Newton's Method: Equations

Newton's method solves a system of nonlinear equations of the form,

$$F(\mathbf{u}) = 0$$

by a sequence of steps including the linear problem

$$\mathbf{J}^{\mathbf{k}}(\mathbf{u}^k)\delta\mathbf{u}^k = -\mathbf{F}(\mathbf{u}^k),$$

where.

$$\mathsf{J}_{i,j}(\mathsf{u}^k) = \frac{\partial \mathsf{F}_i}{\partial \mathsf{u}_i^k},$$

and,

$$\mathbf{u}^{k+1} = \mathbf{u}^k + \delta \mathbf{u}^k.$$

This is continued until

$$\|\mathbf{F}(\mathbf{u}^k)\| < tol_n \|\mathbf{F}(\mathbf{u}^0)\|$$

where tol_n is an input nonlinear tolerance.

Newton's Method: Basics

Newton's method is a first-order Taylor series approximation

$$\mathbf{F}(\mathbf{u} + \delta \mathbf{u}) = 0,$$

$$\mathbf{F}(\mathbf{u}) + \delta \mathbf{u} \frac{\partial \mathbf{F}}{\partial \mathbf{u}} \approx 0$$

$$\mathbf{J} \delta \mathbf{u} = -\mathbf{F}(\mathbf{u}),$$

Forming J can be challenging for many problems. Analytic or numerical evaluation?

$$\mathbf{J}_{i,j} = \frac{\partial \mathbf{F}_i}{\partial \mathbf{u}_j},$$

▶ Radius of convergence can be small. This is a function of the Taylor series approximation. Damping / Globalization is often required. Choosing d?

$$\mathbf{u}^{k+1} = \mathbf{u}^k + d\delta \mathbf{u}^k.$$

Newton Iteration: Implementation

- 1. Define $T^{k=1}$
- 2. Do k = 1, Nnewt
- 3. Evaluate $F(T^k)$ (if $||F(T^k)||_2 < \text{tol then exit}$)
- 4. Evaluate $J(T^k)$
- 5. Solve $J(T^k)\delta T = -F(T^k)$
- 6. Update $T^{k+1} = T^k + \delta T$
- 7. End Do

Standard Newton-Krylov Methods

- Using a Krylov method to solve $\mathbf{J}^{\mathbf{k}}\delta\mathbf{u}^{k}=-\mathbf{F}(\mathbf{u}^{k})$
- ► Form both **J** and **F**(**u**)
- Send J and F(u) to Krylov solver and use standard matrix vector multiply.
- Extract preconditioning directly from J

Inexact Newton Methods

Inexact Newton tolerance on Krylov solver for each Newton iteration limits excessive Krylov iterations. Does prevent quadratic convergence.

$$\parallel \mathbf{J}^{k} \delta \mathbf{u}^{k} + \mathbf{F}(\mathbf{u}^{k}) \parallel_{2} < \gamma \parallel \mathbf{F}(\mathbf{u}^{k}) \parallel_{2}, \tag{1}$$

"Under solving" in early Newton iterations can increase radius of convergence.

Jacobian-Free Newton-Krylov

- ▶ Using a Krylov method to solve $\mathbf{J}^{\mathbf{k}}\delta\mathbf{u}^{k} = -\mathbf{F}(\mathbf{u}^{k})$
- Newton-Krylov methods only need the action of the Jacobian matrix to construct the m^{th} linear iterate of $\delta \mathbf{u}^k$

$$\delta \mathbf{u}_m^k = \beta_0 \mathbf{r}_0 + \beta_1 \mathbf{J} \mathbf{r}_0 + \beta_2 \mathbf{J}^2 \mathbf{r}_0 + \ldots + \beta_m \mathbf{J}^m \mathbf{r}_0$$

where

$$\mathbf{r}_0 = \mathbf{J}^k \delta \mathbf{u}_0^k + \mathbf{F}(\mathbf{u}^k)$$

The Matrix-vector product required by the Krylov method can be approximated with a single function evaluation

$$\mathsf{Jv} pprox rac{\mathsf{F}(\mathsf{u} + \epsilon \mathsf{v}) - \mathsf{F}(\mathsf{u})}{\epsilon}$$

The Jacobian-vector Product

Consider the two coupled nonlinear equations $F_1(u_1, u_2) = 0$, $F_2(u_1, u_2) = 0$. The Jacobian matrix is

$$\mathbf{J} = \begin{bmatrix} \frac{\partial F_1}{\partial u_1} & \frac{\partial F_1}{\partial u_2} \\ \\ \frac{\partial F_2}{\partial u_1} & \frac{\partial F_2}{\partial u_2} \end{bmatrix}.$$

