

Two Loop Matching for Quasi PDF

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1 Renormalization

1.1 One loop diagrams

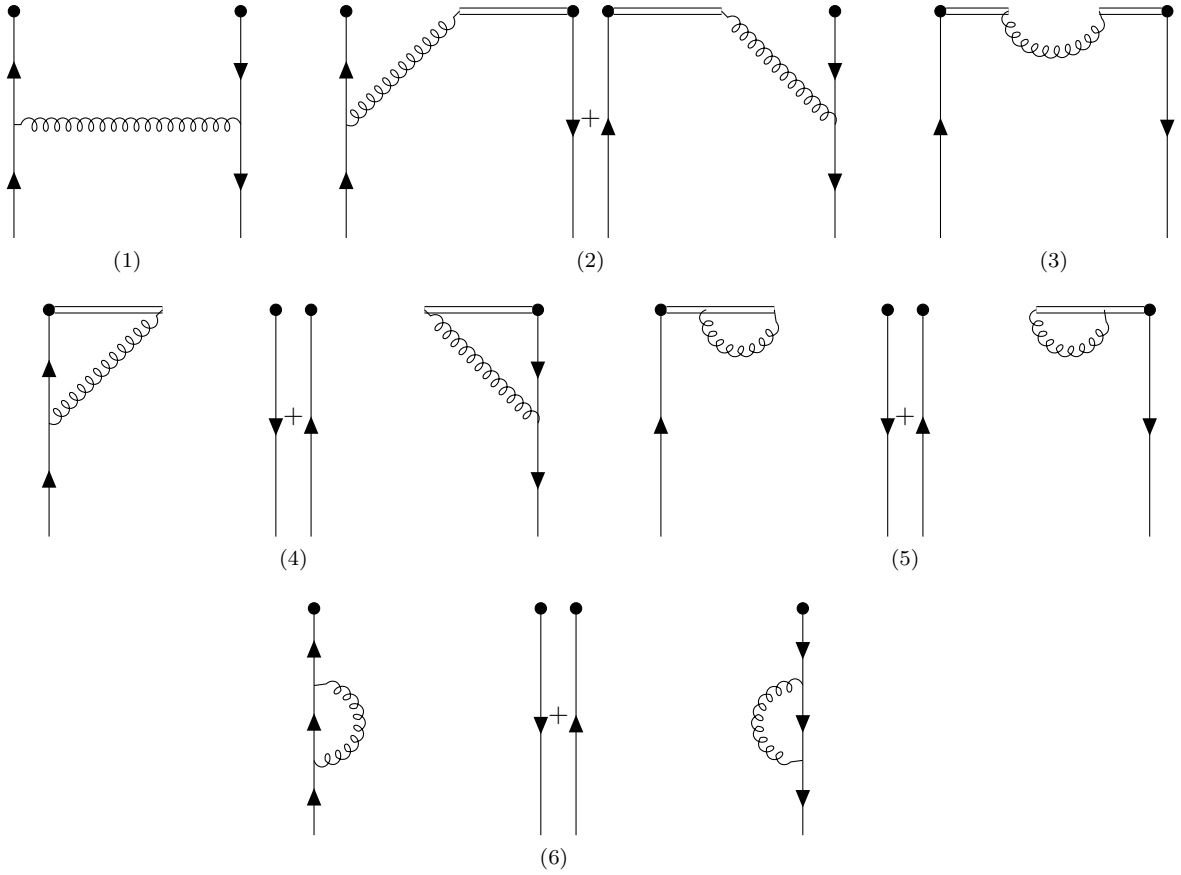
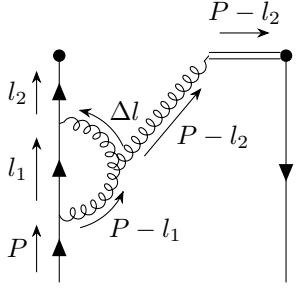


Figure 1: Diagrams of quasi PDF in Feynman gauge.

