Studies of Quantum Dots using Machine Learning



Even Marius Nordhagen

University of Oslo evenmn@fys.uio.no

December 13, 2019

Outline

- Motivation
- Quantum Theory
- ► Machine Learning Theory
- ▶ Methods
- Software
- Results
- ▶ Conclusion



Motivation





Studies of Quantum Dots using Machine Learning





Studies of Quantum Dots using Machine Learning





Machine Learning

Machine learning is the science of getting computers to act without being explicitly programmed¹.





Machine Learning

Machine learning is the science of getting computers to act without being explicitly programmed¹.

► Image recognition



Machine Learning

- Machine learning is the science of getting computers to act without being explicitly programmed¹.
- ► Image recognition
- ▶ Voice commands



Machine Learning + Quantum Mechanics

▶ Neural networks are eminent function approximators





Machine Learning + Quantum Mechanics

▶ Neural networks are eminent function approximators

$$\Psi$$
 = \Rightarrow \wedge





Machine Learning + Quantum Mechanics

▶ Neural networks are eminent function approximators

$$\Psi$$
 = \Rightarrow \wedge

▶ Existing methods are reminiscent of machine learning algorithms





Studies of Quantum Dots using Machine Learning





Studies of Quantum Dots using Machine Learning





Quantum Dots

► Technologically: Quantum dots are expected to be the next big thing in display technology^{2,3}.





Quantum Dots

► Technologically: Quantum dots are expected to be the next big thing in display technology^{2,3}.

Experimentally: Researchers have managed to study two-dimensional quantum dots in the laboratory⁴.





Quantum Dots

► Technologically: Quantum dots are expected to be the next big thing in display technology 2,3 .

► Experimentally: Researchers have managed to study two-dimensional quantum dots in the laboratory⁴.

► Physically: An array of interesting physical phenomena can be observed in quantum dots.





Ethics in Science

► Respect for other's work





Ethics in Science

- ► Respect for other's work
- ► Reproducibility





Quantum Theory





The Schrödinger Equation

 $\hat{\mathcal{H}}\Psi = E\Psi$





The Schrödinger Equation

$$\hat{\mathcal{H}}\Psi = E\Psi$$

$$\Downarrow$$

$$E = \frac{\int dX \Psi^*(X) \hat{\mathcal{H}} \Psi(X)}{\int dX \Psi^*(X) \Psi(X)}$$





The Variational Principle

The variational principle serves as a way of finding the ground state energy. For an arbitrary trial wave function $\Psi_T(X)$, it states that the obtained energy is larger or equal to the ground state,

$$E_0 \leq E = \frac{\int d\mathbf{X} \Psi_T^*(\mathbf{X}) \hat{\mathcal{H}} \Psi_T(\mathbf{X})}{\int d\mathbf{X} \Psi_T^*(\mathbf{X}) \Psi_T(\mathbf{X})}.$$

Thus, by minimizing the obtained energy, *E*, we can estimate the ground state energy.

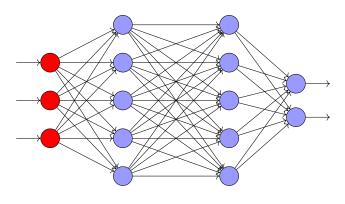




Machine Learning Theory



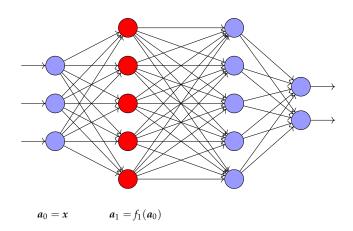




$$a_0 = x$$

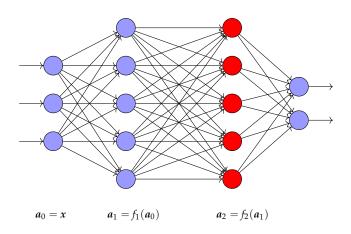






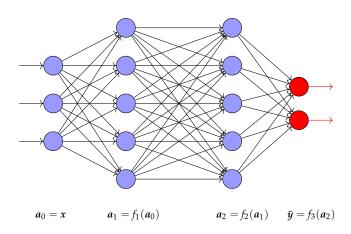
















Cost function

➤ The cost function defines the error





Cost function

- ► The cost function defines the error
- ► Mean square error (MSE):

