

Homework: Quantum Field Theory #5

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1. Complete table at P&S P71.

For the first three column, they're already done in class so I simply give the results. (We use the shorthand in P&S: $(-1)^\mu = 1$ for $\mu = 0$ and $(-1)^\mu = -1$ for $\mu = 1, 2, 3$.¹)

$$\begin{aligned}
 P\bar{\psi}\psi P^{-1} &= +\bar{\psi}\psi(t, -\mathbf{x}) \\
 T\bar{\psi}\psi T^{-1} &= +\bar{\psi}\psi(-t, \mathbf{x}) \\
 C\bar{\psi}\psi C^{-1} &= +\bar{\psi}\psi(t, \mathbf{x}) \\
 CPT\bar{\psi}\psi CPT^{-1} &= +\bar{\psi}\psi(-t, -\mathbf{x}) \\
 P\bar{\psi}\gamma^5\psi P^{-1} &= -\bar{\psi}\gamma^5\psi(t, -\mathbf{x}) \\
 T\bar{\psi}\gamma^5\psi T^{-1} &= -\bar{\psi}\gamma^5\psi(-t, \mathbf{x}) \\
 C\bar{\psi}\gamma^5\psi C^{-1} &= +\bar{\psi}\gamma^5\psi(t, \mathbf{x}) \\
 CPT\bar{\psi}\gamma^5\psi CPT^{-1} &= +\bar{\psi}\gamma^5\psi(-t, -\mathbf{x}) \\
 P\bar{\psi}\gamma^\mu\psi P^{-1} &= (-1)^\mu \bar{\psi}\gamma^\mu\psi(t, -\mathbf{x}) \\
 T\bar{\psi}\gamma^\mu\psi T^{-1} &= (-1)^\mu \bar{\psi}\gamma^\mu\psi(-t, \mathbf{x}) \\
 C\bar{\psi}\gamma^\mu\psi C^{-1} &= -\bar{\psi}\gamma^\mu\psi(t, \mathbf{x}) \\
 CPT\bar{\psi}\gamma^\mu\psi CPT^{-1} &= -\bar{\psi}\gamma^\mu\psi(-t, -\mathbf{x})
 \end{aligned}$$

Now we calculate the rest.

Given

$$\begin{aligned}
 P\psi P^{-1} &= \eta\gamma^0\psi(t, -\mathbf{x}) \\
 P\bar{\psi} P^{-1} &= \eta^*\bar{\psi}(t, -\mathbf{x})\gamma^0
 \end{aligned}$$

we have

$$\begin{aligned}
 P\bar{\psi}\gamma^\mu\gamma^5\psi P^{-1} &= |\eta|^2 \bar{\psi}\gamma^0\gamma^\mu\gamma^5\gamma^0\psi \\
 &= -(-1)^\mu \bar{\psi}\gamma^\mu\gamma^5\psi
 \end{aligned}$$

and

$$\begin{aligned}
 P\bar{\psi}\sigma^{\mu\nu}\psi P^{-1} &= \frac{i}{2} \bar{\psi}\gamma^0[\gamma^\mu, \gamma^\nu]\gamma^0\psi \\
 &= \frac{i}{2} (-1)^\mu (-1)^\nu \bar{\psi}[\gamma^\mu, \gamma^\nu]\psi \\
 &= (-1)^\mu (-1)^\nu \bar{\psi}\sigma^{\mu\nu}\psi
 \end{aligned}$$

¹And I write $(CPT)^{-1}$ as CPT^{-1} for short.

and similarly

$$P\bar{\psi}\partial_\mu\psi P^{-1} = (-1)^\mu\bar{\psi}\partial_\mu\psi$$

Define

$$\mathcal{T} \equiv i\gamma^1\gamma^3$$

and

$$\begin{aligned} T\psi T^{-1} &= \mathcal{T}\psi \\ T\bar{\psi} T^{-1} &= \bar{\psi}\mathcal{T}^{-1} \\ \mathcal{T}(\gamma^\mu)^*\mathcal{T}^{-1} &= \gamma_\mu = (-1)^\mu\gamma^\mu \\ \mathcal{T}(\gamma^5)^*\mathcal{T}^{-1} &= \gamma^5 \\ \mathcal{T} &= \mathcal{T}^{-1} = \mathcal{T}^\dagger \end{aligned}$$

we have

$$\begin{aligned} T\bar{\psi}\gamma^\mu\gamma^5\psi T^{-1} &= \bar{\psi}\mathcal{T}^{-1}(\gamma^\mu\gamma^5)^*\mathcal{T}\psi \\ &= \bar{\psi}\mathcal{T}^{-1}\gamma^\mu*\mathcal{T}^{-1}\mathcal{T}\gamma^5*\mathcal{T}\psi \\ &= \bar{\psi}\gamma_\mu\gamma^5\psi \\ &= (-1)^\mu\bar{\psi}\gamma^\mu\gamma^5\psi \end{aligned}$$

and

$$\begin{aligned} T\bar{\psi}\sigma^{\mu\nu}\psi T^{-1} &= -\frac{i}{2}T\bar{\psi}[\gamma^\mu, \gamma^\nu]\psi T^{-1} \\ &= -\frac{i}{2}\bar{\psi}\mathcal{T}[\gamma^\mu, \gamma^\nu]^*\mathcal{T}^{-1}\psi \\ &= -(-1)^\mu(-1)^\nu\bar{\psi}\sigma^{\mu\nu}\psi \end{aligned}$$

and

$$T\bar{\psi}\partial_\mu\psi T^{-1} = -(-1)^\mu\bar{\psi}\partial_\mu\psi$$

Define

$$\mathcal{C} \equiv i\gamma^2\gamma^0$$

and

$$\begin{aligned} C\psi C^{-1} &= \mathcal{C}\bar{\psi}^T \\ C\bar{\psi} C^{-1} &= \psi^T\mathcal{C} \\ \mathcal{C}(\gamma^\mu)^T\mathcal{C}^{-1} &= -\gamma^\mu \\ \mathcal{C}(\gamma^5)^T\mathcal{C}^{-1} &= \gamma^5 \\ \mathcal{C}^\dagger &= \mathcal{C}^{-1} = -\mathcal{C} = \mathcal{C}^T \end{aligned}$$

thus

$$\begin{aligned} (\mathcal{C}(\gamma^\mu)^T\mathcal{C}^{-1})^\dagger &= -(\gamma^\mu)^\dagger \\ &= (\mathcal{C}^{-1})^\dagger(\gamma^\mu)^*\mathcal{C}^\dagger \\ &= -\mathcal{C}^\dagger(\gamma^\mu)^*\mathcal{C}^{-1} \\ &= \mathcal{C}(\gamma^\mu)^*\mathcal{C}^{-1} = -(\gamma^\mu)^\dagger \end{aligned}$$

and

$$\mathcal{C}\gamma^5\mathcal{C}^{-1} = \gamma^5$$

Then we have

$$\begin{aligned} C\bar{\psi}\gamma^\mu\gamma^5\psi C^{-1} &= \psi^T\mathcal{C}\gamma^\mu\gamma^5\mathcal{C}\bar{\psi}^T \\ &= \psi^T\gamma^{\mu T}\gamma^{5T}\bar{\psi}^T \\ &= -(\bar{\psi}\gamma^5\gamma^\mu\psi)^T \\ &= \bar{\psi}\gamma^\mu\gamma^5\psi \end{aligned}$$

and

$$\begin{aligned} C\bar{\psi}\sigma^{\mu\nu}\psi C^{-1} &= \frac{i}{2}\psi^T\mathcal{C}[\gamma^\mu, \gamma^\nu]\mathcal{C}\bar{\psi}^T \\ &= -\frac{i}{2}\psi^T[\gamma^{\mu T}, \gamma^{\nu T}]\bar{\psi}^T \\ &= \frac{i}{2}(\bar{\psi}[\gamma^\nu, \gamma^\mu]\psi)^T \\ &= -\bar{\psi}\sigma^{\mu\nu}\psi \end{aligned}$$

and

$$C\bar{\psi}\partial_\mu\psi C^{-1} = \bar{\psi}\partial_\mu\psi$$

CPT is to multiply all those coefficients and too trivial to list here.

2. Calculate the Dirac propagator.

$$\begin{aligned} \langle 0|\psi_a(x)\bar{\psi}_b(y)|0\rangle &= \int \frac{d^3p}{(2\pi)^3} \frac{1}{2E_{\mathbf{p}}} \sum_s u_a^s(p)\bar{u}_b^s(p)e^{-ip\cdot(x-y)} = (i\cancel{\partial}_x + m)_{ab} \int \frac{d^3p}{(2\pi)^3} \frac{1}{2E_{\mathbf{p}}} e^{-ip\cdot(x-y)} \\ \langle 0|\bar{\psi}_b(y)\psi_a(x)|0\rangle &= \int \frac{d^3p}{(2\pi)^3} \frac{1}{2E_{\mathbf{p}}} \sum_s v_a^s(p)\bar{v}_b^s(p)e^{-ip\cdot(x-y)} = -(i\cancel{\partial}_x + m)_{ab} \int \frac{d^3p}{(2\pi)^3} \frac{1}{2E_{\mathbf{p}}} e^{ip\cdot(x-y)} \end{aligned}$$

The definition of Dirac propagator

$$S_F(x-y) = \int \frac{d^4p}{(2\pi)^4} \frac{i(\not{p} + m)}{p^2 - m^2 + i\epsilon} e^{-ip\cdot(x-y)}$$

the two poles (if $\epsilon = 0$) are located in $p^0 = \omega_p \equiv \sqrt{\mathbf{p}^2 + m^2}$ and $p^0 = -\omega_p$.

For $x^0 - y^0 > 0$, we have

$$\begin{aligned} \int \frac{dp^0}{2\pi} \frac{i(\not{p} + m)}{p^2 - m^2 + i\epsilon} e^{-ip^0(x^0 - y^0)} &= \int \frac{dp^0}{2\pi} \frac{i(\not{p} + m)}{p_0^2 - \omega_p^2 + i\epsilon} e^{-ip^0(x^0 - y^0)} \\ &= \int \frac{dp^0}{2\pi} \frac{i(\not{p} + m)}{(p_0 - \omega_p + i\epsilon)(p_0 + \omega_p - i\epsilon)} e^{-ip^0(x^0 - y^0)} \\ &= \frac{1}{2(\omega_p - i\epsilon)} \int \frac{dp^0}{2\pi} i(\not{p} + m) \left(\frac{1}{p_0 - \omega_p + i\epsilon} - \frac{1}{p_0 + \omega_p - i\epsilon} \right) e^{-ip^0(x^0 - y^0)} \\ &= \frac{i\cancel{\partial}_x + m}{2(\omega_p - i\epsilon)} \int \frac{dp^0}{2\pi} i \left(\frac{1}{p_0 - \omega_p + i\epsilon} - \frac{1}{p_0 + \omega_p - i\epsilon} \right) e^{-ip^0(x^0 - y^0)} \end{aligned}$$

use residue theorem, and the contour is closed below (only one singularity on the right)

$$\begin{aligned}
&= \frac{i\vec{\phi}_x + m}{2(\omega_p - i\epsilon)} e^{-ip^0(x^0 - y^0)} \\
&= \frac{i\vec{\phi}_x + m}{2\omega_p} e^{-ip^0(x^0 - y^0)}
\end{aligned}$$

which means

$$S_F(x - y) = \langle 0 | \psi_a(x) \bar{\psi}_b(y) | 0 \rangle$$

Similarly, when $x^0 - y^0 < 0$, the contour is closed above and

$$S_F(x - y) = -\frac{i\vec{\phi}_x + m}{2\omega_p} e^{ip^0(x^0 - y^0)} = -\langle 0 | \bar{\psi}_b(y) \psi_a(x) | 0 \rangle$$

Appendix

	γ^0	γ^1	γ^2	γ^3	γ^5
T	1	-1	1	-1	1
-1	1	-1	-1	-1	1
*	1	1	-1	1	1
†	1	-1	-1	-1	1