Hydrogen

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1 Matching

QED Lagrangian is

$$\mathcal{L}_{QED} = \bar{l}(i\slashed{D} - m)l + \bar{N}(iD^0)N - \mathcal{L}_\gamma \tag{1}$$

Set the NRQED Lagrangian as (take large M limit where M is the mass of the proton/hydrogen nucleus)

$$\mathcal{L}_{NRQED} = \psi^\dagger(iD_0 + \frac{\mathbf{D}^2}{2m})\psi + \bar{N}(iD_0)N + \mathcal{L}_{4-fer} + \mathcal{L}_\gamma \tag{2}$$
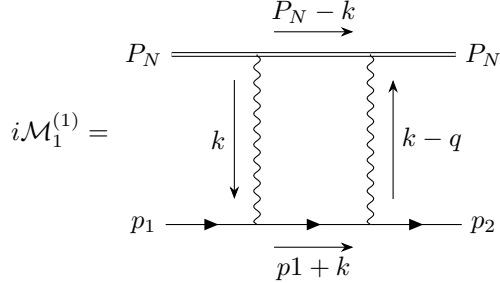
In tree level[1]

$$i\mathcal{M}^{(0)}_{QED} = \qquad\qquad = -e^2 \bar{u}_N(P_N)v^0 u_N(P_N)\frac{i}{\mathbf{q}^2}\bar{u}_e(p_2)\gamma_0 u_e(p1)$$

$$i\mathcal{M}^{(0)}_{NRQED} = \qquad\qquad = -e^2 \bar{u}_N(P_N)v^0 u_N(P_N)\frac{i}{\mathbf{q}^2}\psi^\dagger(p_2)\psi(p1)$$

The box diagram for NRQED process is

$$= e^4 \bar{u}_N(P_N)\frac{1+\gamma^0}{2}u_N(P_N)\psi^\dagger(p_2)\int[\mathrm{d}k]\frac{1}{\mathbf{k}^2(\mathbf{k}-\mathbf{q})^2(-k^0+i\epsilon)(p_1^0+k^0-m-\frac{(\mathbf{p_1}+\mathbf{k})^2}{2m}+i\epsilon)}\psi(p_1)$$

$$= -ie^4 \bar{u}_N(P_N)\frac{1+\gamma^0}{2}u_N(P_N)\psi^\dagger(p_2)\int\frac{\mathrm{d}^3k}{(2\pi)^3}\frac{1}{\mathbf{k}^2(\mathbf{k}-\mathbf{q})^2(E_1-\frac{(\mathbf{p_1}+\mathbf{k})^2}{2m})}\psi(p_1)$$

$$= -ie^4 \bar{u}_N(P_N)\frac{1+\gamma^0}{2}u_N(P_N)\psi^\dagger(p_2)\int\frac{\mathrm{d}^3k}{(2\pi)^3}\frac{1}{(\mathbf{k}-\mathbf{p_1})^2(\mathbf{k}-\mathbf{p_2})^2(E_1-\frac{\mathbf{k}^2}{2m})}\psi(p_1)$$
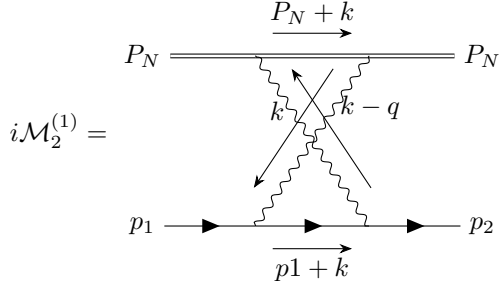
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[1]Note that there's no Gamma matrice in the heavy particle side, they can only appear in the QED side

The box and crossed box diagram for QED process is

$$i\mathcal{M}_1^{(1)} =$$

$$= e^4 \bar{u}_N(P_N)\frac{1+\gamma^0}{2}u_N(P_N)u_e^\dagger(p_2)\int [\mathrm{d}k]\frac{(\not{p}_1+\not{k}+m)\gamma^0}{\mathbf{k}^2(\mathbf{k}-\mathbf{q})^2[(p_1+k)^2-m^2+i\epsilon](-k^0+i\epsilon)}u_e(p_1)$$

