep methods.

A linear s-step method is a formula for calculating the value at a new time step  $u_{n+1}$  using previous time step values  $u_{n-s+1}$ , ...,  $u_n$  and  $f(t_{n-s+1}, u_{n-s+1})$ , ...,  $f(t_n, u_n)$ , where  $s \ge 1$ .

The main idea consists in minimizing the number of evaluations of the right-hand side f(t, u) at each step, because calculation of f(t, u) may be computationally expensive.

Some of the multi-stage methods can be considered as one-step methods, for example

• the Euler method (order 1),

$$u_{n+1} = u_n + h f(t_n, u_n).$$
 (2)

• the trapezoidal rule (order 2)

$$u_{n+1} = u_n + h \frac{f(t_n, u_n) + f(t_{n+1}, u_{n+1})}{2}.$$
 (3)

This scheme is *implicit*, therefore, iteration or a first-order prediction is required to estimate  $u_{n+1}$  in the right-hand side.

• yet another one-step method is the backward Euler method (order 1),

$$u_{n+1} = u_n + hf(t_{n+1}, u_{n+1}). (4)$$

For simplicity, let us write  $f_j$  instead of  $f(t_j, u_j)$ .

Below are some examples of multistep methods of order 2 or higher.

• leap-frog scheme (explicit, order 2)

$$u_{n+1} = u_{n-1} + 2hf_n (5)$$

• fourth-order Adams–Bashforth (AB4) scheme. AB schemes are explicit.

$$u_{n+1} = u_n + \frac{h}{24} \left( 55f_n - 59f_{n-1} + 37f_{n-2} - 9f_{n-3} \right) \tag{6}$$

• fourth-order Adams–Moulton (AM4) scheme. AM schemes are implicit.

$$u_{n+1} = u_n + \frac{h}{24} \left( 9f_{n+1} + 19f_n - 5f_{n-1} + f_{n-2} \right) \tag{7}$$

• third order backward differentiation formula (BDF3). BDF schemes are implicit.

$$u_{n+1} = \frac{18}{11}u_n - \frac{9}{11}u_{n-1} + \frac{2}{11}u_{n-2} + \frac{6}{11}hf_{n+1}$$
 (8)

There is a subtlety about performing the first iterations using multistep methods. For example, suppose we want to use the leapfrog formula. We know the value of  $u_0$ , but it is not sufficient for evaluating  $u_1$ . Where can we find  $u_{-1}$  required for  $u_1$ ? Otherwise, if we start from  $u_2$ , how to calculate  $u_1$ ? There are several practical solutions to this problem.

- Start up using the explicit Euler method with a sufficiently small h.
- Start up a multistep scheme of order p using a Runge-Kutta method of order p-1 to calculate  $u_1, u_2, ..., u_s$ .
- Use analytical formulate for the solution in the neighborhood of the initial condition.

The general form of s-step methods is

$$\sum_{j=0}^{s} \alpha_j u_{n+1+j-s} = h \sum_{j=0}^{s} \beta_j f_{n+1+j-s}, \tag{9}$$

where the coefficients  $\alpha_j$  et  $\beta_j$  are constant,  $\alpha_s = 1$  (the coefficient in front of  $u_{n+1}$ ), and  $\alpha_0 \neq 0$  and/or  $\beta_0 \neq 0$  (coefficients in front of  $u_{n+1-s}$  and  $f_{n+1-s}$ , respectively.). If  $\beta_s = 0$ , the method is *explicit*. If  $\beta_s \neq 0$ , the method is *implicit* ( $\beta_s$  is the coefficient in front of  $f_{n+1}$ ).

Let us define the linear difference operator  $\mathcal{L}_h$  of an s-step scheme as

$$\mathcal{L}_h z(t) = z(t+sh) + \alpha_{s-1} z(t+(s-1)h) + \dots + \alpha_0 z(t) -h \left(\beta_s z'(t+sh) + \beta_{s-1} z'(t+(s-1)h) + \dots + \beta_0 z'(t)\right),$$
(10)

where z(t) is an arbitrary differentiable function (u(t)) is a special case of z(t).  $\mathcal{L}_h$  is a linear operator because

$$\mathcal{L}_h(az(t) + bw(t)) = a\mathcal{L}_h z(t) + b\mathcal{L}_h w(t).$$

**Definition.** The operator  $\mathcal{L}_h$  and the corresponding linear multistep method are consistent of order p if

$$\mathcal{L}_h z(t) = \mathcal{O}(h^{p+1}) \tag{11}$$

with p > 0 for any sufficiently smooth function z.

