# One Loop Gauge Link Self Energy

Yingsheng Huang

January 19, 2020

# 1 One Loop

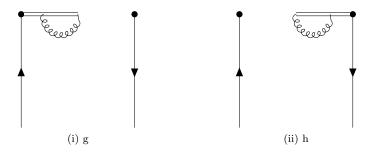


Figure 1: Diagrams of quasi PDF in Feynman gauge.

The definition of the gauge link self energy diagram (diagram g) is

$$\frac{1}{2} \int \frac{\mathrm{d}z}{2\pi} e^{ixP^z z} \langle P, S | \bar{\psi}(z) \gamma^z \frac{\mathcal{P}\left[-ig_s \int_0^\infty \mathrm{d}z' A^{a,z}\left(z'\right) t^a\right] \left[-ig_s \int_0^\infty \mathrm{d}z'' A^{a,z}\left(z''\right) t^a\right]}{2} \psi(0) | P, S \rangle \tag{1}$$

Applying Feynman rule straightaway gives (the overall 1/2 factor has been counted in)

$$\Gamma_g(l) = P$$

$$\downarrow P = -g_s^2 C_F \delta(1-x) \int \frac{\mathrm{d}^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{i\epsilon} \frac{i}{-l^2 + i\epsilon}$$
(2)

#### 1.1 Direct Contraction

## 1.1.1 Left

$$\frac{1}{2!} \mathcal{P} \left[ \int_{0}^{\infty} dz' A^{a,z} (z') \int_{0}^{\infty} dz'' A^{a,z} (z'') \right] 
= \int_{0}^{\infty} dz' A^{a,z} (z') \int_{0}^{\infty} dz'' A^{a,z} (z'') \theta(z' - z'') 
= \int dz' A^{a,z} (z') \int dz'' A^{a,z} (z'') \theta(z' - z'') \theta(z'') 
= \int dz' A^{a,z} (z') \int dz'' A^{a,z} (z'') \theta(z' - z'') \theta(z'') 
= \int dz' \int dz'' \int \frac{d^{4}l}{(2\pi)^{4}} \frac{i}{l^{2} + i\epsilon} e^{-il\cdot(z'' - z')} \theta(z' - z'') \theta(z'')$$

#### 1.1.2 Right

$$\begin{split} &\frac{1}{2!}\mathcal{P}\left[\int_{0}^{\infty}\mathrm{d}z''A^{a,z}\left(z''+z\right)\int_{0}^{\infty}\mathrm{d}z'A^{a,z}\left(z'+z\right)\right]\\ &=\int_{0}^{\infty}\mathrm{d}z''A^{a,z}\left(z''+z\right)\int_{0}^{\infty}\mathrm{d}z'A^{a,z}\left(z'+z\right)\theta(z''-z')\\ &=\int\mathrm{d}z''A^{a,z}\left(z''+z\right)\int\mathrm{d}z'A^{a,z}\left(z'+z\right)\theta(z''-z')\theta(z')\\ &=\int\mathrm{d}z''\overline{A^{a,z}\left(z''+z\right)\int\mathrm{d}z'A^{a,z}\left(z'+z\right)\theta(z''-z')\theta(z')}\\ &=\int\mathrm{d}z'\int\mathrm{d}z''\int\frac{\mathrm{d}^{4}l}{(2\pi)^{4}}\frac{i}{l^{2}+i\epsilon}e^{-il\cdot(z''-z')}\theta(z''-z')\theta(z') \end{split}$$

#### 1.1.3 Summing together

$$\int \mathrm{d}z' \int \mathrm{d}z'' \int \frac{\mathrm{d}^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} e^{-il\cdot(z''-z')} [\theta(z'-z'')\theta(z'') + \theta(z''-z')\theta(z')]$$

