

# Ordinary differential equations

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# Part I

## Explicit and Implicit methods

# Today's outline

- ① Introduction
- ② Euler's method  
Forward Euler
- ③ Rates of convergence
- ④ Runge-Kutta methods  
RK2 methods  
RK4 method
- ⑤ Step size control
- ⑥ Implicit methods  
Backward Euler  
Implicit midpoint method

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  - RK2 methods
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# Overview

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An equation containing a function of one independent variable and its derivatives, in contrast to a *partial differential equation*, which contains derivatives with respect to more independent variables.

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An equation containing a function of one independent variable and its derivatives, in contrast to a *partial differential equation*, which contains derivatives with respect to more independent variables.

## Main question

How to solve

$$\frac{dy}{dx} = f(y(x), x) \quad \text{with} \quad y(x=0) = y_0$$

accurately and efficiently?

# What is an ODE?

- Algebraic equation:

$$f(y(x), x) = 0 \quad \text{e.g. } -\ln(K_{eq}) = (1 - \zeta)$$

- First order ODE:

$$f\left(\frac{dy}{dx}(x), y(x), x\right) = 0 \quad \text{e.g. } \frac{dc}{dt} = -kc^n$$

- Second order ODE:

$$f\left(\frac{d^2y}{dx^2}(x), \frac{dy}{dx}(x), y(x), x\right) = 0 \quad \text{e.g. } \mathcal{D}\frac{d^2c}{dx^2} = -\frac{kc}{1 + Kc}$$

## About second order ODEs

Very often a second order ODE can be rewritten into a system of first order ODEs (whether it is handy depends on the boundary conditions!)



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### Example

Recall:

$$\mathcal{D} \frac{d^2 c}{dx^2} = -\frac{kc}{1 + Kc}$$

Define  $y = -\mathcal{D} \frac{dc}{dx}$ , then  $\frac{dy}{dx} = \frac{kc}{1 + Kc}$ , thus solve system:

$$\begin{aligned} \frac{dc}{dx} &= -\frac{1}{\mathcal{D}} y \\ \frac{dy}{dx} &= \frac{kc}{1 + Kc} \end{aligned}$$

## About second order ODEs

Very often a second order ODE can be rewritten into a system of first order ODEs (whether it is handy depends on the boundary conditions!)

### More general

Consider the second order ODE:

$$\frac{d^2 y}{dx^2} + q(x) \frac{dy}{dx} = r(x)$$

Now define and solve using  $z$  as a new variable:

$$\frac{dy}{dx} = z(x)$$

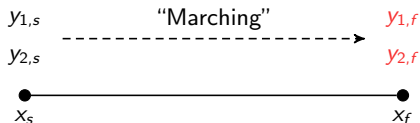
$$\frac{dz}{dx} = r(x) - q(x)z(x)$$

# Importance of boundary conditions

The nature of boundary conditions determines the appropriate numerical method. Classification into 2 main categories:

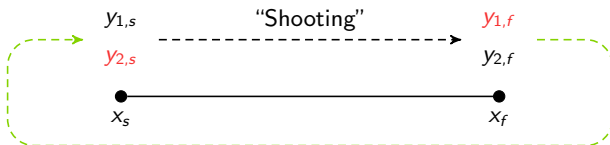
- *Initial value problems (IVP)*

We know the values of all  $y_i$  at some starting position  $x_s$ , and it is desired to find the values of  $y_i$  at some final point  $x_f$ .



- *Boundary value problems (BVP)*

Boundary conditions are specified at more than one  $x$ . Typically, some of the BC are specified at  $x_s$  and the remainder at  $x_f$ .



