# MOLECULAR DYNAMICS SIMULATION OF ARGON

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# **ABSTRACT**

The motion of a collection of 864 particles has been simulated using Molecular Dynamics techniques to compute values for pressure, specific heat and the static pair correlation function in reduced units. The computed properties were compared with experimental results for the properties of argon. The simulation was done with a Lennard-Jones pair-potential and the system was allowed to reach equilibrium. The computed results are in good agreement with the known properties of argon.

#### INTRODUCTION 1

Molecular dynamics is a method to simulate a many-particle system by numerically solving Newton's classical equations of motion for all particles for a period of time. The main limitations are the fact that simulations are only realizable for small systems and times compared to experimental systems in general. Furthermore, the systems to be simulated can only be classical in the usual molecular dynamics approach.

In this report we discuss the results of simulations for a system of argon gas with 864 particles. Molecular dynamics simulations of argon are reported to be in good agreement with experimental results[1]. The particles are placed in a FCC lattice, considering the fact that the ground state configuration is a FCC lattice for argon. A Lennard-Jones Potential is implemented for the calculation of the force between the particles where a pair potential is assumed. The initial velocities of the particles are randomly generated for each velocity component while obeying a Maxwell-Boltzmann distribution for all particles.

After initialization the system's equations of motion are solved numerically after each time step using Verlet's algorithm. The system is allowed to relax to it's equilibrium using a thermostat to rescale the velocities and enforce energy conservation. After the equilibrium has been reached, the simulation continues and collects results for calculating the time average of different properties of the system. The virial's theorem allows the calculation of the pressure, specific heat is evaluated using a formula derived by Lebowitz using the fluctuations of the kinetic energy[2].

2.1 Lennard Jones Potential

2.2 Virial Theorem

2.3 Specific Heat

The specific heat at constant volume  $C_V$  is defined as:

$$C_{\mathbf{V}} = \left(\frac{\delta E}{\delta T}\right)_{\mathbf{V}}$$

In molecular dynamics simulations, a quantity that can be calculated is the ensemble average of the total energy  $\langle E \rangle_{NVT}$ :

$$\langle E \rangle_{\text{NVT}} = \frac{\sum_{X} e^{-\beta \mathfrak{H}(X)} \mathfrak{H}(X)}{\sum_{X} e^{-\beta \mathfrak{H}(X)}} = -\frac{\delta ln(Z)}{\delta \beta}$$

Using the ensemble average for the total energy, the formula for the specific heat can be rewritten as a function of the fluctuations in the total energy:

$$C_{V} = \frac{1}{k_{\rm B}T^2} \frac{\delta^2 \ln(Z)}{\delta \beta^2} \tag{1}$$

$$= \frac{1}{k_{\rm B}T^2} \left( \langle E^2 \rangle_{\rm NVT} - \langle E \rangle_{\rm NVT}^2 \right) \tag{2}$$

This is still difficult to use in a program simulating a system in the microcanonical ensemble as the total energy is kept fixed, but following the derivation by Lebowitz[2] this can be related to the fluctuation in kinetic energy:

$$\frac{\langle \delta K^2 \rangle}{\langle K \rangle^2} = \frac{2}{3N} \left( 1 - \frac{3N}{2C_V} \right)$$

## 2.4 Pair correlation function

The static pair correlation function in the canonical ensemble can be expressed as:

$$g({\bf r},{\bf r}') = V^2 \frac{1}{N! h^{3N} Z} \int_V d^3 r_3 .... d^3 r_N e^{-\beta V_N ({\bf r},{\bf r}',{\bf r}_3 .....{\bf r}_N)}$$

If the system is homogeneous and has a large number of particles N, it can also be written as a function of the difference  $\Delta \mathbf{r} = \mathbf{r} - \mathbf{r}'$ :

$$g(\Delta \mathbf{r}) = \frac{V}{N(N-1)} \left\langle \int d^3 \mathbf{r}' \sum_{i,j i \neq j}^{N} \delta(\mathbf{r} - \mathbf{r}_i) \delta(\mathbf{r}' + \Delta \mathbf{r} - \mathbf{r}_j) \right\rangle$$

