$\bar{c}\gamma^{\mu}c$ matrix element

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1 Kinematics and Conventions

Quark and antiquark momenta are

$$p_1 = P/2 + p = (E, \mathbf{p}) \tag{1}$$

$$p_2 = P/2 - p = (E, -\mathbf{p}) \tag{2}$$

where in rest frame

$$P = (2E(p), 0) \tag{3}$$

$$p = (0, \mathbf{p}) \tag{4}$$

Dirac spinors are normalized as following

$$u(\mathbf{p}) = \sqrt{\frac{E+m}{2E}} \left(\frac{\xi}{\frac{\mathbf{p} \cdot \sigma}{E+m}} \xi \right) \tag{5}$$

$$v(\mathbf{p}) = \sqrt{\frac{E+m}{2E}} \begin{pmatrix} \frac{-\mathbf{p} \cdot \sigma}{E+m} \eta \\ \eta \end{pmatrix} \tag{6}$$

where

$$\xi = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \text{ or } \begin{pmatrix} 0 \\ 1 \end{pmatrix} \tag{7}$$

$$\eta = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \text{ or } \begin{pmatrix} -1 \\ 0 \end{pmatrix} \tag{8}$$

2 State Projection

The bound state is [Weinberg(2015)]

$$|P, E; J, m_j; L; S\rangle = \int d\Omega_{\mathbf{p_1}} \sum_{\substack{\mathbf{s_1, s_2, \dots, m_l}\\ \mathbf{s_1, s_2, s_2, \dots, m_l}}} Y_l^m(\hat{\mathbf{p_1}}) \langle S_1 \lambda_1 S_2 \lambda_2 | S s_z \rangle \langle S s_z L m_l | J m_J \rangle |\mathbf{p_1}, \lambda_1 \rangle |\mathbf{P} - \mathbf{p_1}, \lambda_2 \rangle$$

$$(9)$$

3 Bilinears [Bodwin and Petrelli(2002)]

$$\Pi_0(P, p) \equiv -\sum_{\lambda_1, \lambda_2} u(\mathbf{p}, \lambda_1) \bar{v}(-\mathbf{p}, \lambda_2) \left\langle \frac{1}{2} \lambda_1 \frac{1}{2} \lambda_2 \middle| 00 \right\rangle$$
(10)

$$= \frac{1}{2\sqrt{2}E(E+m)} \left(\frac{1}{2} \not\!\!\!P + m + \not\!\!p\right) \frac{\not\!\!P + 2E}{4E} \gamma_5 \left(\frac{1}{2} \not\!\!\!P - m - \not\!\!p\right) \tag{11}$$

$$\Pi_1(P, p) \equiv \sum_{\lambda_1, \lambda_2} u(\mathbf{p}, \lambda_1) \bar{v}(-\mathbf{p}, \lambda_2) \left\langle \frac{1}{2} \lambda_1 \frac{1}{2} \lambda_2 \middle| 1\epsilon \right\rangle$$
(12)

$$=\frac{-1}{2\sqrt{2}E(E+m)}\left(\frac{1}{2}P+m+p\right)\frac{P+2E}{4E}\not\in\left(\frac{1}{2}P-m-p\right) \tag{13}$$

4 ${}^{3}S_{1}$

Average over **p** with Bodwin's convention (extra $1/(4\pi)$):

$$\left\langle 0 \left| \bar{c} \gamma^{\mu} c \right|^{3} S_{1} \right\rangle^{(0)} = \frac{1}{4\pi} \int d\Omega \operatorname{tr}[\Pi_{1} \gamma^{\mu}] = \sqrt{2} \left(\frac{m}{3E} + \frac{2}{3} \right) \epsilon^{\mu}$$

$5 ^3D_1$

The matrix element reads:

$$\langle 0|\bar{c}\gamma^{\mu}c|^{3}D_{1}\rangle^{(0)} = \int d\Omega \sum_{\lambda_{1}\lambda_{2}s_{z}m_{l}} \operatorname{tr}\{\Pi_{1}\gamma^{\mu}\} \langle 2m_{l}; 1s_{z}|1J_{z}\rangle Y_{2m_{l}}(\theta,\phi)$$

while the trace part is the same as ${}^{3}S_{1}$:

$$\operatorname{tr}\{\Pi_1 \gamma^{\mu}\} = \frac{\sqrt{2}p^{\mu}(p \cdot \epsilon)}{E(E+m)} + \epsilon^{\mu}$$

Chosen polarization vectors:

$$\epsilon^{(-)} = \frac{1}{\sqrt{2}}(0, 1, -i, 0), \epsilon^{(0)} = (0, 0, 0, 1), \epsilon^{(+)} = \frac{1}{\sqrt{2}}(0, -1, -i, 0)$$

Result (the first row and the last are orthogonal):

$$\begin{pmatrix}
\epsilon^{(-)} & 0 & \frac{2\sqrt{2\pi}p^2}{3E(m+E)} & -\frac{2i\sqrt{2\pi}p^2}{3E(m+E)} & 0\\
\hline
\epsilon^{(0)} & 0 & 0 & 0 & \frac{4\sqrt{\pi}p^2}{3E(m+E)} \\
\hline
\epsilon^{(+)} & 0 & -\frac{2\sqrt{2\pi}p^2}{3E(m+E)} & -\frac{2i\sqrt{2\pi}p^2}{3E(m+E)} & 0
\end{pmatrix}$$

and the decay constant is $\frac{4\sqrt{\pi}p^2}{3E(E+m)}$ where $p^2={\bf p}^2=E^2-m^2$.

6 NRQCD 3D_1

$$\langle 0|\frac{1}{2m^2}\chi^{\dagger}D^{\{i}D^{j\}}\sigma^{j}\psi|Q\overline{Q}[^{3}D_{1}(\varepsilon)]\rangle^{(0)} = \int d\Omega \sum_{\lambda_{1}\lambda_{2}s_{z}m_{l}} \frac{-p^{\{i}p^{j\}}}{2m^2}\eta_{\lambda_{2}}^{\dagger}\sigma^{j}\xi_{\lambda_{1}}\langle 2m_{l};1s_{z}|1J_{z}\rangle\left\langle \frac{1}{2}\lambda_{1}\frac{1}{2}\lambda_{2}\Big|1s_{z}\right\rangle Y_{2m_{l}}(\theta,\phi) \qquad (14)$$

$$= \sqrt{\pi}\frac{4p^{2}}{2m^{2}}\varepsilon^{i}$$

if $\partial^{\{i}\partial^{j\}} = 2\partial^i\partial^j$.

References

[Weinberg (2015)] S. Weinberg, Lectures on Quantum Mechanics (Cambridge University Pr., 2015).

[Bodwin and Petrelli(2002)] G. T. Bodwin and A. Petrelli, Phys. Rev. **D66**, 094011 (2002), [Erratum: Phys. Rev. D87,no.3,039902(2013)], arXiv:hep-ph/0205210 [hep-ph].