$$C = \frac{1}{2} \sum_{i=1}^{n} (y - \tilde{y})^2.$$





Cost function

- ▶ The cost function defines the error
- ► Mean square error (MSE):

$$C = \frac{1}{2} \sum_{i=1}^{n} (y - \tilde{y})^2.$$

► Attempt to minimize the cost function



Minimize the cost function





- Minimize the cost function
- ▶ The gradient descent method:

$$\theta^+ = \theta - \frac{\partial \mathcal{C}}{\partial \theta}.$$





- ▶ Minimize the cost function
- ▶ The gradient descent method:

$$\theta^+ = \theta - \frac{\partial \mathcal{C}}{\partial \theta}.$$

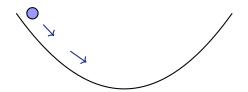






- Minimize the cost function
- ▶ The gradient descent method:

$$\theta^+ = \theta - \frac{\partial \mathcal{C}}{\partial \theta}.$$

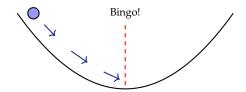






- Minimize the cost function
- ▶ The gradient descent method:

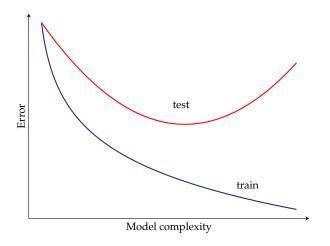
$$\theta^+ = \theta - \frac{\partial \mathcal{C}}{\partial \theta}.$$







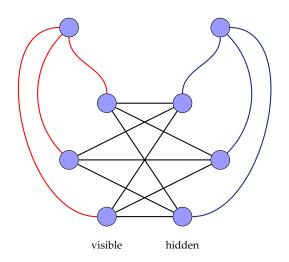
Find Appropriate Complexity







Restricted Boltzmann Machines









Methods





Variational Monte Carlo (VMC)

Exploit the variational principle in order to obtain the ground state energy

$$E_0 < E_{\text{VMC}} = \frac{\int d\mathbf{R} \Psi_T(\mathbf{R})^* \hat{\mathcal{H}} \Psi_T(\mathbf{R})}{\int d\mathbf{R} \Psi_T(\mathbf{R})^* \Psi_T(\mathbf{R})},$$
$$= \int d\mathbf{R} E_L(\mathbf{R}) P(\mathbf{R}),$$

with

$$E_L(\mathbf{R}) = \frac{1}{\Psi_T(\mathbf{R})} \hat{\mathcal{H}} \Psi_T(\mathbf{R}) \quad \wedge \quad P(\mathbf{R}) = \frac{\Psi_T(\mathbf{R})^* \Psi_T(\mathbf{R})}{\int d\mathbf{R} \Psi_T(\mathbf{R})^* \Psi_T(\mathbf{R})}$$





Monte Carlo Integration

We attempt to solve the integral by sampling from the probability density function $P(R) \propto \Psi_T(R)^* \Psi_T(R)$:

$$\begin{split} E_{\text{VMC}} &= \int d\mathbf{R} E_L(\mathbf{R}) P(\mathbf{R}), \\ &\approx \frac{1}{M} \sum_{i=1}^M E_L(\mathbf{R}_i). \end{split}$$



Trial Wave Function Ansatz

The Slater-Jastrow function is the *de facto* standard trial wave function for electronic structure systems,

$$\Psi_T(\mathbf{R}) = |\hat{D}(\mathbf{R})|J(\mathbf{R}),$$

where the Slater matrix,

$$\hat{D}(\mathbf{R}) = \begin{pmatrix} \phi_1(\mathbf{r}_1) & \phi_2(\mathbf{r}_1) & \dots & \phi_N(\mathbf{r}_1) \\ \phi_1(\mathbf{r}_2) & \phi_2(\mathbf{r}_2) & \dots & \phi_N(\mathbf{r}_2) \\ \vdots & \vdots & \ddots & \vdots \\ \phi_1(\mathbf{r}_N) & \phi_2(\mathbf{r}_N) & \dots & \phi_N(\mathbf{r}_N) \end{pmatrix},$$

contains all the single-particle functions.





