## Local Operator Divergence

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## 1 NRQED local matrix elements

Since what we what is an operator equation independent of states, we can choose  $p^{\mu} = (0, \mathbf{p})$  as our outstate electron momentum, and put the nucleus on-shell. This way, we can eliminate possible IR divergence and make the integral much easier (especially at NNLO where the first integral will give a  $\sinh^{-1}\left(\frac{1}{\sqrt{\frac{k_1^2}{2Em}-1}}\right)$  if no constrain condition put at external momentum).

At NLO:

$$\langle 0|\psi_{e}(0)N(0)(-ie)\int d^{4}y \bar{\psi}_{e}\psi_{e}A^{0}(-ie)\int d^{4}z \bar{N}NA^{0}|eN\rangle = P_{N} - k$$

$$= e^{2}u_{N}(v_{N})\left[\int \frac{d^{3}k}{(2\pi)^{3}} \frac{1}{(\mathbf{k} - \mathbf{p})^{2}(E - \frac{\mathbf{k}^{2}}{2m})} (1 - \frac{\mathbf{k}^{4}}{8m^{3}(E - \frac{\mathbf{k}^{2}}{2m})})\right] \frac{i}{E - \frac{\mathbf{p}^{2}}{2m}}$$

$$= -e^{2}u_{N}(v_{N})\left[\frac{\pi}{v} + \mathcal{O}(v^{2})\right] \frac{i}{E - \frac{\mathbf{p}^{2}}{2m}}$$

## Appendices

NRQED matrix element at NLO

$$\langle 0 | \psi_{e}(0) N(0)(-ie) \int d^{4}y \bar{\psi}_{e} \psi_{e} A^{0}(-ie) \int d^{4}z \bar{N} N A^{0} | e N \rangle = \begin{cases} p+k \\ p \\ k_{1} \end{cases} P_{N} - k$$

$$= ie^{2}u_{N}(v_{N}) \left[ \int [dk] \frac{1}{\mathbf{k}^{2}(-k^{0} + i\epsilon)(^{0} + k^{0} - m - \frac{(\mathbf{p} + \mathbf{k})^{2}}{2m} + \frac{(\mathbf{p} + \mathbf{k})^{4}}{8m^{3}} + i\epsilon)} \right] \frac{i}{E - \frac{\mathbf{p}^{2}}{2m}}$$

$$= e^{2}u_{N}(v_{N}) \left[ \int \frac{d^{3}k}{(2\pi)^{3}} \frac{1}{(\mathbf{k} - \mathbf{p})^{2}(E - \frac{\mathbf{k}^{2}}{2m} + \frac{\mathbf{k}^{4}}{8m^{3}})} \right] \frac{i}{E - \frac{\mathbf{p}^{2}}{2m}}$$

$$= e^{2}u_{N}(v_{N}) \left[ \int \frac{d^{3}k}{(2\pi)^{3}} \frac{1}{(\mathbf{k} - \mathbf{p})^{2}(E - \frac{\mathbf{k}^{2}}{2m})} (1 - \frac{\mathbf{k}^{4}}{8m^{3}(E - \frac{\mathbf{k}^{2}}{2m})}) \right] \frac{i}{E - \frac{\mathbf{p}^{2}}{2m}}$$

$$= -2me^{2}u_{N}(v_{N}) \frac{i}{E - \frac{\mathbf{p}^{2}}{2m}} [\frac{\pi}{2p} + \mathcal{O}(p^{2})] = -e^{2}u_{N}(v_{N}) \frac{i}{E - \frac{\mathbf{p}^{2}}{2m}} [\frac{\pi}{v} + \mathcal{O}(v^{2})]$$

At NNLO (where we're only interested in divergent parts)

$$p + k_1 + k_2$$

$$- p + k_1$$

$$p + k_1 + k_2$$

$$P_N - k_1$$

$$P_N - k_1$$

$$P_N = m_N v_N$$

$$= e^4 \left[ \int [dk_1] [dk_2] \frac{1}{|\mathbf{k_1}|^2} \frac{1}{|\mathbf{k_2}|^2} \frac{1}{-k_1^0 - k_2^0 + i\epsilon} \frac{1}{-k_1^0 + i\epsilon} \frac{1}{p^0 + k_1^0 - m - \frac{(\mathbf{p} + \mathbf{k_1})^2}{2m} + i\epsilon} \frac{1}{p^0 + k_1^0 + k_2^0 - m - \frac{(\mathbf{p} + \mathbf{k_1} + \mathbf{k_2})^2}{2m} + i\epsilon} \right] \frac{i}{E - \frac{\mathbf{p}^2}{2m}} u_N(v_N)$$

