

Hyperboloidal neutron star and black hole in spherical symmetry

Alex Vano-Vinuales



IAC3, University of the Balearic Islands

Virtual Infinity Seminar - 16th May 2025

Based on [2505.XXXXXX](#) [gr-qc]



Bondi accretion

Accretion onto a small black hole at the center of a neutron star

Chloe B. Richards,¹ Thomas W. Baumgarte,¹ and Stuart L. Shapiro^{2,3}

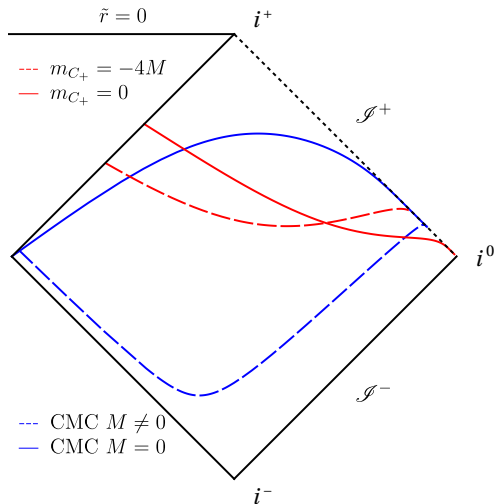
¹*Department of Physics and Astronomy, Bowdoin College, Brunswick, ME 04011*

²*Department of Physics, University of Illinois at Urbana-Champaign, Urbana, IL 61801*

³*Department of Astronomy and NCSA, University of Illinois at Urbana-Champaign, Urbana, IL 61801*

2102.09574 [astro-ph.HE]

Hyperboloidal slice to know mass of system



Peterson et al. *Phys. Rev. D* 110 (2024)

Setup on Cauchy slices

Line element:

$$d\tilde{s}^2 = -e^{\nu(\tilde{r})} d\tilde{t}^2 + e^{\lambda(\tilde{r})} d\tilde{r}^2 + \tilde{r}^2 d\sigma^2$$

Setup on Cauchy slices

Line element:

$$d\tilde{s}^2 = -e^{\nu(\tilde{r})} d\tilde{t}^2 + e^{\lambda(\tilde{r})} d\tilde{r}^2 + \tilde{r}^2 d\sigma^2$$

TOV equations (assuming perfect fluid):

$$\partial_{\tilde{r}} m(\tilde{r}) = 4\pi \tilde{r}^2 \rho(\tilde{r}) \quad \rightarrow \quad e^{\lambda(\tilde{r})} = \left(1 - \frac{2m(\tilde{r})}{\tilde{r}}\right)^{-1}$$

$$\partial_{\tilde{r}} \nu(\tilde{r}) = 2 \left(\frac{m(\tilde{r})}{\tilde{r}^2} + 4\pi \tilde{r} P(\tilde{r}) \right) \left(1 - \frac{2m(\tilde{r})}{\tilde{r}}\right)^{-1}$$

$$\partial_{\tilde{r}} P(\tilde{r}) = -\frac{1}{2} (P(\tilde{r}) + \rho(\tilde{r})) \partial_{\tilde{r}} \nu(\tilde{r})$$

Setup on Cauchy slices

Line element:

$$d\tilde{s}^2 = -e^{\nu(\tilde{r})} d\tilde{t}^2 + e^{\lambda(\tilde{r})} d\tilde{r}^2 + \tilde{r}^2 d\sigma^2$$

TOV equations (assuming perfect fluid):

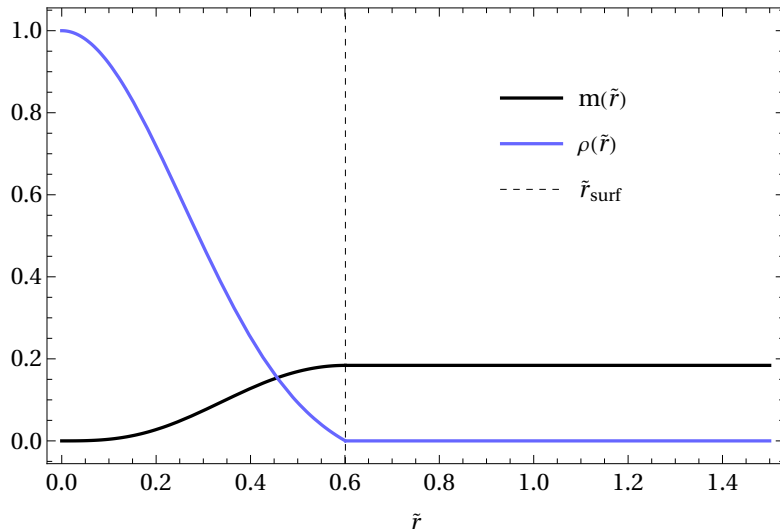
$$\partial_{\tilde{r}} m(\tilde{r}) = 4\pi \tilde{r}^2 \rho(\tilde{r}) \quad \rightarrow \quad e^{\lambda(\tilde{r})} = \left(1 - \frac{2m(\tilde{r})}{\tilde{r}}\right)^{-1}$$

$$\partial_{\tilde{r}} \nu(\tilde{r}) = 2 \left(\frac{m(\tilde{r})}{\tilde{r}^2} + 4\pi \tilde{r} P(\tilde{r}) \right) \left(1 - \frac{2m(\tilde{r})}{\tilde{r}}\right)^{-1}$$

$$\partial_{\tilde{r}} P(\tilde{r}) = -\frac{1}{2} (P(\tilde{r}) + \rho(\tilde{r})) \partial_{\tilde{r}} \nu(\tilde{r})$$