### 1.2 Adding two path order together

$$\begin{split} &\mathcal{P}\left[\int_{0}^{\infty} \mathrm{d}z' A^{a,z}\left(z'\right) \int_{0}^{\infty} \mathrm{d}z'' A^{a,z}\left(z''\right)\right] \\ &= \int_{0}^{\infty} \mathrm{d}z' A^{a,z}\left(z'\right) \int_{0}^{\infty} \mathrm{d}z'' A^{a,z}\left(z''\right) \left[\theta(z'-z'') + \theta(z''-z')\right] \\ &= \int_{0}^{\infty} \mathrm{d}z' A^{a,z}\left(z'\right) \int_{0}^{\infty} \mathrm{d}z'' A^{a,z}\left(z''\right) \\ &= \int \mathrm{d}z' A^{a,z}\left(z'\right) \int \mathrm{d}z'' A^{a,z}\left(z''\right) \int \frac{\mathrm{d}w}{2\pi} \frac{i e^{-iwz'}}{w + i\epsilon} \int \frac{\mathrm{d}h}{2\pi} \frac{i e^{-ihz''}}{h + i\epsilon} \\ &= \int \mathrm{d}z' A^{a,z}\left(z'\right) \int \mathrm{d}z'' A^{a,z}\left(z''\right) \int \frac{\mathrm{d}w}{2\pi} \frac{i e^{-iwz'}}{w + i\epsilon} \int \frac{\mathrm{d}h}{2\pi} \frac{i e^{-ihz''}}{h + i\epsilon} \\ &= \int \mathrm{d}z' \int \mathrm{d}z'' \int \frac{\mathrm{d}^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} e^{-il\cdot(z''-z')} \int \frac{\mathrm{d}w}{2\pi} \frac{i e^{-iwz'}}{w + i\epsilon} \int \frac{\mathrm{d}h}{2\pi} \frac{i e^{-ihz''}}{h + i\epsilon} \\ &= \int \frac{\mathrm{d}^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \int \mathrm{d}z' \int \mathrm{d}z'' \int \frac{\mathrm{d}w}{2\pi} \frac{i}{w + i\epsilon} \int \frac{\mathrm{d}h}{2\pi} \frac{i}{h + i\epsilon} e^{-i(w - l) \cdot z'} e^{-i(h + l) \cdot z''} \\ &= \int \frac{\mathrm{d}^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \end{aligned}$$

The amplitude is

$$P = -\frac{g_s^2 C_F}{2} \delta(1-x) \int \frac{\mathrm{d}^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{l^2 + i\epsilon} \frac{i}{-l^2 + i\epsilon}$$
(3)

## **1.3** Adding $\Gamma_q(l)$ and $\Gamma_q(-l)$

$$\begin{split} \Gamma_g(l) &= \frac{\Gamma_g(l) + \Gamma_g(-l)}{2} = -\frac{1}{2} \bigg[ g_s^2 C_F \delta(1-x) \int \frac{\mathrm{d}^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{i\epsilon} \frac{i}{-l^2 + i\epsilon} + g_s^2 C_F \delta(1-x) \int \frac{\mathrm{d}^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{i\epsilon} \frac{i}{l^2 + i\epsilon} \bigg] \\ &= -\frac{1}{2} g_s^2 C_F \delta(1-x) \int \frac{\mathrm{d}^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{i\epsilon} \bigg[ \frac{i}{-l^2 + i\epsilon} + \frac{i}{l^2 + i\epsilon} \bigg] \\ &= -\frac{1}{2} g_s^2 C_F \delta(1-x) \int \frac{\mathrm{d}^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{i\epsilon} \frac{2\epsilon}{l^2 + \epsilon^2} \\ &= g_s^2 C_F \delta(1-x) \int \frac{\mathrm{d}^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{l^2 + i\epsilon} \frac{i}{l^2 - i\epsilon} \end{split}$$

There's an overall factor of 1/2 missing.

## 1.4 Taking derivatives

Add a small momentum to the gauge link line and consider an actual self energy diagram

$$P \downarrow \downarrow \qquad \qquad \downarrow P = -g_s^2 C_F \delta(1-x) \int \frac{\mathrm{d}^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{k^z - l^z + i\epsilon}$$

$$\tag{4}$$

take the derivative

$$-g_s^2 C_F \delta(1-x) \lim_{k^z \to 0} \frac{\partial}{\partial k^z} \left[ (i) \int \frac{\mathrm{d}^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{k^z - l^z} \right]$$
 (5)

$$=ig_s^2 C_F \delta(1-x) \int \frac{\mathrm{d}^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{(l^2)^2}$$
 (6)