$$= e^4 \bar{u}_N(P_N)\frac{1+\gamma^0}{2}u_N(P_N)u_e^\dagger(p_2)\int [\mathrm{d}k]\frac{2p_1^0+\not{k}\gamma^0}{\mathbf{k}^2(\mathbf{k}-\mathbf{q})^2[(p_1+k)^2-m^2+i\epsilon](-k^0+i\epsilon)}u_e(p_1)$$

$$= ie^4 \bar{u}_N(P_N)\frac{1+\gamma^0}{2}u_N(P_N)u_e^\dagger(p_2)\int \frac{\mathrm{d}^3k}{(2\pi)^3}\frac{p_1^0+k_i\gamma^i\gamma^0+\sqrt{(\mathbf{k}+\mathbf{p_1})^2+m^2}}{2\mathbf{k}^2(\mathbf{k}-\mathbf{q})^2[(\mathbf{k}+\mathbf{p_1})^2+m^2-p_1^0\sqrt{(\mathbf{k}+\mathbf{p_1})^2+m^2}]}u_e(p_1)$$

$$= ie^4 \bar{u}_N(P_N)\frac{1+\gamma^0}{2}u_N(P_N)u_e^\dagger(p_2)\int \frac{\mathrm{d}^3k}{(2\pi)^3}\frac{p_1^0+(k_i-p_{1i})\gamma^i\gamma^0+\sqrt{\mathbf{k}^2+m^2}}{2(\mathbf{k}-\mathbf{p_1})^2(\mathbf{k}-\mathbf{p_2})^2[\mathbf{k}^2+m^2-p_1^0\sqrt{\mathbf{k}^2+m^2}]}u_e(p_1)$$

$i\mathcal{M}_1^{(1)}$ has infrared log divergence and no ultraviolet divergence.

$$i\mathcal{M}_2^{(1)} =$$
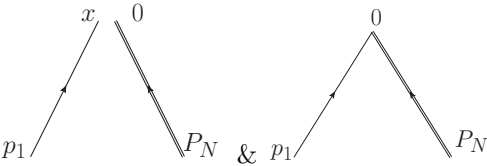
$$= e^4 \bar{u}_N(P_N)\frac{1+\gamma^0}{2}u_N(P_N)u_e^\dagger(p_2)\int [\mathrm{d}k]\frac{(\not{p}_1+\not{k}+m)\gamma^0}{\mathbf{k}^2(\mathbf{k}-\mathbf{q})^2[(p_1+k)^2-m^2+i\epsilon](k^0+i\epsilon)}u_e(p_1)$$

$$= e^4 \bar{u}_N(P_N)\frac{1+\gamma^0}{2}u_N(P_N)u_e^\dagger(p_2)\int [\mathrm{d}k]\frac{2p_1^0+\not{k}\gamma^0}{\mathbf{k}^2(\mathbf{k}-\mathbf{q})^2[(p_1+k)^2-m^2+i\epsilon](k^0+i\epsilon)}u_e(p_1)$$

$$= -ie^4 \bar{u}_N(P_N)\frac{1+\gamma^0}{2}u_N(P_N)u_e^\dagger(p_2)\int \frac{\mathrm{d}^3k}{(2\pi)^3}\frac{p_1^0+k_i\gamma^i\gamma^0-\sqrt{(\mathbf{k}+\mathbf{p_1})^2+m^2}}{2\mathbf{k}^2(\mathbf{k}-\mathbf{q})^2[(\mathbf{k}+\mathbf{p_1})^2+m^2+p_1^0\sqrt{(\mathbf{k}+\mathbf{p_1})^2+m^2}]}u_e(p_1)$$

$$= -ie^4 \bar{u}_N(P_N)\frac{1+\gamma^0}{2}u_N(P_N)u_e^\dagger(p_2)\int \frac{\mathrm{d}^3k}{(2\pi)^3}\frac{p_1^0+(k_i-p_{1i})\gamma^i\gamma^0-\sqrt{\mathbf{k}^2+m^2}}{2(\mathbf{k}-\mathbf{p_1})^2(\mathbf{k}-\mathbf{p_2})^2[\mathbf{k}^2+m^2+p_1^0\sqrt{\mathbf{k}^2+m^2}]}u_e(p_1)$$

$i\mathcal{M}_2^{(1)}$ has no infrared or ultraviolet divergence.