Single-particle Functions

The Hermite functions,

$$\phi_n(\mathbf{r}) \propto H_n(\sqrt{\omega}\mathbf{r}) \exp\left(-\frac{1}{2}\alpha\omega|\mathbf{r}|^2\right),$$

are often used as the single-particle functions for quantum dots. The Gaussian can be factorized out from the Slater determinant,

$$|\hat{D}(\boldsymbol{R};\alpha)| \propto \exp\left(-\frac{1}{2}\alpha\omega|\boldsymbol{R}|^2\right) \begin{vmatrix} H_1(\boldsymbol{r}_1) & H_2(\boldsymbol{r}_1) & \dots & H_N(\boldsymbol{r}_1) \\ H_1(\boldsymbol{r}_2) & H_2(\boldsymbol{r}_2) & \dots & H_N(\boldsymbol{r}_2) \\ \vdots & \vdots & \ddots & \vdots \\ H_1(\boldsymbol{r}_N) & H_2(\boldsymbol{r}_N) & \dots & H_N(\boldsymbol{r}_N) \end{vmatrix}.$$



Restricted Boltzmann Machine

As suggested by Carleo & Troyer⁵, we use the marginal distribution of the visible units as the single-particle functions in the Slater determinant, and see if them can model the correlations

$$\phi_n(\mathbf{r}) \propto H_n(\sqrt{\omega}\mathbf{r})P(\mathbf{r};\boldsymbol{\theta})$$

where $P(\mathbf{r})$ is the marginal distribution of the visible units.

$$|\hat{D}(r;\theta)| \propto P(r;\theta) \begin{vmatrix} H_1(r_1) & H_2(r_1) & \dots & H_N(r_1) \\ H_1(r_2) & H_2(r_2) & \dots & H_N(r_2) \\ \vdots & \vdots & \ddots & \vdots \\ H_1(r_N) & H_2(r_N) & \dots & H_N(r_N) \end{vmatrix}$$



Jastrow Factor

The Jastrow factor is added to account for the correlations Simple Jastrow factor

$$J(\mathbf{r}; \boldsymbol{\beta}) = \exp\left(\sum_{i=1}^{N} \sum_{j>i}^{N} \beta_{ij} r_{ij}\right).$$

Padé-Jastrow factor

$$J(r;\beta) = \exp\left(\sum_{i=1}^{N} \sum_{j>i}^{N} \frac{a_{ij}r_{ij}}{1 + \beta r_{ij}}\right).$$





Our Trial Wave Function Ansätze

$$\blacktriangleright \ \Psi_{\text{RBM}}(\mathbf{R}) = |\hat{D}_{\text{RBM}}(\mathbf{R})|$$

$$\qquad \Psi_{\text{RBM+SJ}}(\mathbf{R}) = |\hat{D}_{\text{RBM}}(\mathbf{R})|J(\mathbf{R};\boldsymbol{\beta})$$

$$\qquad \Psi_{\text{RBM+PJ}}(\mathbf{R}) = |\hat{D}_{\text{RBM}}(\mathbf{R})|J(\mathbf{R};\beta)$$

$$\qquad \qquad \Psi_{\text{VMC}}(\mathbf{R}) = |\hat{D}_{\text{Gauss}}(\mathbf{R})|J(\mathbf{R};\beta)$$





Software





Aims

- ▶ Efficient
- ▶ Flexible
- ► Easy to use





Results





Ground State Energy

Number of electrons: N = 2. Frequency: ω .

ω	RBM	RBM+SJ	RBM+PJ	VMC	HF *	Exact †
1/6	0.7036(1)	0.67684(7)	0.66715(6)	0.66710(1)	0.768675	2/3
0.28 1	1.07050(4) 3.0803(2)	1.03470(7) 3.02108(5)	1.021668(7) 2.999587(5)	1.02192(1) 2.99936(1)	1.14171 3.16190	3

†Semi-analytical ground state energy calculated by Taut, 1993 [7].







^{*}Computation of the Hartree-Fock limit by Mariadason, 2018 [6].

Ground State Energy

Number of electrons: N = 20. Frequency: ω .