Working backwards

$$\frac{\mathbf{F}(\mathbf{u}+\epsilon\mathbf{v})-\mathbf{F}(\mathbf{u})}{\epsilon} = \begin{pmatrix} \frac{F_1(u_1+\epsilon v_1,u_2+\epsilon v_2)-F_1(u_1,u_2)}{\epsilon} \\ \frac{F_2(u_1+\epsilon v_1,u_2+\epsilon v_2)-F_2(u_1,u_2)}{\epsilon} \end{pmatrix}.$$

The Jacobian-vector Product, cont.

▶ Approximating $\mathbf{F}(\mathbf{u} + \epsilon \mathbf{v})$ with a first-order Taylor series

$$\frac{\mathbf{F}(\mathbf{u} + \epsilon \mathbf{v}) - \mathbf{F}(\mathbf{u})}{\epsilon} \approx \begin{pmatrix} \frac{F_1(u_1, u_2) + \epsilon v_1 \frac{\partial F_1}{\partial u_1} + \epsilon v_2 \frac{\partial F_1}{\partial u_2} - F_1(u_1, u_2)}{\epsilon} \\ \frac{F_2(u_1, u_2) + \epsilon v_1 \frac{\partial F_2}{\partial u_1} + \epsilon v_2 \frac{\partial F_2}{\partial u_2} - F_2(u_1, u_2)}{\epsilon} \end{pmatrix},$$

which simplifies to

$$\begin{pmatrix} v_1 \frac{\partial F_1}{\partial u_1} + v_2 \frac{\partial F_1}{\partial u_2} \\ v_1 \frac{\partial F_2}{\partial u_1} + v_2 \frac{\partial F_2}{\partial u_2} \end{pmatrix} = \mathbf{J} v.$$

Jacobian-Free Newton-Krylov: Refs

- Standard "PDE motivated" references:
 P. N. Brown and Y. Saad, SIAM J. Sci. Stat. Comput., 11, pp. 450-481 (1990)
 Tony F. Chan and Kenneth R. Jackson, SIAM J. Sci. Stat. Comput., 5, pp. 533-542 (1984)
- ► Also see the monograph, C. T. Kelley, *Iterative Methods for Linear and Nonlinear Equations*, SIAM, Philadelphia, 1995
- ▶ Recent JFNK review article from the application perspective D.A. Knoll and D.E. Keyes, Jacobian-free Newton-Krylov methods: a survey of approaches and applications, *J. Comput. Phys.*, **193**, pp. 357-397 (2004)

Choosing ϵ

• Choosing ϵ for a numerical derivative of f(x)

$$\frac{\partial f(x)}{\partial x} \approx \frac{f(x+\epsilon) - f(x)}{\epsilon}$$

 $\epsilon = a(1.0 + |x|)$ where a is on order square root of round-off. $a \approx 1.0e - 6$ to 1.0e - 8.

▶ Choosing ϵ for the JFNK matvec. Many algorithms work for choosing ϵ (see refs), one simple one is

$$\epsilon = rac{1}{N||\mathbf{v}||_2}\sum_{i=1}^N \mathsf{a}(1.0+|u_i|)$$

where \mathbf{v} is the Krylov vector and N is the system dimension. Looks like an average of what would be done individually.

▶ Why we use GMRES

Jacobian Times a Vector: Implementation

▶ We need

$$\mathbf{J}\mathbf{v} = \frac{ResT(\mathbf{T} + \epsilon\mathbf{v}) - ResT(\mathbf{T})}{\epsilon}$$

- \blacktriangleright We have $ResT(\mathbf{T})$ and \mathbf{v}
- ▶ Must evaluate an ϵ , then evaluate $ResT(\mathbf{T} + \epsilon \mathbf{v})$, and finally evaluate $\mathbf{J}\mathbf{v}$.

Evaluating y = Jv, **MatrixFreeJv**

- Subroutine MatrixFreeJv(v,y,N)
 - 1. evaluate ϵ
 - 2. evaluate **Tpert** = $\mathbf{T} + \epsilon \mathbf{v}$
 - 3. evaluate $\mathbf{w} = ResT(\mathbf{Tpert})$ (call ResTemp)
 - 4. Do i = 1, N
 - 5. $y(i) = \frac{w(i) ResT(i)}{\epsilon}$
 - 6. End Do
- ▶ This routine / function provides the same service to GMRES as the function AFUN in Matlab. This is a matrix-free AFUN which allows you to do JFNK with GMRES in Matlab.