Polytropic equation of state: $P(\tilde{r}) = K[\rho(\tilde{r})]^\Gamma$

Neutron star profiles



Choice of hyperboloidal slice

Hyperboloidal time: $\tilde{t} = t + h(\tilde{r})$

Choice of hyperboloidal slice

Hyperboloidal time: $\tilde{t} = t + h(\tilde{r})$

Constant-mean curvature slices: $\tilde{K} = -\frac{1}{\sqrt{-\tilde{g}}} \partial_a (\sqrt{-\tilde{g}} \tilde{n}^a) \equiv \text{constant}$

↓

$$\begin{aligned} \frac{e^\nu h'}{\sqrt{e^\lambda - e^\nu (h')^2}} &= -\frac{1}{\tilde{r}^2} \left[\int K_{\text{CMC}} \tilde{r}^2 \sqrt{e^\lambda e^\nu} d\tilde{r} + C_{\text{CMC}} \right] \\ &\equiv \text{int}(\tilde{r}), \end{aligned}$$

Choice of hyperboloidal slice

Hyperboloidal time: $\tilde{t} = t + h(\tilde{r})$

Constant-mean curvature slices: $\tilde{K} = -\frac{1}{\sqrt{-\tilde{g}}} \partial_a (\sqrt{-\tilde{g}} \tilde{n}^a) \equiv \text{constant}$

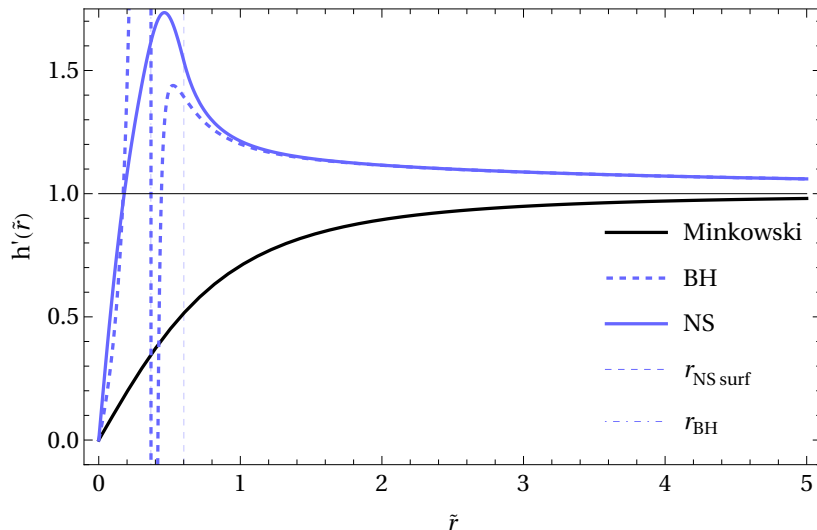
↓

$$\begin{aligned} \frac{e^\nu h'}{\sqrt{e^\lambda - e^\nu (h')^2}} &= -\frac{1}{\tilde{r}^2} \left[\int K_{\text{CMC}} \tilde{r}^2 \sqrt{e^\lambda e^\nu} d\tilde{r} + C_{\text{CMC}} \right] \\ &\equiv \text{int}(\tilde{r}), \end{aligned}$$

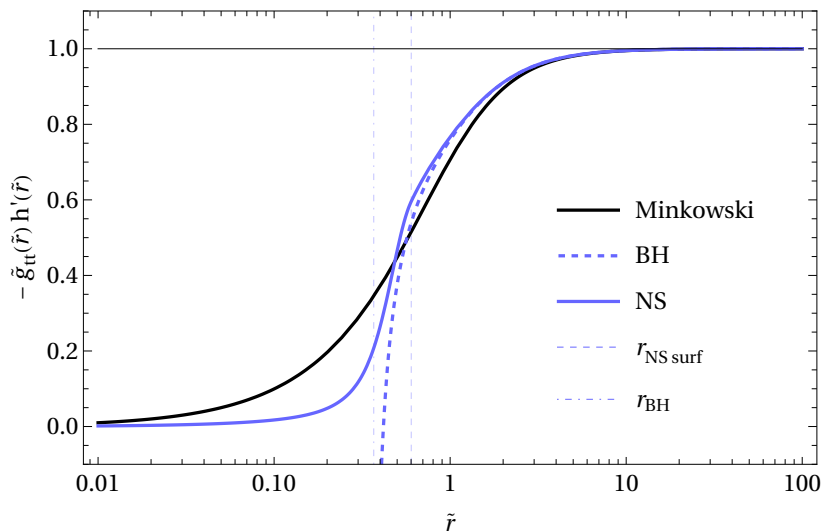
Boost:

$$h'(\tilde{r}) = \pm \text{int}(\tilde{r}) \sqrt{\frac{e^{\lambda(\tilde{r})}}{e^{\nu(\tilde{r})} (e^{\nu(\tilde{r})} + [\text{int}(\tilde{r})]^2)}},$$

