# Studies of Quantum Dots using Machine Learning



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#### Outline

- ▶ Motivation
- ▶ Quantum Theory
- ► Machine Learning Theory
- Methods
- Results
- Conclusion
- ► (Code)





# Motivation





#### Machine Learning

Studies of Quantum Dots using Machine Learning

Machine learning is the science of getting computers to act without being explicitly programmed<sup>1</sup>.





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Image recognition





## **Machine Learning**

Studies of Quantum Dots using Machine Learning

Machine learning is the science of getting computers to act without being explicitly programmed<sup>1</sup>.

- Image recognition
- Voice commands





#### Machine Learning + Quantum Mechanics

Studies of Quantum Dots using Machine Learning

▶ Neural networks are eminent function approximators





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$$\Psi$$
 =  $\Rightarrow$   $\bigwedge$ 





#### Quantum Dots

Studies of Quantum Dots using Machine Learning

► Technologically: Quantum dots are expected to be the next big thing in display technology<sup>2,3</sup>.





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► Experimentally:
Researchers have managed to study two-dimensional quantum dots in the laboratory<sup>4</sup>.





#### **Ouantum Dots**

Studies of Quantum Dots using Machine Learning

► Technologically: Quantum dots are expected to be the next big thing in display technology $^{2,3}$ .

Experimentally: Researchers have managed to study two-dimensional quantum dots in the laboratory<sup>4</sup>.

► Physically: An array of interesting physical phenomena can be observed in quantum dots.





# Quantum Theory





# The Schrödinger Equation

$$\hat{\mathcal{H}}\Psi = E\Psi$$





# The Schrödinger Equation

$$\hat{\mathcal{H}}\Psi = E\Psi$$

$$\Downarrow$$

$$E = \frac{\int dX \Psi^*(X) \hat{\mathcal{H}} \Psi(X)}{\int dX \Psi^*(X) \Psi(X)}$$





# The Variational Principle

The variational principle serves as a way of finding the ground state energy. For an arbitrary trial wave function  $\Psi_T(X)$ , it states that the obtained energy is larger or equal to the ground state,

$$E_0 \leq E = \frac{\int d\mathbf{X} \Psi_T^*(\mathbf{X}) \hat{\mathcal{H}} \Psi_T(\mathbf{X})}{\int d\mathbf{X} \Psi_T^*(\mathbf{X}) \Psi_T(\mathbf{X})}.$$

Thus, by minimizing the obtained energy, *E*, we can estimate the ground state energy.



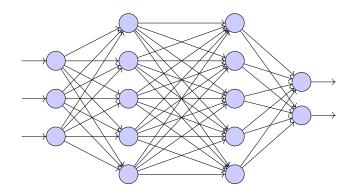


# Machine Learning Theory





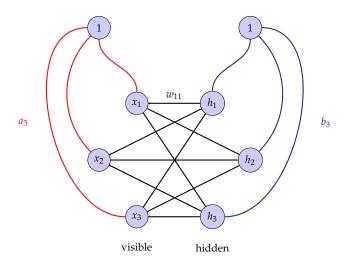
## Feed-forward Neural Network (FNN)







#### Restricted Boltzmann machines







# Methods





#### Variational Monte Carlo (VMC)

Exploit the variational principle in order to obtain the ground state energy

$$E_{0} < E_{VMC} = \frac{\int d\mathbf{R} \Psi_{T}(\mathbf{R})^{*} \hat{\mathcal{H}} \Psi_{T}(\mathbf{R})}{\int d\mathbf{R} \Psi_{T}(\mathbf{R})^{*} \Psi_{T}(\mathbf{R})}$$

$$= \int d\mathbf{R} \underbrace{\frac{\Psi_{T}(\mathbf{R})^{*} \Psi_{T}(\mathbf{R})}{\int d\mathbf{R} \Psi_{T}(\mathbf{R})^{*} \Psi_{T}(\mathbf{R})}}_{P(\mathbf{R})} \cdot \underbrace{\frac{1}{\Psi_{T}(\mathbf{R})} \hat{\mathcal{H}} \Psi_{T}(\mathbf{R})}_{E_{L}(\mathbf{R})}$$