1.2 Vertex corrections

According to [Ji and Zhang(2015)], the vertex correction diagrams in axial gauge (which corresponds to varieties of diagrams in general covariant gauge) don't have total UV divergence. Rather, they only have subdivergence for sub-diagrams. For example the first column (which involves Figure 3), second row of Table 1 in [Ji and Zhang(2015)] is composed of \tilde{q}_{11} and \tilde{q}_{12} , thus we can find some representative diagrams and extract those components ($l \equiv l_1 + l_2, \Delta l \equiv$

$l_1 - l_2$)



$$P \propto \int \frac{d^d l_1}{(2\pi)^d} \frac{d^d l_2}{(2\pi)^d} \frac{1}{[l_1 - m][l_2 - m][(P - l_1)^2][(l_1 - l_2)^2][(P - l_2)^2][n \cdot (P - l_2)]} \quad (1)$$

Take the $l_1 \gg l_2$ limit, the integrand becomes

$$\frac{1}{[l_1 - m][(P - l_1)^2][l_1^2][l_2 - m][(P - l_2)^2][n \cdot (P - l_2)]} \quad (2)$$

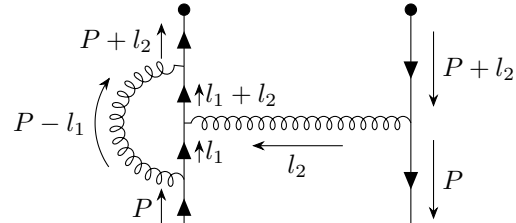
The integral involving l_2 is exactly the integral of \tilde{q}_{12} . By adding the gluon self-interacting vertex we can see that the sub-diagram is logarithmic divergent.

Take the $l_2 \gg l_1$ limit, the integrand becomes

$$\frac{1}{[l_1 - m][(P - l_1)^2][l_2^2][l_2 - m][(P - l_2)^2][n \cdot (P - l_2)]} \quad (3)$$

There's another limit where hard loop momentum flows through all paths except the one that's Δl in our current diagram. This configuration gives a finite integral and a power-divergent integral which happens to be a scaleless integral as well. Thus this configuration won't contribute.

What we extracted above is only the \tilde{q}_{12} part, now we will try on the \tilde{q}_{11} part



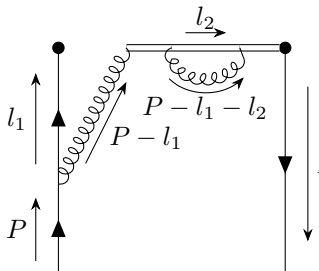
$$\propto \int \frac{d^d l_1}{(2\pi)^d} \frac{d^d l_2}{(2\pi)^d} \frac{1}{[l_1 - m][l_1 + l_2 - m][\not{P} + l_2 - m][\not{P} + l_2 - m][(P - l_1)^2][l_2^2]} \quad (4)$$

In the $l_1 \gg l_2$ limit we have

$$\frac{1}{[l_1 - m][l_1 - m][(P - l_1)^2][\not{P} + l_2 - m][\not{P} + l_2 - m][l_2^2]} \quad (5)$$

and \tilde{q}_{11} is factorized out.

Another example is the sixth row



$$P \propto \int \frac{d^d l_1}{(2\pi)^d} \frac{d^d l_2}{(2\pi)^d} \frac{1}{[l_1 - m][(P - l_1)^2][(P - l_1 - l_2)^2][n \cdot (P - l_1)][n \cdot l_2][n \cdot (P - l_1)]} \quad (6)$$

Take the $l_2 \gg l_1$ limit, the integrand becomes

$$\frac{1}{[l_1 - m] [(P - l_1)^2] [n \cdot (P - l_1)] [n \cdot (P - l_1)] [n \cdot l_2] [(P - l_2)^2]} \quad (7)$$

and the integral involving l_2 should give something proportional to $n \cot(P - l_1)$, thus cancels one eikonal propagator, the remainder is the integral of \tilde{q}_{12} .

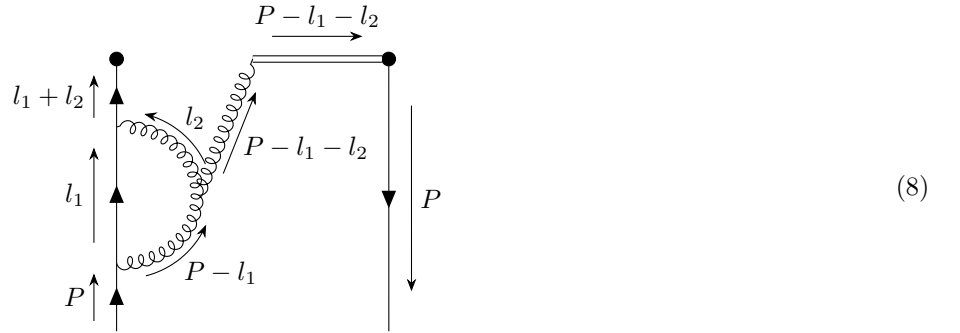
2 Real Diagrams

2.1 All diagrams

Figure 2 lists all self-conjugated real diagrams, and Figure 3 lists all non-self-conjugated diagrams, excluding their conjugates.

