

Day 1

A. preparation

Ref. An Introduction to QFT

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Pauli & Dirac matrices

$$\sigma^1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \sigma^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \sigma^3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

products of sigma matrices

$$\sigma^i \sigma^j = \delta^{ij} + i \epsilon^{ijk} \sigma^k$$

Weyl representation

$$\gamma^0 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \gamma^i = \begin{pmatrix} 0 & \sigma^i \\ -\sigma^i & 0 \end{pmatrix}, \gamma^5 = \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}$$

Anti-commutation relation

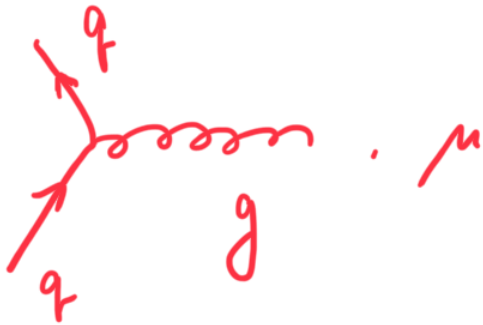
$$\{\gamma^\mu, \gamma^\nu\} = 2g^{\mu\nu}, \{\gamma^\mu, \gamma^5\} = 0$$

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SM Feynman Rules :



$$ig \cdot \gamma^\mu P_L \cdot \delta_{ij} / \sqrt{2}$$



$$ig_s \gamma^\mu T_{ij}^a$$



$$\frac{i}{\not{p} - m + i\epsilon} \delta_{ij}$$



$$\frac{-ig_{\mu\nu}}{p^2 + i\epsilon} \delta_{ab}$$



$$\frac{-i}{p^2 - m_w^2 + i\epsilon} \cdot \left(g^{\mu\nu} - \frac{p^\mu p^\nu}{m_w^2} \right)$$

Feynman Gauge & unitary Gauge.

0..

Physical constants (PDG 2019):

$$m_t = 172.9 \pm 0.4 \text{ GeV}$$

$$m_W = 80.379 \pm 0.012 \text{ GeV}$$

m_b, m_e, m_q set to 0, CKM to diagonal

QCD coupling (\overline{MS} scheme)

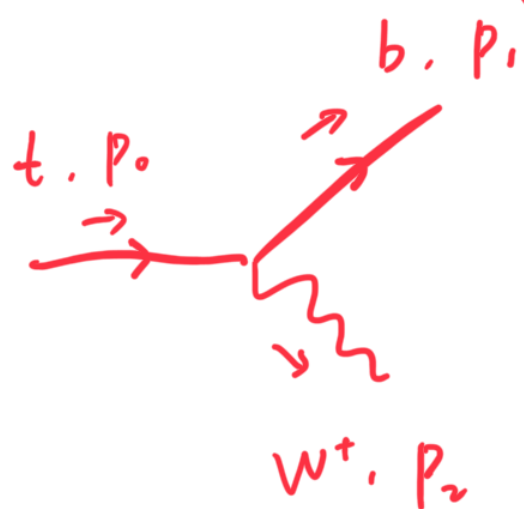
$$\alpha_s(\mu) = \frac{g_s^2(\mu)}{4\pi}, \quad \alpha_s(m_Z) = 0.118 \pm 0.001$$

EW coupling (G_F scheme)

$$G_F = \frac{\sqrt{2} g^2}{8 m_W^2} = 1.16638 \times 10^{-5} \text{ GeV}^{-2}$$

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B. Top quark decay at LO (total width)



$$i\mathcal{M} = \bar{u}(p_1) \gamma^\mu p_\mu u(p_0) \cdot \epsilon_\mu^*(p_2) \cdot ig/\sqrt{2}$$

for unpolarized top and sum over final state spins,

$$\overline{\sum} |\mathcal{M}|^2 = \frac{1}{2N_c} \cdot \text{Tr} [\not{p}_1 \gamma^\mu \not{p}_2 \cdot (\not{p}_0 + m_t) \cdot \gamma^\nu \not{p}_2] \cdot N_c \cdot \sum_{\lambda} \epsilon_\mu^{*,\lambda}(p_2) \epsilon_\nu^\lambda(p_2) \cdot \frac{g^2}{2}$$

recall spin sum for massive vector boson,

$$\sum_{\lambda=\pm,0} \epsilon_\mu^{*,\lambda}(p_2) \epsilon_\nu^\lambda(p_2) = -g_{\mu\nu} + \frac{p_{2\mu} p_{2\nu}}{m_W^2}.$$

thus

$$\overline{\sum} |\mathcal{M}|^2 = \frac{g^2}{8} \left\{ 2 \text{Tr}[\not{p}_1 \not{p}_0] + \frac{1}{m_W^2} \cdot \text{Tr}[\not{p}_1 \not{p}_2 \not{p}_0 \not{p}_2] \right\}$$