$$i\mathcal{M}_1^{(1)}+i\mathcal{M}_2^{(1)} = ie^4 \bar{u}_N(P_N)\frac{1+\gamma^0}{2}u_N(P_N)u_e^\dagger(p_2)\int \frac{\mathrm{d}^3k}{(2\pi)^3}\frac{p_1^{0^2}+k^2+m^2+(k_i-p_{1i})p_1^0\gamma^i\gamma^0}{(\mathbf{k}-\mathbf{p_1})^2(\mathbf{k}-\mathbf{p_2})^2[\mathbf{k}^2+m^2-p_1^{0^2}]\sqrt{\mathbf{k}^2+m^2}}u_e(p_1)$$

$$= ie^4 \bar{u}_N(P_N)\frac{1+\gamma^0}{2}u_N(P_N)u_e^\dagger(p_2)\int \frac{\mathrm{d}^3k}{(2\pi)^3}\frac{p_1^{0^2}+k^2+m^2+(k_i-p_{1i})p_1^0\gamma^i\gamma^0}{(\mathbf{k}-\mathbf{p_1})^2(\mathbf{k}-\mathbf{p_2})^2[\mathbf{k}^2-\mathbf{p_1}^2]\sqrt{\mathbf{k}^2+m^2}}u_e(p_1)$$
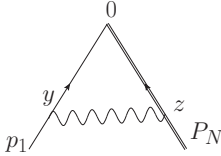
Note that after the expansion over external momentum, k^i can be converted into p^i so it's actually at p^1 order.

Now consider operator product expansion.

Tree level matching: p_1 P_N & p_1 P_N

$$\langle 0|T\psi(x)N(0)|pP_N\rangle = \quad p_1 \qquad P_N \quad = u_e(p)u_N(P_N)e^{-ip\cdot x}$$

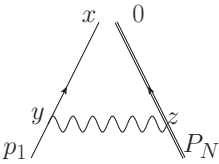$$\langle 0|T\psi_e(x)N(0)|pP_N\rangle = \quad p_1 \qquad P_N \quad = \psi_e(p)u_N(P_N)$$

At leading order $u_e(p) = \begin{pmatrix} \psi_e(p) \\ 0 \end{pmatrix}$.

One loop scenario for NRQED case: p_1 P_N

$$\langle 0|\psi_e(0)N(0)e\int \mathrm{d}^4 y\,\bar{\psi}_e\psi_e A^0\, e\int \mathrm{d}^4 z\,\bar{N}NA^0|eN\rangle = e^2 u_N(v_N)\int[\mathrm{d}k]\frac{1}{\mathbf{k}^2(-k^0+i\epsilon)(p_1^0+k^0-m-\frac{(\mathbf{p_1}+\mathbf{k})^2}{2m}+i\epsilon)}\psi(p_1)$$

$$= -ie^2 u_N(v_N)\int \frac{\mathrm{d}^3 k}{(2\pi)^3}\frac{1}{\mathbf{k}^2(E_1-\frac{(\mathbf{p_1}+\mathbf{k})^2}{2m}+i\epsilon)}\psi(p_1)$$

$$= -ie^2 u_N(v_N)\int \frac{\mathrm{d}^3 k}{(2\pi)^3}\frac{1}{(\mathbf{k}-\mathbf{p_1})^2(E_1-\frac{\mathbf{k}^2}{2m}+i\epsilon)}\psi(p_1)$$

drop p_1

$$= -ie^2 u_N(v_N)\int \frac{\mathrm{d}^3 k}{(2\pi)^3}\frac{1}{\mathbf{k}^2(E_1-\frac{\mathbf{k}^2}{2m}+i\epsilon)}\psi(p_1) = \pi i e^2\sqrt{\frac{2m}{E_1}}u_N(v_N)\psi(p_1)$$
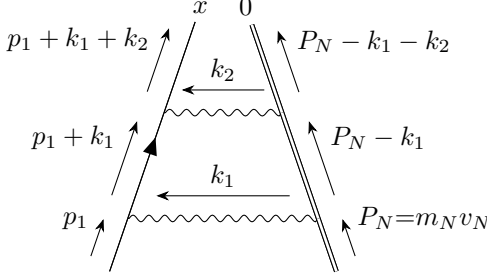
For QED case: p_1 P_N