# Overview

## Initial value problems:

- Explicit methods
  - First order: forward Euler
  - Second order: improved Euler (RK2)
  - Fourth order: Runge-Kutta 4 (RK4)
  - Step size control
- Implicit methods
  - First order: backward Euler
  - Second order: midpoint rule

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## Boundary value problems

- Shooting method

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Easiest solution algorithm: Euler's method, derived here via Taylor series expansion:

$$c(t_0 + \Delta t) \approx c(t_0) + \left. \frac{dc}{dt} \right|_{t_0} \Delta t + \frac{1}{2} \left. \frac{d^2c}{dt^2} \right|_{t_0} (\Delta t)^2 + \mathcal{O}(\Delta t^3)$$

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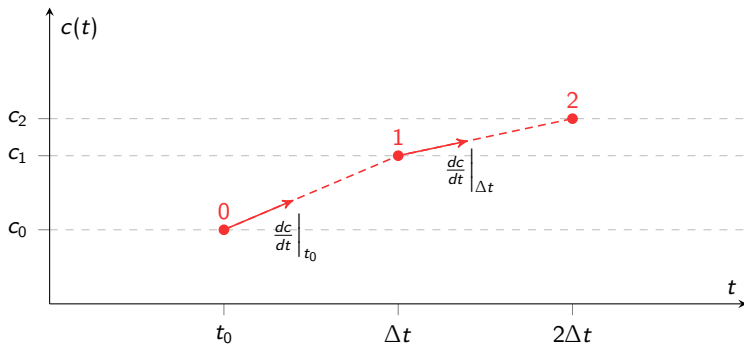
Neglect terms with higher order than two:  $\left. \frac{dc}{dt} \right|_{t_0} = \frac{c(t_0 + \Delta t) - c(t_0)}{\Delta t}$

Substitution:

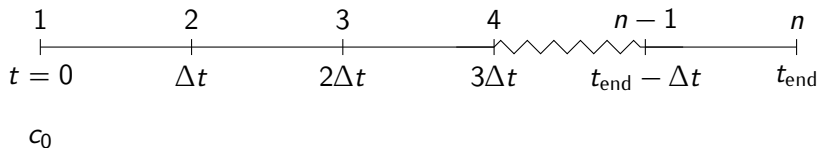
$$\frac{c(t_0 + \Delta t) - c(t_0)}{\Delta t} = f(c_0, t_0) \Rightarrow c(t_0 + \Delta t) = c(t_0) + \Delta t f(c_0, t_0)$$

# Euler's method: graphical example

$$\frac{c(t_0 + \Delta t) - c(t_0)}{\Delta t} = f(c_0, t_0) \Rightarrow c(t_0 + \Delta t) = c(t_0) + \Delta t f(c_0, t_0)$$

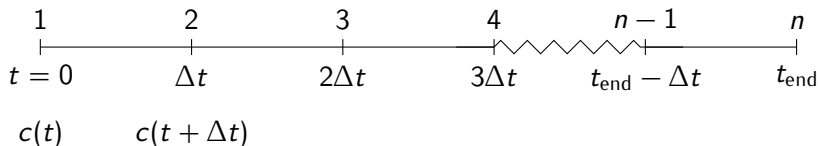


# Euler's method - solution method



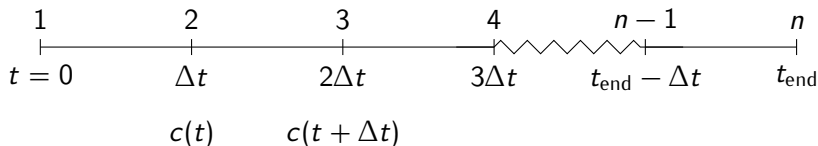
Start with  $t = t_0$ ,  $c = c_0$ , then calculate at discrete points in time:  
 $c(t_1 = t_0 + \Delta t) = c(t_0) + \Delta t f(c_0, t_0)$ .

