Ordinary differential equations

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Chemical Process Intensification, Process Systems Engineering, Eindhoven University of Technology Introduction Explicit methods Implicit methods Boundary value problems Systems of ODEs

Today's outline

- 1 Introduction
- 2 Explicit methods

Forward Euler Convergence rate Runge-Kutta methods Step size control

- 3 Implicit methods Backward Euler Implicit midpoint method
- 4 Boundary value problems
 Shooting method
- **5** Systems of ODEs

Solution methods for systems of ODEs Stiff systems of ODEs Solving systems of ODEs in Matlab

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Overview

Ordinary differential equations

An equation containing a function of one independent variable and its derivatives, in contrast to a *partial differential equation*, which contains derivatives with respect to more independent variables.

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Main question

How to solve

$$\frac{d\mathbf{y}}{dx} = f(\mathbf{y}(x), x) \quad \text{with} \quad \mathbf{y}(x = 0) = \mathbf{y}_0$$

accurately and efficiently?

What is an ODE?

Algebraic equation:

$$f(y(x), x) = 0$$
 e.g. $-\ln(K_{eq}) = (1 - \zeta)$

First order ODE:

$$f\left(\frac{dy}{dx}(x), y(x), x\right) = 0$$
 e.g. $\frac{dc}{dt} = -kc^n$

Second order ODE:

$$f\left(\frac{d^2y}{dx^2}(x), \frac{dy}{dx}(x), y(x), x\right) = 0$$
 e.g. $\mathcal{D}\frac{d^2c}{dx^2} = -\frac{kc}{1 + Kc}$

About second order ODEs

Very often a second order ODE can be rewritten into a system of first order ODEs (whether it is handy depends on the boundary conditions!)

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Example

Recall:

$$\mathcal{D}\frac{d^2c}{dx^2} = -\frac{kc}{1+Kc}$$

Define $y=-\mathcal{D}\frac{dc}{dx}$, then $\frac{dy}{dx}=\frac{kc}{1+Kc}$, thus solve system:

$$\frac{dc}{dx} = -\frac{1}{\mathcal{D}}y$$

$$\frac{dy}{dx} = \frac{kc}{1 + c}$$

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More general

Consider the second order ODE:

$$\frac{d^2y}{dx^2} + q(x)\frac{dy}{dx} = r(x)$$

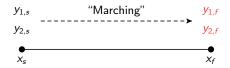
Now define and solve using z as a new variable:

$$\frac{dy}{dx} = z(x)$$

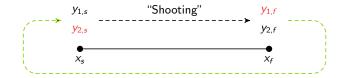
$$\frac{dz}{dx} = r(x) - q(x)z(x)$$

The nature of boundary conditions determines the appropriate numerical method. Classification into 2 main categories:

Initial value problems (IVP) We know the values of all y_i at some starting position x_s , and it is desired to find the values of y_i at some final point x_f .



Boundary value problems (BVP) Boundary conditions are specified at more than one x. Typically, some of the BC are specified at x_s and the remainder at x_f .



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Initial value problems:

- Explicit methods
 - First order: forward Euler
 - Second order: improved Euler (RK2)
 - Fourth order: Runge-Kutta 4 (RK4)
 - Step size control
- Implicit methods
 - First order: backward Euler
 - Second order: midpoint rule

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Boundary value problems

• Shooting method

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Euler's method

Consider the following single initial value problem:

$$\frac{dc}{dt} = f(c(t), t)$$
 with $c(t = 0) = c_0$ (initial value problem)

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Easiest solution algorithm: Euler's method, derived here via Taylor series expansion:

$$c(t_0 + \Delta t) \approx c(t_0) + \left. \frac{dc}{dt} \right|_{t_0} \Delta t + \left. \frac{1}{2} \left. \frac{d^2c}{dt^2} \right|_{t_0} (\Delta t)^2 + \mathcal{O}(\Delta t^3)$$

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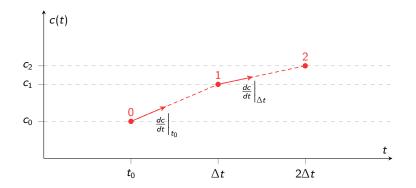
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Neglect terms with higher order than two: $\frac{dc}{dt}\Big|_{t} = \frac{c(t_0 + \Delta t) - c(t_0)}{\Delta t}$ Substitution:

$$\frac{c(t_0 + \Delta t) - c(t_0)}{\Delta t} = f(c_0, t_0) \Rightarrow c(t_0 + \Delta t) = c(t_0) + \Delta t f(c_0, t_0)$$

Euler's method: graphical example

$$\frac{c(t_0 + \Delta t) - c(t_0)}{\Delta t} = f(c_0, t_0) \Rightarrow c(t_0 + \Delta t) = c(t_0) + \Delta t f(c_0, t_0)$$



Euler's method - solution method

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Introduction Explicit methods

Start with $t=t_0$, $c=c_0$, then calculate at discrete points in time: $c(t_1=t_0+\Delta t)=c(t_0)+\Delta t f(c_0,t_0)$.

Pseudo-code Euler's method: $\frac{dy}{dx} = f(x, y)$ and $y(x_0) = y_0$.

- 1 Initialize variables, functions; set $h = \frac{x_1 x_0}{N}$
- 2 Set $x = x_0$, $y = y_0$
- 3 While $x < x_{end}$ do $x_{i+1} = x_i + h$; $y_{i+1} = y_i + hf(x_i, y_i)$

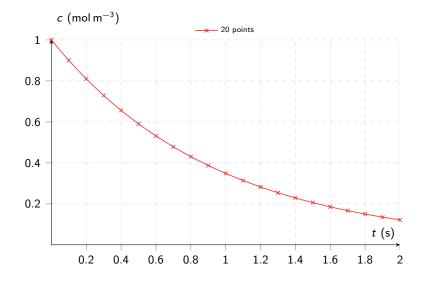
First order reaction in a batch reactor:

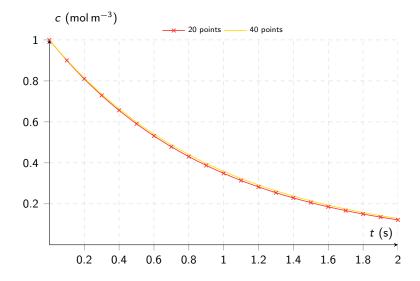
$$\frac{dc}{dt} = -kc$$
 with $c(t=0) = 1 \text{ mol m}^{-3}$, $k = 1 \text{ s}^{-1}$, $t_{\text{end}} = 2 \text{ s}$

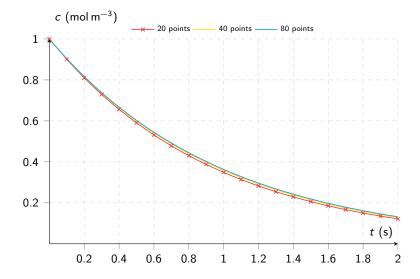
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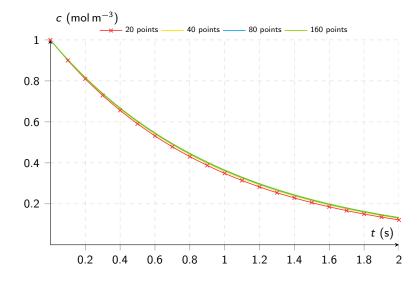
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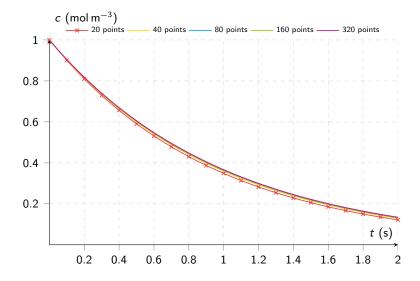
Time [s]	Concentration $[mol m^{-3}]$
$t_0 = 0$	$c_0 = 1.00$
$t_1 = t_0 + \Delta t$	$c_1 = c_0 + \Delta t \cdot (-kc_0)$
= 0 + 0.1 = 0.1	$= 1 + 0.1 \cdot (-1 \cdot 1) = 0.9$
$t_2 = t_1 + \Delta t$	$c_2 = c_1 + \Delta t \cdot (-kc_1)$
= 0.1 + 0.1 = 0.2	$= 0.9 + 0.1 \cdot (-1 \cdot 0.9) = 0.81$
$t_3 = t_2 + \Delta t$	$c_3 = c_2 + \Delta t \cdot (-kc_2)$
= 0.2 + 0.1 = 0.3	$= 0.81 + 0.1 \cdot (-1 \cdot 0.81) = 0.729$
$t_{i+1} = t_i + \Delta t$	$c_{i+1} = c_i + \Delta t \cdot (-kc_i)$
$t_{20} = 2.0$	$c_{20} = c_{19} + \Delta t \cdot (-kc_{19}) = 0.1211577$







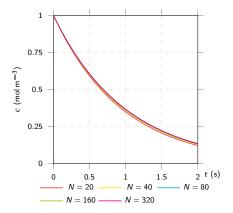




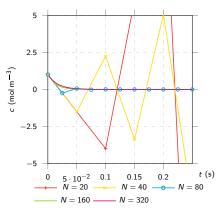
Problems with Euler's method

The question is: What step size, or how many steps to use?

- **1** $Accuracy \Rightarrow$ need information on numerical error!
- 2 Stability \Rightarrow need information on stability limits!



Reaction rate: $k = 1 \text{ s}^{-1}$



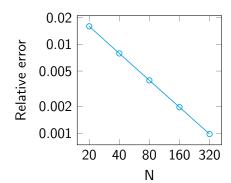
Reaction rate: $k = 50 \text{ s}^{-1}$

Accuracy

Comparison with analytical solution for $k = 1 \text{ s}^{-1}$:

$$c(t) = c_0 \exp(-kt) \Rightarrow \zeta = 1 - \exp(-kt) \Rightarrow \zeta_{\text{analytical}} = 0.864665$$

Ν	ζ	$\frac{\zeta_{\text{numerical}} - \zeta_{\text{analytical}}}{\zeta_{\text{analytical}}}$
20	0.878423	0.015912
40	0.871488	0.007891
80	0.868062	0.003929
160	0.866360	0.001961
320	0.865511	0.000979



For Euler's method: Error halves when the number of grid points is doubled, i.e. error is proportional to Δt : first order method.

Error estimate:

$$\left. \frac{dx}{dt} \right|_{t_0} = \frac{x(t_0 + \Delta t) - x(t_0)}{\Delta t} + \frac{1}{2} \left. \frac{d^2x}{dt^2} \right|_{t_0} (\Delta t) + \mathcal{O}(\Delta t)^2$$

$$\frac{x(t_0 + \Delta t) - x(t_0)}{\Delta t} = f(x_0, t_0) - \frac{1}{2} \left. \frac{d^2 x}{dt^2} \right|_{t_0} (\Delta t) + \mathcal{O}(\Delta t)^2$$

L_2 norm (Euclidean norm)

$$\|\mathbf{v}\|_2 = \sqrt{v_1^2 + v_2^2 + \ldots + v_n^2} = \sqrt{\sum_{i=1}^n v_i^2}$$

L_{∞} norm (maximum norm)

$$\|\mathbf{v}\|_{\infty} = \max(|v_1|,\ldots,|v_n|)$$

Absolute difference

$$\epsilon_{\mathsf{abs}} = \left\| \mathbf{y}_{\mathsf{numerical}} - \mathbf{y}_{\mathsf{analytical}} \right\|_{2,\infty}$$

Relative difference

$$\epsilon_{\mathsf{rel}} = rac{\|\mathbf{y}_{\mathsf{numerical}} - \mathbf{y}_{\mathsf{analytical}}\|_{2,\infty}}{\|\mathbf{y}_{\mathsf{analytical}}\|_{2,\infty}}$$

Errors and convergence rate

Convergence rate (or: order of convergence) r

$$\epsilon = \lim_{\Delta t \to 0} c(\Delta x)^r$$

- A first order method reduces the error by a factor 2 when increasing the number of steps by a factor 2
- A second order method reduces the error by a factor 4 when increasing the number of steps by a factor 2

What to do when there is no analytical solution available?

Errors and convergence rate

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What to do when there is no analytical solution available? Compare to calculations with different number of steps: $\epsilon_1 = c(\Delta x_1)^r$ and $\epsilon_2 = c(\Delta x_2)^r$ and solve for r:

$$\frac{\epsilon_2}{\epsilon_1} = \frac{c(\Delta x_2)^r}{c(\Delta x_1)^r} = \left(\frac{\Delta x_2}{\Delta x_1}\right)^r \Rightarrow \log\left(\frac{\epsilon_2}{\epsilon_1}\right) = \log\left(\frac{\Delta x_2}{\Delta x_1}\right)^r$$

$$\Rightarrow r = \frac{\log\left(\frac{\epsilon_2}{\epsilon_1}\right)}{\log\left(\frac{\Delta x_2}{\Delta x_1}\right)} = \frac{\log\left(\frac{\epsilon_2}{\epsilon_1}\right)}{\log\left(\frac{N_1}{N_2}\right)} \quad \text{in the limit of} \quad \Delta x \to 0 \qquad \text{or} \qquad N \to \infty$$

Example: Euler's method — order of convergence

Ν	ζ	$rac{\zeta_{ ext{numerical}} - \zeta_{ ext{analytical}}}{\zeta_{ ext{analytical}}}$	$r = \frac{\log\left(\frac{\epsilon_i}{\epsilon_{i-1}}\right)}{\log\left(\frac{N_{i-1}}{N_i}\right)}$
20	0.878423	0.015912	_
40	0.871488	0.007891	1.011832
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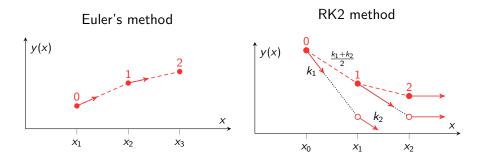
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Wouldn't it be great to have a method that can give the answer using much less steps? ⇒ Higher order methods

Runge-Kutta methods

Propagate a solution by combining the information of several Euler-style steps (each involving one function evaluation) to match a Taylor series expansion up to some higher order.

Euler:
$$y_{i+1} = y_i + hf(x_i, y_i)$$
 with $h = \Delta x$, i.e. slope $= k_1 = f(x_i, y_i)$.