$$\langle 0|\psi(x)N(0)e\int \mathrm{d}^4 y\,\bar{\psi}\gamma^0\psi A^0\, e\int \mathrm{d}^4 z\,\bar{N}NA^0|eN\rangle^{\,2} = e^2 u_N(v_N)\int[\mathrm{d}k]e^{-i(\mathbf{k}+\mathbf{p_1})\cdot\mathbf{x}}\frac{(\not{p}_1+\not{k}+m)\gamma^0}{\mathbf{k}^2[(p_1+k)^2-m^2+i\epsilon](-k^0+i\epsilon)}u_e(p_1)$$

$$= e^2 u_N(v_N)\int[\mathrm{d}k]e^{-i(\mathbf{k}+\mathbf{p_1})\cdot\mathbf{x}}\frac{2p_1^0+\not{k}\gamma^0}{\mathbf{k}^2[(p_1+k)^2-m^2+i\epsilon](-k^0+i\epsilon)}u_e(p_1)$$

$$= ie^2 u_N(v_N)\int \frac{\mathrm{d}^3 k}{(2\pi)^3}e^{-i(\mathbf{k}+\mathbf{p_1})\cdot\mathbf{x}}\frac{p_1^0+k_i\gamma^i\gamma^0+\sqrt{(\mathbf{k}+\mathbf{p_1})^2+m^2}}{2\mathbf{k}^2[(\mathbf{k}+\mathbf{p_1})^2+m^2-p_1^0\sqrt{(\mathbf{k}+\mathbf{p_1})^2+m^2}]}u_e(p_1)$$

$$= ie^2 u_N(v_N)\int \frac{\mathrm{d}^3 k}{(2\pi)^3}e^{-i\mathbf{k}\cdot\mathbf{x}}\frac{p_1^0+(k_i-p_{1i})\gamma^i\gamma^0+\sqrt{\mathbf{k}^2+m^2}}{2(\mathbf{k}-\mathbf{p_1})^2[\mathbf{k}^2+m^2-p_1^0\sqrt{\mathbf{k}^2+m^2}]}u_e(p_1)$$
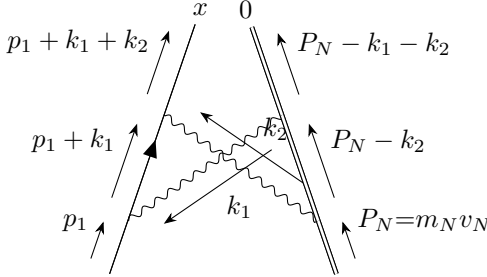
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drop p_1

$$= ie^2 u_N(v_N) \int \frac{\mathrm{d}^3 k}{(2\pi)^3} e^{-i\mathbf{k}\cdot\mathbf{x}} \frac{p_1^0 + \sqrt{\mathbf{k}^2 + m^2}}{2\mathbf{k}^2[\mathbf{k}^2 + m^2 - p_1^0\sqrt{\mathbf{k}^2 + m^2}]} u_e(p_1)$$

Two loop scenario for QED case $\langle 0|T\psi(x)N(0)e \int \mathrm{d}^4 y_1 \bar{\psi}\gamma^0\psi A^0 e \int \mathrm{d}^4 z_1 \bar{N}N A^0 e \int \mathrm{d}^4 y_2 \bar{\psi}\gamma^0\psi A^0 e \int \mathrm{d}^4 z_2 \bar{N}N A^0|eN\rangle$:

$$= e^4 \int [dk_1][dk_2] e^{-i(\mathbf{p_1}+\mathbf{k_1}+\mathbf{k_2})\cdot\mathbf{x}} \frac{1}{|\mathbf{k_1}|^2} \frac{1}{|\mathbf{k_2}|^2} \frac{1}{-k_1^0-k_2^0+i\epsilon} \frac{1}{-k_1^0+i\epsilon} \frac{\not{p}_1+\not{k}_1+\not{k}_2+m}{(p_1+k_1+k_2)^2-m^2+i\epsilon}\gamma^0 \frac{\not{p}_1+\not{k}_1+m}{(p_1+k_1)^2-m^2+i\epsilon}\gamma^0 u_N(v_N)u_e(p_1)$$