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with a bit calculation,

$$|\overline{\mathcal{M}}|^2 = m_t^2 \cdot (1-x) \cdot \left(1 + \frac{1}{2x}\right) \cdot \frac{g^2}{2}, \quad x = \frac{m_W^2}{m_t^2}$$

with the usual formula for decay width,

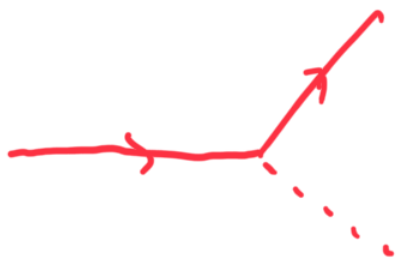
$$\Gamma = \frac{1}{2m_t} \cdot \int d\Pi_2 |\overline{\mathcal{M}}|^2$$

$$= \frac{1}{2m_t} \cdot m_t^2 (1-x) \cdot \left(1 + \frac{1}{2x}\right) \cdot \frac{4m_W^2 G_F}{\sqrt{2}} \cdot \frac{1}{8\pi} \cdot (1-x)$$

$$= \frac{\sqrt{2} G_F m_W^2}{8\pi} \cdot m_t \cdot (1-x)^2 \left(1 + \frac{1}{2x}\right)$$

numerically, $\Gamma \approx 1.49 \text{ GeV}$

note when $m_t \gg m_W$, $\Gamma \propto m_t^3$ due to the longitudinal W / goldstone contribution, e.g.,



$$|\overline{\mathcal{M}_L}|^2 \approx \frac{1}{2} \text{Tr} \left[\frac{\not{p} \not{p}_1}{2} \right] \cdot \left(\frac{\sqrt{2} m_t}{v} \right)^2$$

$$\text{using } v = \frac{2m_W}{g}, \quad |\overline{\mathcal{M}_L}|^2 \approx m_t^2 \cdot \frac{1}{2x} \cdot \frac{g^2}{2}.$$

o.i.

Recall actually width of unstable particles are rather defined as imaginary part of 1PI two-point cor..
It's with optical theorem.

$$\text{Im} \cdot \text{---} \text{---} \text{---} \text{---} \text{---} = \sum_x \left| \text{---} \text{---} \text{---} \text{---} \right|^2 \cdot \frac{1}{2}$$

(unitary condition $|1 + i\hat{T}|^2 = 1 \Rightarrow -i(\hat{T} - \hat{T}^\dagger) = \hat{T}^\dagger \hat{T}$)

That implies for scalar $\Gamma = - \frac{\text{Im}(M(p^2))}{m}$.

For the case of fermion, define

$$\begin{aligned} -i\hat{\Sigma}(p) &= \text{---} \text{---} \text{---} \text{---} \text{---} \cdot \text{1PI} \\ &= -i(\bar{Z}_+(p) p_R + \bar{Z}_-(p) p_L) \end{aligned}$$

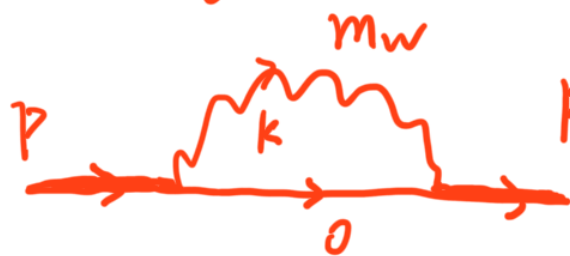
then

$$\Gamma = - \text{Im}(\bar{Z}_+(p) + \bar{Z}_-(p)) \Big|_{p=m}$$

arxiv:0801.0669

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alternative way,



in dimensional regularization
 $d = 4 - 2\epsilon$

$$-i\bar{\Sigma}(p) = \mu^{2\epsilon} \cdot \int \frac{d^d k}{(2\pi)^d} \cdot (-g^2/2) \cdot \frac{1}{(k^2 - m_W^2)(k-p)^2} \cdot \left\{ \gamma^\mu (\not{p} - \not{k}) \gamma_\mu - \frac{1}{m_W^2} \cdot \not{k} (\not{p} - \not{k}) \not{k} \right\} \cdot P_L$$

the Dirac algebra,

$$= -2(1-\epsilon)(\not{p} - \not{k}) - \frac{2p \cdot k}{m_W^2} \not{k} + \frac{k^2}{m_W^2} (\not{k} + \not{p}),$$

let

$$\frac{1}{i2\pi \cdot r_F} \cdot (2\mu^2)^\epsilon \cdot \int d^d k \cdot \frac{1}{(k^2 - m_W^2)(k-p)^2} \equiv B_0(p^2, m_W^2, 0)$$

$$\frac{1}{i2\pi \cdot r_F} \cdot (2\mu^2)^\epsilon \cdot \int d^d k \cdot \frac{k^\mu}{(k^2 - m_W^2)(k-p)^2} \equiv B_1(p^2, m_W^2, 0) \cdot p^\mu$$

(using Passarino - Veltman notation that will

be introduced later, $r_F = \frac{\Gamma^2(1-\epsilon)\Gamma(1+\epsilon)}{\Gamma(1-2\epsilon)}$)

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for the second term,

$$- \frac{2p \cdot k \not{k}}{m_w^2} \cdot \frac{1}{(k^2 - m_w^2)(k-p)^2} \doteq - \left(1 + \frac{p^2}{m_w^2}\right) \cdot \frac{\not{k}}{(k^2 - m_w^2)(k-p)^2}$$

$$\frac{k^2}{m_w^2} \cdot (\not{k} + \not{p}) \cdot \frac{1}{\dots\dots\dots} \doteq \frac{\not{k} + \not{p}}{(k^2 - m_w^2)(k-p)^2}$$

$$\left[\text{e.g., } \int d^d k \cdot \frac{1}{(k-p)^2} = 0, \int d^d k \frac{k^\mu}{k^2 - m_w^2} = 0 \right]$$

thus

$$- i \bar{\Sigma}(p) = i \cdot \frac{(4\pi)^{\epsilon}}{16\pi^2} \cdot g_F \cdot \left(-\frac{g^2}{2}\right) \cdot \left\{ (2\epsilon-1) B_0(p^2, m_w^2, 0) \not{p} \right. \\ \left. + \left(2 - 2\epsilon - \frac{p^2}{m_w^2}\right) \cdot B_1(p^2, m_w^2, 0) \cdot \not{p} \right\} p_L,$$

using

$$B_1(p^2, m_w^2, 0) = \frac{1}{2p^2} \cdot A_0(m_w^2) - \frac{1}{2} \left(1 + \frac{m_w^2}{p^2}\right) B_0(p^2, m_w^2, 0)$$

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also

$$A_0(m_W^2) = m_W^2 \cdot \left(\frac{m^2}{m_W^2 - i\epsilon} \right)^\epsilon \left(\frac{1}{\epsilon} + 1 \right) + \mathcal{O}(\epsilon)$$

$$B_0(p^2, m_W^2, 0) = \left(\frac{m^2}{m_W^2} \right)^\epsilon \left\{ \frac{1}{\epsilon} + 2 + \frac{m_W^2 - p^2}{p^2} \ln \left(\frac{m_W^2 - p^2 - i\epsilon}{m_W^2} \right) \right\}$$

one find the imaginary part,

$$\text{Im } A_0 = 0,$$

$$\text{Im } B_0(p^2, m_W^2, 0) = \frac{p^2 - m_W^2}{p^2} 2 \cdot \theta(p^2 - m_W^2)$$

finally,

$$\text{Im } \bar{\Sigma}_-(p) = \frac{1}{16\pi^2} \cdot \frac{g^2}{2} \cdot \cancel{p} \cdot \left(\frac{m_W^2}{p^2} - \frac{p^2}{2m_W^2} - \frac{1}{2} \right) \cdot$$

$$\left(1 - \frac{m_W^2}{p^2} \right) 2 \cdot \theta(p^2 - m_W^2)$$

So

$$\Gamma = -\text{Im } \bar{\Sigma}_-(p) \Big|_{\cancel{p}=m_t} = \frac{\sqrt{2} G_F m_W^2}{8\pi} \cdot m_t (1-x)^2 \cdot \left(1 + \frac{1}{2x} \right)$$

Proof of optical theory at Feynman diagram level \rightarrow Cutkosky rules !!

Q.E.D.