

Note on Braaten's Paper

Yingsheng Huang

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1 Intro

Hamiltonian [1]:

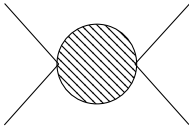
$$\mathcal{H} = \sum_{\sigma} \frac{1}{2m} \nabla \psi_{\sigma}^{\dagger} \cdot \nabla \psi_{\sigma}^{(\Lambda)} + \frac{g(\Lambda)}{m} \psi_1^{\dagger} \psi_2^{\dagger} \psi_3 \psi_4^{(\Lambda)} + \mathcal{V} \quad (1)$$

where the renormalized coupling

$$g(\Lambda) = \frac{4\pi a}{1 - 2a\Lambda/\pi} \quad (2)$$

2 Amplitude

Consider:

$$i\mathcal{A} = \langle 34 | \psi^{\dagger} \psi | 12 \rangle = \text{diagram} \quad (3)$$


Define $P = p_1 + p_2 = (E, \mathbf{0})$, and $E = p^2/m$. The integral equation is

$$i\mathcal{A} = -\frac{ig(\Lambda)}{m} \left(1 + i\mathcal{A} \int [d^4k] \frac{i}{k^0 - \frac{\mathbf{k}^2}{2m} + i\epsilon} \frac{i}{k^0 - p^0 - \frac{|\mathbf{k}-\mathbf{p}|^2}{2m} + i\epsilon} \right) \quad (4)$$

The integral gives (redefine $\epsilon \rightarrow 2m\epsilon$)

$$\mathcal{I} = \frac{im \left(-\Lambda + \sqrt{-mE - i\epsilon} \tan^{-1} \left(\frac{\Lambda}{\sqrt{-mE - i\epsilon}} \right) \right)}{2\pi^2} \quad (5)$$

and

$$i\mathcal{A} = \frac{-1}{\mathcal{I} + \frac{m}{ig(\Lambda)}} = \frac{-1}{\frac{im\sqrt{-mE - i\epsilon} \tan^{-1} \left(\frac{\Lambda}{\sqrt{-mE - i\epsilon}} \right)}{2\pi^2} - \frac{im}{4\pi a}} \quad (6)$$

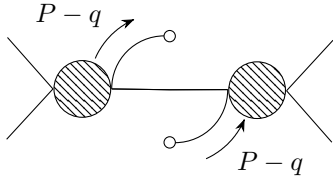
$$\xrightarrow{\Lambda \rightarrow \infty} \frac{4i\pi/m}{-1/a + \sqrt{-mE - i\epsilon}} \quad (7)$$

Question: Why only infinite bubbles in s-channel are considered? What about other channels?

3 OPE

3.1 l.h.s.

Take what we got in the last section as a new non-perturbative vertex, we only need to deal with tree diagram this way. First we have Figure 2(a) in Braaten's paper:



$$= \langle 34 | \psi^\dagger \left(-\frac{\mathbf{r}}{2} \right) \psi \left(\frac{\mathbf{r}}{2} \right) | 12 \rangle \quad (8)$$

$$= \mathcal{A}^2 \int \frac{d^4 q}{(2\pi)^4} \frac{i}{q^0 - \frac{\mathbf{q}^2}{2m} + i\epsilon} \frac{i}{q^0 - \frac{E - \mathbf{q}^2}{2m} + i\epsilon} \frac{i}{E - q^0 - \frac{\mathbf{q}^2}{2m} + i\epsilon} e^{i\mathbf{q} \cdot \mathbf{r}} \quad (9)$$

$$= -\mathcal{A}^2 \int \frac{d^3 \mathbf{q}}{(2\pi)^3} \frac{m^2 e^{i\mathbf{q} \cdot \mathbf{r}}}{(\mathbf{q}^2 - p^2 - i\epsilon)^2} \quad (10)$$

$$= -\frac{im^2 \mathcal{A}^2 e^{ipr}}{8\pi p} \quad (11)$$

References

- [1] Eric Braaten and Lucas Platter. Exact relations for a strongly interacting fermi gas from the operator product expansion. *Phys. Rev. Lett.*, 100(20), may 2008.