


Chiral basis:

$$\gamma^\mu = \begin{pmatrix} 0 & \sigma^\mu \\ \bar{\sigma}^\mu & 0 \end{pmatrix},$$

$$\gamma_5 = \begin{pmatrix} -1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\frac{1-\gamma_5}{2} = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

$$\sigma^\mu_{\alpha\dot{\beta}} = (1, \vec{\sigma})_{\alpha\dot{\beta}}$$

$$(\bar{\sigma}^\mu)^{\dot{\alpha}\beta} = (1, -\vec{\sigma})^{\dot{\alpha}\beta} = \varepsilon^{\dot{\alpha}\dot{\beta}} \varepsilon^{\beta\alpha} \sigma^\mu_{\alpha\dot{\beta}}$$

$$\sigma^1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \sigma^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \sigma^3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

$$\varepsilon^{12} = +1$$

$$\varepsilon_{12} = -1$$

★ Spinor-helicity

Dixon: 1310.5353

Elvang & Huang: 1308.1697

massless fermion

$$\not{p} u(p) = 0 = \not{p} v(p)$$

can take $u(p) = v(p)$

$$\bar{u}(p) \not{p} = 0 = \bar{v}(p) \not{p}$$

$$\bar{u}(p) = \bar{v}(p)$$

$$\frac{1 - \gamma_5}{2} u(p) \equiv \begin{pmatrix} [p]_\alpha \\ 0 \end{pmatrix} \equiv |p]$$

$$\frac{1 + \gamma_5}{2} u(p) \equiv \begin{pmatrix} 0 \\ \underbrace{|p\rangle_{\dot{\alpha}}} \end{pmatrix} \equiv |p\rangle$$

$$\bar{u}(p) \frac{1+\gamma_5}{2} \equiv (0 \mid p|_{\dot{\alpha}}) \equiv \langle p|$$

$$\bar{u}(p) \frac{1-\gamma_5}{2} \equiv (\bar{p}|^{\dot{\alpha}} \mid 0) \equiv [\bar{p}|$$

$$\not{p} = p^\mu \gamma_\mu = \begin{pmatrix} 0 & p \cdot \sigma \\ p \cdot \bar{\sigma} & 0 \end{pmatrix} = \begin{pmatrix} 0 & \underline{p_{\dot{\alpha}\beta}} \\ p^{\dot{\alpha}\beta} & 0 \end{pmatrix}$$

$$p_\mu \sigma^\mu_{\dot{\alpha}\beta} \equiv p_{\dot{\alpha}\beta} \quad , \quad p^\mu \bar{\sigma}_\mu^{\dot{\alpha}\beta} \equiv \underline{p^{\dot{\alpha}\beta}} = \begin{pmatrix} p^0 + p^3 & p^1 - ip^2 \\ p^1 + ip^2 & p^0 - p^3 \end{pmatrix}$$

$$\det(p^{\dot{\alpha}\beta}) = p^2 = 0$$

$$\text{Dirac eq. : } p_{\alpha\dot{\beta}} |p\rangle^{\dot{\beta}} = 0$$

$$p^{\dot{\alpha}\beta} |p]_{\beta} = 0$$

$$\langle p|_{\dot{\alpha}} p^{\dot{\alpha}\beta} = 0$$

$$[\bar{p}]^{\alpha} p_{\alpha\dot{\beta}} = 0$$

7月16日

Spinor product

$$|p_i\rangle \equiv |i\rangle, \quad |p_i] \equiv [i]$$

$$\langle i j \rangle \equiv \langle i|_{\dot{\alpha}} |j\rangle^{\dot{\alpha}} = \varepsilon_{\dot{\alpha}\dot{\beta}} |i\rangle^{\dot{\beta}} |j\rangle^{\dot{\alpha}} = -\langle j i \rangle$$

$$[i j] \equiv [i|^{\alpha} |j]_{\alpha} = -[j i]$$

$$\langle i | \gamma^{\mu} | j] = \langle i|_{\dot{\alpha}} (\bar{\sigma}^{\mu})^{\dot{\alpha}\beta} |j]_{\beta}$$

$$[i | \gamma^{\mu} | j \rangle = [i|^{\alpha} \sigma_{\alpha\dot{\beta}}^{\mu} |j\rangle^{\dot{\beta}}$$

* Momentum

$$\not{p} = \begin{pmatrix} 0 & p_{\alpha\dot{\beta}} \\ p^{\dot{\alpha}\beta} & 0 \end{pmatrix}$$