$$= e^4 \int [dk_1][dk_2] e^{-i(\mathbf{p_1}+\mathbf{k_1}+\mathbf{k_2})\cdot\mathbf{x}} \frac{1}{|\mathbf{k_1}|^2} \frac{1}{|\mathbf{k_2}|^2} \frac{1}{-k_1^0-k_2^0+i\epsilon} \frac{1}{-k_1^0+i\epsilon} \frac{\not{k}_2(\not{k}_1+2m)+4m^2+2k_1\cdot p_1}{[(p_1+k_1+k_2)^2-m^2+i\epsilon][(p_1+k_1)^2-m^2+i\epsilon]} u_N(v_N)u_e(p_1){}^{[3]}$$

$$= e^4 \int [dk_1][dk_2] e^{-i(\mathbf{p_1}+\mathbf{k_1}+\mathbf{k_2})\cdot\mathbf{x}} \frac{1}{|\mathbf{k_1}|^2} \frac{1}{|\mathbf{k_2}|^2} \frac{1}{-k_1^0-k_2^0+i\epsilon} \frac{1}{-k_2^0+i\epsilon} \frac{\not{p}_1+\not{k}_1+\not{k}_2+m}{(p_1+k_1+k_2)^2-m^2+i\epsilon}\gamma^0 \frac{\not{p}_1+\not{k}_1+m}{(p_1+k_1)^2-m^2+i\epsilon}\gamma^0 u_N(v_N)u_e(p_1)$$

2 HSET

2.1 Lagrangian

For scalar QED

$$\mathcal{L} = (D_\mu\phi)^\dagger D^\mu\phi - m^2\phi^\dagger\phi$$

where

$$D_\mu = \partial_\mu + ieA_\mu$$

In Schwartz's QFT (Chap. 35) he mentioned a choice of χ_v and $\tilde{\chi}_v$:

$$\phi(x) = e^{-imv\cdot x}\frac{1}{\sqrt{2m}}(\chi_v(x) + \tilde{\chi}_v(x)) \tag{3}$$

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[2] $\langle 0|\psi(x)N(0)e \int \mathrm{d}^4 y \bar{\psi}\gamma^0\psi A^0 e \int \mathrm{d}^4 z \bar{N}N A^0|eN\rangle = e^2 \int \mathrm{d}^4 y \int \mathrm{d}^4 z \int \frac{\mathrm{d}^4 k}{(2\pi)^4} \frac{i}{\mathbf{k}^2} e^{-ik\cdot(z-y)} \int \frac{\mathrm{d}^4 k_1}{(2\pi)^4} \tilde{S}_e(k_1)e^{-ik_1\cdot(y-x)} \int \frac{\mathrm{d}^4 k_2}{(2\pi)^4} \tilde{S}_N(k_2)u_N(v_N)u_e(p)e^{-ip_1\cdot y}$.

[3] With a $u_e(p_1)$ on the right hand side, $(\not{p}_1+\not{k}_1+\not{k}_2+m)\gamma^0(\not{p}_1+\not{k}_1+m)\gamma^0 = (\not{p}_1+\not{k}_1+\not{k}_2+m)[2(p_1^0+k_1^0)\gamma^0-(\not{p}_1+\not{k}_1-m)] = 2(p_1^0+k_1^0)(\not{p}_1+\not{k}_1+\not{k}_2+m)\gamma^0-(\not{p}_1+\not{k}_1+\not{k}_2+m)(\not{p}_1+\not{k}_1-m) = 2(p_1^0+k_1^0)(\not{p}_1+\not{k}_1+\not{k}_2+m)\gamma^0-(\not{p}_1+\not{k}_1+\not{k}_2+m)\not{k}_1 = 2(p_1^0+k_1^0)(\not{p}_1+\not{k}_1+\not{k}_2+m)\gamma^0-(2p_1\cdot k_1-\not{k}_1\not{p}_1+\not{k}_2\not{k}_1+\not{k}_1 m) = 2(p_1^0+k_1^0)(\not{p}_1+\not{k}_1+\not{k}_2+m)\gamma^0 - (2p_1\cdot k_1+\not{k}_2\not{k}_1) =.$

$$\chi_v(x) = e^{imv \cdot x} \frac{1}{\sqrt{2m}}(iv \cdot D + m)\phi(x), \ \tilde{\chi}_v(x) = e^{imv \cdot x} \frac{1}{\sqrt{2m}}(-iv \cdot D + m)\phi(x) \tag{4}$$

Put (3) into (4), a simple relation is derived:

$$(-iv \cdot D)\chi_v(x) = (2m + iv \cdot D)\tilde{\chi}_v(x)$$

It can also be writen as

$$2m\tilde{\chi}_v = (-iv \cdot D)(\chi_v + \tilde{\chi}_v)$$

Use this result

$$\mathcal{L} = \frac{1}{2m}\left\{\{[D^\mu(\chi_v + \tilde{\chi}_v)]^\dagger + imv^\mu(\chi_v + \tilde{\chi}_v)^\dagger\}\{[D_\mu(\chi_v + \tilde{\chi}_v)] - imv_\mu(\chi_v + \tilde{\chi}_v)\} - m^2(\chi_v + \tilde{\chi}_v)^\dagger(\chi_v + \tilde{\chi}_v)\right\}$$

$$= (\chi_v + \tilde{\chi}_v)^\dagger(iv \cdot D)(\chi_v + \tilde{\chi}_v) + \frac{1}{2m}[D^\mu(\chi_v + \tilde{\chi}_v)]^\dagger D_\mu(\chi_v + \tilde{\chi}_v) \tag{5}$$

$$= (\chi_v(x) + \tilde{\chi}_v(x))^\dagger(iv \cdot D)(\chi_v(x) + \tilde{\chi}_v(x)) + \mathcal{O}(\frac{1}{m}) \tag{6}$$

(note that $D_\mu\phi = e^{-imv \cdot x}[D_\mu(\chi_v + \tilde{\chi}_v) - imv_\mu(\chi_v + \tilde{\chi}_v)]$ and $-imv^\mu[D_\mu(\chi_v + \tilde{\chi}_v)]^\dagger(\chi_v + \tilde{\chi}_v) = imv^\mu(\chi_v + \tilde{\chi}_v)^\dagger D_\mu(\chi_v + \tilde{\chi}_v) - total \ derivative \ term$)

Use the leading order of (5)

$$\begin{aligned}
\mathcal{L}^{(0)} &= (\chi_v + \tilde{\chi}_v)^\dagger(iv \cdot D)(\chi_v + \tilde{\chi}_v)\\
&= \chi_v^\dagger iv \cdot D\chi_v + \tilde{\chi}_v^\dagger iv \cdot D(\chi_v + \tilde{\chi}_v) + \chi_v^\dagger iv \cdot D\tilde{\chi}_v\\
&= \chi_v^\dagger iv \cdot D\chi_v - 2m\tilde{\chi}_v^\dagger\tilde{\chi}_v + (iv \cdot D\chi_v)^\dagger\tilde{\chi}_v\\
&= \chi_v^\dagger iv \cdot D\chi_v - 2m\tilde{\chi}_v^\dagger\tilde{\chi}_v + [(-2m - iv \cdot D)\tilde{\chi}_v]^\dagger\tilde{\chi}_v\\
&= \chi_v^\dagger iv \cdot D\chi_v - \tilde{\chi}_v^\dagger(iv \cdot D + 4m)\tilde{\chi}_v
\end{aligned}$$

We can have the final form[4]

$$\mathcal{L} = \chi_v^\dagger iv \cdot D\chi_v - \tilde{\chi}_v^\dagger(iv \cdot D + 4m)\tilde{\chi}_v + \mathcal{O}(\frac{1}{m})$$

2.2 Quantization

2.2.1 HQET as an example

The leading term of HQET Lagrangian is

$$\mathcal{L} = \bar{Q}_v(iv \cdot D)Q_v$$

$$Q_v(x) = e^{imv \cdot x}\frac{1 + \not{v}}{2}\psi(x)$$

In free Dirac fermion theory we know that

$$\left\{\psi_a(\mathbf{x}), \psi_b^\dagger(\mathbf{y})\right\} = \delta^{(3)}(\mathbf{x} - \mathbf{y})\delta_{ab}$$

$$\{a_\mathbf{p}, a_{\mathbf{p}'}^\dagger\} = (2\pi)^3\delta^{(3)}(\mathbf{p} - \mathbf{p}')$$

also the plane wave expansion of ψ is

$$\begin{aligned}
\psi(x) &= \int \frac{\mathrm{d}^3p}{(2\pi)^3}\frac{1}{\sqrt{2E_p}}a_\mathbf{p}u(p)e^{-ip \cdot x}\\
&= \int \frac{\mathrm{d}^3k}{(2\pi)^3}\frac{1}{\sqrt{2mv^0}}a_v\sqrt{m}u(v)e^{-imv \cdot x - ik \cdot x}
\end{aligned}$$