2.2 A first attempt

Let's first take a look at the following diagram



And

3 Virtual Diagrams (Excluding Gauge Link Self-Energy Diagrams)

3.1 All diagrams

4 Gauge Link Self-Energy Diagrams

4.1 All diagrams

5 HQET Correspondence

References

[Ji and Zhang(2015)] X. Ji and J.-H. Zhang, *Phys. Rev.* **D92**, 034006 (2015), [arXiv:1505.07699 \[hep-ph\]](#) .

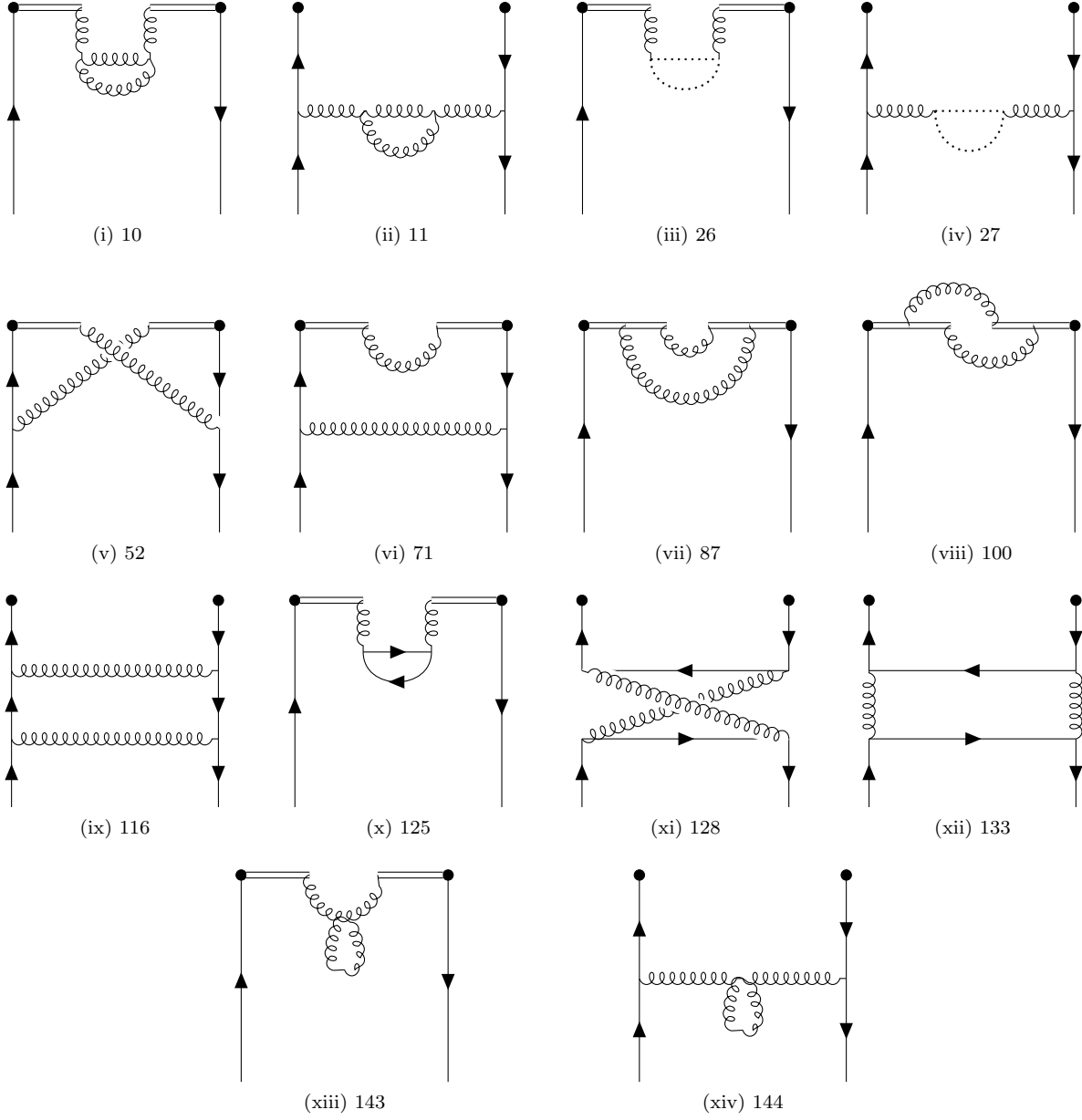


Figure 2: All self-conjugated diagrams.

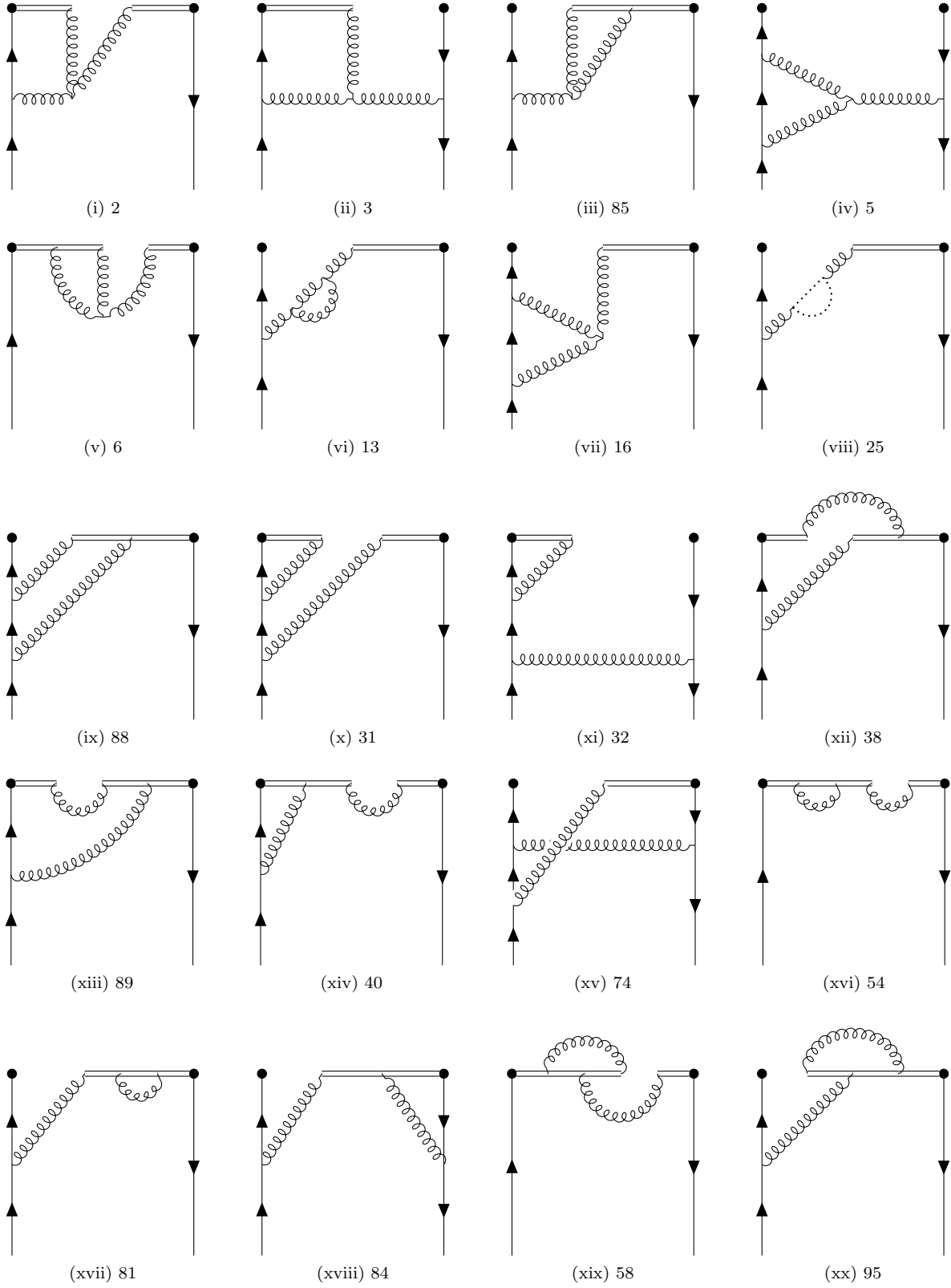


Figure 3: All real diagrams (excluding conjugated diagrams).