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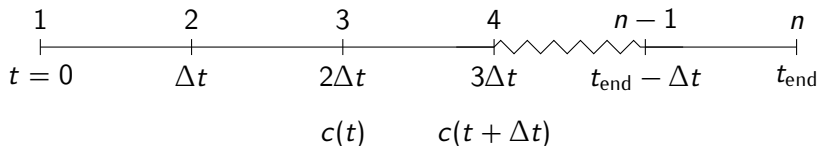
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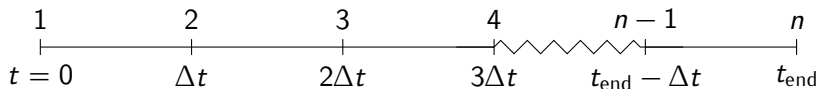


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 $c(t_1 = t_0 + \Delta t) = c(t_0) + \Delta t f(c_0, t_0)$ .

Pseudo-code Euler's method:  $\frac{dy}{dx} = f(x, y)$  and  $y(x_0) = y_0$ .

- 1 Initialize variables, functions; set  $h = \frac{x_1 - x_0}{N}$
- 2 Set  $x = x_0$ ,  $y = y_0$
- 3 While  $x < x_{\text{end}}$  do  
 $x_{i+1} = x_i + h$ ;  $y_{i+1} = y_i + hf(x_i, y_i)$

## Euler's method - example

First order reaction in a batch reactor:

$$\frac{dc}{dt} = -kc \quad \text{with} \quad c(t=0) = 1 \text{ mol m}^{-3}, \quad k = 1 \text{ s}^{-1}, \quad t_{\text{end}} = 2 \text{ s}$$

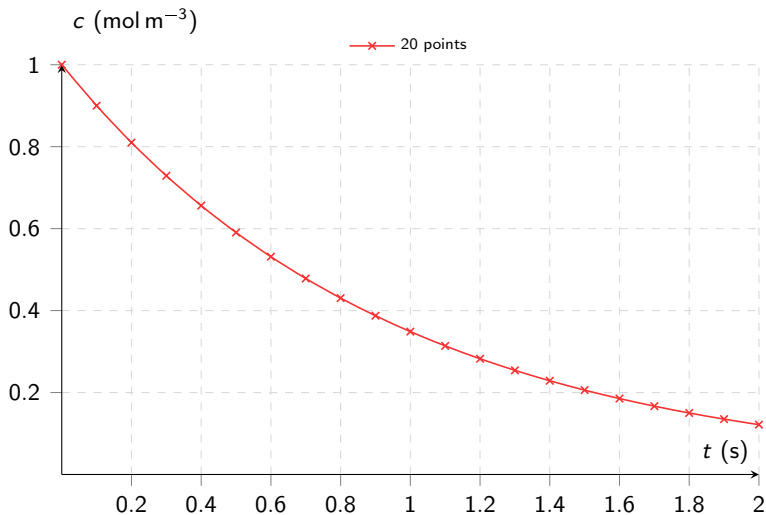
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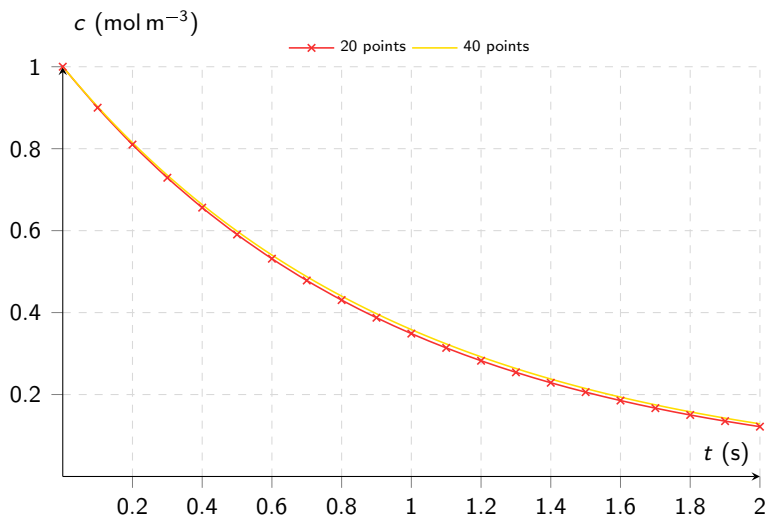
$$\frac{dc}{dt} = -kc \quad \text{with} \quad c(t=0) = 1 \text{ mol m}^{-3}, \quad k = 1 \text{ s}^{-1}, \quad t_{\text{end}} = 2 \text{ s}$$