Standard Newton-Krylov vs JFNK

- ▶ Standard: Using a Krylov method to solve $\mathbf{J}^{\mathbf{k}}\delta\mathbf{u}^{k} = -\mathbf{F}(\mathbf{u}^{k})$
 - Form both J and F(u)
 - Send J and F(u) to Krylov solver and use standard matrix vector multiply.
 - Extract preconditioning directly from J
- ▶ JFNK: Using a Krylov method to solve $\mathbf{J}^{\mathbf{k}} \delta \mathbf{u}^{k} = -\mathbf{F}(\mathbf{u}^{k})$
 - ► Form only **F**(**u**)
 - Send only F(u) to Krylov solver and use Matrix-free matrix vector multiply.
 - Extract preconditioning from something less costly than J

Preconditioning: Concept and Choices

In left preconditioning the Krylov method is solving the altered system

$$M^{-1}Ax = M^{-1}b$$

- ▶ What is the "meaning" of M^{-1} vs that of A^{-1} ?
- We have two choices to make
 - 1. Do we use the same discrete operator to form M as we are using to solve Ax = b?
 - 2. How do we approximately invert M?
- ▶ GOAL: Reduce the number of Krylov iterations from N_1 to N_2 with a preconditioning process which cost less than $(N_1 N_2)$ Krylov iterations.

Preconditioning JFNK (1 of 2)

- Preconditioning is the KEY for efficient application to multiphysics engineering problems.
- Using right preconditioning one solves

$$(\mathsf{J}\mathsf{M}^{-1})(\mathsf{M}\delta\mathsf{u}) = -\mathsf{F}(\mathsf{u}). \tag{1}$$

 ${\bf M}$ symbolically represents the preconditioning matrix (or process) and ${\bf M}^{-1}$ the inverse of preconditioning matrix.

▶ Actually realized through a two-step process. First solve

$$(\mathbf{J}\mathbf{M}^{-1})\mathbf{w} = -\mathbf{F}(\mathbf{u}),\tag{2}$$

for **w** with the Krylov methods. Then solve for $\delta \mathbf{u}$,

$$\delta \mathbf{u} = \mathbf{M}^{-1} \mathbf{w}. \tag{3}$$

Preconditioning JFNK (2 of 2)

Right-preconditioned matrix-free version is:

$$\mathsf{J}\mathsf{M}^{-1}\mathsf{v}\approx\left[\mathsf{F}(\mathsf{u}+\epsilon\mathsf{M}^{-1}\mathsf{v})-\mathsf{F}(\mathsf{u})\right]/\epsilon. \tag{4}$$

- Required in step 1 of previous slide
- Actually done in two steps (v is given):
 - 1. Preconditioning: Solve (approximately) for \mathbf{y} in $\mathbf{y} = \mathbf{M}^{-1}\mathbf{v}$.
 - 2. Perform matrix-free product $\mathbf{J}\mathbf{y} \approx \left[\mathbf{F}(\mathbf{u} + \epsilon \mathbf{y}) \mathbf{F}(\mathbf{u})\right] / \epsilon$.
- ▶ Only the matrix elements required for the action of M^{-1} are formed.

Preconditioning JFNK: Options (1 of 2)

- ▶ Form **J** every Newton iteration, $\mathbf{M} \equiv \mathbf{J}$, and approximately invert \mathbf{M} as preconditioning.
 - ▶ There is NO advantage to this !!
 - Should perform the same as standard Newton Krylov
- Form J every N Newton iterations, M ≡ J, and approximately invert M as preconditioning.
 - ► This can significantly reduce the cost of forming **J**.
 - Strong Newton convergence maintained since GMRES is using most current version of J in matrix-vector multiply.
 - Use of same J for a number of Newton iterations is often referred to as Modified Newton Krylov (MNK).
 - ▶ We will consider this as a solver and a preconditioner.

Preconditioning JFNK: Options (2 of 2)

- ▶ Do not form **J**. Form **M** from a Picard linearization, and approximately invert **M** as preconditioning.
 - Here M can typically be formed with less effort.
 - Often the matrix from Picard linearization has better properties with respect to linear algebra.
 - A key concept in physics-based preconditioning.
- ▶ Do not form J. Form M from operators which control numerical stiffness (spread in eigenvalues) and approximately invert M as preconditioning.
 - Here M is formed from targeted physical processes such as diffusion, advection, reaction ...
 - This can often result optimal preconditioning cost-benefit.