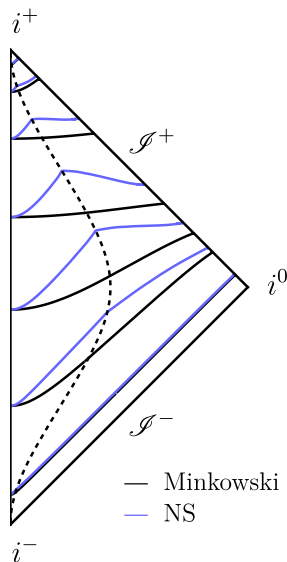
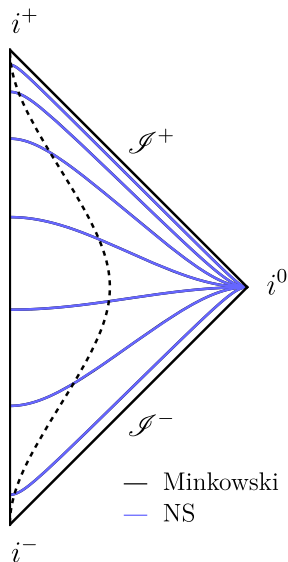
Boost function



Rescaled boost function



Penrose diagrams



Compactification by imposing conformal flatness

Compactification: $\tilde{r} = \frac{r}{\Omega(r)}$

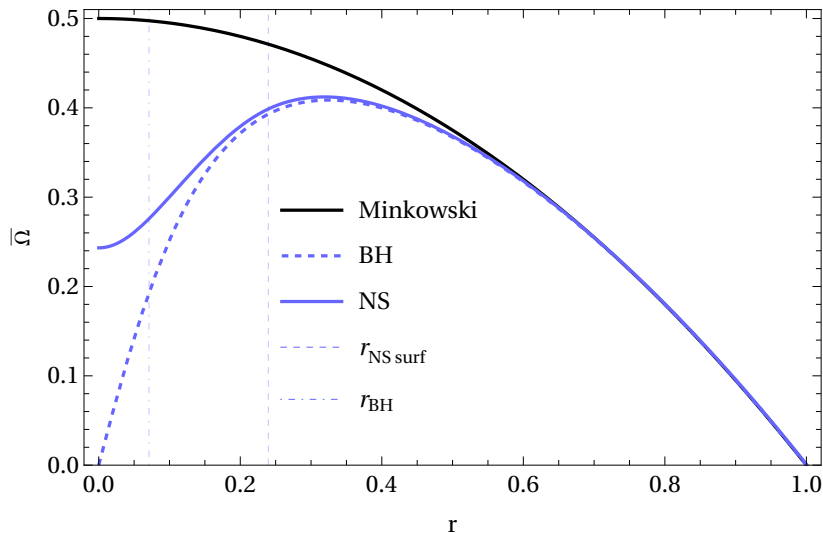
Compactification by imposing conformal flatness

Compactification: $\tilde{r} = \frac{r}{\bar{\Omega}(r)}$

Impose:

$$\gamma_{rr} = \left[\left(1 - \frac{2m(\frac{r}{\bar{\Omega}})\bar{\Omega}}{r} \right)^{-1} - e^{\nu(\frac{r}{\bar{\Omega}})} \left(h'(\frac{r}{\bar{\Omega}}) \right)^2 \right] \left(\frac{\bar{\Omega} - r\bar{\Omega}'}{\bar{\Omega}} \right)^2 = 1$$

Compactification factors



Procedure

- ① Solve TOV for NS on a Cauchy slice
- ② Express it in isotropic form
- ③ Add BH in isotropic form
- ④ Solve Hamiltonian constraint
- ⑤ Hyperboloidalize
- ⑥ Compactify

Transforming to isotropic radius

Want to add contributions as

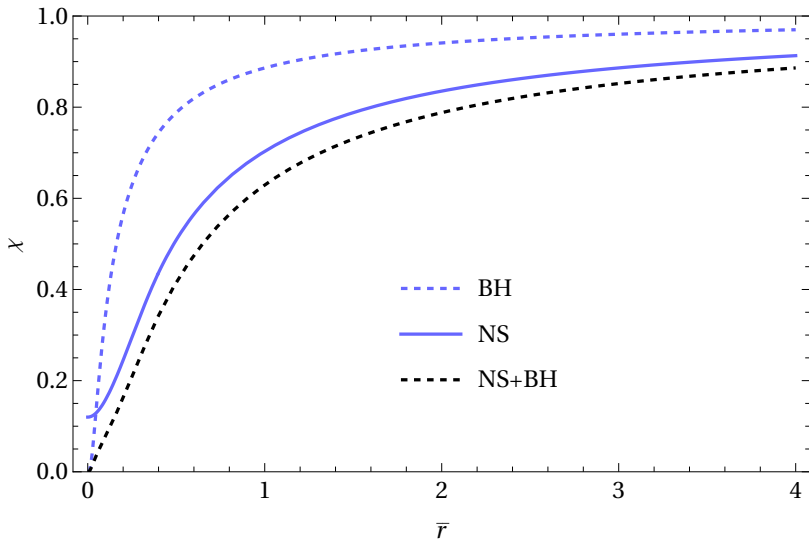
$$\psi = \psi_{\text{NS}} + \psi_{\text{BH}} + \delta\psi$$