### Monte Carlo Integration

We attempt to solve the integral by sampling from the probability density function  $P(\mathbf{R})$ 

$$E_{VMC} = \int d\mathbf{R} E_L(\mathbf{R}) P(\mathbf{R})$$

$$\approx \frac{1}{M} \sum_{i=1}^{M} E_L(\mathbf{R}_i)$$





#### **Trial Wave Function**

$$P(\mathbf{R}) \propto \Psi_T(\mathbf{R})^* \Psi_T(\mathbf{R})$$

Use the Slater-Jastrow function as our trial wave function

$$\Psi_T(\mathbf{R}) = |\hat{D}(\mathbf{R})|J(\mathbf{R})$$

where the Slater matrix,  $\hat{D}(R)$ , contains all the single-particle functions

$$\hat{D}(\mathbf{R}) = \begin{pmatrix} \phi_1(\mathbf{r}_1) & \phi_2(\mathbf{r}_1) & \dots & \phi_N(\mathbf{r}_1) \\ \phi_1(\mathbf{r}_2) & \phi_2(\mathbf{r}_2) & \dots & \phi_N(\mathbf{r}_2) \\ \vdots & \vdots & \ddots & \vdots \\ \phi_1(\mathbf{r}_N) & \phi_2(\mathbf{r}_N) & \dots & \phi_N(\mathbf{r}_N) \end{pmatrix}$$





## Single-particle Functions

The Hermite functions,

$$\phi_n(\mathbf{r}) \propto H_n(\mathbf{r}) \exp\left(-\frac{1}{2}\alpha\omega|\mathbf{r}|^2\right),$$

are used as the single-particle functions for quantum dots in standard VMC. The Gaussian can be factorized out from the Slater determinant.

$$|\hat{D}(\boldsymbol{R};\alpha)| \propto \exp\left(-\frac{1}{2}\alpha\omega|\boldsymbol{R}|^2\right) \begin{vmatrix} H_1(\boldsymbol{r}_1) & H_2(\boldsymbol{r}_1) & \dots & H_N(\boldsymbol{r}_1) \\ H_1(\boldsymbol{r}_2) & H_2(\boldsymbol{r}_2) & \dots & H_N(\boldsymbol{r}_2) \\ \vdots & \vdots & \ddots & \vdots \\ H_1(\boldsymbol{r}_N) & H_2(\boldsymbol{r}_N) & \dots & H_N(\boldsymbol{r}_N) \end{vmatrix}$$





#### Restricted Boltzmann Machine

We use the marginal distribution of the visible units as the single-particle functions in the Slater determinant, and see if them can model the correlations

$$\phi_n(\mathbf{r}) \propto H_n(\mathbf{r}) P(\mathbf{r}; \mathbf{a}, \mathbf{b}, \mathbf{W})$$

where P(r) is the marginal distribution of the visible units.

$$|\hat{D}(\textbf{r};\textbf{a},\textbf{b},\textbf{W})| \propto P(\textbf{r};\textbf{a},\textbf{b},\textbf{W}) \begin{vmatrix} H_1(\textbf{r}_1) & H_2(\textbf{r}_1) & \dots & H_N(\textbf{r}_1) \\ H_1(\textbf{r}_2) & H_2(\textbf{r}_2) & \dots & H_N(\textbf{r}_2) \\ \vdots & \vdots & \ddots & \vdots \\ H_1(\textbf{r}_N) & H_2(\textbf{r}_N) & \dots & H_N(\textbf{r}_N) \end{vmatrix}$$







#### Jastrow Factor

The Jastrow factor is added to account for the correlations Simple Jastrow factor

$$J(\mathbf{r}; \boldsymbol{\beta}) = \exp\left(\sum_{i=1}^{N} \sum_{j>i}^{N} \beta_{ij} r_{ij}\right).$$

Padé-Jastrow factor

$$J(r;\beta) = \exp\left(\sum_{i=1}^{N} \sum_{j>i}^{N} \frac{a_{ij}r_{ij}}{1 + \beta r_{ij}}\right).$$





# Results





## **Ground State Energy**

Number of electrons: N = 2. Frequency:  $\omega$ .