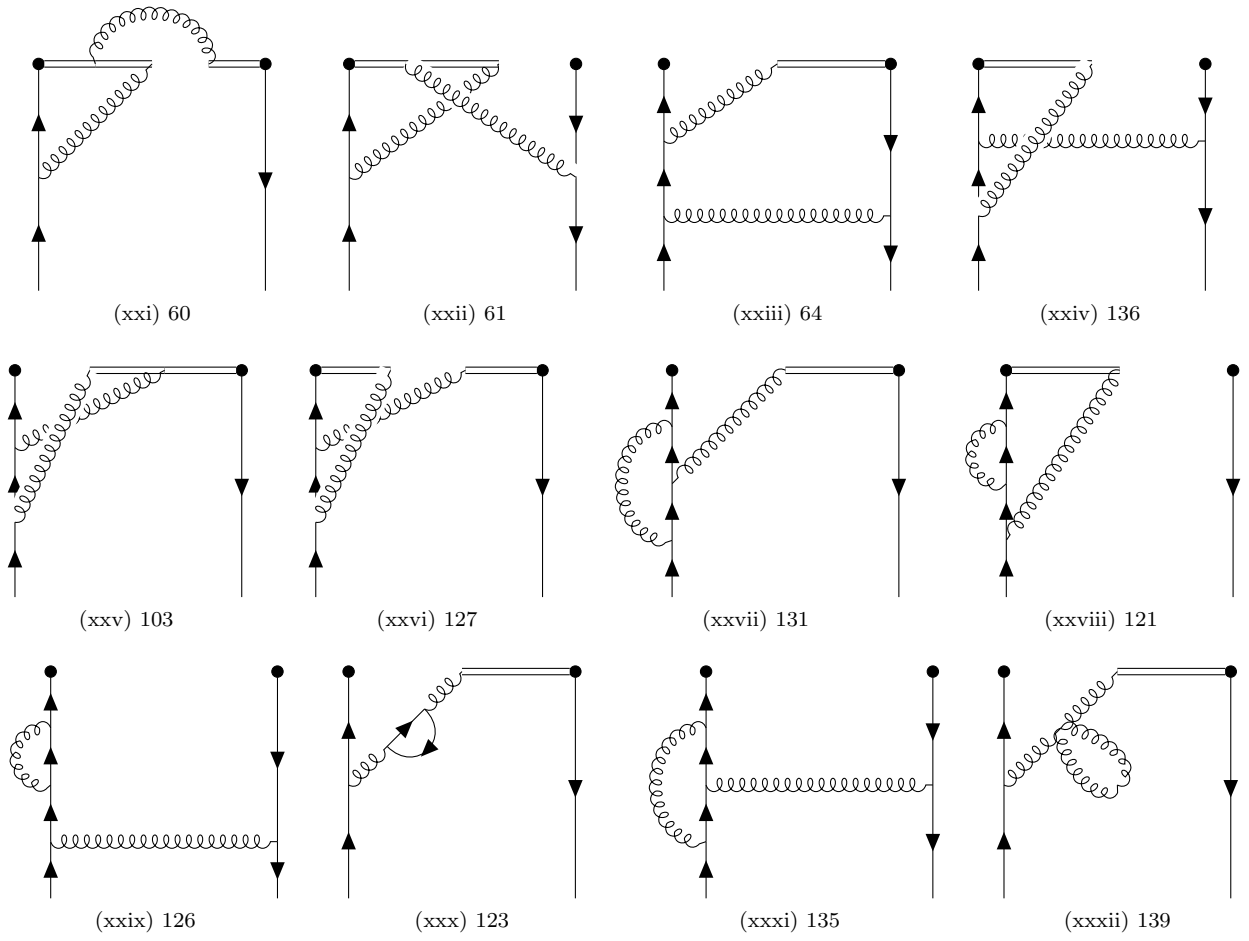


Figure 3: All real diagrams (excluding conjugated and self-conjugated diagrams).