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[4]With one problem: if we can tolerate coupled particle-anti particle pair, we can trade $iv \cdot D$ for mass term, so the leading part is the same but the anti-particle part could be different with the mixing?

using normalization of states $u(k) = \sqrt{m}u(v)$[5], $\langle p'|p\rangle = 2E_p(2\pi)^3\delta^{(3)}(\mathbf{p}' - \mathbf{p})$ and $\langle v', k'|v, k\rangle = 2v^0\delta_{vv'}(2\pi)^3\delta^{(3)}(\mathbf{k}' - \mathbf{k})$ we have $|p\rangle = \sqrt{m}|v\rangle$ ($|p\rangle = \sqrt{2E_p}a_{\mathbf{p}}^\dagger|0\rangle$ while $|v, k\rangle = \sqrt{2v^0}a_{v,\mathbf{k}}^\dagger|0\rangle$)

$$= \int \frac{\mathrm{d}^3k}{(2\pi)^3} \frac{1}{\sqrt{2v^0}} a_v u(v) e^{-imv\cdot x - ik\cdot x}$$

Using the definition of $Q_v(x)$

$$Q_v(x) = e^{imv\cdot x}\frac{1 + \not{v}}{2}\psi(x)$$

$$= \int \frac{\mathrm{d}^3k}{(2\pi)^3} \frac{1}{\sqrt{2v^0}} a_v \frac{1 + \not{v}}{2} u(v) e^{-ik\cdot x}$$

$$= \int \frac{\mathrm{d}^3k}{(2\pi)^3} \frac{1}{\sqrt{2v^0}} a_v u(v) e^{-ik\cdot x}$$

The commutation relation should be

$$\{Q_{va}(\mathbf{x}), Q_{v'b}(\mathbf{x}')\} = \int \frac{\mathrm{d}^3k}{(2\pi)^3} \frac{\mathrm{d}^3k'}{(2\pi)^3} \frac{1}{\sqrt{4v^0v'^0}} \{a_v, a_{v'}^\dagger\} u_a(v) u_b^\dagger(v') e^{-ik\cdot x + ik'\cdot x'}$$

using $\sum_s u_a(v) u_b^\dagger(v) = \frac{1}{m}\sum_s u_a(p) u_b^\dagger(p) = [(\not{v} + 1)\gamma^0]_{ab}$

$$= \int \frac{\mathrm{d}^3k}{(2\pi)^3} \frac{\mathrm{d}^3k'}{(2\pi)^3} \frac{1}{\sqrt{4v^0v'^0}} \{a_v, a_{v'}^\dagger\}[(\not{v} + 1)\gamma^0]_{ab} e^{-ik\cdot x + ik'\cdot x'}$$

assuming $\{a_v, a_{v'}\} = (2\pi)^3\delta_{vv'}\delta^{(3)}(\mathbf{k} - \mathbf{k}')$

$$= \int \frac{\mathrm{d}^3k}{(2\pi)^3} \frac{1}{2v^0}[(\not{v} + 1)\gamma^0]_{ab} e^{-ik\cdot(x - x')}\delta_{vv'}$$

$$= [\frac{(\not{v} + 1)\gamma^0}{2v^0}]_{ab}\delta_{vv'}\delta^{(3)}(\mathbf{x} - \mathbf{x}')$$

2.2.2 HSET

The anti-particle field is decoupled so we don't have to consider that for now. The equation-of-motion is

$$\begin{cases} v \cdot D\chi_v^\dagger = 0 \\ v \cdot D\chi_v = 0 \end{cases}$$

By definition

$$\chi_v(x) = \frac{e^{imv\cdot x}}{\sqrt{2m}}(iv \cdot D + m)\phi(x)$$

$$= \frac{1}{\sqrt{2m}}(iv \cdot D + 2m)e^{imv\cdot x}\phi(x)$$

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[5]The relation $\bar{u}^s(p)\gamma^\mu u^s(p) = 2p^\mu$ can be derived using Gordon identity, same for $\bar{u}^s(v)\gamma^\mu u^s(v) = 2v^\mu$, but it's actually $\bar{u}u$.