This method is also called Heun's method, or improved Euler method:

- **1** Approximate the slope at x_i : $k_1 = f(x_i, y_i)$
- 2 Approximate the slope at x_{i+1} : $k_2 = f(x_{i+1}, y_{i+1})$ where we use Euler's method to approximate $y_{i+1} = y_i + hf(x_i, y_i) = y_i + hk_1$
- 3 Perform an Euler step with the average of the slopes: $y_{i+1} = y_i + h\frac{1}{2}(k_1 + k_2)$

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In pseudocode:

```
x = x_0, y = y_0

while x < x_{end} do

x_{i+1} = x_i + h

k_1 = f(x_i, y_i)

k_2 = f(x_i + h, y_i + hk_1)

y_{i+1} = y_i + h\frac{1}{2}(k_1 + k_2)

end while
```

$$\frac{dy}{dx} = f(x, y(x))$$

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Using Taylor series expansion: $y_{i+1} = y_i + h \left. \frac{dy}{dx} \right|_i + \frac{h^2}{2} \frac{d^2y}{dx^2} \right| + \mathcal{O}(h^3)$

$$\left. \frac{dy}{dx} \right|_i = f(x_i, y_i) \equiv f_i$$

$$\frac{d^2y}{dx^2}\bigg| = \frac{d}{dx}f(x,y(x))\bigg|_i = \frac{\partial f}{\partial x}\bigg|_i + \frac{\partial f}{\partial y}\bigg|_i \frac{\partial y}{\partial x}\bigg|_i = \frac{\partial f}{\partial x}\bigg|_i + \frac{\partial f}{\partial y}\bigg|_i f_i \quad \text{(chain rule)}$$

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Substitution gives:

$$y_{i+1} = y_i + hf_i + \frac{h^2}{2} \left(\frac{\partial f}{\partial x} \Big|_i + \frac{\partial f}{\partial y} \Big|_i f_i \right) + \mathcal{O}(h^3)$$

$$y_{i+1} = y_i + \frac{h}{2} f_i + \frac{h}{2} \left(f_i + h \frac{\partial f}{\partial x} \Big|_i + hf_i \frac{\partial f}{\partial y} \Big|_i \right) + \mathcal{O}(h^3)$$

Note multivariate Taylor expansion:

$$f(x_i + h, y_i + k) = f_i + h \frac{\partial f}{\partial x}\Big|_i + k \frac{\partial f}{\partial y}\Big|_i + \mathcal{O}(h^2)$$

$$\Rightarrow \frac{h}{2} \left(f_i + h \frac{\partial f}{\partial x}\Big|_i + h f_i \frac{\partial f}{\partial y}\Big|_i \right) = \frac{h}{2} f(x_i + h, y_i + k f_i) + \mathcal{O}(h^3)$$

Concluding:

$$y_{i+1} = y_i + \frac{h}{2}f_i + \frac{h}{2}f(x_i + h, y_i + kf_i) + \mathcal{O}(h^3)$$

Rewriting:

$$k_1 = f(x_i, y_i)$$

$$k_2 = f(x_i + h, y_i + hk_1)$$

$$\Rightarrow y_{i+1} = y_i + \frac{h}{2}(k_1 + k_2)$$

```
Generalization: y_{i+1} = y_i + h(b_1k_1 + b_2k_2) + \mathcal{O}(h^3) with k_1 = f_i, k_2 = f(x_i + c_2h, y_1 + a_{2,1}hk_1) (Note that classical RK2: b_1 = b_2 = \frac{1}{2} and c_2 = a_{2,1} = 1.)
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Bivariate Taylor expansion:

$$f(x_{i} + c_{2}h, y_{i} + a_{2,1}hk_{1}) = f_{i} + c_{2}h \frac{\partial f}{\partial x}\Big|_{i} + a_{2,1}hk_{1} \frac{\partial f}{\partial y}\Big|_{i} + \mathcal{O}(h^{2})$$

$$y_{i+1} = y_{i} + h(b_{1}k_{1} + b_{2}k_{2}) + \mathcal{O}(h^{3})$$

$$= y_{i} + h\left[b_{1}f_{i} + b_{2}f(x_{i} + c_{2}h, y_{1} + a_{2,1}hk_{1})\right] + \mathcal{O}(h^{3})$$

$$= y_{i} + h\left[b_{1}f_{i} + b_{2}\left\{f_{i} + c_{2}h \frac{\partial f}{\partial x}\Big|_{i} + a_{2,1}hk_{1} \frac{\partial f}{\partial y}\Big|_{i} + \mathcal{O}(h^{2})\right\}\right] + \mathcal{O}(h^{3})$$

$$= y_{i} + h(b_{1} + b_{2})f_{i} + h^{2}b_{2}\left(c_{2} \frac{\partial f}{\partial x}\Big|_{i} + a_{2,1}f_{i} \frac{\partial f}{\partial y}\Big|_{i}\right) + \mathcal{O}(h^{3})$$

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Bivariate Taylor expansion:

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$$y_{i+1} = y_{i} + h(b_{1}k_{1} + b_{2}k_{2}) + \mathcal{O}(h^{3})$$

$$= y_{i} + h\left[b_{1}f_{i} + b_{2}f(x_{i} + c_{2}h, y_{1} + a_{2,1}hk_{1})\right] + \mathcal{O}(h^{3})$$

$$= y_{i} + h\left[b_{1}f_{i} + b_{2}\left\{f_{i} + c_{2}h \frac{\partial f}{\partial x}\Big|_{i} + a_{2,1}hk_{1} \frac{\partial f}{\partial y}\Big|_{i} + \mathcal{O}(h^{2})\right\}\right] + \mathcal{O}(h^{3})$$

$$= y_{i} + h(b_{1} + b_{2})f_{i} + h^{2}b_{2}\left(c_{2} \frac{\partial f}{\partial x}\Big|_{i} + a_{2,1}f_{i} \frac{\partial f}{\partial y}\Big|_{i}\right) + \mathcal{O}(h^{3})$$

Comparison with Taylor:

$$y_{i+1} = y_i + hf_i + \frac{h^2}{2} \left(\frac{\partial f}{\partial x} \Big|_i + \frac{\partial f}{\partial y} \Big|_i f_i \right) + \mathcal{O}(h^3)$$

Generalization:
$$y_{i+1} = y_i + h(b_1k_1 + b_2k_2) + \mathcal{O}(h^3)$$
 with $k_1 = f_i$, $k_2 = f(x_i + c_2h, y_1 + a_{2,1}hk_1)$ (Note that classical RK2: $b_1 = b_2 = \frac{1}{2}$ and $c_2 = a_{2,1} = 1$.)

Bivariate Taylor expansion:

$$f(x_{i} + c_{2}h, y_{i} + a_{2,1}hk_{1}) = f_{i} + c_{2}h \frac{\partial f}{\partial x}\Big|_{i} + a_{2,1}hk_{1} \frac{\partial f}{\partial y}\Big|_{i} + \mathcal{O}(h^{2})$$

$$y_{i+1} = y_{i} + h(b_{1}k_{1} + b_{2}k_{2}) + \mathcal{O}(h^{3})$$

$$= y_{i} + h\left[b_{1}f_{i} + b_{2}f(x_{i} + c_{2}h, y_{1} + a_{2,1}hk_{1})\right] + \mathcal{O}(h^{3})$$

$$= y_{i} + h\left[b_{1}f_{i} + b_{2}\left\{f_{i} + c_{2}h \frac{\partial f}{\partial x}\Big|_{i} + a_{2,1}hk_{1} \frac{\partial f}{\partial y}\Big|_{i} + \mathcal{O}(h^{2})\right\}\right] + \mathcal{O}(h^{3})$$

$$= y_{i} + h(b_{1} + b_{2})f_{i} + h^{2}b_{2}\left(c_{2} \frac{\partial f}{\partial x}\Big|_{i} + a_{2,1}f_{i} \frac{\partial f}{\partial y}\Big|_{i}\right) + \mathcal{O}(h^{3})$$

Comparison with Taylor:

$$y_{i+1} = y_i + hf_i + \frac{h^2}{2} \left(\frac{\partial f}{\partial x} \Big|_i + \frac{\partial f}{\partial y} \Big|_i f_i \right) + \mathcal{O}(h^3)$$

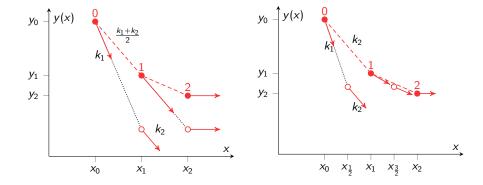
$$y_{i+1} = y_i + h(b_1 + b_2)f_i + h^2 b_2 \left(c_2 \frac{\partial f}{\partial x} \Big|_i + a_{2,1} f_i \frac{\partial f}{\partial y} \Big|_i \right) + \mathcal{O}(h^3)$$

$$y_{i+1} = y_i + h f_i + \frac{h^2}{2} \left(\frac{\partial f}{\partial x} \Big|_i + \frac{\partial f}{\partial y} \Big|_i f_i \right) + \mathcal{O}(h^3)$$

- \Rightarrow 3 eqns and 4 unknowns \Rightarrow multiple possibilities!
 - 1 Classical RK2: $b_1 = b_2 = \frac{1}{2}$ and $c_2 = a_{2,1} = 1$
 - ② Midpoint rule (modified Euler): $b_1 = 0, b_2 = 1, c_2 = a_{2,1} = \frac{1}{2}$

Second order Runge-Kutta methods

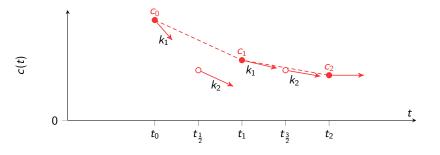
Classical RK2 method (= Heun's method, improved Euler method)	Explicit midpoint rule (modified Euler method)
$k_1 = f_i$	$k_1 = f_i$
$k_2 = f(x_i + h, y_i + hk_1)$	$k_2 = f(x_i + \frac{1}{2}h, y_i + \frac{1}{2}hk_1)$
$y_{i+1} = y_i + \frac{1}{2}h(k_1 + k_2)$	$y_{i+1} = y_i + hk_2$



Second order Runge-Kutta method — Example

First order reaction in a batch reactor: $\frac{dc}{dt} = -kc$ with $c(t=0) = 1 \text{ mol m}^{-3}$, $k=1 \text{ s}^{-1}$, $t_{\text{end}} = 2 \text{ s}$.

Time [s]	C [mol m ⁻³]	$k_1 = hf(x_i, y_i)$	$k_2 = hf(x_i + \frac{1}{2}h, y_n + \frac{1}{2}k_1)$
0	1.00	$0.1 \cdot (-1 \cdot 1) = -0.1$	$0.1 \cdot (-1 \cdot (1 - 0.5 \cdot 0.1)) = -0.095$
0.1	1 - 0.095 = 0.905	$\begin{array}{l} 0.1 \cdot (-1 \cdot 0.0905) = \\ -0.0905 \end{array}$	$0.1 \cdot (-1 \cdot (0.905 - 0.5 \cdot 0.0905)) = 0.085975$
2	0.1358225	-0.0135822	-0.0129031



RK2 method — order of convergence

N	ζ	$\frac{\zeta_{\text{numerical}} - \zeta_{\text{analytical}}}{\zeta_{\text{analytical}}}$	$r = \frac{\log\left(\frac{\epsilon_i}{\epsilon_{i-1}}\right)}{\log\left(\frac{N_{i-1}}{N_i}\right)}$
20	0.864178	5.634×10^{-4}	_
40	0.864548	1.355×10^{-4}	2.056
80	0.864636	3.323×10^{-5}	2.028
160	0.864658	8.229×10^{-6}	2.014
320	0.864663	2.048×10^{-6}	2.007

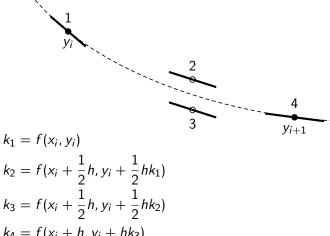
RK2 method — order of convergence

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 \Rightarrow RK2 is a second order method. Doubling the number of cells reduces the error by a factor 4!

Can we do even better?

RK4 method (classical fourth order Runge-Kutta method)



$$k_4 = f(x_i + h, y_i + hk_3)$$

$$y_{i+1} = y_i + h\left(\frac{1}{6}k_1 + \frac{1}{3}(k_2 + k_3) + \frac{1}{6}k_4\right)$$

RK4 method — order of convergence

Ν	ζ	$rac{\zeta_{ ext{numerical}} - \zeta_{ ext{analytical}}}{\zeta_{ ext{analytical}}}$	$r = \frac{\log\left(\frac{\epsilon_i}{\epsilon_{i-1}}\right)}{\log\left(\frac{N_{i-1}}{N_i}\right)}$
20	0.864664472	2.836×10^{-7}	_
40	0.864664702	1.700×10^{-8}	4.060
80	0.864664716	1.040×10^{-9}	4.030
160	0.864664717	6.435×10^{-11}	4.015
320	0.864664717	4.001×10^{-12}	4.007

RK4 method — order of convergence

Ν	ζ	$rac{\zeta_{ ext{numerical}} - \zeta_{ ext{analytical}}}{\zeta_{ ext{analytical}}}$	$r = \frac{\log\left(\frac{\epsilon_i}{\epsilon_{i-1}}\right)}{\log\left(\frac{N_{i-1}}{N_i}\right)}$
20	0.864664472	2.836×10^{-7}	_
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320	0.864664717	4.001×10^{-12}	4.007

 \Rightarrow RK4 is a fourth order method: Doubling the number of cells reduces the error by a factor 16!

Can we do even better?

Adaptive step size control

The step size (be it either position, time or both (PDEs)) cannot be decreased indefinitely to favour a higher accuracy, since each additional grid point causes additional computation time. It may be wise to adapt the step size according to the computation requirements.

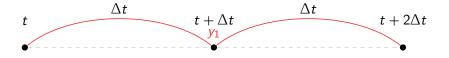
Globally two different approaches can be used:

- Step doubling: compare solutions when taking one full step or two consecutive halve steps
- 2 Embedded methods: Compare solutions when using two approximations of different order

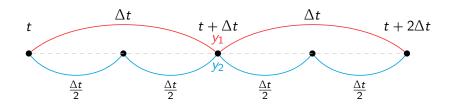
Adaptive step size control: step doubling

$$t$$
 $t + \Delta t$ $t + 2\Delta t$

Adaptive step size control: step doubling



• RK4 with one large step of h: $y_{i+1} = y_1 + ch^5 + \mathcal{O}(h^6)$



- RK4 with one large step of h: $y_{i+1} = y_1 + ch^5 + \mathcal{O}(h^6)$
- RK4 with two steps of $\frac{1}{2}h$: $y_{i+1} = y_2 + 2c(\frac{1}{2}h)^5 + \mathcal{O}(h^6)$

Adaptive step size control: step doubling

- Estimation of truncation error by comparing y₁ and y₂: $\Delta = y_2 - y_1$
- If ∆ too large, reduce step size for accuracy
- If Δ too small, increase step size for efficiency.
- Ignoring higher order terms and solving for c: $\Delta = \frac{15}{16}ch^5 \Rightarrow ch^5 = \frac{16}{15}\Delta \Rightarrow y_{i+1} = y_2 + \frac{\Delta}{15} + \mathcal{O}(h^6)$ (local Richardson extrapolation)

Note that when we specify a tolerance tol, we can estimate the maximum allowable step size as: $h_{\text{new}} = \alpha h_{\text{old}} \left| \frac{\text{tol}}{\Delta} \right|^{\frac{1}{5}}$ with α a safety factor (typically $\alpha = 0.9$).

Adaptive step size control: embedded methods

Use a special fourth and a fifth order Runge Kutta method to approximate y_{i+1}

- The fourth order method is special because we want to use the same positions for the evaluation for computational efficiency.
- RK45 is there preferred method (minimum number of function evaluations) (this is built in Matlab as ode45).

Today's outline

- Introduction
- Explicit methods

Forward Euler Convergence rate Runge-Kutta methods Step size control

- 3 Implicit methods Backward Euler Implicit midpoint method
- 4 Boundary value problems Shooting method
- **5** Systems of ODEs

Solution methods for systems of OD Stiff systems of ODEs Solving systems of ODEs in Matlab

6 Conclusion

Problems with Euler's method: instability

Consider the ODE:

$$\frac{dy}{dx} = f(x, y(x)) \quad \text{with} \quad y(x = 0) = y_0$$

Problems with Euler's method: instability

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First order approximation of derivative: $\frac{dy}{dx} = \frac{y_{i+1} - y_i}{\Delta x}$.

Where to evaluate the function f?

Problems with Euler's method: instability

Consider the ODE:

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First order approximation of derivative: $\frac{dy}{dx} = \frac{y_{i+1} - y_i}{\Delta x}$.

Where to evaluate the function f?

- 1 Evaluation at x_i : Explicit Euler method (forward Euler)
- 2 Evaluation at x_{i+1} : Implicit Euler method (backward Euler)

Explicit Euler method (forward Euler):

Use values at x_i:

$$\frac{y_{i+1}-y_i}{\Delta x}=f(x_i,y_i)\Rightarrow y_{i+1}=y_i+hf(x_i,y_i).$$

- This is an explicit equation for y_{i+1} in terms of y_i.
- It can give instabilities with large function values.

Explicit Euler method (forward Euler):

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- This is an explicit equation for y_{i+1} in terms of y_i .
- It can give instabilities with large function values.

Consider the first order batch reactor:

$$\frac{dc}{dt} = -kc \Rightarrow c_{i+1} = c_i - kc_i \Delta t \Rightarrow \frac{c_{i+1}}{c_i} = 1 - k\Delta t$$

Problems with Euler's method: instability – forward Euler

Explicit Euler method (forward Euler):

- Use values at x_i: $\frac{y_{i+1}-y_i}{\wedge y}=f(x_i,y_i)\Rightarrow y_{i+1}=y_i+hf(x_i,y_i).$
- This is an explicit equation for y_{i+1} in terms of y_i .
- It can give instabilities with large function values.

Consider the first order batch reactor:

$$\frac{dc}{dt} = -kc \Rightarrow c_{i+1} = c_i - k \frac{c_i}{\Delta t} \Rightarrow \frac{c_{i+1}}{c_i} = 1 - k \Delta t$$

It follows that unphysical results are obtained for $k\Delta t > 1!!$

Stability requirement

$$k\Delta t < 1$$

(but probably accuracy requirements are more stringent here!)

Problems with Euler's method: instability – backward Euler

Implicit Euler method (backward Euler):

- Use values at x_{i+1}: $\frac{y_{i+1}-y_i}{\Lambda x} = f(x_{i+1}, y_{i+1}) \Rightarrow y_{i+1} = y_i + hf(x_{i+1}, y_{i+1}).$
- This is an implicit equation for y_{i+1} , because it also depends on terms of y_{i+1} .

Problems with Euler's method: instability – backward Euler

Implicit Euler method (backward Euler):

- Use values at x_{i+1}: $\frac{y_{i+1}-y_i}{\Delta x} = f(x_{i+1}, y_{i+1}) \Rightarrow y_{i+1} = y_i + hf(x_{i+1}, y_{i+1}).$
- This is an implicit equation for y_{i+1} , because it also depends on terms of y_{i+1} .

Consider the first order batch reactor:

$$\frac{dc}{dt} = -kc \Rightarrow c_{i+1} = c_i - k \frac{c_{i+1}}{c_i} \Delta t \Rightarrow \frac{c_{i+1}}{c_i} = \frac{1}{1 + k\Delta t}$$

Problems with Euler's method: instability – backward Euler

Implicit Euler method (backward Euler):

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Consider the first order batch reactor:

$$\frac{dc}{dt} = -kc \Rightarrow c_{i+1} = c_i - kc_{i+1}\Delta t \Rightarrow \frac{c_{i+1}}{c_i} = \frac{1}{1 + k\Delta t}$$

This equation does never give unphysical results! The implicit Euler method is unconditionally stable (but maybe not very accurate or efficient).

Usually f is a non-linear function of y, so that linearization is required (recall Newton's method).

$$\frac{dy}{dx} = f(y) \Rightarrow y_{i+1} = y_i + hf(y_{i+1}) \quad \text{using} \quad f(y_{i+1}) = f(y_i) + \frac{df}{dy} \Big|_i (y_{i+1} - y_i) + \dots$$

$$\Rightarrow y_{i+1} = y_i + h \left[f(y_i) + \frac{df}{dy} \Big|_i (y_{i+1} - y_i) \right]$$

$$\Rightarrow \left(1 - h \frac{df}{dy} \Big|_i \right) y_{i+1} = \left(1 - h \frac{df}{dy} \Big|_i \right) y_i + hf(y_i)$$

$$\Rightarrow y_{i+1} = y_i + h \left(1 - h \frac{df}{dy}\Big|_{i}\right)^{-1} f(y_i)$$

Semi-implicit Euler method

Usually f is a non-linear function of y, so that linearization is required (recall Newton's method).

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$$\Rightarrow y_{i+1} = y_i + h \left(1 - h \frac{df}{dy}\Big|_i\right)^{-1} f(y_i)$$

For the case that f(x, y(x)) we could add the variable x as an additional variable $y_{n+1} = x$. Or add one fully implicit Euler step (which avoids the computation of $\frac{\partial f}{\partial x}$):

$$y_{i+1} = y_i + hf(x_{i+1}, y_{i+1}) \Rightarrow y_{i+1} = y_i + h\left(1 - h\left(\frac{df}{dy}\right)_i\right)^{-1} f(x_{i+1}, y_i)$$

Semi-implicit Euler method - example

$$\frac{dc}{dt}=-kc^2$$
 with $c_0=1~{\rm mol\,m^{-3}},~k=1~{\rm m^3\,mol^{-1}\,s^{-1}},~t_{\rm end}=2~{\rm s}$ Analytical solution: $c(t)=\frac{c_0}{1+kc_0t}$

Semi-implicit Euler method - example

$$\frac{dc}{dt}=-kc^2$$
 with $c_0=1$ mol m⁻³, $k=1$ m³ mol⁻¹ s⁻¹, $t_{\rm end}=2$ s Analytical solution: $c(t)=\frac{c_0}{1+kc_0t}$

Define
$$f = -kc^2$$
, then $\frac{df}{dc} = -2kc \Rightarrow c_{i+1} = c_i - \frac{hkc_i^2}{1+2hkc_i}$.

Semi-implicit Euler method - example

$$\frac{dc}{dt}=-kc^2$$
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Ν	ζ	$rac{\zeta_{ ext{numerical}} - \zeta_{ ext{analytical}}}{\zeta_{ ext{analytical}}}$	$r = \frac{\log\left(rac{\epsilon_i}{\epsilon_{i-1}} ight)}{\log\left(rac{N_{i-1}}{N_i} ight)}$
20	0.654066262	1.89×10^{-2}	_
40	0.660462687	9.31×10^{-3}	1.02220
80	0.663589561	4.62×10^{-3}	1.01162
160	0.665134433	2.30×10^{-3}	1.00594
320	0.665902142	1.15×10^{-3}	1.00300

Second order implicit method: Implicit midpoint method

Implicit midpoint rule	Explicit midpoint rule
(second order)	(modified Euler method)
$y_{i+1} = y_i + hf\left(x_i + \frac{1}{2}h, \frac{1}{2}(y_i + y_{i+1})\right)$	$y_{i+1} = y_i + hf(x_i + \frac{1}{2}h, y_i + \frac{1}{2}hk_1)$

in case f(y) then:

$$f\left(\frac{1}{2}(y_i+y_{i+1})\right) = f_i + \frac{df}{dy}\Big|_i \left(\frac{1}{2}(y_i+y_{i+1}) - y_i\right) = f_i + \frac{1}{2} \frac{df}{dy}\Big|_i (y_{i+1}-y_i)$$

Implicit midpoint rule Explicit midpoint rule (second order) (modified Euler method)
$$y_{i+1} = y_i + hf\left(x_i + \frac{1}{2}h, \frac{1}{2}(y_i + y_{i+1})\right) \quad y_{i+1} = y_i + hf(x_i + \frac{1}{2}h, y_i + \frac{1}{2}hk_1)$$

in case f(y) then:

$$f\left(\frac{1}{2}(y_i+y_{i+1})\right) = f_i + \frac{df}{dy}\Big|_i \left(\frac{1}{2}(y_i+y_{i+1}) - y_i\right) = f_i + \frac{1}{2} \frac{df}{dy}\Big|_i (y_{i+1}-y_i)$$

Implicit midpoint rule reduces to:

$$y_{i+1} = y_i + hf_i + \frac{h}{2} \left. \frac{df}{dy} \right|_i (y_{i+1} - y_i)$$

$$\Rightarrow \left(1 - \frac{h}{2} \left. \frac{df}{dy} \right|_i \right) y_{i+1} = \left(1 - \frac{h}{2} \left. \frac{df}{dy} \right|_i \right) y_i + hf_i$$

$$\Rightarrow y_{i+1} = y_i + h \left(1 - \frac{h}{2} \frac{df}{dy}\Big|_i\right)^{-1} f_i$$

$$\frac{dc}{dt}=-kc^2$$
 with $c_0=1$ mol m⁻³, $k=1$ m³ mol⁻¹ s⁻¹, $t_{\rm end}=2$ s (Analytical solution: $c(t)=\frac{c_0}{1+kc_0t}$).

$$\frac{dc}{dt}=-kc^2$$
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Define
$$f = -kc^2$$
, then $\frac{df}{dc} = -2kc$.

Second order reaction in a batch reactor:

$$\frac{dc}{dt}=-kc^2$$
 with $c_0=1$ mol m⁻³, $k=1$ m³ mol⁻¹ s⁻¹, $t_{\rm end}=2$ s (Analytical solution: $c(t)=\frac{c_0}{1+kc_0t}$).

Define
$$f = -kc^2$$
, then $\frac{df}{dc} = -2kc$.

Substitution:

$$c_{i+1} = c_i + h \left(1 - \frac{h}{2} \cdot (-2kc_i) \right)^{-1} \cdot (-kc_i^2)$$

$$= c_i - \frac{hkc_i^2}{1 + hkc_i} = \frac{c_i + hkc_i^2 - hkc_i^2}{1 + hkc_i} \Rightarrow c_{i+1} = \frac{c_i}{1 + hkc_i}$$

Second order reaction in a batch reactor:

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 with $c_0=1~\text{mol}\,\text{m}^{-3},~k=1~\text{m}^3\,\text{mol}^{-1}\,\text{s}^{-1},~t_{\text{end}}=2~\text{s}$ (Analytical solution: $c(t)=\frac{c_0}{1+kc_0t}$).

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You will find that this method is exact for all step sizes h because of the quadratic source term!

$$\frac{dc}{dt}=-kc^2$$
 with $c_0=1~\rm mol\,m^{-3},~k=1~m^3\,mol^{-1}\,s^{-1},~t_{\rm end}=2~\rm s$ Analytical solution: $c(t)=\frac{c_0}{1+kc_0t}$

$$c_{i+1} = \frac{c_i}{1 + hkc_i}$$

$$\frac{dc}{dt}=-kc^2$$
 with $c_0=1~{\rm mol\,m^{-3}},~k=1~{\rm m^3\,mol^{-1}\,s^{-1}},~t_{\rm end}=2~{\rm s}$ Analytical solution: $c(t)=\frac{c_0}{1+kc_0t}$

$$c_{i+1} = \frac{c_i}{1 + hkc_i}$$

Ν	ζ	$rac{\zeta_{ ext{numerical}} - \zeta_{ ext{analytical}}}{\zeta_{ ext{analytical}}}$	$r = \frac{\log\left(\frac{\epsilon_i}{\epsilon_{i-1}}\right)}{\log\left(\frac{N_{i-1}}{N_i}\right)}$
20	0.6666666667	1.665×10^{-16}	_
40	0.6666666667	0	
80	0.6666666667	0	_
160	0.6666666667	0	-
320	0.6666666667	0	_

Third order reaction in a batch reactor Analytical solution: $c(t) = \frac{c_0}{\sqrt{1+2kc_0^2t}}$

$$c_{i+1} = c_i - \frac{hkc_i^3}{1 + \frac{3}{2}hkc_i^2}$$

Third order reaction in a batch reactor Analytical solution: $c(t) = \frac{c_0}{\sqrt{1+2kc_0^2t}}$

$$c_{i+1} = c_i - \frac{hkc_i^3}{1 + \frac{3}{2}hkc_i^2}$$

Ν	ζ	$\frac{\zeta_{numerical} - \zeta_{analytical}}{\zeta_{analytical}}$	$r = \frac{\log\left(\frac{\epsilon_i}{\epsilon_{i-1}}\right)}{\log\left(\frac{N_{i-1}}{N_i}\right)}$
20	0.5526916174	1.71×10^{-4}	_
40	0.5527633731	4.17×10^{-5}	2.041
80	0.5527807304	1.03×10^{-5}	2.021
160	0.5527849965	2.55×10^{-6}	2.011
320	0.5527860538	6.34×10^{-7}	2.005

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Forward Euler Convergence rate Runge-Kutta methods Step size control

- Mackward Euler Implicit midpoint method
- 4 Boundary value problems Shooting method
- **5** Systems of ODEs

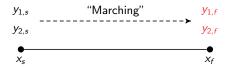
Stiff systems of ODEs
Solving systems of ODEs in Matlab

Conclusion

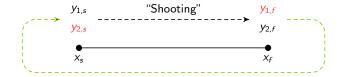
Importance of boundary conditions

The nature of boundary conditions determines the appropriate numerical method. Classification into 2 main categories:

Initial value problems (IVP) We know the values of all y_i at some starting position x_s , and it is desired to find the values of y_i at some final point x_f .

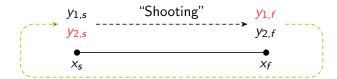


Boundary value problems (BVP) Boundary conditions are specified at more than one x. Typically, some of the BC are specified at x_s and the remainder at x_f .



Shooting method

How to solve a BVP using the shooting method:



- Define the system of ODEs
- Provide an initial guess for the unknown boundary condition
- Solve the system and compare the resulting boundary condition to the expected value
- Adjust the guessed boundary value, and solve again. Repeat until convergence.
 - Of course, you can subtract the expected value from the computed value at the boundary, and use a non-linear root finding method

Consider a chemical reaction in a liquid film layer of thickness δ :

$$\mathcal{D} \frac{d^2c}{dx^2} = k_Rc \text{ with } egin{array}{c} c(x=0) = C_{A,i,L} = 1 \ c(x=\delta) = 0 \end{array} \qquad ext{(interface concentration)}$$

Question: compute the concentration profile in the film layer.

Consider a chemical reaction in a liquid film layer of thickness δ :

$$\mathcal{D} \frac{d^2c}{dx^2} = k_Rc$$
 with $c(x=0) = C_{A,i,L} = 1$ (interface concentration) $c(x=\delta) = 0$ (bulk concentration)

Question: compute the concentration profile in the film layer.

Step 1: Define the system of ODEs

This second-order ODE can be rewritten as a system of first-order ODEs, if we define the flux q as:

$$q = -\mathcal{D}\frac{dc}{dx}$$

Now, we find:

$$\frac{dc}{dx} = -\frac{1}{\mathcal{D}}q$$

$$\frac{dq}{dx} = -k_R c$$

Consider a chemical reaction in a liquid film layer of thickness δ :

$$\mathcal{D}\frac{d^2c}{dx^2} = k_Rc \text{ with } \begin{cases} c(x=0) = C_{A,i,L} = 1 \\ c(x=\delta) = 0 \end{cases} \qquad \text{(interface concentration)}$$

Question: compute the concentration profile in the film layer.

Step 2: Set the boundary conditions

The boundary conditions for the concentrations at x=0 and $x=\delta$ are known.

The flux at the interface, however, is not known, and should be solved for.

$$\frac{dc}{dx} = -\frac{1}{\mathcal{D}}q$$

$$\frac{dq}{dx} = -k_R c$$

Solving the two first-order ODEs in Excel. First, the cells with constants:

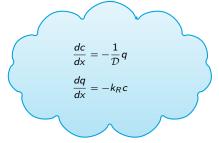
	Α	В	С
1	CAiL	1	ml/m3
2	D	1e-8	m2/s
3	kR	10	1/s
4	delta	1e-4	m
5	N	100	
6	dx	=B4/B5	

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Now, we program the forward Euler (explicit) schemes for $\it c$ and $\it q$ below:

	А	В	С
10	×	С	q
11	0	=B1	10
12	=A11+\$B\$6	=B11+\$B\$6*(-1/\$B\$2*C11)	=C11+\$B\$6*(-\$B\$3*B11)
13	=A12+\$B\$6	=B12+\$B\$6*(-1/\$B\$2*C12)	=C12+\$B\$6*(-\$B\$3*B12)
111	=A110+\$B\$6	=B110+\$B\$6*(-1/\$B\$2*C110)	=C110+\$B\$6*(-\$B\$3*B110)

- We now have profiles for c and q as a function of position x.
- The concentration $c(x = \delta)$ depends (eventually) on the boundary condition at the interface q(x = 0)
- We can use the solver to change q(x=0) such that the concentration at the bulk meets our requirement: $c(x=\delta)=0$

We first program the system of ODEs in a separate function:

$$\frac{d}{dx} = -\frac{1}{D}q$$

$$\frac{dq}{dx} = -k_R c$$

$$function [dxdt] = BVPODE(t,x,ps)$$

$$dxdt(1) = -1/ps.D*x(2);$$

$$dxdt(2) = -ps.kR*x(1);$$

$$dxdt = dxdt';$$

$$return$$

We first program the system of ODEs in a separate function:

$$\begin{aligned} \frac{dc}{dx} &= -\frac{1}{\mathcal{D}}q \\ \frac{dq}{dx} &= -k_R c \\ \\ \text{function [dxdt] = BVPODE(t,x,ps)} \\ \text{dxdt(1) = -1/ps.D*x(2);} \\ \text{dxdt(2) = -ps.kR*x(1);} \\ \text{dxdt=dxdt';} \\ \text{return} \end{aligned}$$

Note that we pass a variable (type: struct) that contains required parameters: ps.

The ODE function is solved via ode45, after setting a number of initial and boundary conditions:

```
function f = RunBVP(bcq,ps)
[x,cq] = ode45(@BVPODE,[0 ps.delta],[1 bcq], [], ps);
f = cq(end,1) - 0;
plotyy(x,cq(:,1),x,cq(:,2));
return;
```

The ODE function is solved via ode45, after setting a number of initial and boundary conditions:

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return;
```

Note the following:

- We use the interval $0 \le x \le \delta$
- Boundary conditions are given as: c(x = 0) = 1 and q(x = 0) = bcq, which is given as an argument to the function (i.e. changable from 'outside'!)
- The function returns f, the difference between the computed and desired concentration at $x = \delta$.

Finally, we should solve the system so that we obtain the right boundary condition q = bcq such that $c(x = \delta) = 0$. We can use the built-in function fzero to do this

```
% Parameter definition
ps.D=1e-8;
ps.kR=10;
ps.delta=1e-4;

% Solve for flux boundary condition (initial guess: 0)
opt = optimset('Display','iter');
flux = fzero(@RunBVP,0,opt,ps);
```

BVP example: analytical solution

Compare with the analytical solution:

$$q = k_L E_A C_{A,i,L} \quad \text{with}$$

$$E_A = \frac{Ha}{\tanh Ha} \qquad \text{(Enhancement factor)}$$

$$Ha = \frac{\sqrt{k_R \mathcal{D}}}{k_L} \qquad \text{(Hatta number)}$$

$$k_L = \frac{\mathcal{D}}{\delta} \qquad \text{(mass transfer coefficient)}$$

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A system of ODEs is specified using vector notation:

$$\frac{d\mathbf{y}}{dx} = \mathbf{f}(x, \mathbf{y}(x))$$

for

$$\frac{dy_1}{dx} = f_1(x, y_1(x), y_2(x))$$
 or $f_1(x, y_1, y_2)$

$$\frac{dy_1}{dx} = f_1(x, y_1(x), y_2(x)) \quad \text{or} \quad f_1(x, y_1, y_2)$$

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The solution techniques discussed before can also be used to solve systems of equations.

Forward Euler method

$$\mathbf{y}_{i+1} = \mathbf{y}_i + h\mathbf{f}(x_i, \mathbf{y}_i)$$

Improved Euler method (classical RK2)

$$\mathbf{y}_{i+1} = \mathbf{y}_i + \frac{h}{2}(\mathbf{k}_1 + \mathbf{k}_2)$$
 using $\mathbf{k}_1 = \mathbf{f}(x_i, \mathbf{y}_i)$
 $\mathbf{k}_2 = \mathbf{f}(x_i + h, \mathbf{y}_i + h\mathbf{k}_1)$

Modified Euler method (midpoint rule)

$$\mathbf{y}_{i+1} = \mathbf{y}_i + h\mathbf{k}_2$$
 using $\mathbf{k}_1 = \mathbf{f}(x_i, \mathbf{y}_i)$
 $\mathbf{k}_2 = \mathbf{f}(x_i + \frac{h}{2}, \mathbf{y}_i + \frac{h}{2}\mathbf{k}_1)$

Classical fourth order Runge-Kutta method (RK4)

$$\mathbf{y}_{i+1} = \mathbf{y}_i + h\left(\frac{\mathbf{k}_1}{6} + \frac{1}{3}(\mathbf{k}_2 + \mathbf{k}_3) + \frac{\mathbf{k}_4}{6}\right)$$

$$\mathbf{k}_1 = \mathbf{f}(x_i, \mathbf{y}_i)$$

$$\mathbf{k}_2 = \mathbf{f}(x_i + \frac{h}{2}, \mathbf{y}_i + \frac{h}{2}\mathbf{k}_1)$$

using

$$\mathbf{k}_3 = \mathbf{f}(x_i + \frac{h}{2}, \mathbf{y}_i + \frac{h}{2}\mathbf{k}_2)$$

$$\mathbf{k}_4 = \mathbf{f}(x_i + h, \mathbf{y}_i + h\mathbf{k}_3)$$

Backward Euler method

$$\mathbf{y}_{i+1} = \mathbf{y}_i + h \left(1 - h \left. \frac{d\mathbf{f}}{d\mathbf{y}} \right|_i \right)^{-1} \mathbf{f}(\mathbf{y}_i)$$

Implicit midpoint method

$$\mathbf{y}_{i+1} = \mathbf{y}_i + h \left(1 - \frac{h}{2} \left. \frac{d\mathbf{f}}{d\mathbf{y}} \right|_i \right)^{-1} \mathbf{f}(\mathbf{y}_i)$$

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Stiff systems of ODEs

A system of ODEs can be stiff and require a different solution method.

Stiff systems of ODEs

A system of ODEs can be stiff and require a different solution method. For example:

$$\frac{dc_1}{dt} = 998c_1 + 1998c_2 \qquad \frac{dc_2}{dt} = -999c_1 - 1999c_2$$

with boundary conditions $c_1(t=0)=1$ and $c_2(t=0)=0$. The analytical solution is:

$$c_1 = 2e^{-t} - e^{-1000t}$$
 $c_2 = -e^{-t} + e^{-1000t}$

For the explicit method we require $\Delta t < 10^{-3}$ despite the fact that the term is completely negligible, but essential to keep stability.