Time [s]	Concentration [mol m <sup>-3</sup> ]
$t_0 = 0$	$c_0 = 1.00$
$t_1 = t_0 + \Delta t$	$c_1 = c_0 + \Delta t \cdot (-kc_0)$
$= 0 + 0.1 = 0.1$	$= 1 + 0.1 \cdot (-1 \cdot 1) = 0.9$
$t_2 = t_1 + \Delta t$	$c_2 = c_1 + \Delta t \cdot (-kc_1)$
$= 0.1 + 0.1 = 0.2$	$= 0.9 + 0.1 \cdot (-1 \cdot 0.9) = 0.81$
$t_3 = t_2 + \Delta t$	$c_3 = c_2 + \Delta t \cdot (-kc_2)$
$= 0.2 + 0.1 = 0.3$	$= 0.81 + 0.1 \cdot (-1 \cdot 0.81) = 0.729$
...	...
$t_{i+1} = t_i + \Delta t$	$c_{i+1} = c_i + \Delta t \cdot (-kc_i)$
...	...
$t_{20} = 2.0$	$c_{20} = c_{19} + \Delta t \cdot (-kc_{19}) = 0.1211577$

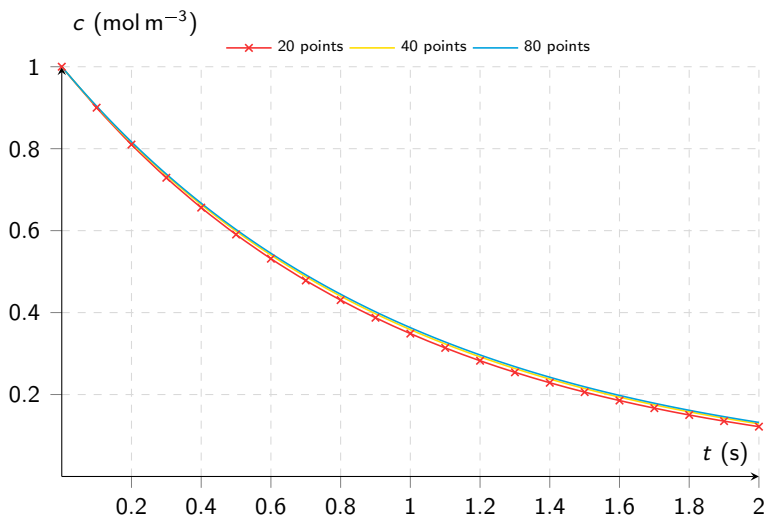
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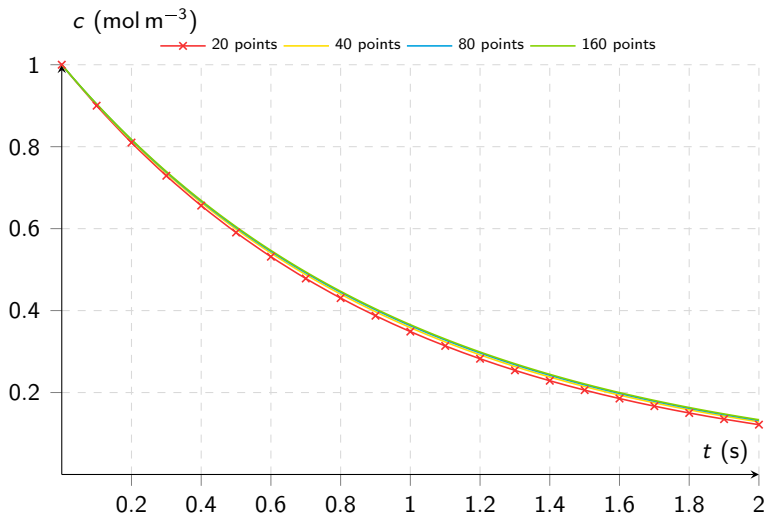
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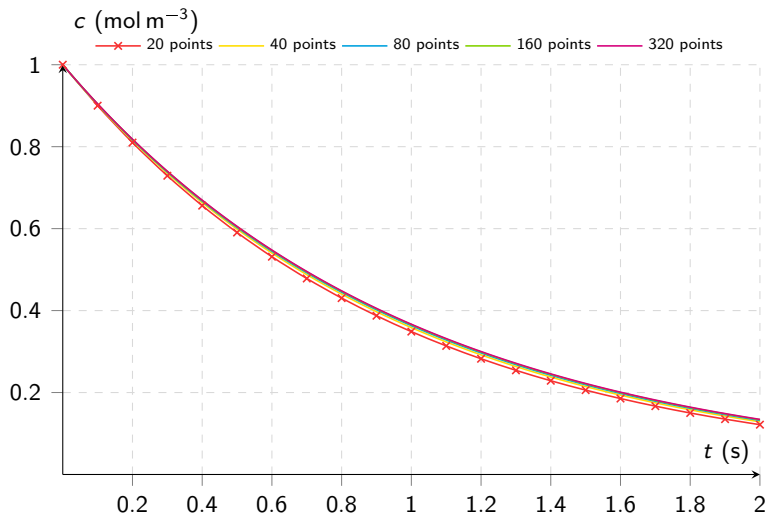
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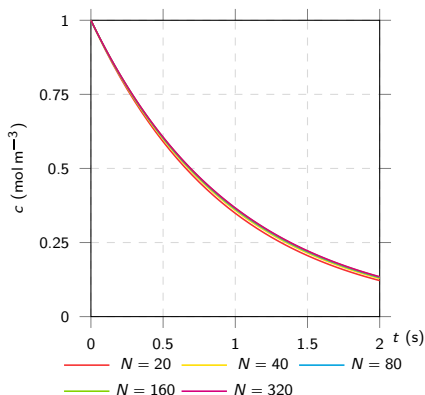