do the shift as above

$$= -e^{4} \left[ \int \frac{\mathrm{d}^{3}\mathbf{k_{1}}}{(2\pi)^{3}} \frac{\mathrm{d}^{3}\mathbf{k_{2}}}{(2\pi)^{3}} \frac{1}{|\mathbf{k_{1}} - \mathbf{p}|^{2}} \frac{1}{|\mathbf{k_{2}} - \mathbf{k_{1}}|^{2}} \frac{1}{E - \frac{|\mathbf{k_{1}}|^{2}}{2m} + 2i\epsilon} \frac{1}{E - \frac{|\mathbf{k_{2}}|^{2}}{2m} + 2i\epsilon} \right] \frac{i}{E - \frac{\mathbf{p}^{2}}{2m}} u_{N}(v_{N})$$

drop p

$$= -e^{4} \left[ \int \frac{\mathrm{d}^{3} \mathbf{k_{1}}}{(2\pi)^{3}} \frac{\mathrm{d}^{3} \mathbf{k_{2}}}{(2\pi)^{3}} \frac{1}{|\mathbf{k_{1}}|^{2}} \frac{1}{|\mathbf{k_{2}} - \mathbf{k_{1}}|^{2}} \frac{1}{-\frac{|\mathbf{k_{1}}|^{2}}{2m} + 2i\epsilon} \frac{1}{-\frac{|\mathbf{k_{2}}|^{2}}{2m} + 2i\epsilon} \right] \frac{i}{E - \frac{\mathbf{p}^{2}}{2m}} u_{N}(v_{N})$$

if we add higher reltivistic correction

$$= -e^{4} \left[ \int \frac{d^{3}\mathbf{k_{1}}}{(2\pi)^{3}} \frac{d^{3}\mathbf{k_{2}}}{(2\pi)^{3}} \frac{1}{|\mathbf{k_{1}}|^{2}} \frac{1}{|\mathbf{k_{2}} - \mathbf{k_{1}}|^{2}} \frac{2m}{|\mathbf{k_{1}}|^{2} - \frac{|\mathbf{k_{1}}|^{4}}{4m^{2}}} \frac{2m}{|\mathbf{k_{2}}|^{2} - \frac{|\mathbf{k_{2}}|^{4}}{4m^{2}}} \right] \frac{i}{E - \frac{\mathbf{p}^{2}}{2m}} u_{N}(v_{N})$$

$$= -4m^{2}e^{4} \left[ \int \frac{d^{3}\mathbf{k_{1}}}{(2\pi)^{3}} \frac{d^{3}\mathbf{k_{2}}}{(2\pi)^{3}} \frac{1}{|\mathbf{k_{1}}|^{2}} \frac{1}{|\mathbf{k_{2}} - \mathbf{k_{1}}|^{2}} \frac{1}{|\mathbf{k_{1}}|^{2}} (1 + \frac{|\mathbf{k_{1}}|^{2}}{4m^{2}}) \frac{1}{|\mathbf{k_{2}}|^{2}} (1 + \frac{|\mathbf{k_{2}}|^{2}}{4m^{2}}) \right] \frac{i}{E - \frac{\mathbf{p}^{2}}{2m}} u_{N}(v_{N})$$

The integral (we can verify that  $\mathbf{k}^0$  part is not UV divergent, and we don't care about  $\mathbf{k}^8$  term for now) for

$$p + k_{1} + k_{2} + k_{1} + k_{2} +$$

where 
$$\Delta_2 = (1 - x) \left( |\mathbf{k_1}|^2 x - 2Em \right)$$

$$\begin{split} &= \frac{1}{(4\pi)^2} \int_0^1 \mathrm{d}x \int \frac{\mathrm{d}^3\mathbf{k_1}}{(2\pi)^3} \frac{1}{|\mathbf{k_1} - \mathbf{p}|^2} \frac{|\mathbf{k_1}|^4 / 4m^2}{[|\mathbf{k_1}|^2 - 2mE]^2} \frac{1}{(|\mathbf{k_1}|^2 - 2mE/x)^{2-d/2}} \left(\frac{4\pi}{x(1-x)}\right)^{2-d/2} \Gamma(2-d/2) \\ &= \frac{1}{(4\pi)^2} \int_0^1 \mathrm{d}x \int_0^1 \mathrm{d}y \mathrm{d}z \mathrm{d}t \delta(y+z+t-1) \int \frac{\mathrm{d}^3\mathbf{k_1}}{(2\pi)^3} \frac{zt^{1-d/2} |\mathbf{k_1}|^4 / 4m^2}{[|\mathbf{k_1}|^2 + \Delta_1]^{5-d/2}} \frac{\Gamma(5-d/2)}{\Gamma(2-d/2)} \left(\frac{4\pi}{x(1-x)}\right)^{2-d/2} \Gamma(2-d/2) \end{split}$$