with

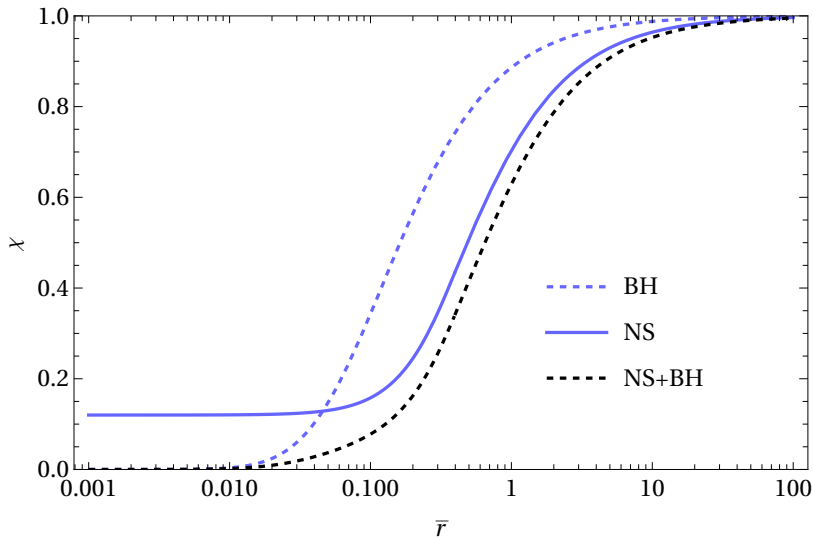
$$d\tilde{s}^2 = -f d\tilde{t}^2 + \psi^4 (d\bar{r}^2 + \bar{r}^2 d\sigma^2)$$

- ψ_{NS} determined from numerical solution
- $\psi_{\text{BH}} = \frac{m_{\text{BH}}}{2\bar{r}}$
- $\delta\psi$ to solve the Hamiltonian constraint for

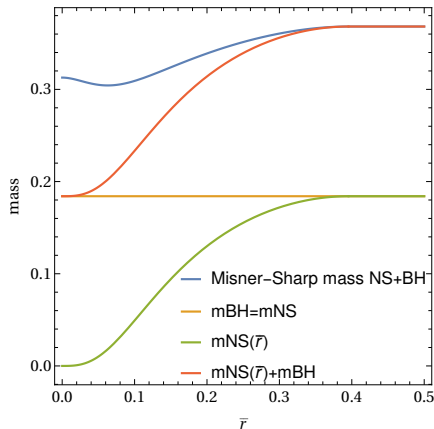
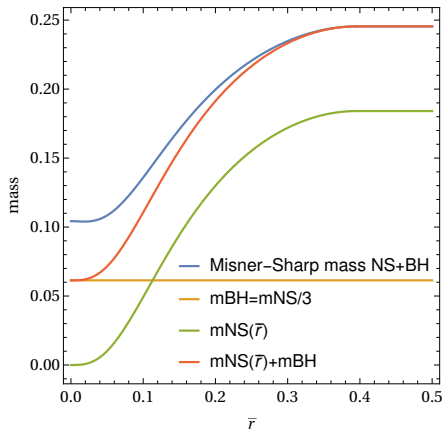
Addition of conformal factors



Addition of conformal factors (logscale)



Misner-Sharp mass



Left: $m_{\text{BH}} = m_{\text{NS}}/3$, right: $m_{\text{BH}} = m_{\text{NS}}$.

Setup

Following [3]

$$\Delta\psi = -2\pi\psi^5\rho \quad \text{using} \quad \rho = \psi^m\bar{\rho}$$