$\nearrow p \cdot \sigma$

$\searrow p \cdot \bar{\sigma}$

$$\sigma_{\alpha\dot{\beta}}^{\mu} (\bar{\sigma}^{\nu})^{\dot{\beta}\alpha} = \text{Tr}(\sigma^{\mu} \bar{\sigma}^{\nu}) = 2g^{\mu\nu}$$

$$\Rightarrow p^{\dot{\alpha}\beta} \sigma_{\beta\dot{\alpha}}^{\mu} = \text{Tr}(p \sigma^{\mu}) \stackrel{?}{=} 2p^{\mu}$$

$$p_{\alpha\dot{\beta}} (\bar{\sigma}^{\mu})^{\dot{\beta}\alpha} =$$

$$\begin{aligned}
 * \quad S_{ij} &\equiv 2P_i \cdot P_j = \underbrace{\sigma_{\alpha\dot{\beta}}^{\mu}}_{\substack{\text{metric} \\ \text{tensor}}} P_{i\mu} \underbrace{\bar{\sigma}^{\nu\dot{\beta}\alpha}}_{\substack{\text{metric} \\ \text{tensor}}} P_{j\nu} = \underbrace{(P_i)_{\alpha\dot{\beta}}}_{\substack{\text{spinor} \\ \text{index}}} \underbrace{(P_j)^{\dot{\beta}\alpha}}_{\substack{\text{spinor} \\ \text{index}}} \\
 &= [i]_{\alpha} \langle i |_{\dot{\beta}} | j \rangle^{\dot{\beta}} [j]^{\alpha} = \langle i j \rangle [j i]
 \end{aligned}$$

$$* \quad \text{For real momenta} \quad [P]^{\alpha} = (|P\rangle^{\dot{\alpha}})^* \quad (\alpha = \dot{\alpha})$$

$$\langle P |_{\dot{\alpha}} = (|P]_{\alpha})^*$$

$$\Rightarrow \langle i j \rangle = [j i]^* = \sqrt{S_{ij}} \times e^{i\phi_{ij}}$$

☆ Back to $e^+e^- \rightarrow f\bar{f}$

$$i\mathcal{M}_4(\text{LRRL}) = 2ie^2 Q_e Q_f \mathbb{I} A_4(\text{LRRL})$$

$$A_4(\text{LRRL}) = \frac{1}{2S_{12}} \langle 2 | \gamma^\mu | 1 \rangle [3 | \gamma_\mu | 4 \rangle$$

$$= \frac{1}{2S_{12}} 2 \langle 2 4 \rangle [3 1]$$

$$= \frac{\langle 2 4 \rangle [3 1]}{\langle 1 2 \rangle [2 1]}$$

* Other helicity configurations

$$P: L \leftrightarrow R, \quad | \rangle \leftrightarrow |]$$

$$\Rightarrow A_4(RLLR) = \frac{\langle 13 \rangle^2}{\langle 12 \rangle \langle 34 \rangle} = \frac{[24]^2}{[12][34]}$$

C on 1-2 fermion-line: $1 \leftrightarrow 2$

$$\Rightarrow A_4(RLRL) = - \frac{\langle 14 \rangle^2}{\langle 12 \rangle \langle 34 \rangle}$$

C on 3-4 line:

$$\Rightarrow A_4(LRLR) = - \frac{\langle 23 \rangle^2}{\langle 12 \rangle \langle 34 \rangle}$$

* Squared-amplitude

$$\overline{|M_4|^2} = e^4 Q_e^2 Q_f^2 N_c \sum_{\text{hel}} |A_4|^2$$

number of colors (=3)

$$= 2e^4 Q_e^2 Q_f^2 N_c \left[\left| \frac{\langle 24 \rangle^2}{\langle 12 \rangle \langle 34 \rangle} \right|^2 + \left| \frac{\langle 14 \rangle^2}{\langle 12 \rangle \langle 34 \rangle} \right|^2 \right]$$

$$= 2e^4 Q_e^2 Q_f^2 N_c \frac{S_{24}^2 + S_{13}^2}{S_{12}^2}$$

$$\hookrightarrow \frac{1}{4}(1-\cos\theta)^2 + \frac{1}{4}(1+\cos\theta)^2$$
$$= \frac{1}{2}(1+\cos^2\theta)$$

* Cross section

$$S = S_{12}$$

$$\sigma = \frac{1}{2S} \int d\Phi_2 \overline{|M_4|^2}$$

↳ 2-body phase space

$$d\Phi_2 = \frac{d^3\vec{p}_3}{(2\pi)^3 2E_3} \frac{d^3\vec{p}_4}{(2\pi)^3 2E_4} (2\pi)^4 \delta^{(4)}(p_1 + p_2 - p_3 - p_4)$$

$$\frac{d^3\vec{p}_4}{2E_4} = d^4p_4 \delta(p_4^2) \Theta(E_4)$$

↳ $|\vec{p}_4|$

$$\int d\Phi_2 = \frac{1}{(2\pi)^2} \int \frac{d^3\vec{P}_3}{2\bar{E}_3} \underbrace{\delta((p_1+p_2-p_3)^2)}_{\substack{\text{c.o.m. frame} \\ \delta(s-2\sqrt{s}\bar{E}_3)}} \underbrace{\theta(E_1+E_2-\bar{E}_3)}_{\theta(\sqrt{s}-\bar{E}_3)}$$