#### METHODS FOR SIMULATION 3

## Initialization

The system is initialized by creating an fcc lattice structure and assigning each position of the particles to the lattice points. This is done for 864 particles but can in principle be easily implemented for a number of particles N that can be written as 4M<sup>3</sup> with M an integer and representing the number of fcc cells in one dimension.

```
subroutine cubic_fcc_lattice(N_cell_dim, L_side, pos)
 do i = 1,N_cell_dim
   do j = 1,N_cell_dim
     do k = 1,N_cell_dim
       do l = 1,4
         n = n + 1
         pos(:,n) = L_side/N_cell_dim*((/ i-1, j-1, k-1 /) + fcc_cell(l))
     end do
   end do
 end do
end subroutine
```

Listing 1: Constructing the fcc lattice

Above can you see the code with i, j and k for the x, y and z component and a function fcc\_cell is called that takes an argument l to add the relative position of each particle in a particular cell.

## 3.1.1 Velocity Distribution

The particles are then assigned an initial velocity that is randomly distributed but the distribution in each direction obeys a Maxwell-Boltzmann distribution:

$$f(\nu_i) = \sqrt{\frac{m}{2\pi k_B T}} e^{\frac{-m\nu_i^2}{2k_B T}}$$

```
subroutine init_vel(T, Kb, m, N_part, vel)
 do n = 1, N_part
   do i = 1,3
     CALL RANDOM_NUMBER(xs(1))
     CALL RANDOM_NUMBER(xs(2))
     vel(i,n) = sqrt(Kb*T/m) * sqrt(-2d0*log(xs(1)))*cos(2*pi*xs(2)) !sigma
           * box muller
   end do
 end do
end subroutine
```

**Listing 2:** Generating initial velocities

```
subroutine init_vel(T, Kb, m, N_part, vel)
  !Set center of mass velocity to zero
```

```
do i = 1,3
   vel(i,:) = vel(i,:) - sum(vel(i,:))/N_part
 end do
end subroutine
```

**Listing 3:** Removing the center of mass degree of freedom

#### 3.1.1.1BOX MULLER ALGORITHM

## 3.2 Dynamics

## 3.2.1 Boundary Conditions

As the system to be simulated is very small compared to real experimental systems, boundaries would probably have a larger effect in the simulation. To counter this periodic boundary conditions are used and this is implemented by copying the system in all directions creating a large box of 9 identical smaller boxes with the box in the middle as our box of interest. As we calculate the forces the Lennard Jones pair potential is calculated for all the particles inside the box and also all the other 8 copies that are closer than a predefined cut off distance.

```
subroutine calc_dynamics(calc_quant, N_part, L_side, time_step, m, e, s,
     r_cut, pos, kin_energy, pot_energy, virial, vel, num_intervals,
    histogram_vector)
 do n = 1, N_part
   do\ i = 1, N_part ! integrate over all particles inside box except <math>i = n
     do j = -1, 1
     do k = -1, 1 !periodic boundary condition
     do l = -1, 1
        if (n/=i) then
          r_{\text{vec}} = (/pos(1,n) - pos(1,i), pos(2,n) - pos(2,i), pos(3,n) - pos(3,i)
               /) + L_side*(/j,k,l/)
          r = sqrt(dot_product(r_vec, r_vec))
          !force calculation
          if (r<r_cut) then
            dF = e*(48*s**12/r**14 - 24*s**6/r**8) * r_vec
            F = F + dF
          end if
        end if
      end do
      end do
     end do
   end do
 end do
end subroutine
```

Listing 4: Creating periodic boundary conditions

- 3.2.2 Lennard Jones Potential
- 3.2.3 Force Calculation
- 3.2.3.1LEAP FROG/VERLET ALGORITHM
- 3.2.4 Pressure Calculation
- 3.2.4.1 VIRIAL THEOREM