Example. The linear difference operator of the Euler method is

$$\mathcal{L}_h z(t) = z(t+h) - z(t) - hz'(t),$$

which gives, by Taylor series expansion,

$$\mathcal{L}_h z(t) = \frac{1}{2} h^2 z''(t) + \mathcal{O}(h^3),$$

hence  $\mathcal{L}_h z(t) = \mathcal{O}(h^2)$  and the method is consistent of order 1.

Example. The linear difference operator of the BDF3 scheme:

$$\mathcal{L}_h z(t) = z(t+3h) - \frac{18}{11}z(t+2h) + \frac{9}{11}z(t+h) - \frac{2}{11}z(t) - \frac{6}{11}hz'(t+3h).$$

Its Taylor series expansion:

$$\begin{split} z(t+3h) &= z(t) + 3hz'(t) + \tfrac{9}{2}h^2z''(t) + \tfrac{9}{2}h^3z'''(t) + \tfrac{27}{8}h^4z''''(t) + \mathcal{O}(h^5), \\ z(t+2h) &= z(t) + 2hz'(t) + 2h^2z''(t) + \tfrac{4}{3}h^3z'''(t) + \tfrac{2}{3}h^4z''''(t) + \mathcal{O}(h^5), \\ z(t+h) &= z(t) + hz'(t) + \tfrac{1}{2}h^2z''(t) + \tfrac{1}{6}h^3z'''(t) + \tfrac{1}{24}h^4z''''(t) + \mathcal{O}(h^5), \\ z'(t+3h) &= z'(t) + 3hz''(t) + \tfrac{9}{2}h^2z'''(t) + \tfrac{9}{2}h^3z''''(t) + \mathcal{O}(h^4). \end{split}$$

$$\mathcal{L}_{h}z(t) = \left(1 - \frac{18}{11} + \frac{9}{11} - \frac{2}{11}\right)z(t) + h\left(3 - 2 \cdot \frac{18}{11} + \frac{9}{11} - \frac{6}{11}\right)z'(t) + h^{2}\left(\frac{9}{2} - 2 \cdot \frac{18}{11} + \frac{1}{2} \cdot \frac{9}{11} - 3 \cdot \frac{6}{11}\right)z''(t) + h^{3}\left(\frac{9}{2} - \frac{4}{3} \cdot \frac{18}{11} + \frac{1}{6} \cdot \frac{9}{11} - \frac{9}{2} \cdot \frac{6}{11}\right)z'''(t) + h^{4}\left(\frac{27}{8} - \frac{2}{3} \cdot \frac{18}{11} + \frac{1}{24} \cdot \frac{9}{11} - \frac{9}{2} \cdot \frac{6}{11}\right)z''''(t) + \mathcal{O}(h^{5}) = -\frac{3}{22}h^{4}z''''(t) + \mathcal{O}(h^{5}) = \mathcal{O}(h^{4}).$$

Conclusion: the method is consistent of order 3.

Same approach to determine the coefficients of any method.

$$\mathcal{L}_h z(t) = \sum_{j=0}^s \alpha_j z(t+jh) - h \sum_{j=0}^s \beta_j z'(t+jh)$$
 (12)

Taylor series:

$$z(t+jh) = z(t) + jhz'(t) + \frac{1}{2}(jh)^2 z''(t) + ...,$$
  

$$z'(t+jh) = z'(t) + jhz''(t) + \frac{1}{2}(jh)^2 z'''(t) + ...$$
(13)