Adding  $i\epsilon$  by hand and we got

$$g_s^2 C_F \delta(1-x) \int \frac{\mathrm{d}^4 l}{(2\pi)^4} \frac{i}{l^2 + i\epsilon} \frac{i}{l^2 + i\epsilon} \frac{i}{l^2 - i\epsilon}$$
 (7)

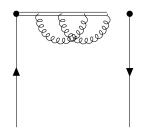
# 2 Two Loop

In coordinate space

$$\langle P, S | \bar{\psi}(z) \gamma^{z} \mathcal{P} \frac{\left[ -ig_{s} n_{\mu} \int_{0}^{\infty} \mathrm{d}z_{1} A^{a,\mu}(z_{1}) t^{a} \right] \left[ -ig_{s} n_{\nu} \int_{0}^{\infty} \mathrm{d}z_{2} A^{b,\nu}(z_{2}) t^{b} \right] \left[ -ig_{s} n_{\rho} \int_{0}^{\infty} \mathrm{d}z_{3} A^{c,\rho}(z_{3}) t^{c} \right] \left[ -ig_{s} n_{\sigma} \int_{0}^{\infty} \mathrm{d}z_{4} A^{d,\sigma}(z_{4}) t^{d} \right]}{4!} \psi(0) |P, S\rangle \tag{8}$$

$$\begin{split} &\frac{1}{4!} \mathcal{P} \bigg[ \int_0^\infty \mathrm{d}z_1 A^{a,\mu}(z_1) \bigg] \bigg[ \int_0^\infty \mathrm{d}z_2 A^{b,\nu}(z_2) \bigg] \bigg[ \int_0^\infty \mathrm{d}z_3 A^{c,\rho}(z_3) \bigg] \bigg[ \int_0^\infty \mathrm{d}z_4 A^{d,\sigma}(z_4) \bigg] \\ &= \bigg[ \int_0^\infty \mathrm{d}z_1 A^{a,\mu}(z_1) \bigg] \bigg[ \int_0^\infty \mathrm{d}z_2 A^{b,\nu}(z_2) \bigg] \bigg[ \int_0^\infty \mathrm{d}z_3 A^{c,\rho}(z_3) \bigg] \bigg[ \int_0^\infty \mathrm{d}z_4 A^{d,\sigma}(z_4) \bigg] \theta(z_1 - z_2) \theta(z_2 - z_3) \theta(z_3 - z_4) \end{split}$$

The amplitude for



is related to the color ordering  $t^a t^b t^a t^b$ .

$$\begin{split} & \int_{0}^{\infty} \mathrm{d}z_{1} \int_{0}^{\infty} \mathrm{d}z_{2} \int_{0}^{\infty} \mathrm{d}z_{3} \int_{0}^{\infty} \mathrm{d}z_{4} A^{a,\mu}(z_{1}) A^{b,\nu}(z_{2}) A^{c,\rho}(z_{3}) A^{d,\sigma}(z_{4}) \theta(z_{1}-z_{2}) \theta(z_{2}-z_{3}) \theta(z_{3}-z_{4}) \\ &= \int_{0}^{\infty} \mathrm{d}z_{1} \int_{0}^{\infty} \mathrm{d}z_{2} \int_{0}^{\infty} \mathrm{d}z_{3} \int_{0}^{\infty} \mathrm{d}z_{4} \int \frac{\mathrm{d}^{4}l_{1}}{(2\pi)^{4}} \frac{-ig^{\mu\rho} \delta^{ac}}{l_{1}^{2}+i\epsilon} e^{-il_{1}\cdot(z_{3}-z_{1})} \int \frac{\mathrm{d}^{4}l_{2}}{(2\pi)^{4}} \frac{-ig^{\nu\sigma} \delta^{bd}}{l_{2}^{2}+i\epsilon} e^{-il_{2}\cdot(z_{4}-z_{2})} \\ &= \int \mathrm{d}z_{1} \int \mathrm{d}z_{2} \int \mathrm{d}z_{3} \int \mathrm{d}z_{4} \int \frac{\mathrm{d}^{4}l_{1}}{(2\pi)^{4}} \frac{-ig^{\mu\rho} \delta^{ac}}{l_{1}^{2}+i\epsilon} e^{-il_{1}\cdot(z_{3}-z_{1})} \int \frac{\mathrm{d}^{4}l_{2}}{(2\pi)^{4}} \frac{-ig^{\nu\sigma} \delta^{bd}}{l_{2}^{2}+i\epsilon} e^{-il_{2}\cdot(z_{4}-z_{2})} \\ &\theta(z_{1}-z_{2})\theta(z_{2}-z_{3})\theta(z_{3}-z_{4})\theta(z_{4}) \end{split}$$