Thermostat 3.2.5

Information Processing

#### Specific Heat 3.3.1

The implementation in our simulation is inside the algorithm for the dynamics of the particles. For each time step the new velocities are calculated, the kinetic energy is also calculated and the relevant sums of the kinetic energy are updated in a subroutine "calc\_specific\_heat":

```
subroutine calc_specific_heat(end_of_routine,N_part, kin_energy,
    sum_kin_energy, sum_kin_energy_sqr)
  sum_kin_energy_sqr = sum_kin_energy_sqr + kin_energy**2
  sum_kin_energy = sum_kin_energy + kin_energy
end subroutine
```

Listing 5: Updating the relevant sums of the kinetic energy

As discussed earlier, the specific heat is related to the fluctuations in kinetic energy using a formula derived by Lebowitz[2]:

$$\frac{\langle \delta K^2 \rangle}{\langle K \rangle^2} = \frac{2}{3N} \left( 1 - \frac{3N}{2C_V} \right)$$

Rewriting this to get an expression for the specific heat:

$$C_{V} = \left(\frac{2}{3N} - \frac{\langle \delta K^{2} \rangle}{\langle K \rangle^{2}}\right)^{-1} \tag{3}$$

$$= \left(\frac{2}{3N} - \frac{\langle K^2 \rangle - \langle K \rangle^2}{\langle K \rangle^2}\right)^{-1} \tag{4}$$

Using the time average instead of the ensemble average, the averages can be expressed as a sum over all time steps  $n_t$  divided by the time steps:

$$\langle K^2 \rangle = \frac{\sum_{i=1}^{n_t} K(i)^2}{n_t}$$

$$\langle K \rangle^2 = \left(\frac{\sum_{i=1}^{n_t} K(i)}{n_t}\right)^2$$

Now the specific heat can be calculated as:

$$C_V = \left(\frac{2}{3N} - \frac{\langle K^2 \rangle - \langle K \rangle^2}{\langle K \rangle^2}\right)^{-1} \tag{5}$$

$$= \left(\frac{2}{3N} - \frac{\sum_{i=1}^{n_t} K(i)^2 * n_t - \left(\sum_{i=1}^{n_t} K(i)\right)^2}{\left(\sum_{i=1}^{n_t} K(i)\right)^2}\right)^{-1}$$
(6)

This is implemented in our code within the same subroutine "calc\_specific\_heat" after an if statement is enabled at the end of the simulation:

```
subroutine calc_specific_heat(end_of_routine,N_part, kin_energy,
    sum_kin_energy, sum_kin_energy_sqr, step)
  if (end_of_routine .eqv. .true.) then
   specific_heat = ((2d0/(3d0*N_part)) - (((sum_kin_energy_sqr * step) -
        sum_kin_energy**2) / (sum_kin_energy**2)))**(-1)
   print *, "The specific heat is ", specific_heat
end subroutine
```

Listing 6: Calculating the specific heat

## Pair Correlation Function

The pair correlation function can be evaluated in molecular dynamics systems if a histogram is recorded of the number of pairs of particles n(r) for each distance r:

$$g(r) = \frac{2V}{N(N-1)} \left[ \frac{\langle n(r) \rangle}{4\pi r^2 \Delta r} \right]$$

In our simulation a number of intervals n<sub>i</sub> is defined which together with the length in each direction L sets the small  $\Delta r$ :

$$\Delta r = \frac{L}{n_i}$$

Again a time average is used and the average number of pairs  $\langle n(r) \rangle$  is calculated as a sum of all the pair occurences N(r,t) in a time step t and in a distance r over all time steps  $n_t$ :

$$\langle n(r) \rangle = \frac{\sum_{t}^{n_t} N(r,t)}{n_t}$$

For a large enough number of intervals the distance r can be approximated by  $i\Delta r$  with i the index of the corresponding interval. The number of pair occurences N(r,t) in a time step t and in a distance r can be expressed now as function of the interval you are in N(i, t). This way we are able to count all pair occurences in our program and calculate the pair correlation function:

$$g(r) \approx \frac{2V}{N(N-1)} \left[ \frac{\langle n(r) \rangle}{4\pi r^2 \Delta r} \right]$$
 (7)

$$=\frac{2L^3}{N(N-1)}\left[\frac{\sum_{t}^{n_t}N(i,t)}{n_t4\pi i^2\Delta r^2\Delta r}\right] \tag{8}$$

$$= \frac{2L^3}{N(N-1)} \frac{\sum_{t}^{n_t} N(i,t)}{n_t 4\pi i^2} \frac{n_i^3}{L^3}$$
 (9)

$$= \frac{2}{N(N-1)} \frac{n_i^3}{n_t 4\pi} \frac{\sum_t^{n_t} N(i,t)}{i^2} \tag{10}$$