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For the explicit method we require $\Delta t < 10^{-3}$ despite the fact that the term is completely negligible, but essential to keep stability.

> The "disease" of stiff equations: we need to follow the solution on the shortest length scale to maintain stability of the integration, although accuracy requirements would allow a much larger time step.

Forward Euler (explicit)

$$\begin{aligned} \frac{dc_{1,i+1} - c_{1,i}}{dt} &= 998c_{1,i} + 1998c_{2,i} \\ \frac{dc_{2,i+1} - c_{2,i}}{dt} &= -999c_{1,i} - 1999c_{2,i} \\ &\Rightarrow \frac{c_{1,i+1} = (1 + 998\Delta t) c_{1,i} + 1998\Delta t c_{2,i}}{c_{2,i+1} = -999\Delta t c_{1,i} + (1 - 1999\Delta t) c_{2,i}} \end{aligned}$$

Backward Euler (implicit)

$$\begin{split} \frac{dc_{1,i+1} - c_{1,i+1}}{dt} &= 998c_{1,i+1} + 1998c_{2,i+1} \\ \frac{dc_{2,i+1} - c_{2,i+1}}{dt} &= -999c_{1,i+1} - 1999c_{2,i+1} \\ &\Rightarrow \frac{(1 - 998\Delta t)c_{1,i+1} - 1998\Delta tc_{2,i} = c_{1,i}}{999\Delta tc_{1,i+1} + (1 + 999\Delta t)c_{2,i+1} = c_{2,i}} \end{split}$$

Backward Euler (implicit)

$$\frac{dc_{1,i+1} - c_{1,i+1}}{dt} = 998c_{1,i+1} + 1998c_{2,i+1}$$

$$\frac{dc_{2,i+1} - c_{2,i+1}}{dt} = -999c_{1,i+1} - 1999c_{2,i+1}$$

$$\Rightarrow \frac{(1 - 998\Delta t) c_{1,i+1} - 1998\Delta t c_{2,i} = c_{1,i}}{999\Delta t c_{1,i+1} + (1 + 999\Delta t) c_{2,i+1} = c_{2,i}}$$

$$Ac_{i+1} = c_i \text{ with } A = \begin{pmatrix} 1 - 998\Delta t & -1998\Delta t \\ 999\Delta t & 1 + 1999\Delta t \end{pmatrix} \text{ and } \mathbf{b} = \begin{pmatrix} c_{1,i} \\ c_{2,i} \end{pmatrix}$$

Backward Euler (implicit)
$$A\mathbf{c}_{i+1} = \mathbf{c}_i$$
 with $A = \begin{pmatrix} 1 - 998\Delta t & -1998\Delta t \\ 999\Delta t & 1 + 1999\Delta t \end{pmatrix}$ and $\mathbf{b} = \begin{pmatrix} c_{1,i} \\ c_{2,i} \end{pmatrix}$

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Cramers rule:

$$c_{1,i+1} = \frac{\begin{vmatrix} c_{1,i} & -1998\Delta t \\ c_{2,i} & 1+1999\Delta t \end{vmatrix}}{\det(A)} = \frac{\frac{(1+1999\Delta t)c_{1,i}+1998\Delta tc_{2,i}}{(1-998\Delta t)(1+1999\Delta t)+1998\cdot999\Delta t^2}}{\begin{vmatrix} 1-998\Delta t & c_{1,i} \\ 999\Delta t & c_{2,i} \end{vmatrix}}{\det(A)} = \frac{\frac{-999\Delta tc_{1,i}+(1-998\Delta t)c_{2,i}}{(1-998\Delta t)(1+1999\Delta t)+1998\cdot999\Delta t^2}}$$

Forward Euler: $\Delta t \leq 0.001$ for stability Backward Euler: always stable, even for $\Delta t > 100$ (but then not very accurate!)

Cure for stiff problems: use implicit methods! To find out whether your system is stiff: check whether one of the eigenvalues have an imaginary part

Solving systems of ODEs in Matlab

Matlab provides convenient procedures to solve (systems of) ODEs automatically.

The procedure is as follows:

- 1 Create a function that specifies the ODEs. Specifically, this function returns the $\frac{d\mathbf{y}}{d\mathbf{x}}$ vector.
- 2 Initialise solver variables and settings (e.g. step size, initial conditions, tolerance), in a separate script
- 3 Call the ODE solver function, using a function handle to the ODE function described in point 1.
 - The ODE solver will return the vector for the independent variable, and a solution vector (matrix for systems of ODEs).

We solve the system: $\frac{dx_1}{dt} = -x_1 - x_2$, $\frac{dx_2}{dt} = x_1 - 2x_2$

Create an ODE function

```
function [dxdt] = myODEFunction(t,x)
dxdt(1) = -x(1) - x(2);
dxdt(2) = x(1) - 2*x(2);
dxdt=dxdt'; % Transpose to column vector
return
```

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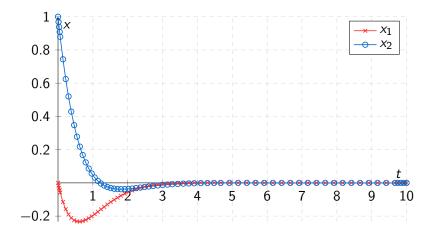
Create a solution script

```
options = odeset('RelTol',1e-4,'AbsTol',[1e-4 1e-4]);
[t,x] = ode45(@myODEfunction,tspan,x_init,options);
```

Plot the solution:

```
plot(t,x(:,1),'r-x',t,x(:,2),'b-o')
```

Plot the solution:



A few notes on working with ode45 and other solvers. If we want to give additional arguments (e.g. a, b and c) to our ODE function, we can list them in the function line:

```
function [dxdt] = myODE(t,x,a,b,c)
```

The additional arguments can now be set in the solver script by adding them after the options:

```
[t,x] = ode45(@myODE,tspan,x_0,options,a,b,c);
```

A few notes on working with ode45 and other solvers. If we want to give additional arguments (e.g. a, b and c) to our ODE function, we can list them in the function line:

```
function [dxdt] = myODE(t,x,a,b,c)
```

The additional arguments can now be set in the solver script by adding them after the options:

```
[t,x] = ode45(@myODE,tspan,x_0,options,a,b,c);
```

 Of course, in the solver script, the variables do not need to be called a, b and c:

```
[t,x] = ode45 (@myODE, tspan, x_0, options, k1, phi, V);
```

 These variables may be of any type (vectors, matrix, struct). Especially a struct is useful to carry many values in 1 variable.

You may have noticed that the step size in t varied. This is because we have given the begin and end times of our time span:

```
tspan = [0 10];
```

You may have noticed that the step size in t varied. This is because we have given the begin and end times of our time span:

```
tspan = [0 10];
```

You can also solve at specific steps, by supplying all steps explicitly, e.g.:

```
tspan = linspace(0,10,101);
```

This example provides 101 explicit time steps between 0 and 10 seconds.

Note that you may affect the efficiency and accuracy of the solver algorithm by doing this!

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Other methods

Other explicit methods:

 Burlisch-Stoer method (Richardson extrapolation + modified midpoint method)

Other implicit methods:

- Rosenbrock methods (higher order implicit Runge-Kutta methods)
- Predictor-corrector methods

- Several solution methods and their derivation were discussed:
 - Explicit solution methods: Euler, Improved Euler, Midpoint method, RK45
 - Implicit methods: Implicit Euler and Implicit midpoint method
 - A few examples of their spreadsheet implementation were shown
- We have paid attention to accuracy and instability, rate of convergence and step size
- Systems of ODEs can be solved by the same algorithms. Stiff problems should be treated with care.
- An example of solving ODEs with Matlab was demonstrated.