The exponent is (for simplicity we assume vectors  $z_i = (0, 0, 0, z_i)$  and  $k_i = (0, 0, 0, -k_i)$ )

$$-il_1 \cdot (z_3 - z_1) - il_2 \cdot (z_4 - z_2) - ik_1 \cdot (z_1 - z_2) - ik_2 \cdot (z_2 - z_3) - ik_3 \cdot (z_3 - z_4) - ik_4 \cdot z_4$$

$$= -iz_3 \cdot (l_1 + k_3 - k_2) - iz_1 \cdot (k_1 - l_1) - iz_4 \cdot (l_2 + k_4 - k_3) - iz_2 \cdot (k_2 - k_1 - l_2)$$

which gives 4 delta functions. The propagators involved are then

$$\int \frac{\mathrm{d}^4 l_1}{(2\pi)^4} \int \frac{\mathrm{d}^4 l_2}{(2\pi)^4} \frac{-ig^{\mu\rho}\delta^{ac}}{l_1^2 + i\epsilon} \frac{-ig^{\nu\sigma}\delta^{bd}}{l_2^2 + i\epsilon} \frac{i}{-l_1^z + i\epsilon} \frac{i}{-l_1^z + i\epsilon} \frac{i}{-l_2^z + i\epsilon} \frac{i}{i\epsilon}$$

$$\int dz_{1} \int dz_{2} \int dz_{3} \int dz_{4} A^{a,\mu}(z_{1}) A^{b,\nu}(z_{2}) A^{c,\rho}(z_{3}) A^{d,\sigma}(z_{4}) \theta(z_{1}-z_{2}) \theta(z_{2}-z_{3}) \theta(z_{3}-z_{4}) \theta(z_{4})$$

$$= \frac{1}{2} \int dz_{1} \int dz_{2} \int dz_{3} \int dz_{4} A^{a,\mu}(z_{1}) A^{b,\nu}(z_{2}) A^{c,\rho}(z_{3}) A^{d,\sigma}(z_{4}) [\theta(z_{1}-z_{2})\theta(z_{2}-z_{3})\theta(z_{3}-z_{4})\theta(z_{4})$$

$$+ \theta(z_{3}-z_{4})\theta(z_{4}-z_{1})\theta(z_{1}-z_{2})\theta(z_{2})]$$

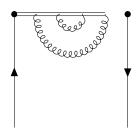
$$= \int \frac{d^{4}l_{1}}{(2\pi)^{4}} \int \frac{d^{4}l_{2}}{(2\pi)^{4}} \frac{-ig^{\mu\rho}\delta^{ac}}{l_{1}^{2}+i\epsilon} \frac{-ig^{\nu\sigma}\delta^{bd}}{l_{2}^{2}+i\epsilon} \frac{-(l_{1}^{z^{2}}+3l_{1}^{z}l_{2}^{z}+l_{2}^{z^{2}}-\epsilon^{2})}{(l_{1}^{z}-i\epsilon)(l_{1}^{z}+i\epsilon)(l_{2}^{z}-i\epsilon)(l_{1}^{z}+l_{2}^{z}-i\epsilon)(l_{1}^{z$$

Taking the derivative of the former expression, we can also arrive at a divergence-free form

$$\frac{i}{2}\lim_{p\to 0}\frac{\partial}{\partial p}\left[\frac{i}{p+l_1^z}\frac{i}{p+l_1^z+l_2^z}\frac{i}{p+l_2^z}\right] = \frac{-\left(l_1^{z^2}+3l_1^zl_2^z+l_2^{z^2}\right)}{l_1^{z^2}l_2^{z^2}(l_1^z+l_2^z)^2}$$

which is equivalent to above expression.