This is implemented in the code as follows. For each interval i the number of pair occurences in that interval is registered in a particular time interval. This is done while evaluating the forces of the particles:

```
subroutine calc_dynamics(calc_quant, N_part, L_side, time_step, m, e, s,
     r\_cut, pos, kin\_energy, pot\_energy,&
                                & virial, vel, num_intervals,
                                      histogram_{-}vector)
  histogram_vector = 0
 do n = 1, N_part
   do\ i = 1, N_part ! integrate over all particles inside box except <math>i = n
      do j = -1, 1
      do k = -1, 1 !periodic boundary condition
     do l = -1, 1
        if (n/=i) then
          r_{-}vec = \dots !Calculation of the difference vector
          r = sqrt(dot_product(r_vec, r_vec))
          ! histogram for the pair correlation function
          if ((n > i) .and. (calc_quant .eqv. .true.)) then
            hist_i = 1 + floor(r/delta_r_hist)
            if (hist_i < num_intervals + 1) then ! defines a cut off</pre>
              histogram_vector(hist_i) = histogram_vector(hist_i) + 1
            end if
          end if
       end if
      end do
      end do
      end do
    end do
 end do
end subroutine
```

**Listing 7**: Updating the histogram for each interval

As the final time step is done and the routine for the dynamics is finished, the sum for each interval is passed to a routine for calculating the pair correlation function that is called at the end:

```
subroutine write_histogram_file(average_number, num_intervals, N_part, step
 temp_factor = 2d0 * num_intervals / N_part
  temp_factor2 = 1d0 * num_intervals / (N_part - 1)
 temp_factor3 = 1d0 * num_intervals / step
 constant_factor = (temp_factor * temp_factor2 * temp_factor3) / ( 4 * abs(
      atan(1d0)) * 4)
 open (unit=6, file="histogram.dat",action="write")
 do i=1,num_intervals
    write (6,"(I3, 4F18.6)") i, (constant_factor * average_number(i) )/ (i
        **2)
 end do
end subroutine
```

Listing 8: Factors are needed to prevent overflow errors

#### RESULTS AND DISCUSSION 4

- 4.1 Pressure
- 4.2 Temperature
- 4.3 Specific Heat

Some text with a citation [1] The other citation [3] and another [4]

- 4.4 Pair Correlation function
- 4.5 Phase Transitions

# REFERENCES

- [1] L. Verlet. 'Computer "Experiments" on Classical Fluids. I. Thermodynamical Properties of Lennard-Jones Molecules'. Physical Review, 159, 1967.
- [2] S. Duane. 'Stochastic quantization versus the microcanonical ensemble – getting the best of both worlds'. Nucl. Phys., 275:398–420, 1985.
- [3] Glosser C. 'ICCP Coding Manual'. 2015.
- [4] J. Thijssen. 'Computational Physics'. 2013.

## MAIN FORTRAN SOURCE CODE

```
!argon gas in a box simulation, molecular dynamics.
!the cubic geometry sides of length = L
!initial positions initialized according to fcc lattice structure
!number of fcc cells per cartesian dimension = Ncell,
!number of particles is N, (4 particles per cube)
!velocity verlet method: v' = v+1/2*F(x)/m*dt; x = x+v'*dt; v = v'+1/2*F/m*dt
! (coverted from initially an implementation of the semi iplicit euler
    method)
!time evolution for particles in lennard jones potential: U = 4*e*((s/r)
    **12-(s/r)**6),
!Fij = -du/dx = -du/dr*dr/dx = e*(48*s**12/r**13 - 6*s**6/r**7) * x/r,
!r = sqrt(x**2+y**2+z**2)
program argon_box
 use argon_box_init
 use argon_box_dynamics
 use argon_box_results
! use md_plot
 implicit none
  integer, parameter :: N_cell_dim = 6, velocity_rescale_steps = 50,
       equilibration_steps = 120
  real(8), parameter :: dt = 0.004_8, T_initial = 1d0, rho = 0.88_8, t_stop
      = 1d0
  integer, parameter :: N_cell = N_cell_dim**3, N_part = N_cell*4
  real(8), parameter :: L_side = (N_part/rho)**(1._8/3)
  real(8), parameter :: s = 1d0, e = 1d0, r_cut = 5d-1*L_side ! lennard
       jones potential
  real(8), parameter :: m = 1d0, Kb = 1d0 !mass and boltzman constant
  integer, parameter :: hist_num_intervals = 500
  integer, dimension(1:hist_num_intervals) :: histogram_vector,
       tot_histogram_vector = 0
  real(8), dimension(1:3, 1:N_part) :: pos, vel
  integer :: step = 0
  real(8) :: time = 0, kin_energy, pot_energy, virial
  real(8) :: Pressure, Temperature, tot_energy
  real(8) :: sum_kin_energy_sqr = 0, sum_kin_energy = 0, sum_virial = 0,
       sum_potential_energy = 0
  ! Create initial state
  call cubic_fcc_lattice(N_cell_dim, L_side, pos)
  call init_random_seed
  call init_vel(T_initial, Kb, m, N_part, vel)
! call plot_init(0d0, L_side,0d0, L_side,0d0, L_side)
  do while (time < t_stop)</pre>
   time = time + dt
    step = step + 1
    !velocity verlet integration method, .true. triggers the calculation of
         thermodynamic quantities.
    call calc_dynamics(.false., N_part, L_side, dt, m, e, s, r_cut, pos,
        kin_energy, pot_energy, &
                                   & virial, vel, hist_num_intervals,
                                        histogram_vector)
```

```
call new_pos(N_part, L_side, dt, vel, pos)
        call calc_dynamics(.true., N_part, L_side, dt, m, e, s, r_cut, pos,
            kin_energy, pot_energy, &
                                       & virial, vel, hist_num_intervals,
                                            histogram_vector)
        !Temperature control
        if (step < velocity_rescale_steps) then</pre>
          call rescale_vel(T_initial, kin_energy, Kb, N_part, vel)
       end if
       tot_energy = pot_energy + kin_energy
       Temperature = 2*kin_energy/(3* (N_part-1) *Kb) !Center of mass degrees
65
            of freedom substracted..
       Pressure = 1 + 1/(3*Kb*Temperature*N_part)* virial !P/(Kb T rho) + TODO:
             correction cuttoff
        tot_energy = tot_energy/N_part
        pot_energy = pot_energy/N_part
        kin_energy = kin_energy/N_part
       call plot_points(pos)
        print *, step, "t=", time, "H=", tot_energy, "K=", kin_energy, "U=",
            pot_energy, "virial=", virial, &
             & "T=", Temperature, "P=", Pressure
        !print *, histogram_vector
        if (equilibration_steps < step) then</pre>
         tot_histogram_vector = tot_histogram_vector + histogram_vector
         sum_virial = sum_virial + virial
          sum_potential_energy = sum_potential_energy + pot_energy
          ! variables sum_kin_energy and sum_kin_energy_sqr calculated in
              subroutine calc_specific_heat
          call calc_specific_heat(.false., N_part, kin_energy, sum_kin_energy,
              sum_kin_energy_sqr, step - equilibration_steps)
       end if
       call write_energy_file(kin_energy, pot_energy, virial, time, step)
     end do
    ! call plot_end
      pot_energy = sum_potential_energy/(step - equilibration_steps)
      kin_energy = sum_kin_energy/(step - equilibration_steps)
      tot_energy = kin_energy + pot_energy
     virial = sum_virial/(step - equilibration_steps)
     Temperature = 2*kin_energy*N_part/(3* (N_part-1) *Kb) !Center of mass
          degrees of freedom substracted ...
     Pressure = 1 + 1/(3*Kb*Temperature*N_part)* virial !P/(Kb T rho) + TODO:
          correction cuttoff
     print *, "H=", tot_energy, "K=", kin_energy, "U=", pot_energy, "virial=",
          virial, "T=", Temperature, "P=", Pressure
      call calc_specific_heat(.true., N_part, kin_energy, sum_kin_energy,
           sum_kin_energy_sqr, step - equilibration_steps)
      call write_histogram_file(tot_histogram_vector, hist_num_intervals, N_part
          , step - equilibration_steps)
   end program
```

Listing 9: argon\_box.f90

## FORTRAN SOURCE CODE FOR SETTING THE INI-В TIAL CONDITIONS

```
module argon_box_init
 implicit none
  private
 public cubic_fcc_lattice, init_vel, init_random_seed
contains
 function fcc_cell(i) result(output)
    implicit none
   integer, intent(in) :: i
   real(8) :: output(3)
    ! face centered cubic unit cell with basis particle positions:
   real(8), dimension(1:3), parameter :: &
    fcc_part1 = (/0d0, 0d0, 0d0/), &
   fcc_part2 = (/0d0, 5d-1, 5d-1/), &
   fcc_part3 = (/5d-1, 0d0, 5d-1/), &
   fcc_part4 = (/5d-1, 5d-1, 0d0/)
    real(8), dimension(1:3,1:4), parameter :: &
   R_cell = reshape( (/fcc_part1, fcc_part2, fcc_part3, fcc_part4/),
        (/3,4/))
   output = R_cell(:,i)
    ! print *, "face centered cubic unit cell"
    ! do i = 1,3 ! print *, (R_cell(i,j), j=1,4)
   ! end do
 end function fcc_cell
 subroutine cubic_fcc_lattice(N_cell_dim, L_side, pos)
    !initial positions of all particles according to an fcc lattice
        structure
   integer :: i,j,k,l,n
   integer, intent(in) :: N_cell_dim
   real(8), intent(in) :: L_side
   real(8), intent(out), dimension(1:3, 1:(4*N_cell_dim**3)) :: pos
   n = 0
   do i = 1,N_cell_dim
     do j = 1,N_cell_dim
       do k = 1,N_cell_dim
         do l = 1,4
           n = n + 1
           pos(:,n) = L_side/N_cell_dim*((/ i-1, j-1, k-1 /) + fcc_cell(l))
                 print *, "particle:", n, "/", N_part, "position:", pos(:,
         end do
       end do
     end do
   end do
 end subroutine
  subroutine init_vel(T, Kb, m, N_part, vel)
    !initial particles velocities according to the maxwell distribution
    ! in the maxwell boltzman distribution each velocity component is
        normally distributed:
    ! Box muller transform used for converting uniform dist to normal dist
   ! f(v) = sqrt(m/(2*PI*Kb*T)) * exp(-(v**2)/2 *m/(*Kb*T))
    ! sigma**2 = Kb*T/m, and zero mean
   real(8), intent(in) :: T, Kb, m
   integer, intent(in) :: N_part
   real(8), intent(out), dimension(1:3, 1:N_part) :: vel
   real(8), parameter :: pi = 4*atan(1d0)
```

```
real(8) :: xs(2) !two random numbers
   integer :: n, i
   do n = 1, N_part
     do i = 1,3
       CALL RANDOM_NUMBER(xs(1))
       CALL RANDOM_NUMBER(xs(2))
       vel(i,n) = sqrt(Kb*T/m) * sqrt(-2d0*log(xs(1)))*cos(2*pi*xs(2)) !
            sigma * box_muller
     end do
   end do
   !Set center of mass velocity to zero
   do i = 1,3
     vel(i,:) = vel(i,:) - sum(vel(i,:))/N_part
   end do
 end subroutine
  ! copied from ICCP coding-notes
  subroutine init_random_seed()
   implicit none
   integer, allocatable :: seed(:)
   integer :: i, n, un, istat, dt(8), pid, t(2), s
   integer(8) :: count, tms
    call random_seed(size = n)
    allocate(seed(n))
   open(newunit=un, file="/dev/urandom", access="stream",&
   form="unformatted", action="read", status="old", &
    iostat=istat)
   if (istat == 0) then
      read(un) seed
      close (un)
    else
      call system_clock(count)
      if (count /= 0) then
       t = transfer(count, t)
      else
        call date_and_time(values=dt)
       tms = (dt(1) - 1970)*365_8 * 24 * 60 * 60 * 1000 &
       + dt(2) * 31_8 * 24 * 60 * 60 * 1000 &
       + dt(3) * 24 * 60 * 60 * 60 * 1000 &
       + dt(5) * 60 * 60 * 1000 &
       + dt(6) * 60 * 1000 + dt(7) * 1000 &
       + dt(8)
       t = transfer(tms, t)
     end if
      s = ieor(t(1), t(2))
     pid = getpid() + 1099279 ! Add a prime
      s = ieor(s, pid)
      if (n >= 3) then
       seed(1) = t(1) + 36269
        seed(2) = t(2) + 72551
        seed(3) = pid
        if (n > 3) then
         seed(4:) = s + 37 * (/ (i, i = 0, n - 4) /)
       end if
      else
       seed = s + 37 * (/ (i, i = 0, n - 1) /)
     end if
   end if
    call random_seed(put=seed)
 end subroutine init_random_seed
end module
```

Listing 10: argon\_box\_init.f90

## FORTRAN SOURCE CODE FOR THE DYNAMICS OF C THE SYSTEM

```
module argon_box_dynamics
     implicit none
     private
     public calc_dynamics, rescale_vel, new_pos
    contains
     subroutine calc_dynamics(calc_quant, N_part, L_side, time_step, m, e, s,
           r_cut, pos, kin_energy, pot_energy,&
                                    & virial, vel, num_intervals,
                                         histogram_vector)
       logical, intent(in) :: calc_quant !for improving efficiency with
             velocity verlet method
       integer, intent(in) :: N_part, num_intervals
        real(8), intent(in) :: e, s, r_cut !Lennard Jones
        real(8), intent(in) :: m, time_step, L_side
        real(8), intent(inout), dimension(1:3, 1:N_part) :: vel
        real(8), intent(in), dimension(1:3, 1:N_part) :: pos
        integer :: i,j,k,l,n
        real(8), intent(out) :: kin_energy, pot_energy, virial
        real(8) :: sum_v_2, F(3), dF(3), r, r_vec(3)
        integer, intent(out), dimension(1:num_intervals) :: histogram_vector
       integer :: hist_i
        real(8) :: delta_r_hist
       delta_r_hist = L_side / num_intervals
       virial = 0
       pot_energy = 0
        sum_{-}v_{-}2 = 0
       histogram_vector = 0
       do n = 1, N_part
         do\ i = 1, N\_part ! integrate over all particles inside box except <math>i = n
            do j = -1, 1
            do \ k = -1, 1 !periodic boundary condition
           do l = -1, 1
              if (n/=i) then
                r_{\text{vec}} = (/pos(1,n) - pos(1,i), pos(2,n) - pos(2,i), pos(3,n) - pos(3,i)
                     )/) + L_side*(/j,k,l/)
                r = sqrt(dot_product(r_vec, r_vec))
                !force calculation
                if (r<r_cut) then
                  dF = e*(48*s**12/r**14 - 24*s**6/r**8) * r_vec
                  F = F + dF
                  ! calculate other quantities:
45
                  if (calc_quant .eqv. .true.) then
                    if (n > i) then
                      pot\_energy = pot\_energy + 4*e*((s/r)**12-(s/r)**6)
                      virial = virial + dot_product(r_vec, dF)
                    end if
                  end if
                end if
                ! histogram for the pair correlation function
                if ((n > i) .and. (calc_quant .eqv. .true.)) then
                  hist_i = 1 + floor(r/delta_r_hist)
                  if (hist_i < num_intervals + 1) then ! defines a cut off</pre>
                    histogram_vector(hist_i) = histogram_vector(hist_i) + 1
                  !else
```

```
! histogram_vector(num_intervals) = histogram_vector(
                  num_intervals) + 1
             end if
           end if
         end if
       end do
       end do
       end do
     end do
     vel(:,n) = vel(:,n) + F/m*time_step/2 ! velocity verlet method ->
          factor 1/2 !
     if (calc_quant .eqv. .true.) then
       sum_v_2 = sum_v_2 + dot_product(vel(:,n),vel(:,n))
     end if
   end do
   kin\_energy = m/2*sum\_v\_2
 end subroutine
 subroutine rescale_vel(T_intended, kin_energy, Kb, N_part, Vel)
   !rescale velocities in order to move system to desired temperature
   real(8), intent(in) :: T_intended, kin_energy, kb
   integer, intent(in) :: N_part
   real(8), intent(inout), dimension(1:3, 1:N_part) :: vel
   real(8) :: scaling_factor
   scaling_factor = sqrt((N_part - 1)*3/2*kb*T_intended/kin_energy)
   vel = scaling_factor*vel
 end subroutine
 subroutine new_pos(N_part, L_side, dt, vel, pos)
   ! Postion calculation
   real(8), intent(in) :: L_side, dt
   integer, intent(in) :: N_part
   real(8), intent(in), dimension(1:3, 1:N_part) :: vel
   real(8), intent(inout), dimension(1:3, 1:N_part) :: pos
   integer :: n, i
   do n = 1, N_part
     pos(:,n) = pos(:,n) + vel(:,n)*dt
     do i = 1,3 !implements periodic boundary conditions
       if (pos(i,n) < 0d0) then
         pos(i,n) = pos(i,n) + L_side
       else if (pos(i,n) > L_side) then
         pos(i,n) = pos(i,n) - L_side
       end if
     end do
   end do
 end subroutine
end module
```