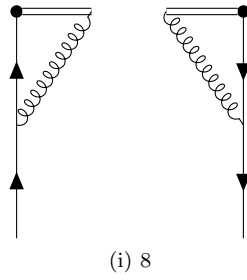


Figure 4: All self-conjugated virtual diagrams (actually there's only one).

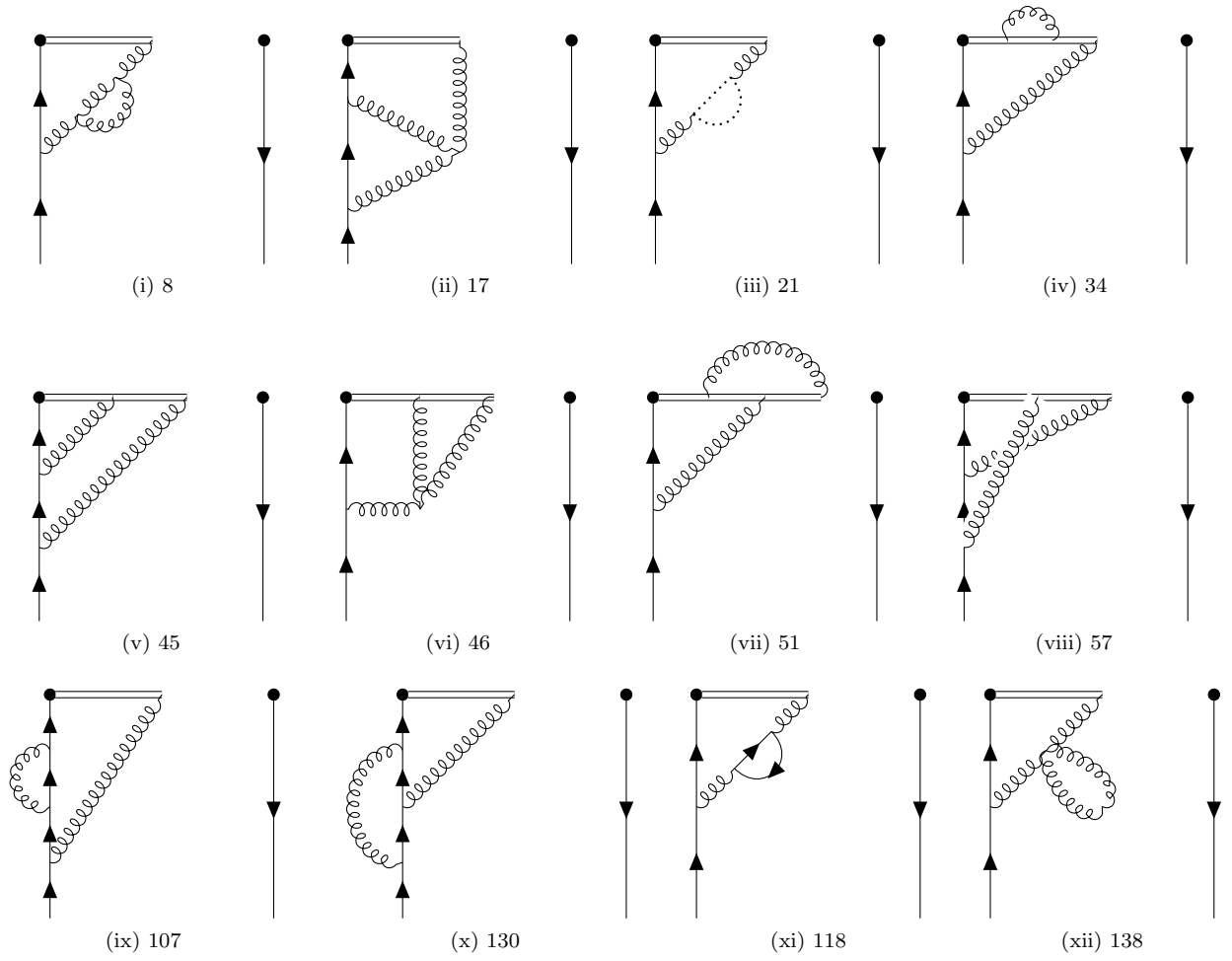


Figure 5: All virtual diagrams (excluding conjugated and self-conjugated diagrams).