The amplitude for



is related to the color ordering  $t^a t^b t^b t^a$ .

$$\int_{0}^{\infty} dz_{1} \int_{0}^{\infty} dz_{2} \int_{0}^{\infty} dz_{3} \int_{0}^{\infty} dz_{4} A^{a,\mu}(z_{1}) A^{b,\nu}(z_{2}) A^{c,\rho}(z_{3}) A^{d,\sigma}(z_{4}) \theta(z_{1} - z_{2}) \theta(z_{2} - z_{3}) \theta(z_{3} - z_{4})$$

$$= \int dz_{1} \int dz_{2} \int dz_{3} \int dz_{4} \int \frac{d^{4}l_{1}}{(2\pi)^{4}} \frac{-ig^{\mu\sigma} \delta^{ad}}{l_{1}^{2} + i\epsilon} e^{-il_{1} \cdot (z_{3} - z_{1})} \int \frac{d^{4}l_{2}}{(2\pi)^{4}} \frac{-ig^{\nu\rho} \delta^{bc}}{l_{2}^{2} + i\epsilon} e^{-il_{2} \cdot (z_{4} - z_{2})}$$

$$\theta(z_{1} - z_{2}) \theta(z_{2} - z_{3}) \theta(z_{3} - z_{4}) \theta(z_{4})$$

The propagators involved are

$$\int \frac{\mathrm{d}^{4}l_{1}}{(2\pi)^{4}} \int \frac{\mathrm{d}^{4}l_{2}}{(2\pi)^{4}} \frac{-ig^{\mu\sigma}\delta^{ad}}{l_{1}^{2} + i\epsilon} \frac{-ig^{\nu\rho}\delta^{bc}}{l_{2}^{2} + i\epsilon} \frac{i}{-l_{1}^{z} + i\epsilon} \frac{i}{-l_{1}^{z} - l_{2}^{z} + i\epsilon} \frac{i}{-l_{1}^{z} + i\epsilon} \frac{i}{i\epsilon}$$

$$\frac{1}{2} \int \mathrm{d}z_{1} \int \mathrm{d}z_{2} \int \mathrm{d}z_{3} \int \mathrm{d}z_{4} A^{a,\mu}(z_{1}) A^{b,\nu}(z_{2}) A^{c,\rho}(z_{3}) A^{d,\sigma}(z_{4}) [\theta(z_{1} - z_{2})\theta(z_{2} - z_{3})\theta(z_{3} - z_{4})\theta(z_{4}) + \theta(z_{4} - z_{3})\theta(z_{3} - z_{2})\theta(z_{2} - z_{1})\theta(z_{1})]$$

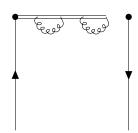
$$= \int \frac{\mathrm{d}^{4}l_{1}}{(2\pi)^{4}} \int \frac{\mathrm{d}^{4}l_{2}}{(2\pi)^{4}} \frac{-ig^{\mu\sigma}\delta^{ad}}{l_{1}^{2} + i\epsilon} \frac{-ig^{\nu\rho}\delta^{bc}}{l_{2}^{2} + i\epsilon} \frac{-3l_{1}^{z^{2}} - 2l_{1}^{z}l_{2}^{z} + \epsilon^{2}}{(l_{1}^{z^{2}} + 2l_{1}^{z}l_{2}^{z} + l_{2}^{z}^{z} + \epsilon^{2})}$$

Taking the derivative of the former expression, we can also arrive at a divergence-free form

$$\frac{i}{2} \lim_{p \to 0} \frac{\partial}{\partial p} \left[ \frac{i}{p + l_1^z} \frac{i}{p + l_1^z + l_2^z} \frac{i}{p + l_1^z} \right] = \frac{-(3l_1^z + 2l_2^z)}{l_1^z}$$

which is equivalent to above expression.

The amplitude for



is related to the color ordering  $t^a t^a t^b t^b$ .

$$\int_{0}^{\infty} dz_{1} \int_{0}^{\infty} dz_{2} \int_{0}^{\infty} dz_{3} \int_{0}^{\infty} dz_{4} A^{a,\mu}(z_{1}) A^{b,\nu}(z_{2}) A^{c,\rho}(z_{3}) A^{d,\sigma}(z_{4}) \theta(z_{1}-z_{2}) \theta(z_{2}-z_{3}) \theta(z_{3}-z_{4})$$

$$= \int dz_{1} \int dz_{2} \int dz_{3} \int dz_{4} \int \frac{d^{4}l_{1}}{(2\pi)^{4}} \frac{-ig^{\mu\sigma} \delta^{ad}}{l_{1}^{2}+i\epsilon} e^{-il_{1}\cdot(z_{3}-z_{1})} \int \frac{d^{4}l_{2}}{(2\pi)^{4}} \frac{-ig^{\nu\rho} \delta^{bc}}{l_{2}^{2}+i\epsilon} e^{-il_{2}\cdot(z_{4}-z_{2})}$$

$$\theta(z_{1}-z_{2})\theta(z_{2}-z_{3})\theta(z_{3}-z_{4})\theta(z_{4})$$

The propagators involved are

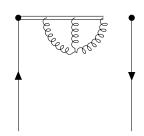
$$\int \frac{\mathrm{d}^{4}l_{1}}{(2\pi)^{4}} \int \frac{\mathrm{d}^{4}l_{2}}{(2\pi)^{4}} \frac{-ig^{\mu\sigma}\delta^{ad}}{l_{1}^{2} + i\epsilon} \frac{-ig^{\nu\rho}\delta^{bc}}{l_{2}^{2} + i\epsilon} \frac{i}{-l_{1}^{z} + i\epsilon} \frac{i}{i\epsilon} \frac{i}{-l_{2}^{z} + i\epsilon} \frac{i}{i\epsilon}$$

$$\frac{1}{4} \int \mathrm{d}z_{1} \int \mathrm{d}z_{2} \int \mathrm{d}z_{3} \int \mathrm{d}z_{4} \overline{A^{a,\mu}(z_{1})} \overline{A^{b,\nu}(z_{2})} \overline{A^{c,\rho}(z_{3})} \overline{A^{d,\sigma}(z_{4})} [\theta(z_{1} - z_{2})\theta(z_{2} - z_{3})\theta(z_{3} - z_{4})\theta(z_{4})$$

$$+ \theta(z_{2} - z_{1})\theta(z_{1} - z_{3})\theta(z_{3} - z_{4})\theta(z_{4}) + \theta(z_{1} - z_{2})\theta(z_{2} - z_{4})\theta(z_{4} - z_{3})\theta(z_{3}) + \theta(z_{2} - z_{1})\theta(z_{1} - z_{4})\theta(z_{4} - z_{3})\theta(z_{3}) ]$$

$$= \int \frac{\mathrm{d}^{4}l_{1}}{(2\pi)^{4}} \int \frac{\mathrm{d}^{4}l_{2}}{(2\pi)^{4}} \frac{-ig^{\mu\sigma}\delta^{ad}}{l_{1}^{2} + i\epsilon} \frac{-ig^{\nu\rho}\delta^{bc}}{l_{2}^{2} + i\epsilon} \frac{1}{(l_{1}^{z^{2}} + \epsilon^{2})(l_{2}^{z^{2}} + \epsilon^{2})}$$

The amplitude for



is related to

$$\langle P, S | \bar{\psi}(z) \gamma^z V_3 \mathcal{P} \frac{\left[ -ig_s n_\mu \int_0^\infty \mathrm{d}z_1 A^{a,\mu}(z_1) t^a \right] \left[ -ig_s n_\nu \int_0^\infty \mathrm{d}z_2 A^{b,\nu}(z_2) t^b \right] \left[ -ig_s n_\rho \int_0^\infty \mathrm{d}z_3 A^{c,\rho}(z_3) t^c \right]}{3!} \psi(0) | P, S \rangle \tag{9}$$

where

$$V_3 = -\frac{g_s}{2} f^{def} \int d^4t \left( \partial^\alpha A_d^\beta - \partial^\beta A_d^\alpha \right) A_\alpha^e A_\beta^f$$

The gluon related contraction (one out of three, others can be obtained by exchanging d, e, f.)