Set

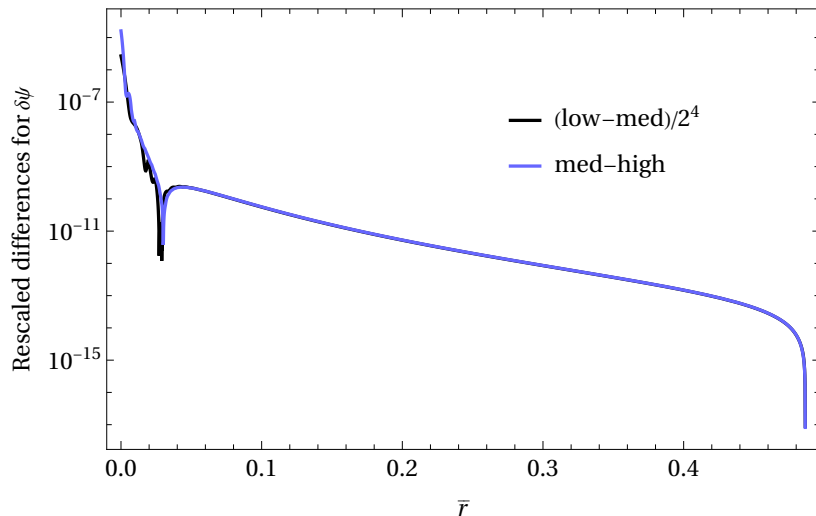
$$\Delta\psi_{\text{NS}} = -2\pi\psi_{\text{NS}}^5\rho_{\text{NS}}, \quad \Delta\psi_{\text{BH}} = 0, \quad \bar{\rho} = \psi_{\text{NS}}^{-m}\rho_{\text{NS}}$$

Solve with $m = -6$

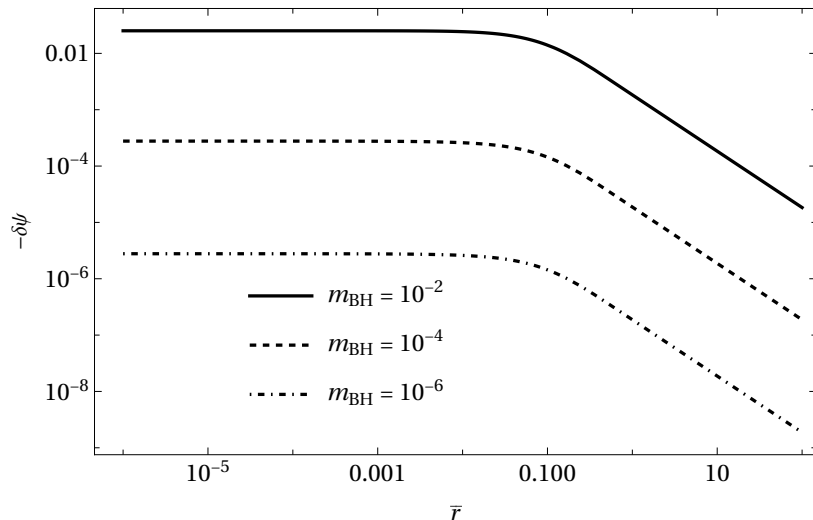
$$\begin{aligned} \Delta\delta\psi &= 2\pi\rho_{\text{NS}} \left(\psi_{\text{NS}}^5 - \frac{(\psi_{\text{NS}} + \psi_{\text{BH}} + \delta\psi)^{5+m}}{\psi_{\text{NS}}^m} \right), \\ \partial_{\bar{r}}^2\delta\psi + \frac{2}{\bar{r}}\partial_{\bar{r}}\delta\psi &= 2\pi\rho_{\text{NS}} \left(\psi_{\text{NS}}^5 - \frac{\psi_{\text{NS}}^6}{(\psi_{\text{NS}} + \psi_{\text{BH}} + \delta\psi)} \right), \end{aligned}$$

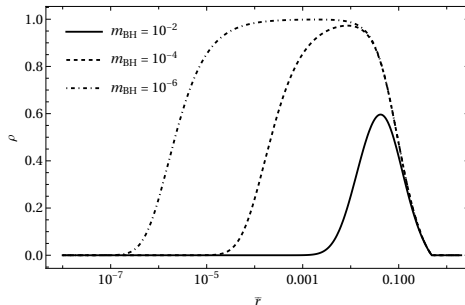
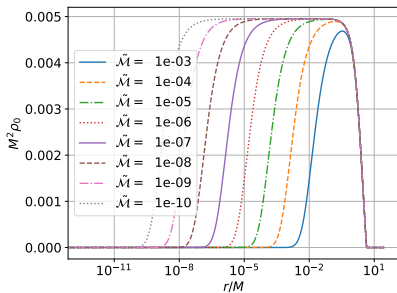
assuming $\delta\psi \sim A/\bar{r}$ at NS's surface.

Convergence



$\delta\psi$ solutions for small BH





Left: Richards, Baumgarte and Shapiro. *Phys. Rev. D* 103.10 (2021), right: this work.