Listing 11: argon\_box\_dynamics.f90

## FORTRAN SOURCE CODE FOR OUTPUT OF RE-D **SULTS**

```
module argon_box_results
            implicit none
                private
            public calc_specific_heat, write_energy_file, write_histogram_file
        contains
            subroutine write_energy_file(kin_energy, pot_energy, virial, time, cnt)
                real(8), intent(in) :: virial, kin_energy, time, pot_energy
                integer, intent(in) :: cnt
                open (unit=1, file="energy_matrix.dat", action="write")
                write (1,"(I6, 4F18.6)") cnt, time, kin_energy, pot_energy, virial
           end subroutine
            subroutine write_histogram_file(average_number, num_intervals, N_part,
                      step )
                integer, intent(in) :: num_intervals, N_part, step
                integer, intent(in), dimension(1:num_intervals) :: average_number
                real(8) :: constant_factor, temp_factor, temp_factor2, temp_factor3
                integer :: i
                temp_factor = 2d0 * num_intervals / N_part
                temp_factor2 = 1d0 * num_intervals / (N_part - 1)
                temp_factor3 = 1d0 * num_intervals / step
                {\tt constant\_factor = (temp\_factor * temp\_factor2 * temp\_factor3) / (~4 *}
                           abs(atan(1d0)) * 4)
                open (unit=6, file="histogram.dat", action="write")
                do i=1,num_intervals
                     write (6,"(I3, 4F18.6)") i, (constant_factor * average_number(i) )/ (
                end do
            end subroutine
            subroutine calc_specific_heat(end_of_routine,N_part, kin_energy,
                      sum_kin_energy, sum_kin_energy_sqr, step)
                logical, intent(in) :: end_of_routine
                real(8), intent(in) :: kin_energy
                real(8), intent(out) :: sum_kin_energy, sum_kin_energy_sqr
                integer, intent(in) :: N_part, step
                real(8) :: specific_heat
                sum_kin_energy_sqr = sum_kin_energy_sqr + kin_energy**2
                sum_kin_energy = sum_kin_energy + kin_energy
45
                if (end_of_routine .eqv. .true.) then
                     specific_heat = ((2d0/(3d0*N_part)) - (((sum_kin_energy_sqr * step) - ((sum_kin_energy_sqr * step) - ((sum_kin_energy_sqr * step) - (sum_kin_energy_sqr * 
                               sum_kin_energy**2) / (sum_kin_energy**2)))**(-1)
                    print *, "The specific heat is ", specific_heat
                end if
            end subroutine
       end module
```

Listing 12: argon\_box\_results.f90