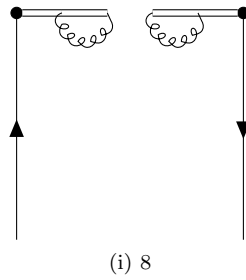


Figure 6: All self-conjugated gauge link self-energy diagrams (actually there's only one).

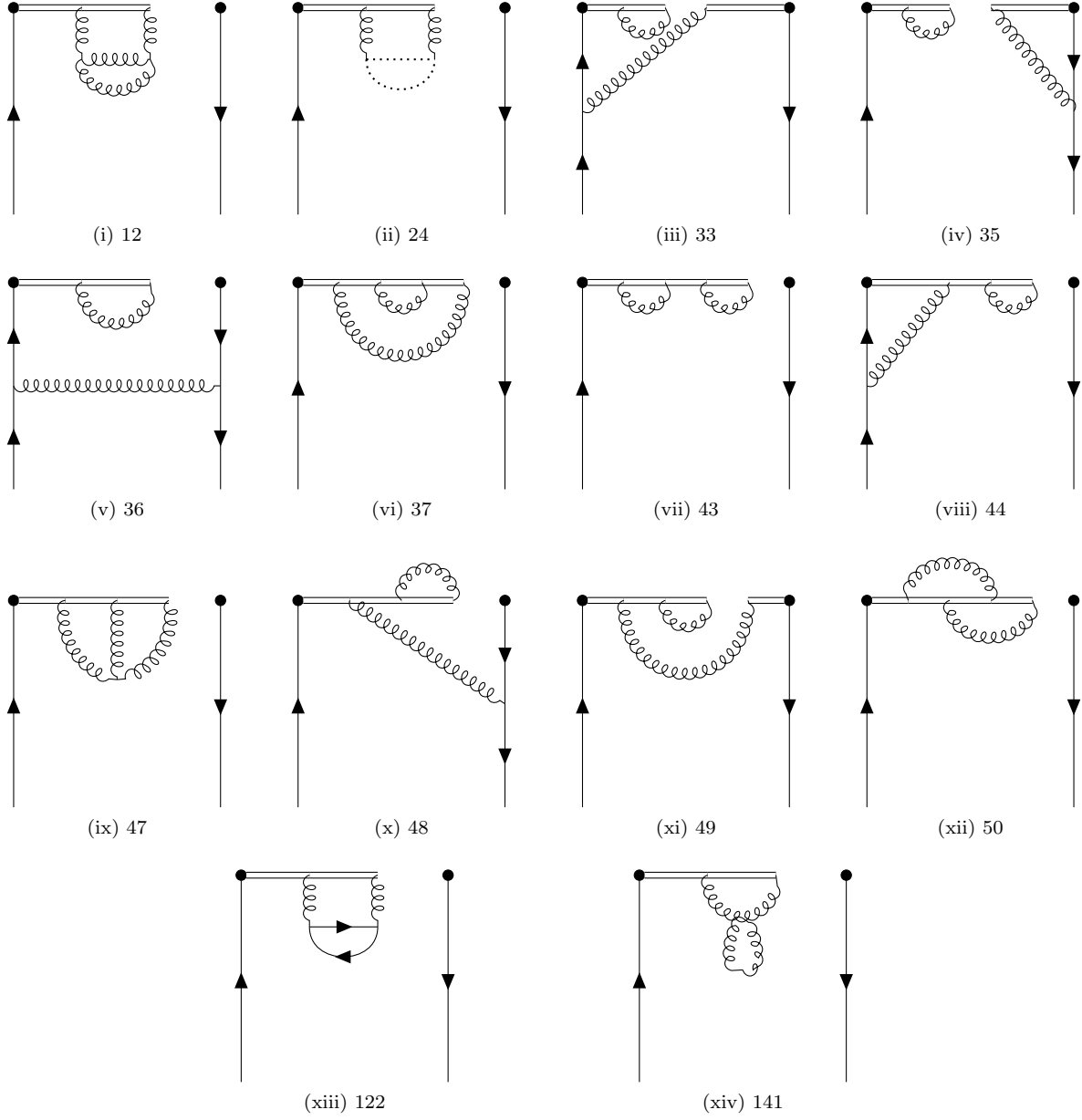


Figure 7: