

# Homework: Gauge Field Theory #3

Yingsheng Huang

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1. The standard model Lagrangian without fermion part and gluon part (since we only need to consider Faddeev-Popov determinant, the effect of Higgs field can also be neglected, unless we wish to include the effect in  $R_\xi$  gauge fixing term):

$$\mathcal{L} = -\frac{1}{2} \text{Tr} W^{\mu\nu} W_{\mu\nu} - \frac{1}{4} B^{\mu\nu} B_{\mu\nu} + |D^\mu \phi|^2 - V(\phi)$$

note that  $W$  has 3 colors, the covariant derivatives are

$$D^\mu = \partial^\mu - igW^{a,\mu}T^a + ig_B B^\mu$$

and the gauge transforms are:

U(1):

$$B^\mu \rightarrow B^\mu - \frac{1}{g_B} \partial^\mu \beta$$

$$\phi \rightarrow e^{i\beta(x)} \phi$$

SU(2):

$$W^{a,\mu} \rightarrow W^{a,\mu} + \frac{1}{g} \partial^\mu \alpha^a - f^{abc} \alpha^b W^{c,\mu}$$

$$\phi \rightarrow e^{i\alpha^a(x)T^a} \phi$$

Make  $\phi = \frac{1}{\sqrt{2}}(v + h(x))$ , With  $R_\xi$  gauge, the gauge fixing term is

$$\mathcal{L}_{GF} = -\frac{1}{2\xi} (\partial_\mu W^{a,\mu} - \xi g T_{ij}^a v_j h_i)^2 - \frac{1}{2\xi} (\partial_\mu B^{a,\mu} - \xi g_B v_i h_i)^2$$

2. QED Lagrangian:

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \bar{\psi}(i\not{D} + m)\psi$$

We can add gauge fixing term

$$\mathcal{L}_{GF} = -\frac{1}{2\xi} (\partial_\mu A^\mu)^2$$

and ignore the fermion part, the generating functional is then

$$Z[J] = \int D[A] e^{i \int d^4x (-\frac{1}{4} F_{\mu\nu} F^{\mu\nu} - \frac{1}{2\xi} (\partial_\mu A^\mu)^2 - J^\mu A_\mu)}$$

Note that the kinetic term can be rewrite as

$$-\frac{1}{4} F_{\mu\nu} F^{\mu\nu} = -\frac{1}{2} (\partial_\mu A^\nu)^2 + \frac{1}{2} (\partial_\mu A^\mu)^2$$

so

$$Z[J] = \int D[A] e^{i \int d^4x (-\frac{1}{2} (\partial_\mu A^\nu)^2 + \frac{1-\xi^{-1}}{2} (\partial_\mu A^\mu)^2 - J^\mu A_\mu)}$$

$$= \int D[A] e^{i \int d^4x (A_\mu (\frac{1}{2} g^{\mu\nu} \partial^2 + \frac{1-\xi^{-1}}{2} \partial^\mu \partial^\nu) A_\nu - J^\mu A_\mu)}$$

The propagator is then

$$(g^{\mu\nu} \partial^2 - (1 - \xi^{-1}) \partial^\mu \partial^\nu) \Delta_{\mu\nu}(x - y) = i\delta^4(x - y)$$

$$\Delta^{\mu\nu}(x - y) = \int \frac{d^4k}{(2\pi)^4} (-i) \left( \frac{g^{\mu\nu}}{k^2 + i\epsilon} - \frac{(1 - \xi) k^\mu k^\nu / k^2}{k^2 + i\epsilon} \right) e^{ik \cdot (x - y)}$$

### 3. BRST symmetry.

We have

$$\delta_B \psi = -igc^a T^a \psi, \delta_B \bar{\psi} = \bar{\psi}(-igc^a T^a)$$

$$\delta_B G^{a,\mu} = (D^\mu)^{ab} c^b, \delta_B c^a = \frac{1}{2} g f^{abc} c^b c^c$$

$$\delta_B \bar{c}^a = B^a(x), \delta_B B^a = 0, (D^\mu)^{ab} = \partial^\mu \delta^{ab} + g f^{cab} G^{c,\mu}$$

$$\text{so } (T^a T^b = i f^{abc} T^c + T^b T^a = i f^{abc} T^c + \frac{1}{2} \delta^{ab} - T^a T^b = \frac{i}{2} f^{abc} T^c)$$

$$\delta_B(\delta_B \psi) = -ig(\delta_B c^a) T^a \psi + g^2 c^a T^a c^b T^b \psi = -\frac{ig^2}{2} f^{abc} c^b c^c T^a \psi + g^2 c^a T^a c^b T^b \psi = 0$$

$$\delta_B^2 c^a = \frac{1}{2} g f^{abc} (\frac{1}{2} g f^{bde} c^d c^e c^c - \frac{1}{2} g f^{cde} c^b c^d c^e) = \frac{g^2}{4} (f^{eac} f^{cbd} c^b c^d c^e - f^{abc} f^{cde} c^b c^d c^e) = -\frac{g^2}{4} f^{adc} f^{ceb} c^b c^d c^e = 0$$

$$\delta_B^2 \bar{c}^a = 0$$

$$\delta_B^2 \bar{\psi} = \bar{\psi}(-igc^a T^a)(-igc^b T^b) - \bar{\psi}(-ig\frac{1}{2} g f^{abc} c^b c^c T^a) = \bar{\psi} \frac{ig^2}{2} f^{abc} c^b c^c T^a - \bar{\psi} g^2 c^a T^a c^b T^b = 0$$

$$\delta_B^2 G^{a,\mu} = \delta_B(\partial^\mu c^a + g f^{cab} G^{c,\mu} c^b) = \frac{1}{2} g f^{abc} \partial^\mu (c^b c^c) + g f^{cab} (\partial^\mu c^c + g f^{dce} G^{d,\mu} c^e) c^b + g f^{cab} G^{c,\mu} \frac{1}{2} g f^{bde} c^d c^e = \frac{1}{2} g f^{abc} \partial^\mu (c^b c^c)$$

$$- g f^{abc} c^b (\partial^\mu c^c) + g^2 f^{cab} f^{dce} G^{d,\mu} c^e c^b - \frac{g^2}{2} f^{cab} f^{bde} G^{c,\mu} c^d c^e = g^2 f^{cab} f^{dce} G^{d,\mu} c^e c^b + \frac{g^2}{2} f^{cab} f^{bde} G^{c,\mu} c^d c^e$$

$$= \frac{g^2}{2} (f^{aeb} f^{bdc} - f^{adb} f^{bec} + f^{cab} f^{bde}) G^{c,\mu} c^d c^e = 0$$