

Notes on 10d Super Yang-Mills

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1 Minimal Super Yang-Mills

Our starting point will be the theory complexifying the usual 10d super Yang-Mills theory. Fix a complex reductive gauge group G with Lie algebra \mathfrak{g} . The ordinary fields of super Yang-Mills theory on \mathbb{R}^{10} consist of a boson: a connection A on the trivial G -bundle, and a fermion: a \mathfrak{g} -valued section λ of the Weyl spinor bundle associated to the spinor representation S_+^{-1} . These fields are acted upon by the group of gauge transformations – G -valued functions on \mathbb{R}^{10} .

We can model the stack of fields modulo gauge transformations infinitesimally near the point 0 by the corresponding BRST complex. This is the local super Lie algebra

$$L_{\text{BRST}} = \Omega^0(\mathbb{R}^{10}; \mathfrak{g}) \rightarrow \Omega^1(\mathbb{R}^{10}; \mathfrak{g}) \oplus \Omega^0(\mathbb{R}^{10}; \Pi S_+ \otimes \mathfrak{g})$$

with the de Rham differential, placed in cohomological degrees 0 and 1, with bracket induced from the Lie bracket on \mathfrak{g} .

The action functional in 10d super Yang-Mills is given by

$$S(A, \lambda) = \int_{\mathbb{R}^{10}} \langle F_A \wedge *F_A + (\lambda, \not{D}_A \lambda) \rangle,$$

where $\langle - \rangle$ denotes an invariant pairing on \mathfrak{g} , and $(,)$ denotes a scalar-valued pairing $S_+ \otimes S_- \rightarrow \mathbb{C}$ (there will be a unique such pairing, up to rescaling, characterized by the condition that $(\rho(v)\lambda_1, \rho(v)\lambda_2) = (\lambda_1, \lambda_2)$ for each $v \in \mathbb{C}^{10}$, where ρ denotes Clifford multiplication).

We can re-encode this data in terms of the classical BV complex (Phil and I wrote this down in [EY18, Section 3.1]). This is the L_∞ -algebra whose underlying cochain complex takes the form

$$\Omega^0(\mathbb{R}^{10}; \mathfrak{g}) \xrightarrow{d} \Omega^1(\mathbb{R}^{10}; \mathfrak{g}) \xrightarrow{d \circ d} \Omega^2(\mathbb{R}^{10}; \mathfrak{g}) \xrightarrow{d} \Omega^3(\mathbb{R}^{10}; \mathfrak{g}) \xrightarrow{d} \Omega^4(\mathbb{R}^{10}; \mathfrak{g})$$

$$\Omega^0(\mathbb{R}^{10}; \Pi S_- \otimes \mathfrak{g}) \xrightarrow{*d} \Omega^1(\mathbb{R}^{10}; \Pi S_- \otimes \mathfrak{g}),$$

with degree -3 invariant pairing induced by the invariant pairing on \mathfrak{g} and the pairing $(,)$ between S_+ and S_- , and

¹If we didn't complexify we would instead consider $G_{\mathbb{R}}$ a compact connected Lie group, and a section of the Majorana-Weyl spinor bundle, which necessitates working in Lorentzian signature. I think for our purposes it's interesting enough to just consider the complexified theory and avoid signature issues. The complexified theory twists to holomorphic Chern-Simons theory with complex gauge group.

with degree 2 and 3 brackets given by the action of $\Omega^0(\mathbb{R}^{10}; \mathfrak{g})$ on everything along with

$$\begin{aligned} \ell_2^{\text{Bos}}: \Omega^1(\mathbb{R}^{10}; \mathfrak{g}) \otimes \Omega^1(\mathbb{R}^{10}; \mathfrak{g}) &\rightarrow \Omega^9(\mathbb{R}^{10}; \mathfrak{g}) \\ (A \otimes B) &\mapsto [A \wedge *dB] + [*dA \wedge B] + d * [A \wedge B] \\ \ell_2^{\text{Fer}}: \Omega^1(\mathbb{R}^{10}; \mathfrak{g}) \otimes \Omega^0(\mathbb{R}^{10}; S_+ \otimes \mathfrak{g}) &\rightarrow \Omega^{10}(\mathbb{R}^{10}; S_- \otimes \mathfrak{g}) \\ (A \otimes \lambda) &\mapsto *A\lambda \end{aligned}$$

in degree 2, and the map

$$\begin{aligned} \ell_3: \Omega^1(\mathbb{R}^{10}; \mathfrak{g}) \otimes \Omega^1(\mathbb{R}^{10}; \mathfrak{g}) \otimes \Omega^1(\mathbb{R}^{10}; \mathfrak{g}) &\rightarrow \Omega^9(\mathbb{R}^{10}; \mathfrak{g}) \\ (A \otimes B \otimes C) &\mapsto [A \wedge *[B \wedge C]] + [B \wedge *[C \wedge A]] + [C \wedge *[A \wedge B]] \end{aligned}$$

in degree 3.

1.1 On-Shell Supersymmetry Action

We can define an action of the 10d $\mathcal{N} = 1$ supersymmetry algebra on the complex of BRST fields in this minimal super Yang-Mills theory. The Poincaré action is clear, and the action of the supersymmetry Q is generated – in the usual Physics notation – by the transformation [BSS77]

$$\begin{aligned} \delta_Q A &= \Gamma(Q, \lambda) \\ \delta_Q \lambda &= \frac{1}{4} \not{F}_A Q, \end{aligned}$$

where the notation \not{F}_A stands for the iterated Clifford multiplication $\not{F}_A = F^{ij} \gamma_i \gamma_j$. To check that this defines a supersymmetry action we need to check it's compatible with the brackets in the supersymmetry algebra. So, we calculate

$$\begin{aligned} [\delta_{Q_1}, \delta_{Q_2}] A &= \frac{1}{4} (\Gamma(Q_2, \not{F}_A Q_1) - \Gamma(Q_1, \not{F}_A Q_2)) \\ &= \\ \text{and } [\delta_{Q_1}, \delta_{Q_2}] \lambda &= \frac{1}{4} (\not{F}_{\Gamma(Q_2, \lambda)} Q_1 - \not{F}_{\Gamma(Q_1, \lambda)} Q_2) \\ &= \end{aligned}$$

2 Baulieu's 10 Model

2.1 Off-Shell Action of a Scalar Supersymmetry

3 Homotopy Data

3.1 Homotopy Transfer of the Scalar Supersymmetry

References

- [BSS77] L. Brink, J. Schwarz, and J. Scherk. “Supersymmetric Yang-Mills theories”. *Nuclear Phys. B* 121.1 (1977), pp. 77–92. URL: [https://doi.org/10.1016/0550-3213\(77\)90328-5](https://doi.org/10.1016/0550-3213(77)90328-5).
- [EY18] C. Elliott and P. Yoo. “Geometric Langlands Twists of $N = 4$ Gauge Theory from Derived Algebraic Geometry”. *Advances in Theoretical and Mathematical Physics* 22.3 (2018). arXiv: 1507.03048 [math-ph].