where 
$$\Delta_1 = y(1-y)\mathbf{p}^2 - 2mE(z+t/x)$$

$$\begin{split} &= \frac{1}{4m^2(4\pi)^2} \int_0^1 \mathrm{d}x \int_0^1 \mathrm{d}y \mathrm{d}z \mathrm{d}t \\ \delta(y+z+t-1)zt^{1-d/2} \frac{d(d+2)}{4} \frac{\Gamma\left(3-d\right)}{(4\pi)^{3-d/2}} \left(\frac{4\pi}{\Delta_1}\right)^{3-d} \left(\frac{4\pi}{x(1-x)}\right)^{2-d/2} \\ &= -\frac{1}{128\pi^2(d-3)m^2} + \text{finite terms} \end{split}$$

$$p + k_{1} + k_{2} + k_{1} + k_{2} +$$

where  $\Delta_2 = x(1-x)|\mathbf{k_1}|^2 - 2mE(1-x)$ 

$$= \frac{1}{4m^2} \int_0^1 \mathrm{d}x \int_0^1 \mathrm{d}y \mathrm{d}z \mathrm{d}t \frac{t^{-d/2}}{\left[\left|\mathbf{k_1}\right|^2 + \Delta_1\right]^{3-d/2}} \frac{\Gamma(3-d/2)}{\Gamma(1-d/2)} \delta(y+z+t-1) \frac{\Gamma(1-d/2)}{8\pi} \left(\frac{4\pi}{x(1-x)}\right)^{1-d/2} \frac{d(d+2)x}{4\pi} \delta(y+z+t-1) \frac{(d+2)x}{4\pi} \delta(y+z+t-1)$$

where  $\Delta_1 = y(1-y)\mathbf{p}^2 - 2mEz - 2mE\frac{t}{x}$ 

$$\begin{split} &= \frac{1}{4m^2} \int_0^1 \mathrm{d}x \int_0^1 \mathrm{d}y \mathrm{d}z \mathrm{d}t \delta(y+z+t-1) \frac{1}{(4\pi)^{3-d/2}} \left(\frac{4\pi}{\Delta_1}\right)^{3-d} \frac{\Gamma\left(3-d\right)}{8\pi} \left(\frac{4\pi}{x(1-x)}\right)^{1-d/2} \frac{d(d+2)x}{4} t^{-d/2} \\ &= \frac{15}{8192\pi^2(d-3)m^2} + \text{finite terms} \end{split}$$

Check the contour integral

$$p + k_{1} + k_{2}$$

$$- p + k_{1}$$

$$p$$

$$P_{N} - k_{1}$$

$$P_{N} = m_{N}v_{N}$$

$$= e^{4} \left[ \int [dk_{1}][dk_{2}] \frac{1}{|\mathbf{k}_{1}|^{2}} \frac{1}{|\mathbf{k}_{2}|^{2}} \frac{1}{-k_{1}^{0} - k_{2}^{0} + i\epsilon} \frac{1}{-k_{1}^{0} + i\epsilon} \frac{(\mathbf{p} + \mathbf{k})^{4}/8m^{3}}{[p^{0} + k_{1}^{0} - m - \frac{(\mathbf{p} + \mathbf{k}_{1})^{2}}{2m} + i\epsilon]^{2}} \frac{1}{p^{0} + k_{1}^{0} + k_{2}^{0} - m - \frac{(\mathbf{p} + \mathbf{k}_{1} + \mathbf{k}_{2})^{2}}{2m} + i\epsilon} \right] \frac{i}{E - \frac{\mathbf{p}^{2}}{2m}} u_{N}(v_{N})$$

$$= -e^{4} \left[ \int \frac{d^{3}\mathbf{k}_{1}}{(2\pi)^{3}} \frac{d^{3}\mathbf{k}_{2}}{(2\pi)^{3}} \frac{1}{|\mathbf{k}_{1} - \mathbf{p}|^{2}} \frac{|\mathbf{k}_{1}|^{4}/8m^{3}}{|\mathbf{k}_{2} - \mathbf{k}_{1}|^{2}} \frac{1}{E - \frac{|\mathbf{k}_{1}|^{2}}{2m} + 2i\epsilon} \frac{i}{E - \frac{|\mathbf{k}_{2}|^{2}}{2m} + 2i\epsilon} \right] \frac{i}{E - \frac{\mathbf{p}^{2}}{2m}} u_{N}(v_{N})$$

which is the same as previous (by expansion) one.

Contact (Darwin) term

if the dispersion relation is up to  $\mathbf{k}^4$  then

$$=\frac{4e^4m}{3}$$