ω	RBM	RBM+SJ	RBM+PJ	VMC	HF *	Exact †
1/6	0.7036(1)	0.67684(7)	0.66715(6)	0.66710(1)	0.768675	2/3
1	3.0803(2)	3.02108(5)	2.999587(5)	2.99936(1)	3.16190	3

†Semi-analytical ground state energy calculated by Taut, 1993 [6].





<sup>\*</sup>Computation of the Hartree-Fock limit by Mariadason, 2018 [5].

#### **Ground State Energy**

Number of electrons: N = 20. Frequency:  $\omega$ .

ω	RBM	RBM+SJ	RBM+PJ	VMC	HF <sup>‡</sup>	DMC §
0.1	30.824(2)	30.567(3)	30.1553(9)	30.0403(2)	31.1902	29.9779(1)
1.0	159.428(3)	156.816(4)	156.104(1)	155.8900(4)	158.004	155.8822(1)

<sup>§</sup>Ground state energy estimate using the diffusion Monte Carlo method. By Høgberget, 2013 [7].



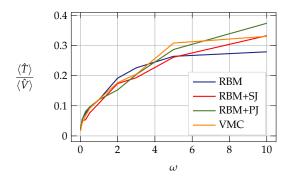


<sup>&</sup>lt;sup>‡</sup>Computation of the Hartree-Fock limit by Mariadason, 2018 [5].

### Energy distribution

Number of electrons: N = 20. Frequency:  $\omega$ .

Ratio between the kinetic energy,  $\langle \hat{T} \rangle$ , and the total potential energy,  $\langle \hat{V} \rangle$ .

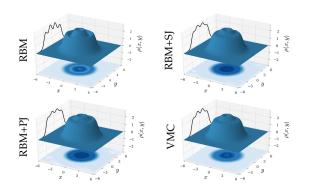






#### One-body density

Number of electrons: N = 20. Frequency:  $\omega = 1.0$ .

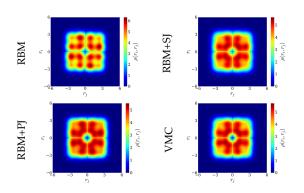






#### Two-body density

Number of electrons: N = 20. Frequency:  $\omega = 1.0$ .

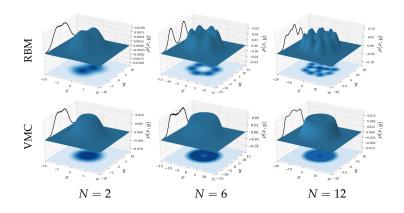






#### Low-frequency dots

Number of electrons: *N*. Frequency:  $\omega = 0.1$ .

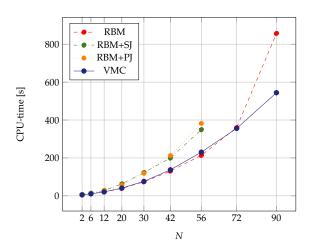






### Computational Cost

Number of electrons: *N*.







# Conclusion







#### **Conclusions**

- ▶ RBM is able to account for most of the correlations
- RBM+PJ implies to give a lower ground state energy and model the correlations better than a traditional VMC
- RBM+SJ is both more expensive and less accurate than its fellow methods, and we see no reason to choose it





#### **Future Work**

- Repeat the exercise using spherical coordinates interactions are easier to model in spherical coordinates
- Check the ability of modeling the three-body correlations, considering nuclear systems
- Reduce the computational cost





# Thank you!





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