\downarrow $|\vec{P}_3|$

$$d^3\vec{P}_3 = |\vec{P}_3|^2 d|\vec{P}_3| d\cos\theta d\phi$$

$$\alpha = \frac{e^2}{4\pi}$$

$$\frac{1}{16\pi} \int_{-1}^{+1} d\cos\theta$$

$$\sigma = \frac{e^4 Q_f^2 N_c}{32\pi s} \int_{-1}^1 d\cos\theta (1 + \cos^2\theta) = \frac{4\pi\alpha^2}{3s} \underline{N_c Q_f^2}$$

* R ratio

$$R = \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)} \approx N_c \sum_f Q_f^2$$

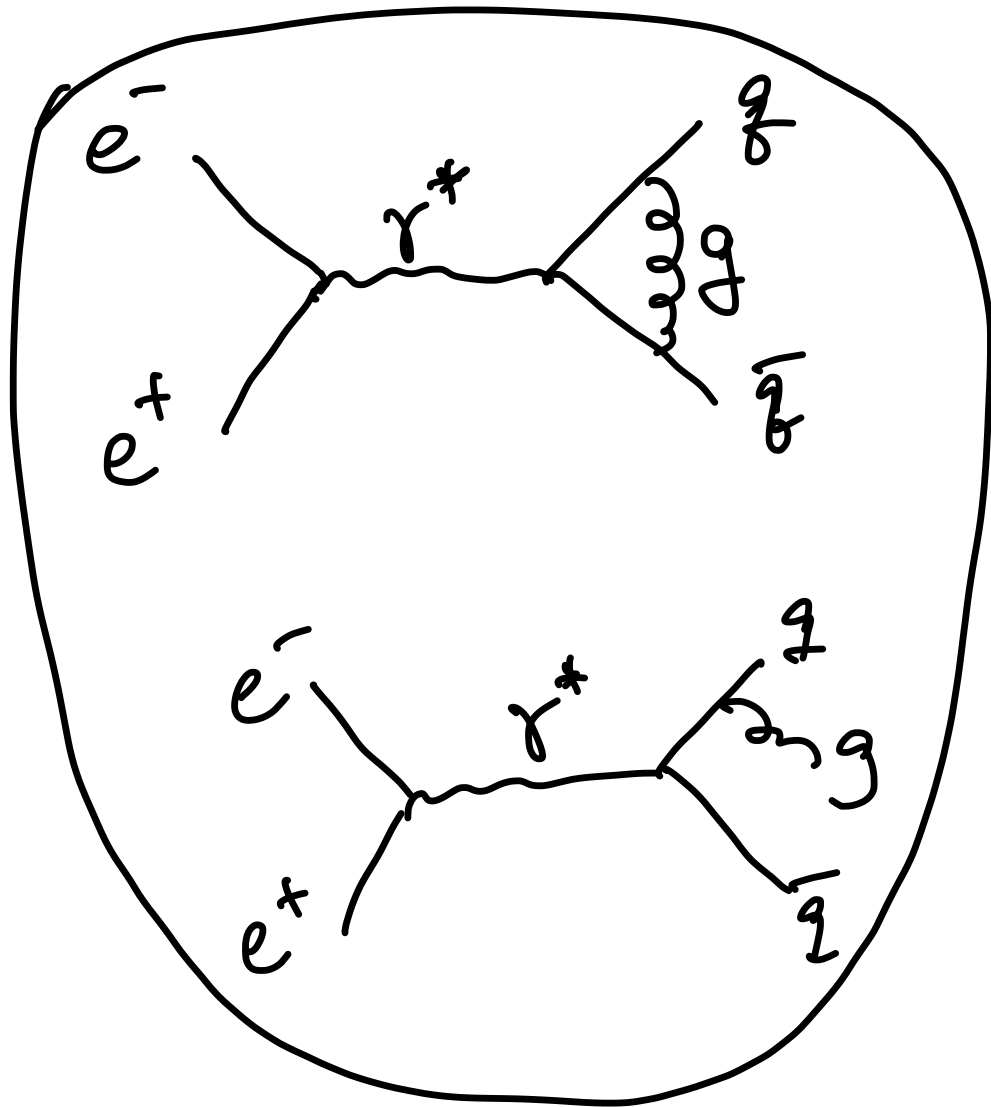
sum over "active" flavors

$$\sqrt{s} > 2m_f$$

Evidence for 3 colors

★ Next-to-leading order

$$\sigma = \sigma_0 + \boxed{\alpha_s \text{ (loop)}}$$



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