Radial transformation

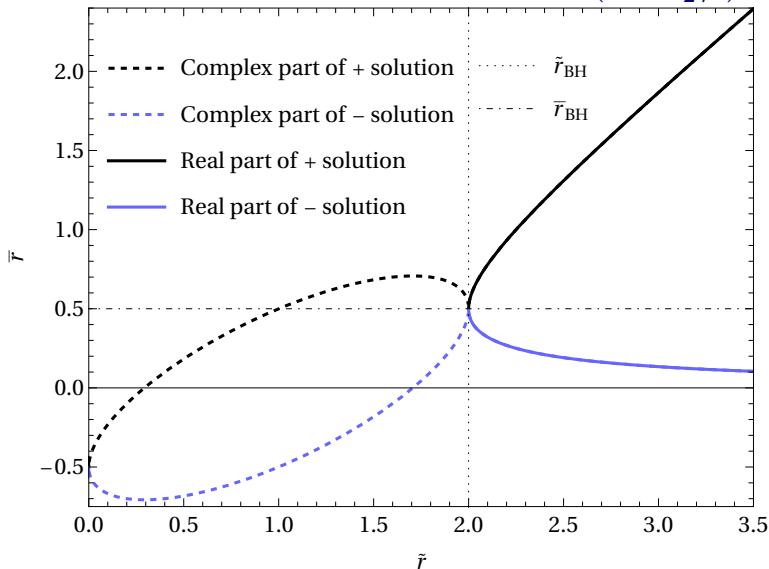
$$d\tilde{t}^2 = \psi^4 (d\bar{r}^2 + \bar{r}^2 d\sigma^2) = g_{\tilde{r}\tilde{r}} d\tilde{r}^2 + \tilde{r}^2 d\sigma^2,$$

Radial transformation

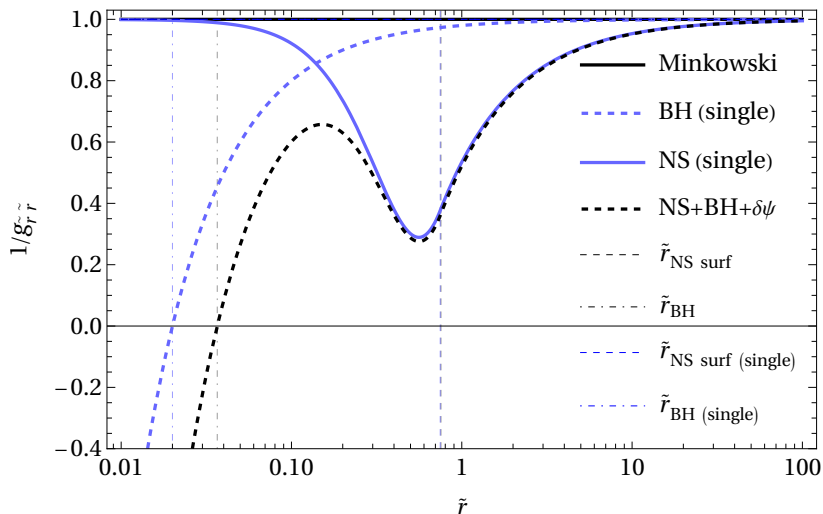
$$d\tilde{t}^2 = \psi^4 (d\bar{r}^2 + \bar{r}^2 d\sigma^2) = g_{\tilde{r}\tilde{r}} d\tilde{r}^2 + \tilde{r}^2 d\sigma^2,$$

Difficulty: $g_{\tilde{r}\tilde{r}}$ changes sign at the horizon.

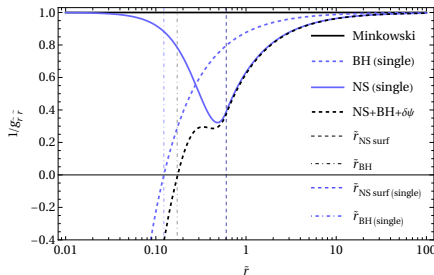
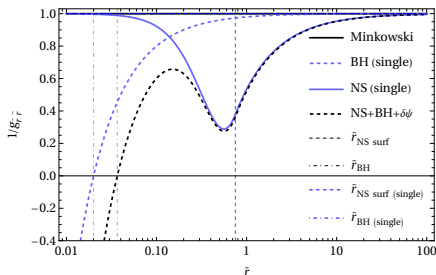
Outside/inside of horizon of BH – $\tilde{r} = \left(1 + \frac{m_{\text{BH}}}{2\tilde{r}}\right)^2 \bar{r}$



Metric deformation



Comparison of effect of BH masses



Left: $m_{\text{BH}} = 10^{-2}$, right: $m_{\text{BH}} = m_{\text{NS}}/3$.

Determine boost

Choose

$$g_{\tilde{t}\tilde{t}} \doteq -\frac{1}{g_{\tilde{r}\tilde{r}}} \equiv -f(\tilde{r})$$

Determine boost

Choose

$$g_{\tilde{t}\tilde{t}} \doteq -\frac{1}{g_{\tilde{r}\tilde{r}}} \equiv -f(\tilde{r})$$

Introduce in

$$h'(\tilde{r}) = -\frac{\left(\frac{K_{\text{CMC}}\tilde{r}}{3} + \frac{C_{\text{CMC}}}{\tilde{r}^3}\right)}{f(\tilde{r})\sqrt{f(\tilde{r}) + \left(\frac{K_{\text{CMC}}\tilde{r}}{3} + \frac{C_{\text{CMC}}}{\tilde{r}^3}\right)^2}}$$