Substitute (13) in (12) to obtain

$$\mathcal{L}_{h}z(t) = C_{0}z(t) + C_{1}hz'(t) + C_{2}h^{2}z''(t) + ..., \quad \text{où}$$

$$C_{0} = \alpha_{0} + ... + \alpha_{s},$$

$$C_{1} = (\alpha_{1} + 2\alpha_{2} + ... + s\alpha_{s}) - (\beta_{0} + ... + \beta_{s}),$$

$$C_{2} = \frac{1}{2}(\alpha_{1} + 4\alpha_{2} + ... + s^{2}\alpha_{s}) - (\beta_{1} + 2\beta_{2} + ... + s\beta_{s}),$$

$$\vdots$$

$$C_{m} = \sum_{j=0}^{s} \frac{j^{m}}{m!} \alpha_{j} - \sum_{j=0}^{s} \frac{j^{m-1}}{(m-1)!} \beta_{j}$$
(14)

We conclude that a method is consistent of order p if  $C_0 = C_1 = ... = C_p = 0$  et  $C_{p+1} \neq 0$ . This condition defines a system of p+1 linear equations with 2s+1 unknowns at most:  $\alpha_0,...,\alpha_{s-1}$  and  $\beta_0,...,\beta_s$  (remember that  $\alpha_s=1$ ). One may think it could be possible to construct an s-step method of order 2s. In reality, it turns out that such a method is not convergent and, therefore, useless. We will see that a convergent s-step method can reach the convergence order s+2 at the maximum (Dahlquist barrier). For this reason, in the most used schemes, a large part of the coefficients  $\alpha_j$  and  $\beta_j$  are equal to zero. For example, Adams–Bashforth and Adams–Moulton s-step schemes have order s. Their general form is

$$u_{n+1} = u_n + h \sum_{j=0}^{s} \beta_j f_{n+1+j-s}.$$
 (15)

The coefficients  $\alpha$  are  $\alpha_s = 1$ ,  $\alpha_{s-1} = -1$  and  $\alpha_{s-2} = \dots = \alpha_0 = 0$ . The coefficients  $\beta_0...\beta_{s-1}$  are non-zero.  $\beta_s = 0$  in the Adams–Bashforth schemes and  $\beta_s \neq 0$  in the Adams–Moulton schemes.

There exist an alternative approach to determine the coefficients  $\beta_j$  of these schemes. Function values  $f_{n+1-s}$ , ...,  $f_{n+1}$  are treated as F(t) = f(t, u(t)) (u(t) being the exact solution) evaluated at  $t_{n+1-s}$ , ...,  $t_{n+1}$ . Then we rewrite the differential equation as

$$u(t_{n+1}) - u(t_n) = \int_{t_n}^{t_{n+1}} F(t) dt.$$
 (16)

Let q(t) be a polynomial of degree s-1 (AB) or s (AM) that interpolates the data points  $\{f_j\}$ . Then we can write the AB and AM schemes as

$$u_{n+1} - u_n = \int_{t_n}^{t_{n+1}} q(t) dt.$$
 (17)

As q(t) is a linear combination of the  $\{f_j\}$ , we can take them out of the integral and calculate the weights – coefficients of the numerical scheme.

Example. Calculation of the second-order Adams–Bashforth (AB2) coefficients. We have s = 2, we need to determine the coefficients  $\beta_0$  and  $\beta_1$ . We use the values of  $f_{n-1}$  and  $f_n$  to construct a linear polynomial,

$$q(t) = f_n - \frac{f_n - f_{n-1}}{h}(t_n - t). \tag{18}$$

Using (17), we obtain

$$u_{n+1} - u_n = \int_{t_n}^{t_{n+1}} \left[ f_n - \frac{f_n - f_{n-1}}{h} (t_n - t) \right] dt$$

$$= h f_n - \frac{f_n - f_{n-1}}{h} \int_{t_n}^{t_{n+1}} (t_n - t) dt$$

$$= h f_n - \frac{f_n - f_{n-1}}{h} \left( -\frac{h^2}{2} \right)$$

$$= \frac{3}{2} h f_n - \frac{1}{2} h f_{n-1}.$$
(19)