ω	RBM	RBM+SJ	RBM+PJ	VMC	HF [‡]	DMC §
0.1	30.824(2)	30.567(3)	30.1553(9)	30.0403(2)	31.1902	29.9779(1)
1.0	159.428(3)	156.816(4)	156.104(1)	155.8900(4)	158.004	155.8822(1)

[§]Ground state energy estimate using the diffusion Monte Carlo method. By Høgberget, 2013 [8].



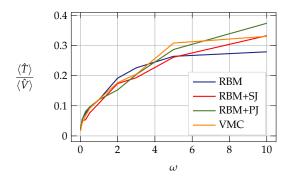


[‡]Computation of the Hartree-Fock limit by Mariadason, 2018 [6].

Energy distribution

Number of electrons: N = 20. Frequency: ω .

Ratio between the kinetic energy, $\langle \hat{T} \rangle$, and the total potential energy, $\langle \hat{V} \rangle$.

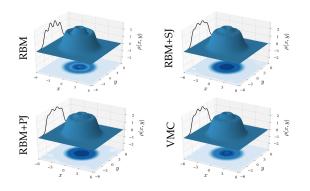






One-body density

Number of electrons: N = 20. Frequency: $\omega = 1.0$.

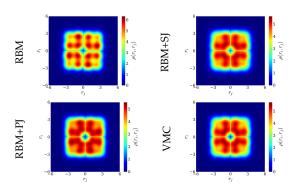






Two-body density

Number of electrons: N = 20. Frequency: $\omega = 1.0$.

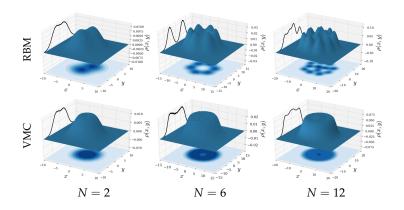






Low-frequency dots

Number of electrons: *N*. Frequency: $\omega = 0.1$.

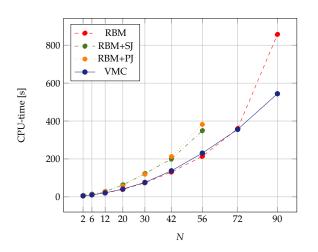






Computational Cost

Number of electrons: *N*.







Conclusions





Conclusions

- ▶ The RBM ansatz is able to account for most of the correlations
- The RBM+PJ ansatz might give a better ground state estimate of small quantum dots, compared to the traditional VMC ansatz
- For larger quantum dots, RBM+PJ gives slightly larger energy than VMC
- The energy distribution is different for the different methods, indicating different electron configurations.
- The ground state energy might not be the best way to evaluate the various ansätze





Future Work

- ▶ Investigate restricted Boltzmann machines with other architectures
- ▶ Pass more information to the restricted Boltzmann machine
- ► Apply the method on more complex systems





Thank you!





References

- Machine Learning Stanford University. http://mlclass.stanford.edu/.
- Samsung QLED TV | The Frame | Premium UHD TV. http://www.samsung.com/global/tv/blog/why-are-quantum-dot-displays-so-good/.
- Manders, J. R. et al. 8.3: Distinguished Paper: Next-Generation Display Technology: Quantum-Dot LEDs. SID Symposium Digest of Technical Papers 46, 73 (2015).
- Brunner, K., Abstreiter, G., Böhm, G., Tränkle, G. & Weimann, G. Sharp-Line Photoluminescence and Two-Photon Absorption of Zero-Dimensional Biexcitons in a GaAs/AlGaAs Structure. *Physical Review Letters* 73, 1138 (1994).
- Carleo, G. & Troyer, M. Solving the Quantum Many-Body Problem with Artificial Neural Networks. Science 355, 602 (2017).
- 6. Mariadason, A. A. Quantum Many-Body Simulations of Double Dot System. MA thesis (2018).
- Taut, M. Two electrons in an external oscillator potential: Particular analytic solutions of a Coulomb correlation problem. *Physical Review A* 48, 3561 (1993).
- 8. Høgberget, J. Quantum Monte-Carlo Studies of Generalized Many-body Systems. MA thesis (2013).





Test

Test

TEST





Test

Test

TEST





Test

Test

TEST



