

## One Loop Gauge Link Self Energy

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## 1 One Loop

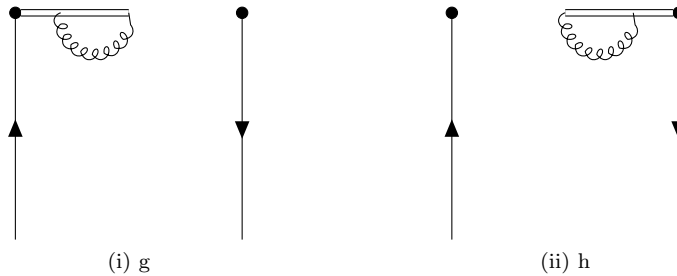



Figure 1: Diagrams of quasi PDF in Feynman gauge.

The definition of the gauge link self energy diagram (diagram g) is

$$\frac{1}{2} \int \frac{dz}{2\pi} e^{ixPz} \langle P, S | \bar{\psi}(z) \gamma^z \frac{\mathcal{P}[-ig_s \int_0^\infty dz' A^{a,z}(z') t^a]}{2} [-ig_s \int_0^\infty dz'' A^{a,z}(z'') t^a] \psi(0) | P, S \rangle \quad (1)$$

Applying Feynman rule straightaway gives (the overall  $1/2$  factor has been counted in)



$$\Gamma_g(l) = P \uparrow \downarrow P = -g_s^2 C_F \delta(1-x) \int \frac{d^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{i\epsilon} \frac{i}{-l^2 + i\epsilon} \quad (2)$$

### 1.1 Direct Contraction

### 1.1.1 Left

$$\begin{aligned}
& \frac{1}{2!} \mathcal{P} \left[ \int_0^\infty dz' A^{a,z}(z') \int_0^\infty dz'' A^{a,z}(z'') \right] \\
&= \int_0^\infty dz' A^{a,z}(z') \int_0^\infty dz'' A^{a,z}(z'') \theta(z' - z'') \\
&= \int dz' A^{a,z}(z') \int dz'' A^{a,z}(z'') \theta(z' - z'') \theta(z'') \\
&= \int dz' \overbrace{A^{a,z}(z') \int dz'' A^{a,z}(z'') \theta(z' - z'') \theta(z'')} \\
&= \int dz' \int dz'' \int \frac{d^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} e^{-il \cdot (z'' - z')} \theta(z' - z'') \theta(z'')
\end{aligned}$$

### 1.1.2 Right

$$\begin{aligned}
& \frac{1}{2!} \mathcal{P} \left[ \int_0^\infty dz'' A^{a,z}(z''+z) \int_0^\infty dz' A^{a,z}(z'+z) \right] \\
&= \int_0^\infty dz'' A^{a,z}(z''+z) \int_0^\infty dz' A^{a,z}(z'+z) \theta(z''-z') \\
&= \int dz'' A^{a,z}(z''+z) \int dz' A^{a,z}(z'+z) \theta(z''-z') \theta(z') \\
&= \int dz'' \overline{A^{a,z}(z''+z)} \int dz' A^{a,z}(z'+z) \theta(z''-z') \theta(z') \\
&= \int dz' \int dz'' \int \frac{d^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} e^{-il \cdot (z''-z')} \theta(z''-z') \theta(z')
\end{aligned}$$

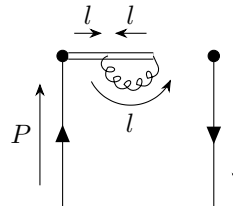
### 1.1.3 Summing together

$$\int dz' \int dz'' \int \frac{d^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} e^{-il \cdot (z''-z')} [\theta(z'-z'') \theta(z'') + \theta(z''-z') \theta(z')]$$

## 1.2 Adding two path order together

$$\begin{aligned}
& \mathcal{P} \left[ \int_0^\infty dz' A^{a,z}(z') \int_0^\infty dz'' A^{a,z}(z'') \right] \\
&= \int_0^\infty dz' A^{a,z}(z') \int_0^\infty dz'' A^{a,z}(z'') [\theta(z'-z'') + \theta(z''-z')] \\
&= \int_0^\infty dz' A^{a,z}(z') \int_0^\infty dz'' A^{a,z}(z'') \\
&= \int dz' A^{a,z}(z') \int dz'' A^{a,z}(z'') \int \frac{dw}{2\pi} \frac{ie^{-iwz'}}{w+i\epsilon} \int \frac{dh}{2\pi} \frac{ie^{-ihz''}}{h+i\epsilon} \\
&= \int dz' \overline{A^{a,z}(z')} \int dz'' A^{a,z}(z'') \int \frac{dw}{2\pi} \frac{ie^{-iwz'}}{w+i\epsilon} \int \frac{dh}{2\pi} \frac{ie^{-ihz''}}{h+i\epsilon} \\
&= \int dz' \int dz'' \int \frac{d^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} e^{-il \cdot (z''-z')} \int \frac{dw}{2\pi} \frac{ie^{-iwz'}}{w+i\epsilon} \int \frac{dh}{2\pi} \frac{ie^{-ihz''}}{h+i\epsilon} \\
&= \int \frac{d^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \int dz' \int dz'' \int \frac{dw}{2\pi} \frac{i}{w+i\epsilon} \int \frac{dh}{2\pi} \frac{i}{h+i\epsilon} e^{-i(w-l) \cdot z'} e^{-i(h+l) \cdot z''} \\
&= \int \frac{d^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{lz + i\epsilon} \frac{i}{l^z + i\epsilon} \frac{i}{-lz + i\epsilon}
\end{aligned}$$

The amplitude is



$$P = -\frac{g_s^2 C_F}{2} \delta(1-x) \int \frac{d^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{lz + i\epsilon} \frac{i}{l^z + i\epsilon} \frac{i}{-lz + i\epsilon} \quad (3)$$

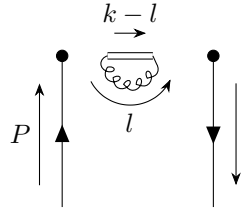
### 1.3 Adding $\Gamma_g(l)$ and $\Gamma_g(-l)$

$$\begin{aligned}
\Gamma_g(l) &= \frac{\Gamma_g(l) + \Gamma_g(-l)}{2} = -\frac{1}{2} \left[ g_s^2 C_F \delta(1-x) \int \frac{d^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{i\epsilon} \frac{i}{-lz + i\epsilon} + g_s^2 C_F \delta(1-x) \int \frac{d^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{i\epsilon} \frac{i}{lz + i\epsilon} \right] \\
&= -\frac{1}{2} g_s^2 C_F \delta(1-x) \int \frac{d^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{i\epsilon} \left[ \frac{i}{-lz + i\epsilon} + \frac{i}{lz + i\epsilon} \right] \\
&= -\frac{1}{2} g_s^2 C_F \delta(1-x) \int \frac{d^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{i\epsilon} \frac{2\epsilon}{lz^2 + \epsilon^2} \\
&= g_s^2 C_F \delta(1-x) \int \frac{d^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{lz + i\epsilon} \frac{i}{lz - i\epsilon}
\end{aligned}$$

There's an overall factor of 1/2 missing.

### 1.4 Taking derivatives

Add a small momentum to the gauge link line and consider an actual self energy diagram



$$P = -g_s^2 C_F \delta(1-x) \int \frac{d^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{kz - lz + i\epsilon} \quad (4)$$

take the derivative

$$-g_s^2 C_F \delta(1-x) \lim_{kz \rightarrow 0} \frac{\partial}{\partial kz} \left[ (i) \int \frac{d^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{kz - lz} \right] \quad (5)$$

$$= i g_s^2 C_F \delta(1-x) \int \frac{d^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{(lz)^2} \quad (6)$$

Adding  $i\epsilon$  by hand and we got

$$g_s^2 C_F \delta(1-x) \int \frac{d^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{lz + i\epsilon} \frac{i}{lz - i\epsilon} \quad (7)$$

## 2 Two Loop

In coordinate space

$$\langle P, S | \bar{\psi}(z) \gamma^z \mathcal{P} \frac{[-ig_s n_\mu \int_0^\infty dz_1 A^{a,\mu}(z_1) t^a] [-ig_s n_\nu \int_0^\infty dz_2 A^{b,\nu}(z_2) t^b] [-ig_s n_\rho \int_0^\infty dz_3 A^{c,\rho}(z_3) t^c] [-ig_s n_\sigma \int_0^\infty dz_4 A^{d,\sigma}(z_4) t^d]}{4!} \psi(0) | P, S \rangle \quad (8)$$

$$\begin{aligned}
&\frac{1}{4!} \mathcal{P} \left[ \int_0^\infty dz_1 A^{a,\mu}(z_1) \right] \left[ \int_0^\infty dz_2 A^{b,\nu}(z_2) \right] \left[ \int_0^\infty dz_3 A^{c,\rho}(z_3) \right] \left[ \int_0^\infty dz_4 A^{d,\sigma}(z_4) \right] \\
&= \left[ \int_0^\infty dz_1 A^{a,\mu}(z_1) \right] \left[ \int_0^\infty dz_2 A^{b,\nu}(z_2) \right] \left[ \int_0^\infty dz_3 A^{c,\rho}(z_3) \right] \left[ \int_0^\infty dz_4 A^{d,\sigma}(z_4) \right] \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3 - z_4)
\end{aligned}$$

The coefficient of the above expression is

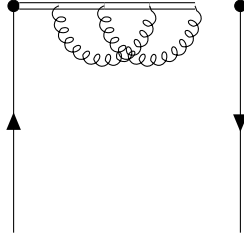
$$\langle P | \bar{\psi}(z) \gamma^z \psi(0) | P \rangle (-ig_s n^\mu) (-ig_s n^\nu) (-ig_s n^\rho) (-ig_s n^\sigma) t^a t^b t^c t^d \quad (9)$$

take a trace

$$e^{-iPz} \text{Tr}\{(\not{P} + m) \gamma^z\} \text{Tr}\{t^a t^b t^c t^d\} g_s^4 n^\mu n^\nu n^\rho n^\sigma \quad (10)$$

## 2.1 Diag. 50

The amplitude for



is related to the color ordering  $t^a t^b t^a t^b$ .

$$\begin{aligned}
& \int_0^\infty dz_1 \int_0^\infty dz_2 \int_0^\infty dz_3 \int_0^\infty dz_4 \overbrace{A^{a,\mu}(z_1) A^{b,\nu}(z_2) A^{c,\rho}(z_3) A^{d,\sigma}(z_4)}^{\text{gluon loop}} \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3 - z_4) \\
&= \int_0^\infty dz_1 \int_0^\infty dz_2 \int_0^\infty dz_3 \int_0^\infty dz_4 \int \frac{d^4 l_1}{(2\pi)^4} \frac{-ig^{\mu\rho} \delta^{ac}}{l_1^2 + i\epsilon} e^{-il_1 \cdot (z_3 - z_1)} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-ig^{\nu\sigma} \delta^{bd}}{l_2^2 + i\epsilon} e^{-il_2 \cdot (z_4 - z_2)} \\
&\quad \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3 - z_4) \\
&= \int dz_1 \int dz_2 \int dz_3 \int dz_4 \int \frac{d^4 l_1}{(2\pi)^4} \frac{-ig^{\mu\rho} \delta^{ac}}{l_1^2 + i\epsilon} e^{-il_1 \cdot (z_3 - z_1)} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-ig^{\nu\sigma} \delta^{bd}}{l_2^2 + i\epsilon} e^{-il_2 \cdot (z_4 - z_2)} \\
&\quad \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3 - z_4) \theta(z_4)
\end{aligned}$$

The exponent is (for simplicity we assume vectors  $z_i = (0, 0, 0, z_i)$  and  $k_i = (0, 0, 0, -k_i)$ )

$$\begin{aligned}
& -il_1 \cdot (z_3 - z_1) - il_2 \cdot (z_4 - z_2) - ik_1 \cdot (z_1 - z_2) - ik_2 \cdot (z_2 - z_3) - ik_3 \cdot (z_3 - z_4) - ik_4 \cdot z_4 \\
&= -iz_3 \cdot (l_1 + k_3 - k_2) - iz_1 \cdot (k_1 - l_1) - iz_4 \cdot (l_2 + k_4 - k_3) - iz_2 \cdot (k_2 - k_1 - l_2)
\end{aligned}$$

which gives 4 delta functions. The propagators involved are then

$$\begin{aligned}
& \int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-ig^{\mu\rho} \delta^{ac}}{l_1^2 + i\epsilon} \frac{-ig^{\nu\sigma} \delta^{bd}}{l_2^2 + i\epsilon} \frac{i}{-l_1^z + i\epsilon} \frac{i}{-l_1^z - l_2^z + i\epsilon} \frac{i}{-l_2^z + i\epsilon} \frac{i}{i\epsilon} \\
& \int dz_1 \int dz_2 \int dz_3 \int dz_4 \overbrace{A^{a,\mu}(z_1) A^{b,\nu}(z_2) A^{c,\rho}(z_3) A^{d,\sigma}(z_4)}^{\text{gluon loop}} \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3 - z_4) \theta(z_4) \\
&= \frac{1}{2} \int dz_1 \int dz_2 \int dz_3 \int dz_4 \overbrace{A^{a,\mu}(z_1) A^{b,\nu}(z_2) A^{c,\rho}(z_3) A^{d,\sigma}(z_4)}^{\text{gluon loop}} [\theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3 - z_4) \theta(z_4) \\
&\quad + \theta(z_3 - z_4) \theta(z_4 - z_1) \theta(z_1 - z_2) \theta(z_2)] \\
&= \int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-ig^{\mu\rho} \delta^{ac}}{l_1^2 + i\epsilon} \frac{-ig^{\nu\sigma} \delta^{bd}}{l_2^2 + i\epsilon} \frac{-(l_1^{z^2} + 3l_1^z l_2^z + l_2^{z^2} - \epsilon^2)}{(l_1^z - i\epsilon)(l_1^z + i\epsilon)(l_2^z + i\epsilon)(l_2^z - i\epsilon)(l_1^z + l_2^z - i\epsilon)(l_1^z + l_2^z + i\epsilon)}
\end{aligned}$$

Taking the derivative of the former expression, we can also arrive at a divergence-free form

$$\frac{i}{2} \lim_{p \rightarrow 0} \frac{\partial}{\partial p} \left[ \frac{i}{p + l_1^z} \frac{i}{p + l_1^z + l_2^z} \frac{i}{p + l_2^z} \right] = \frac{-(l_1^{z^2} + 3l_1^z l_2^z + l_2^{z^2})}{l_1^{z^2} l_2^{z^2} (l_1^z + l_2^z)^2}$$

which is equivalent to above expression.

Adding  $\epsilon$  in the definition so that the definition is  $\mathcal{P} e^{-ig_s \int_0^\infty e^{-z\epsilon} n \cdot A^a(z) t^a}$ , the expression becomes

$$\begin{aligned}
& \int dz_1 \int dz_2 \int dz_3 \int dz_4 \overbrace{A^{a,\mu}(z_1) A^{b,\nu}(z_2) A^{c,\rho}(z_3) A^{d,\sigma}(z_4)}^{\text{gluon loop}} e^{-(z_1 + z_2 + z_3 + z_4)\epsilon} \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3 - z_4) \theta(z_4) \\
&= \frac{1}{2} \int dz_1 \int dz_2 \int dz_3 \int dz_4 \overbrace{A^{a,\mu}(z_1) A^{b,\nu}(z_2) A^{c,\rho}(z_3) A^{d,\sigma}(z_4)}^{\text{gluon loop}} e^{-(z_1 + z_2 + z_3 + z_4)\epsilon} [\theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3 - z_4) \theta(z_4) \\
&\quad + \theta(z_3 - z_4) \theta(z_4 - z_1) \theta(z_1 - z_2) \theta(z_2)]
\end{aligned}$$

$$\begin{aligned}
& + \theta(z_3 - z_4)\theta(z_4 - z_1)\theta(z_1 - z_2)\theta(z_2)] \\
& = \int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-ig^{\mu\rho}\delta^{ac}}{l_1^2 + i\epsilon} \frac{-ig^{\nu\sigma}\delta^{bd}}{l_2^2 + i\epsilon} \frac{-3l_1^{z^2} - 6l_1^z l_2^z - l_2^{z^2}}{4(l_1^z - i\epsilon)(l_1^z + i\epsilon)(3\epsilon - il_2^z)(3\epsilon + il_2^z)(l_1^z + l_2^z - 2i\epsilon)(l_1^z + l_2^z + 2i\epsilon)}
\end{aligned}$$

The full amplitude in coordinate space is

$$4P^z e^{-iP^z z} g_s^4 \text{Tr}\{t^a t^b t^c t^d\} n^\mu n^\nu n^\rho n^\sigma \quad (11)$$

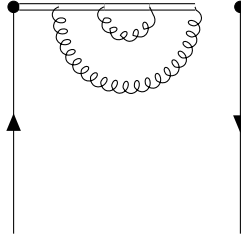
$$\begin{aligned}
& \int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-ig^{\mu\rho}\delta^{ac}}{l_1^2 + i\epsilon} \frac{-ig^{\nu\sigma}\delta^{bd}}{l_2^2 + i\epsilon} \frac{-3l_1^{z^2} - 6l_1^z l_2^z - l_2^{z^2}}{4(l_1^z - i\epsilon)(l_1^z + i\epsilon)(3\epsilon - il_2^z)(3\epsilon + il_2^z)(l_1^z + l_2^z - 2i\epsilon)(l_1^z + l_2^z + 2i\epsilon)} \\
& = 4P^z e^{-iP^z z} g_s^4 \text{Tr}\{t^a t^b t^a t^b\} \quad (12)
\end{aligned}$$

$$\begin{aligned}
& \int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-in^2}{l_1^2 + i\epsilon} \frac{-in^2}{l_2^2 + i\epsilon} \frac{-3l_1^{z^2} - 6l_1^z l_2^z - l_2^{z^2}}{4(l_1^z - i\epsilon)(l_1^z + i\epsilon)(3\epsilon - il_2^z)(3\epsilon + il_2^z)(l_1^z + l_2^z - 2i\epsilon)(l_1^z + l_2^z + 2i\epsilon)} \\
& = 4P^z e^{-iP^z z} g_s^4 \text{Tr}\{t^a t^b t^a t^b\} \quad (13)
\end{aligned}$$

$$\int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \frac{i}{l_1^2 + i\epsilon} \frac{i}{l_2^2 + i\epsilon} \frac{-3l_1^{z^2} - 6l_1^z l_2^z - l_2^{z^2}}{4(l_1^z - i\epsilon)(l_1^z + i\epsilon)(l_2^z + 3i\epsilon)(l_2^z - 3i\epsilon)(l_1^z + l_2^z - 2i\epsilon)(l_1^z + l_2^z + 2i\epsilon)}$$

## 2.2 Diag. 37

The amplitude for



is related to the color ordering  $t^a t^b t^b t^a$ .

$$\begin{aligned}
& \int_0^\infty dz_1 \int_0^\infty dz_2 \int_0^\infty dz_3 \int_0^\infty dz_4 \overline{A^{a,\mu}(z_1) A^{b,\nu}(z_2) A^{c,\rho}(z_3) A^{d,\sigma}(z_4)} \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3 - z_4) \\
& = \int dz_1 \int dz_2 \int dz_3 \int dz_4 \int \frac{d^4 l_1}{(2\pi)^4} \frac{-ig^{\mu\sigma}\delta^{ad}}{l_1^2 + i\epsilon} e^{-il_1 \cdot (z_3 - z_1)} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-ig^{\nu\rho}\delta^{bc}}{l_2^2 + i\epsilon} e^{-il_2 \cdot (z_4 - z_2)} \\
& \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3 - z_4) \theta(z_4)
\end{aligned}$$

The propagators involved are

$$\begin{aligned}
& \int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-ig^{\mu\sigma}\delta^{ad}}{l_1^2 + i\epsilon} \frac{-ig^{\nu\rho}\delta^{bc}}{l_2^2 + i\epsilon} \frac{i}{-l_1^z + i\epsilon} \frac{i}{-l_1^z - l_2^z + i\epsilon} \frac{i}{-l_1^z + i\epsilon} \frac{i}{i\epsilon} \\
& \frac{1}{2} \int dz_1 \int dz_2 \int dz_3 \int dz_4 \overline{A^{a,\mu}(z_1) A^{b,\nu}(z_2) A^{c,\rho}(z_3) A^{d,\sigma}(z_4)} [\theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3 - z_4) \theta(z_4) \\
& \quad + \theta(z_4 - z_3) \theta(z_3 - z_2) \theta(z_2 - z_1) \theta(z_1)] \\
& = \int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-ig^{\mu\sigma}\delta^{ad}}{l_1^2 + i\epsilon} \frac{-ig^{\nu\rho}\delta^{bc}}{l_2^2 + i\epsilon} \frac{-3l_1^{z^2} - 2l_1^z l_2^z + \epsilon^2}{(l_1^{z^2} + \epsilon^2)^2 (l_1^z + 2l_1^z l_2^z + l_2^{z^2} + \epsilon^2)}
\end{aligned}$$

Taking the derivative of the former expression, we can also arrive at a divergence-free form

$$\frac{i}{2} \lim_{p \rightarrow 0} \frac{\partial}{\partial p} \left[ \frac{i}{p + l_1^z} \frac{i}{p + l_1^z + l_2^z} \frac{i}{p + l_1^z} \right] = \frac{-(3l_1^z + 2l_2^z)}{l_1^{z^3} (l_1^z + l_2^z)^2}$$

which is equivalent to above expression.

$$\frac{-3l_1^{z^2} - 2l_1^z l_2^z + 3\epsilon^2}{2(l_1^z - i\epsilon)(l_1^z + i\epsilon)(l_1^z - 3i\epsilon)(l_1^z + 3i\epsilon)(l_1^z + l_2^z - 2i\epsilon)(l_1^z + l_2^z + 2i\epsilon)}$$

The full amplitude in coordinate space is

$$4P^z e^{-iP^z z} g_s^4 \text{Tr}\{t^a t^b t^c t^d\} n^\mu n^\nu n^\rho n^\sigma \quad (14)$$

$$\int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-ig^{\mu\sigma}\delta^{ad}}{l_1^2 + i\epsilon} \frac{-ig^{\nu\rho}\delta^{bc}}{l_2^2 + i\epsilon} \frac{-3l_1^{z^2} - 2l_1^z l_2^z + 3\epsilon^2}{2(l_1^z - i\epsilon)(l_1^z + i\epsilon)(l_1^z - 3i\epsilon)(l_1^z + 3i\epsilon)(l_1^z + l_2^z - 2i\epsilon)(l_1^z + l_2^z + 2i\epsilon)}$$

$$= 4P^z e^{-iP^z z} g_s^4 \text{Tr}\{t^a t^b t^b t^a\} \quad (15)$$

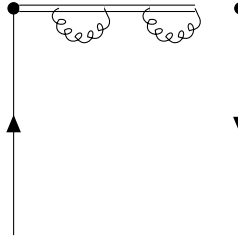
$$\int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-in^2}{l_1^2 + i\epsilon} \frac{-in^2}{l_2^2 + i\epsilon} \frac{-3l_1^{z^2} - 2l_1^z l_2^z + 3\epsilon^2}{2(l_1^z - i\epsilon)(l_1^z + i\epsilon)(l_1^z - 3i\epsilon)(l_1^z + 3i\epsilon)(l_1^z + l_2^z - 2i\epsilon)(l_1^z + l_2^z + 2i\epsilon)}$$

$$= 4P^z e^{-iP^z z} g_s^4 \text{Tr}\{t^a t^b t^b t^a\} \quad (16)$$

$$\int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \frac{i}{l_1^2 + i\epsilon} \frac{i}{l_2^2 + i\epsilon} \frac{-3l_1^{z^2} - 2l_1^z l_2^z}{2(l_1^z - i\epsilon)(l_1^z + i\epsilon)(l_1^z - 3i\epsilon)(l_1^z + 3i\epsilon)(l_1^z + l_2^z - 2i\epsilon)(l_1^z + l_2^z + 2i\epsilon)}$$

### 2.3 Diag. 43

The amplitude for



is related to the color ordering  $t^a t^a t^b t^b$ .

$$\int_0^\infty dz_1 \int_0^\infty dz_2 \int_0^\infty dz_3 \int_0^\infty dz_4 \overline{A^{a,\mu}(z_1)} A^{b,\nu}(z_2) \overline{A^{c,\rho}(z_3)} A^{d,\sigma}(z_4) \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3 - z_4)$$

$$= \int dz_1 \int dz_2 \int dz_3 \int dz_4 \int \frac{d^4 l_1}{(2\pi)^4} \frac{-ig^{\mu\sigma}\delta^{ad}}{l_1^2 + i\epsilon} e^{-il_1 \cdot (z_2 - z_1)} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-ig^{\nu\rho}\delta^{bc}}{l_2^2 + i\epsilon} e^{-il_2 \cdot (z_4 - z_3)}$$

$$\theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3 - z_4) \theta(z_4)$$

The propagators involved are

$$\int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-ig^{\mu\sigma}\delta^{ad}}{l_1^2 + i\epsilon} \frac{-ig^{\nu\rho}\delta^{bc}}{l_2^2 + i\epsilon} \frac{i}{-l_1^z + i\epsilon} \frac{i}{i\epsilon} \frac{i}{-l_2^z + i\epsilon} \frac{i}{i\epsilon}$$

$$\frac{1}{4} \int dz_1 \int dz_2 \int dz_3 \int dz_4 \overline{A^{a,\mu}(z_1)} A^{b,\nu}(z_2) \overline{A^{c,\rho}(z_3)} A^{d,\sigma}(z_4) [\theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3 - z_4) \theta(z_4)$$

$$+ \theta(z_2 - z_1) \theta(z_1 - z_3) \theta(z_3 - z_4) \theta(z_4) + \theta(z_1 - z_2) \theta(z_2 - z_4) \theta(z_4 - z_3) \theta(z_3) + \theta(z_2 - z_1) \theta(z_1 - z_4) \theta(z_4 - z_3) \theta(z_3)]$$

$$= \int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-ig^{\mu\sigma}\delta^{ab}}{l_1^2 + i\epsilon} \frac{-ig^{\nu\rho}\delta^{cd}}{l_2^2 + i\epsilon} \frac{1}{(l_1^{z^2} + \epsilon^2)(l_2^{z^2} + \epsilon^2)}$$

$$\int_0^\infty dz_1 \int_0^\infty dz_2 \int_0^\infty dz_3 \int_0^\infty dz_4 \overline{A^{a,\mu}(z_1)} A^{b,\nu}(z_2) \overline{A^{c,\rho}(z_3)} A^{d,\sigma}(z_4) e^{-(z_1 + z_2 + z_3 + z_4)\epsilon} \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3 - z_4)$$

$$= \int dz_1 \int dz_2 \int dz_3 \int dz_4 \int \frac{d^4 l_1}{(2\pi)^4} \frac{-ig^{\mu\sigma}\delta^{ab}}{l_1^2 + i\epsilon} e^{-il_1 \cdot (z_2 - z_1)} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-ig^{\nu\rho}\delta^{cd}}{l_2^2 + i\epsilon} e^{-il_2 \cdot (z_4 - z_3)} e^{-(z_1 + z_2 + z_3 + z_4)\epsilon}$$

$$\theta(z_1 - z_2)\theta(z_2 - z_3)\theta(z_3 - z_4)\theta(z_4)$$

and the eikonal part is

$$\frac{3}{8(l_1^z - i\epsilon)(l_1^z + i\epsilon)(3\epsilon - il_2^z)(3\epsilon + il_2^z)}$$

The full amplitude in coordinate space is

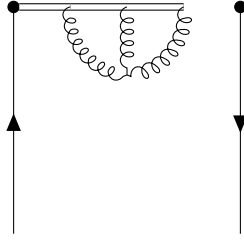
$$4P^z e^{-iP^z z} g_s^4 \text{Tr}\{t^a t^b t^c t^d\} n^\mu n^\nu n^\rho n^\sigma \int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-ig^{\mu\sigma}\delta^{ab}}{l_1^2 + i\epsilon} \frac{-ig^{\nu\rho}\delta^{cd}}{l_2^2 + i\epsilon} \frac{3}{8(l_1^z - i\epsilon)(l_1^z + i\epsilon)(3\epsilon - il_2^z)(3\epsilon + il_2^z)} \quad (17)$$

$$= 4P^z e^{-iP^z z} g_s^4 \text{Tr}\{t^a t^a t^b t^b\} \int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-in^2}{l_1^2 + i\epsilon} \frac{-in^2}{l_2^2 + i\epsilon} \frac{3}{8(l_1^z - i\epsilon)(l_1^z + i\epsilon)(3\epsilon - il_2^z)(3\epsilon + il_2^z)} \quad (18)$$

$$= 4P^z e^{-iP^z z} g_s^4 \text{Tr}\{t^a t^a t^b t^b\} \int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \frac{i}{l_1^2 + i\epsilon} \frac{i}{l_2^2 + i\epsilon} \frac{3}{8(l_1^z - i\epsilon)(l_1^z + i\epsilon)(l_2^z - 3i\epsilon)(l_2^z + 3i\epsilon)} \quad (19)$$

## 2.4 Diag. 47

The amplitude for



is related to

$$\langle P, S | \bar{\psi}(z) \gamma^z V_3 \mathcal{P} \frac{[-ig_s n_\mu \int_0^\infty dz_1 A^{a,\mu}(z_1) t^a] [-ig_s n_\nu \int_0^\infty dz_2 A^{b,\nu}(z_2) t^b] [-ig_s n_\rho \int_0^\infty dz_3 A^{c,\rho}(z_3) t^c]}{3!} \psi(0) | P, S \rangle \quad (20)$$

where

$$V_3 = -\frac{g_s}{2} f^{def} \int d^4 t \left( \partial^\alpha A_d^\beta - \partial^\beta A_d^\alpha \right) A_\alpha^e A_\beta^f$$

The gluon related contraction (one out of three, others can be obtained by exchanging  $d, e, f$ .)

$$\begin{aligned} & \int dz_1 \int dz_2 \int dz_3 \int d^4 t \left( \partial^\alpha A^{d,\beta} - \partial^\beta A^{d,\alpha} \right) A_\alpha^e A_\beta^f A^{a,\mu}(z_1) A^{b,\nu}(z_2) A^{c,\rho}(z_3) \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3) \\ &= \int dz_1 \int dz_2 \int dz_3 \int d^4 t \int \frac{d^4 l_1}{(2\pi)^4} \frac{-i\delta^{dc}}{l_1^2 + i\epsilon} \left[ g^{\beta\rho} \partial^\alpha e^{-il_1 \cdot (z_3 - t)} - g^{\alpha\rho} \partial^\beta e^{-il_1 \cdot (z_3 - t)} \right] \int \frac{d^4 l_2}{(2\pi)^4} \frac{-ig_\alpha^\mu \delta^{ea}}{l_2^2 + i\epsilon} e^{-il_2 \cdot (z_1 - t)} \\ & \int \frac{d^4 l_3}{(2\pi)^4} \frac{-ig_\beta^\nu \delta^{fb}}{l_3^2 + i\epsilon} e^{-il_3 \cdot (z_2 - t)} \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3) \\ &= \int dz_1 \int dz_2 \int dz_3 \int d^4 t \int \frac{d^4 l_1}{(2\pi)^4} \frac{-i\delta^{dc}}{l_1^2 + i\epsilon} \left[ g^{\beta\rho} il_1^\alpha - g^{\alpha\rho} il_1^\beta \right] e^{-il_1 \cdot (z_3 - t)} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-ig_\alpha^\mu \delta^{ea}}{l_2^2 + i\epsilon} e^{-il_2 \cdot (z_1 - t)} \\ & \int \frac{d^4 l_3}{(2\pi)^4} \frac{-ig_\beta^\nu \delta^{fb}}{l_3^2 + i\epsilon} e^{-il_3 \cdot (z_2 - t)} \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3) \\ &= \int dz_1 \int dz_2 \int dz_3 \int d^4 t \int \frac{d^4 l_1}{(2\pi)^4} \frac{-i\delta^{dc}}{l_1^2 + i\epsilon} \left[ g^{\nu\rho} il_1^\mu - g^{\mu\rho} il_1^\nu \right] e^{-il_1 \cdot (z_3 - t)} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-i\delta^{ea}}{l_2^2 + i\epsilon} e^{-il_2 \cdot (z_1 - t)} \\ & \int \frac{d^4 l_3}{(2\pi)^4} \frac{-i\delta^{fb}}{l_3^2 + i\epsilon} e^{-il_3 \cdot (z_2 - t)} \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3) \end{aligned}$$

Multiplied by  $n$

$$n_\mu n_\nu n_\rho [g^{\nu\rho} i l_1^\mu - g^{\mu\rho} i l_1^\nu] = i n^2 [l_1^z - l_1^z] = 0$$

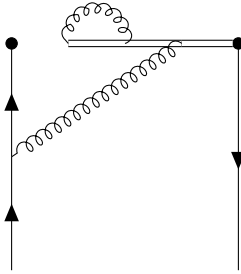
Contracting with different fields in the parenthesis

$$\begin{aligned} & \partial^\alpha \overbrace{A^{d,\beta} A_\alpha^e A_\beta^f A^{a,\mu}(z_1) A^{b,\nu}(z_2) A^{c,\rho}(z_3)} - \partial^\beta \overbrace{A^{d,\alpha} A_\alpha^e A_\beta^f A^{a,\mu}(z_1) A^{b,\nu}(z_2) A^{c,\rho}(z_3)} \\ & \int dz_1 \int dz_2 \int dz_3 \int d^4 t \int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \int \frac{d^4 l_3}{(2\pi)^4} \frac{-i\delta^{dc}}{l_3^2 + i\epsilon} \frac{-ig_\alpha^\mu \delta^{ea}}{l_1^2 + i\epsilon} e^{-il_1 \cdot (z_1 - t)} \frac{-ig_\beta^\nu \delta^{fb}}{l_2^2 + i\epsilon} \\ & \left[ e^{-il_2 \cdot (z_2 - t)} g^{\beta\rho} \partial^\alpha e^{-il_3 \cdot (z_3 - t)} - e^{-il_3 \cdot (z_3 - t)} g^{\alpha\rho} \partial^\beta e^{-il_2 \cdot (z_2 - t)} \right] \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3) \\ & = \int dz_1 \int dz_2 \int dz_3 \int d^4 t \int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \int \frac{d^4 l_3}{(2\pi)^4} \frac{-i\delta^{dc}}{l_3^2 + i\epsilon} \frac{-i\delta^{ea}}{l_1^2 + i\epsilon} \frac{-i\delta^{fb}}{l_2^2 + i\epsilon} e^{-il_1 \cdot (z_1 - t)} e^{-il_2 \cdot (z_2 - t)} e^{-il_3 \cdot (z_3 - t)} \\ & [g^{\nu\rho} i l_3^\mu - g^{\mu\rho} i l_2^\nu] \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3) \end{aligned}$$

Make  $z_1 \rightarrow -z_1, t \rightarrow -t, l_2 \rightarrow -l_2$ ,

$$\begin{aligned} & \int dz_1 \int dz_2 \int dz_3 \int d^4 t \int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \int \frac{d^4 l_3}{(2\pi)^4} \frac{-i\delta^{dc}}{l_3^2 + i\epsilon} \frac{-i\delta^{ea}}{l_1^2 + i\epsilon} \frac{-i\delta^{fb}}{l_2^2 + i\epsilon} e^{-il_1 \cdot (z_1 - t)} e^{-il_2 \cdot (z_2 - t)} e^{-il_3 \cdot (z_3 - t)} \\ & [g^{\nu\rho} i l_3^\mu - g^{\mu\rho} i l_2^\nu] \theta(-z_1 - z_2) \theta(z_2 - z_3) \theta(z_3) \\ & = \int \frac{d^4 l_2}{(2\pi)^4} \int \frac{d^4 l_3}{(2\pi)^4} \frac{-i\delta^{dc}}{l_3^2 + i\epsilon} \frac{-i\delta^{ea}}{(l_3 + l_2)^2 + i\epsilon} \frac{-i\delta^{fb}}{l_2^2 + i\epsilon} \frac{i}{-l_3^z + i\epsilon} \frac{i}{-2l_3^z + i\epsilon} \frac{i}{-2l_3^z - l_2^z + i\epsilon} [g^{\nu\rho} i l_3^\mu - g^{\mu\rho} i l_2^\nu] \end{aligned}$$

$$\begin{aligned} & \partial^\alpha \overbrace{A^{d,\beta} A_\alpha^e A_\beta^f A^{a,\mu}(z_1) A^{b,\nu}(z_2) A^{c,\rho}(z_3)} - \partial^\beta \overbrace{A^{d,\alpha} A_\alpha^e A_\beta^f A^{a,\mu}(z_1) A^{b,\nu}(z_2) A^{c,\rho}(z_3)} \\ & \int dz_1 \int dz_2 \int dz_3 \int d^4 t \int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \int \frac{d^4 l_3}{(2\pi)^4} \frac{-i\delta^{dc}}{l_3^2 + i\epsilon} \frac{-ig_\alpha^\mu \delta^{ea}}{l_1^2 + i\epsilon} e^{-il_1 \cdot (z_1 - t)} \frac{-ig_\beta^\nu \delta^{fb}}{l_2^2 + i\epsilon} \\ & \left[ e^{-il_2 \cdot (z_2 - t)} g^{\beta\rho} \partial^\alpha e^{-il_3 \cdot (z_3 - t)} - e^{-il_3 \cdot (z_3 - t)} g^{\alpha\rho} \partial^\beta e^{-il_2 \cdot (z_2 - t)} \right] \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3) \end{aligned}$$



is related to

$$\langle P, S | \bar{\psi}(z) \gamma^z Q_3 \mathcal{P} \frac{[-ig_s n_\mu \int_0^\infty dz_1 A^{a,\mu}(z_1) t^a] [-ig_s n_\nu \int_0^\infty dz_2 A^{b,\nu}(z_2) t^b] [-ig_s n_\rho \int_0^\infty dz_3 A^{c,\rho}(z_3) t^c]}{3!} \psi(0) | P, S \rangle \quad (21)$$

where

$$\begin{aligned} Q_3 &= (-ig_s \gamma_\sigma) \int d^4 t \bar{\psi} \psi A^{d,\sigma} \\ & \int dz_1 \int dz_2 \int dz_3 \overbrace{A^{d,\sigma}(t) A^{a,\mu}(z_1) A^{b,\nu}(z_2) A^{c,\rho}(z_3)} \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3) \\ & = \int dz_1 \int dz_2 \int dz_3 \int \frac{d^4 l_1}{(2\pi)^4} \int \frac{d^4 l_2}{(2\pi)^4} \frac{-ig^{\mu\nu} \delta^{ab}}{l_1^2 + i\epsilon} \frac{-ig^{\sigma\rho} \delta^{dc}}{l_2^2 + i\epsilon} e^{-il_1 \cdot (z_2 - z_1)} e^{-il_2 \cdot (z_3 - t)} \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3) \end{aligned}$$



The eikonal propagators are

$$\frac{i}{-l_1^z + i\epsilon} \frac{i}{i\epsilon} \frac{i}{l_2^z + i\epsilon}$$

Flip  $z_1$  and  $z_2$

$$-\frac{i}{-l_1^z + i\epsilon} \frac{i}{l_1^z + i\epsilon} \frac{i}{l_2^z + i\epsilon} \quad (22)$$

This covers for all diagrams involved one gauge field contracting with the would-be-divergent three gauge link fields, and the extra eikonal line behaves exactly as in one loop level.