Determine boost

Choose

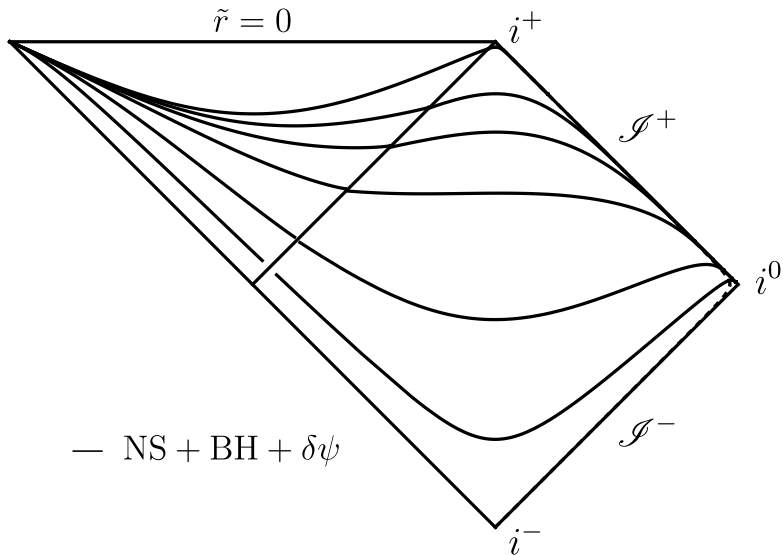
$$g_{\tilde{t}\tilde{t}} \doteq -\frac{1}{g_{\tilde{r}\tilde{r}}} \equiv -f(\tilde{r})$$

Introduce in

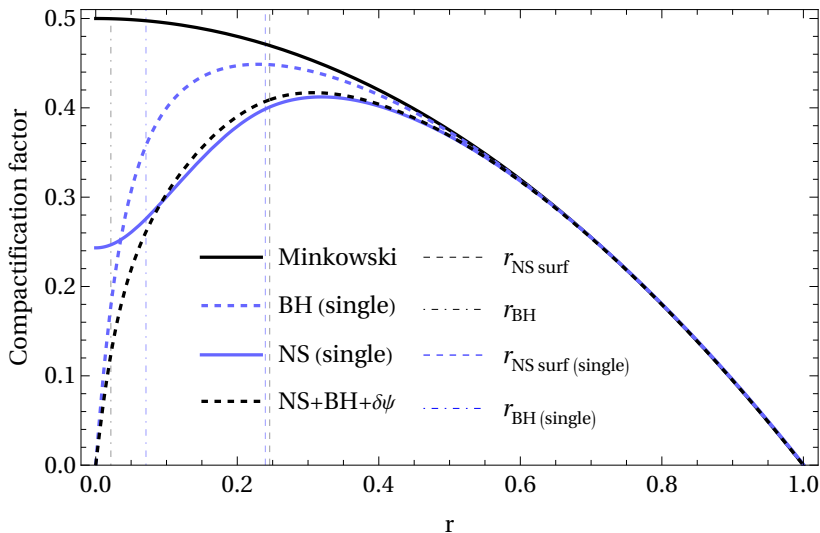
$$h'(\tilde{r}) = -\frac{\left(\frac{K_{\text{CMC}}\tilde{r}}{3} + \frac{C_{\text{CMC}}}{\tilde{r}^3}\right)}{f(\tilde{r})\sqrt{f(\tilde{r}) + \left(\frac{K_{\text{CMC}}\tilde{r}}{3} + \frac{C_{\text{CMC}}}{\tilde{r}^3}\right)^2}}$$

Tune value of C_{CMC} for boost to diverge at trumpet.