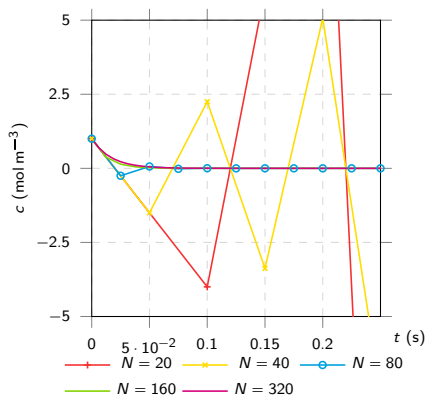
# Problems with Euler's method

The question is: What step size, or how many steps to use?

- 1 *Accuracy*  $\Rightarrow$  need information on numerical error!
- 2 *Stability*  $\Rightarrow$  need information on stability limits!



Reaction rate:  $k = 1 \text{ s}^{-1}$



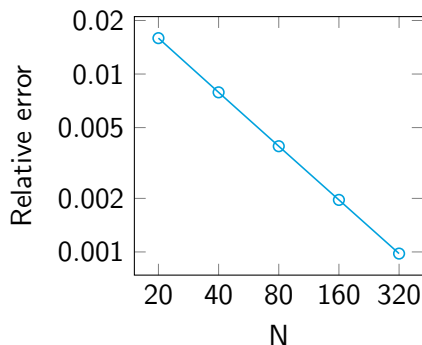
Reaction rate:  $k = 50 \text{ s}^{-1}$

# Accuracy

Comparison with analytical solution for  $k = 1 \text{ s}^{-1}$ :

$$c(t) = c_0 \exp(-kt) \Rightarrow \zeta = 1 - \exp(-kt) \Rightarrow \zeta_{\text{analytical}} = 0.864665$$

$N$	$\zeta$	$\frac{\zeta_{\text{numerical}} - \zeta_{\text{analytical}}}{\zeta_{\text{analytical}}}$
20	0.878423	0.015912
40	0.871488	0.007891
80	0.868062	0.003929
160	0.866360	0.001961
320	0.865511	0.000979



# Accuracy

For Euler's method: Error halves when the number of grid points is doubled, i.e. error is proportional to  $\Delta t$ : first order method.