$$\begin{split} &\int \mathrm{d}z_{1} \int \mathrm{d}z_{2} \int \mathrm{d}z_{3} \int \mathrm{d}^{4}t \Big( \partial^{\alpha}A^{d,\beta} - \partial^{\beta}A^{d,\alpha} \Big) A_{\alpha}^{e,A} A_{\beta}^{f}A^{a,\mu}(z_{1}) A^{b,\nu}(z_{2}) A^{c,\rho}(z_{3}) \theta(z_{1} - z_{2}) \theta(z_{2} - z_{3}) \theta(z_{3}) \\ &= \int \mathrm{d}z_{1} \int \mathrm{d}z_{2} \int \mathrm{d}z_{3} \int \mathrm{d}^{4}t \int \frac{\mathrm{d}^{4}l_{1}}{(2\pi)^{4}} \frac{-i\delta^{dc}}{l_{1}^{2} + i\epsilon} \Big[ g^{\beta\rho}\partial^{\alpha}e^{-il_{1}\cdot(z_{3} - t)} - g^{\alpha\rho}\partial^{\beta}e^{-il_{1}\cdot(z_{3} - t)} \Big] \int \frac{\mathrm{d}^{4}l_{2}}{(2\pi)^{4}} \frac{-ig_{\alpha}^{\mu}\delta^{ea}}{l_{2}^{2} + i\epsilon} e^{-il_{2}\cdot(z_{1} - t)} \\ &\int \frac{\mathrm{d}^{4}l_{3}}{(2\pi)^{4}} \frac{-ig_{\beta}^{\nu}\delta^{fb}}{l_{3}^{2} + i\epsilon} e^{-il_{3}\cdot(z_{2} - t)} \theta(z_{1} - z_{2}) \theta(z_{2} - z_{3}) \theta(z_{3}) \\ &= \int \mathrm{d}z_{1} \int \mathrm{d}z_{2} \int \mathrm{d}z_{3} \int \mathrm{d}^{4}t \int \frac{\mathrm{d}^{4}l_{1}}{(2\pi)^{4}} \frac{-i\delta^{dc}}{l_{1}^{2} + i\epsilon} \Big[ g^{\beta\rho}il_{1}^{\alpha} - g^{\alpha\rho}il_{1}^{\beta} \Big] e^{-il_{1}\cdot(z_{3} - t)} \int \frac{\mathrm{d}^{4}l_{2}}{(2\pi)^{4}} \frac{-ig_{\alpha}^{\mu}\delta^{ea}}{l_{2}^{2} + i\epsilon} e^{-il_{2}\cdot(z_{1} - t)} \\ &\int \frac{\mathrm{d}^{4}l_{3}}{(2\pi)^{4}} \frac{-ig_{\beta}^{\nu}\delta^{fb}}{l_{3}^{2} + i\epsilon} e^{-il_{3}\cdot(z_{2} - t)} \theta(z_{1} - z_{2}) \theta(z_{2} - z_{3}) \theta(z_{3}) \\ &= \int \mathrm{d}z_{1} \int \mathrm{d}z_{2} \int \mathrm{d}z_{3} \int \mathrm{d}^{4}t \int \frac{\mathrm{d}^{4}l_{1}}{(2\pi)^{4}} \frac{-i\delta^{dc}}{l_{1}^{2} + i\epsilon} \Big[ g^{\nu\rho}il_{1}^{\mu} - g^{\mu\rho}il_{1}^{\nu} \Big] e^{-il_{1}\cdot(z_{3} - t)} \int \frac{\mathrm{d}^{4}l_{2}}{(2\pi)^{4}} \frac{-i\delta^{ea}}{l_{2}^{2} + i\epsilon} e^{-il_{2}\cdot(z_{1} - t)} \\ &\int \frac{\mathrm{d}^{4}l_{3}}{(2\pi)^{4}} \frac{-i\delta^{fb}}{l_{2}^{2} + i\epsilon} e^{-il_{3}\cdot(z_{2} - t)} \theta(z_{1} - z_{2}) \theta(z_{2} - z_{3}) \theta(z_{3}) \end{aligned}$$

Multiplied by n

$$n_{\mu}n_{\nu}n_{\rho}[g^{\nu\rho}il_{1}^{\mu}-g^{\mu\rho}il_{1}^{\nu}]=in^{2}[l_{1}^{z}-l_{1}^{z}]=0$$