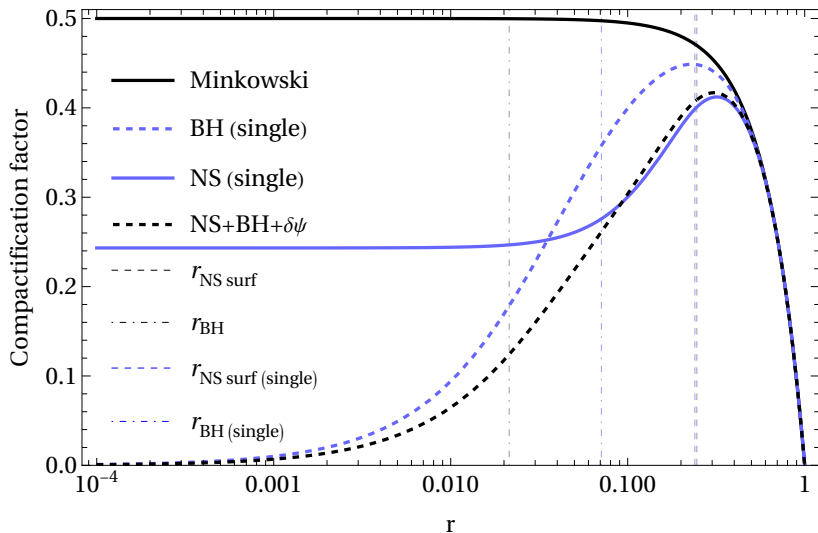
Penrose diagram



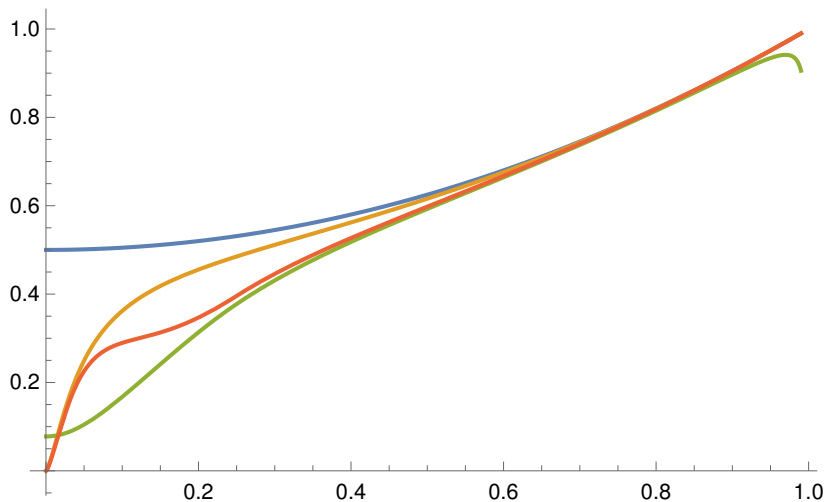
Compactification factor



Compactification factor – log scale



Ready as initial data for evolutions



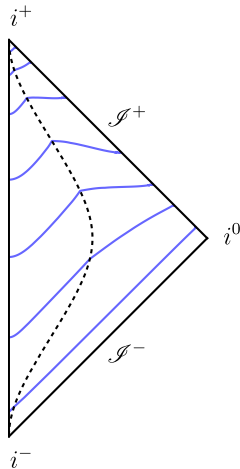
Future plans

In spherical symmetry:

- Evolve NS initial data (Einstein + relativistic Euler)
- Evolve perturbed NS
- Bondi accretion: evolve NS + small BH initial data

Beyond spherical symmetry:

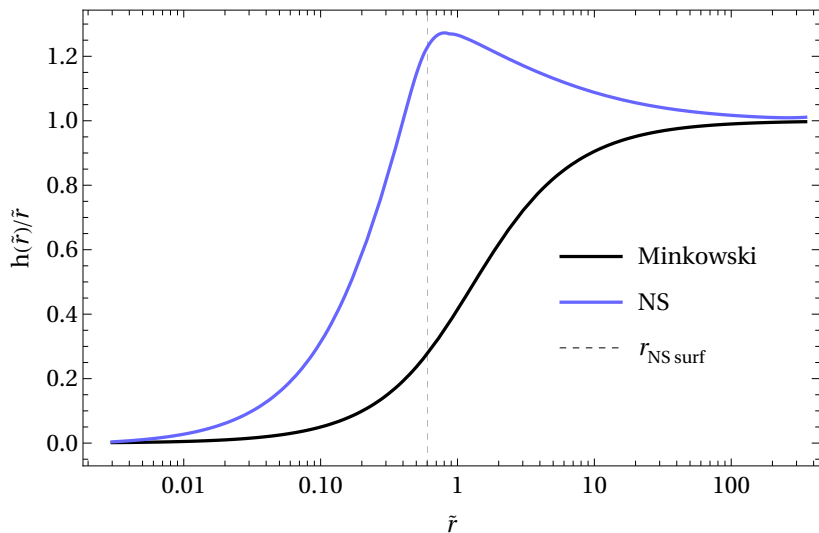
- Hyperboloidalize superposition of bodies in different locations
- **Evolve 3D perturbed NS**
- ...



Thanks for listening! Questions?

Backup slides

Integration of the height function



Tortoise-like coordinate

Express metric as: $d\tilde{s}^2 = \Xi^2 (-d\tilde{t}^2 + d\tilde{r}_*^2) \equiv -\Xi^2 d\tilde{u} d\tilde{v}$

For NS: $d\tilde{s}^2 = e^{\nu(\tilde{r})} (-d\tilde{t}^2 + d\tilde{r}_*^2)$ with $d\tilde{r}_* = \sqrt{\frac{e^{\lambda(\tilde{r})}}{e^{\nu(\tilde{r})}}} d\tilde{r}$

Tortoise-like coordinate

Express metric as: $d\tilde{s}^2 = \Xi^2 (-d\tilde{t}^2 + d\tilde{r}_*^2) \equiv -\Xi^2 d\tilde{u} d\tilde{v}$

For NS: $d\tilde{s}^2 = e^{\nu(\tilde{r})} (-d\tilde{t}^2 + d\tilde{r}_*^2)$ with $d\tilde{r}_* = \sqrt{\frac{e^{\lambda(\tilde{r})}}{e^{\nu(\tilde{r})}}} d\tilde{r}$

Usual compactification along null directions:

$$\begin{aligned}\tilde{U} &= \tilde{t} - \tilde{r}_*, & \tilde{V} &= \tilde{t} + \tilde{r}_*, \\ U &= \arctan \tilde{U}, & V &= \arctan \tilde{V}, \\ T &= \frac{V + U}{2}, & R &= \frac{V - U}{2}.\end{aligned}$$

Tortoise-like coordinate $d\tilde{r}_* = \frac{1}{f(\tilde{r})}d\tilde{r}$