We derived the AB2 scheme:

$$u_{n+1} = u_n + h\left(\frac{3}{2}f_n - \frac{1}{2}f_{n-1}\right). \tag{20}$$

One can use the Newtonian interpolation to calculate the coefficients of the AB and AM schemes of any order. We introduce the backward difference operator,

$$\nabla z(t) = z(t) - z(t - h). \tag{21}$$

Applied to discrete data, (21) yields

$$\nabla z_n = z_n - z_{n-1}. (22)$$

We will use the powers of this operator,  $\nabla^{j}$ . For example,

$$\nabla^2 z_n = \nabla(\nabla z_n) = \nabla(z_n - z_{n-1})$$
  
=  $(z_n - z_{n-1}) - (z_{n-1} - z_{n-2}) = z_n - 2z_{n-1} + z_{n-2}.$  (23)

• For any  $s \ge 1$ , the s-step Adams–Bashforth scheme is consistent of order s. It can be written in the following form:

$$u_{n+1} = u_n + h \sum_{j=0}^{s-1} \gamma_j \nabla^j f_n.$$
 (24)

The coefficients  $\gamma_j$ ,  $j \geq 0$ , are calculated using the recurrent relation

$$\gamma_j + \frac{1}{2}\gamma_{j-1} + \frac{1}{3}\gamma_{j-2} + \dots + \frac{1}{j+1}\gamma_0 = 1.$$
 (25)

• For all  $s \ge 0$ , the Adams–Moulton scheme with s steps (with 1 step if s = 0) is consistent of order s + 1.

$$u_{n+1} = u_n + h \sum_{j=0}^{s} \gamma_j^* \nabla^j f_{n+1}.$$
 (26)

The coefficients  $\gamma_j$ ,  $j \geq 1$ , are calculated using the relation

$$\gamma_j^* + \frac{1}{2}\gamma_{j-1}^* + \frac{1}{3}\gamma_{j-2}^* + \dots + \frac{1}{j+1}\gamma_0^* = 0$$
 (27)

with  $\gamma_0^* = 1$ . The case s = 0 corresponds to the backward Euler scheme.

ullet For any  $s\geq 1$ , the s steps backward differentiation formula (BDF) is consistent of order s.

$$\sum_{j=1}^{s} \frac{1}{j} \nabla^{j} v_{n+1} = h f_{n+1}. \tag{28}$$

We must divide (28) by the coefficient in front of  $u_n + 1$  to write this scheme in its classical form.

Using the same procedure, we can construct <u>adaptive schemes</u>. For example, an AB2 method with variable time step:

$$u_{n+1} = u_n + \beta_{1a} f_n + \beta_{0a} f_{n-1} \tag{29}$$

The coefficients need to be recalculated on each time step:

$$\beta_{1a} = \frac{1}{2} \frac{h_{n+1}}{h_n} (h_{n+1} + 2h_n), \qquad \beta_{0a} = -\frac{1}{2} \frac{h_{n+1}^2}{h_n}, \tag{30}$$

where  $h_n = t_n - t_{n-1}$ ,  $h_{n+1} = t_{n+1} - t_n$ .

Implicit schemes, like AM or BDF, result in a nonlinear equation (or system of equations). We can use a fixed point method to solve this equation at each time step. For example, for the trapezoidal rule, at all  $t_n$ , n = 1, ..., N, we calculate, for k = 0, 1, ... until convergence,

$$u_{n+1}^{(k+1)} = u_n + \frac{h}{2}(f_n + f(t_n, u_{n+1}^{(k)})).$$

Predictor-corrector methods avoid iteration. Suppose we want to compute  $u_{n+1}$  using an implicit method of order p. To estimate the value of  $f(t_n, u_{n+1})$ , we can use the 'prediction'  $\tilde{u}_{n+1}$  calculated by a explicit method of order  $\geq p-1$ . Example. Improved Euler method (Heun) where  $\tilde{u}$  is calculated using the first-order explicit Euler method,

$$\tilde{u}_{n+1} = u_n + h f(t_n, u_n) 
u_{n+1} = u_n + \frac{h}{2} (f(t_n, u_n) + f(t_{n+1}, \tilde{u}_{n+1}))$$
(31)

## Comparison between two different fourth-order methods.

We consider again the example of a mathematical pendulum. A program that calculates the right side of the equation:

```
function f = rhs_p(u)
global nop;
f = [u(2) -sin(u(1))];
nop = nop + 1;
return
```

We are going to compare the accuracy of Gill's RK4 scheme and of the AB4 scheme, for a given computational cost. We use RK4 as a startup method for AB4.