Contracting with different fields in the parenthesis

$$\partial^{\alpha} A^{d,\beta} A^{e}_{\alpha} A^{f}_{\beta} A^{a,\mu}(z_{1}) A^{b,\nu}(z_{2}) A^{c,\rho}(z_{3}) - \partial^{\beta} A^{d,\alpha} A^{e}_{\alpha} A^{f}_{\beta} A^{a,\mu}(z_{1}) A^{b,\nu}(z_{2}) A^{c,\rho}(z_{3})$$

$$\int dz_{1} \int dz_{2} \int dz_{3} \int d^{4}t \int \frac{d^{4}l_{1}}{(2\pi)^{4}} \int \frac{d^{4}l_{2}}{(2\pi)^{4}} \int \frac{d^{4}l_{3}}{(2\pi)^{4}} \frac{-i\delta^{dc}}{l_{1}^{2} + i\epsilon} \frac{-ig_{\alpha}^{\mu} \delta^{ea}}{l_{2}^{2} + i\epsilon} e^{-il_{2} \cdot (z_{1} - t)} \frac{-ig_{\beta}^{\nu} \delta^{fb}}{l_{3}^{2} + i\epsilon}$$

$$\left[ e^{-il_{3} \cdot (z_{2} - t)} g^{\beta\rho} \partial^{\alpha} e^{-il_{1} \cdot (z_{3} - t)} - e^{-il_{1} \cdot (z_{3} - t)} g^{\alpha\rho} \partial^{\beta} e^{-il_{3} \cdot (z_{2} - t)} \right] \theta(z_{1} - z_{2}) \theta(z_{2} - z_{3}) \theta(z_{3})$$

$$= \int dz_{1} \int dz_{2} \int dz_{3} \int d^{4}t \int \frac{d^{4}l_{1}}{(2\pi)^{4}} \int \frac{d^{4}l_{2}}{(2\pi)^{4}} \int \frac{d^{4}l_{3}}{(2\pi)^{4}} \frac{-i\delta^{dc}}{l_{1}^{2} + i\epsilon} \frac{-i\delta^{fb}}{l_{2}^{2} + i\epsilon} e^{-il_{2} \cdot (z_{1} - t)} e^{-il_{3} \cdot (z_{2} - t)} e^{-il_{1} \cdot (z_{3} - t)}$$

$$\left[ g^{\nu\rho} i l_{1}^{\mu} - g^{\mu\rho} i l_{3}^{\nu} \right] \theta(z_{1} - z_{2}) \theta(z_{2} - z_{3}) \theta(z_{3})$$

Make  $z_1 \to -z_1, t \to -t, l_2 \to -l_2,$ 

$$\int dz_{1} \int dz_{2} \int dz_{3} \int d^{4}t \int \frac{d^{4}l_{1}}{(2\pi)^{4}} \int \frac{d^{4}l_{2}}{(2\pi)^{4}} \int \frac{d^{4}l_{3}}{(2\pi)^{4}} \frac{-i\delta^{dc}}{l_{1}^{2} + i\epsilon} \frac{-i\delta^{fb}}{l_{2}^{2} + i\epsilon} e^{-il_{2}\cdot(z_{1} - t)} e^{-il_{3}\cdot(z_{2} - t)} e^{-il_{1}\cdot(z_{3} - t)} \\
[g^{\nu\rho}il_{1}^{\mu} - g^{\mu\rho}il_{3}^{\nu}]\theta(-z_{1} - z_{2})\theta(z_{2} - z_{3})\theta(z_{3}) \\
= \int \frac{d^{4}l_{1}}{(2\pi)^{4}} \int \frac{d^{4}l_{3}}{(2\pi)^{4}} \frac{-i\delta^{dc}}{l_{1}^{2} + i\epsilon} \frac{-i\delta^{fb}}{(l_{1} + l_{3})^{2} + i\epsilon} \frac{i}{l_{3}^{2} + i\epsilon} \frac{i}{-l_{1}^{z} + i\epsilon} \frac{i}{-2l_{1}^{z} + i\epsilon} \frac{i}{-2l_{1}^{z} - l_{3}^{z} + i\epsilon}$$