#### Gill RK4 method:

```
function [uend,h] = func_pendulum_rk4gill(theta0,omega0,tmax,nt)
sqrt2 = sqrt(2);
sqrt05 = sqrt(0.5);
h = tmax/nt;
u = [theta0 omega0];
% t = 0:
for i = 1:nt
k = h * rhs_p(u);
u = u + 0.5*k;
q = k;
k = h * rhs_p(u);
u = u + (1-sqrt05)*(k-q);
q = (2-sqrt2)*k + (-2+3*sqrt05)*q;
k = h * rhs_p(u);
u = u + (1+sqrt05)*(k-q);
q = (2+sqrt2)*k + (-2-3*sqrt05)*q;
k = h * rhs_p(u);
u = u + k/6 - q/3;
% t = t + h;
end;
uend = u;
return
```

#### AB4 method:

```
function [uend,h] = func_pendulum_ab4(theta0,omega0,tmax,nt)
sqrt2 = sqrt(2); sqrt05 = sqrt(0.5);
h = tmax/nt;
u = [theta0 omega0];
f3 = [0\ 0]; f2 = [0\ 0]; f1 = [0\ 0];
% t = 0;
for i = 1:nt
f = rhs_p(u);
if (i < 4)
k = h * f; u = u + 0.5*k; q = k;
k = h * rhs_p(u); u = u + (1-sqrt05)*(k-q);
q = (2-sqrt2)*k + (-2+3*sqrt05)*q;
k = h * rhs_p(u); u = u + (1+sqrt05)*(k-q);
q = (2+sqrt2)*k + (-2-3*sqrt05)*q;
k = h * rhs_p(u); u = u + k/6 - q/3;
else
u = u + h/24 * (55*f - 59*f1 + 37*f2 - 9*f3);
end;
% t = t + h;
f3 = f2; f2 = f1; f1 = f;
end;
uend = u;
return
```

### Comparison:

```
clear all; set(0,'DefaultaxesFontSize',15); set(0,'DefaulttextFontsize',15);
global nop;
tmax = 2.697200567700154e+001; \% tmax = 4 periodes
theta0 = pi/3;
omega0 = 0;
% METHODE RK4
vecnt = 2.^[4:18]; vecnop = zeros(length(vecnt),1);
for it = 1:length(vecnt)
nop = 0;
[X,h] = func_pendulum_rk4gill(theta0,omega0,tmax,vecnt(it));
vecnop(it) = nop;
err_theta = abs(X(1)-theta0);
err_omega = abs(X(2)-omega0);
vecerr(it) = sqrt((err_theta).^2+(err_omega).^2);
end;
figure(1); clf;
loglog(vecnop, vecerr, 'r-', 'LineWidth', 2); hold on;
```

```
% METHODE AB4
vecnt = 2.^[4:18]; vecnop = zeros(length(vecnt),1);
for it = 1:length(vecnt)
nop = 0;
[X,h] = func_pendulum_ab4(theta0,omega0,tmax,vecnt(it));
vecnop(it) = nop;
err_theta = abs(X(1)-theta0);
err_omega = abs(X(2)-omega0);
vecerr(it) = sqrt((err_theta).^2+(err_omega).^2);
end:
figure(1);
loglog(vecnop, vecerr, 'b--', 'LineWidth', 2); hold on;
loglog(vecnop,1e8*vecnop.^(-4),'k:','LineWidth',2);
hold on:
figure(1); axis([1e1 1e6 1e-15 1e1]);
legend('RK4 (Gill)','AB4','O(nop^4)','Location','NorthEast');
xlabel('nop : nombre d''appels rhs_p');
title('((\theta(T)-\theta_{ref})^2+(\omega(T)-\omega_{ref})^2)^\{1/2\}');
set(gca,'XTick',10.^[1:1:6]); set(gca,'YTick',10.^[-15:2:1]); grid on;
print('-depsc','fig_err_nop.eps');
```