Error estimate:

$$\left. \frac{dx}{dt} \right|_{t_0} = \frac{x(t_0 + \Delta t) - x(t_0)}{\Delta t} + \frac{1}{2} \left. \frac{d^2x}{dt^2} \right|_{t_0} (\Delta t) + \mathcal{O}(\Delta t)^2$$

$$\frac{x(t_0 + \Delta t) - x(t_0)}{\Delta t} = f(x_0, t_0) - \frac{1}{2} \left. \frac{d^2x}{dt^2} \right|_{t_0} (\Delta t) + \mathcal{O}(\Delta t)^2$$

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# Errors and convergence rate

## $L_2$ norm (Euclidean norm)

$$\|\mathbf{v}\|_2 = \sqrt{v_1^2 + v_2^2 + \dots + v_n^2} = \sqrt{\sum_{i=1}^n v_i^2}$$

## $L_\infty$ norm (maximum norm)

$$\|\mathbf{v}\|_\infty = \max(|v_1|, \dots, |v_n|)$$

## Absolute difference

$$\epsilon_{\text{abs}} = \left\| \mathbf{y}_{\text{numerical}} - \mathbf{y}_{\text{analytical}} \right\|_{2,\infty}$$

## Relative difference

$$\epsilon_{\text{rel}} = \frac{\left\| \mathbf{y}_{\text{numerical}} - \mathbf{y}_{\text{analytical}} \right\|_{2,\infty}}{\left\| \mathbf{y}_{\text{analytical}} \right\|_{2,\infty}}$$

# Errors and convergence rate

Convergence rate (or: order of convergence)  $r$

$$\epsilon = \lim_{\Delta x \rightarrow 0} c(\Delta x)^r$$

- A first order method reduces the error by a factor 2 when increasing the number of steps by a factor 2
- A second order method reduces the error by a factor 4 when increasing the number of steps by a factor 2

# Computing the rate of convergence

When the analytical solution is available, choose ① or ② for a particular number of grid points  $N$ :

- ① Compute the relative or absolute error vector  $\bar{\epsilon}$ . Take the norm to compute a single error value  $\epsilon$  following:

- Based on  $L_1$ -norm:  $\epsilon = \frac{\|\bar{\epsilon}\|_1}{N}$

- Based on  $L_2$ -norm:  $\epsilon = \frac{\|\bar{\epsilon}\|_2}{\sqrt{N}}$

- Based on  $L_\infty$ -norm:  $\epsilon = \|\bar{\epsilon}\|_\infty$

- ② Compute the relative or absolute error at a single indicative points (e.g. middle of domain, outlet).

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- Based on  $L_\infty$ -norm:  $\epsilon = \|\bar{\epsilon}\|_\infty$

- ② Compute the relative or absolute error at a single indicative points (e.g. middle of domain, outlet).

Compare to calculations with different number of steps:  $\epsilon_1 = c(\Delta x_1)^r$  and  $\epsilon_2 = c(\Delta x_2)^r$  and solve for  $r$ :

$$\frac{\epsilon_2}{\epsilon_1} = \frac{c(\Delta x_2)^r}{c(\Delta x_1)^r} = \left( \frac{\Delta x_2}{\Delta x_1} \right)^r \Rightarrow \log \left( \frac{\epsilon_2}{\epsilon_1} \right) = \log \left( \frac{\Delta x_2}{\Delta x_1} \right)^r$$

$$\Rightarrow r = \frac{\log \left( \frac{\epsilon_2}{\epsilon_1} \right)}{\log \left( \frac{\Delta x_2}{\Delta x_1} \right)} = \frac{\log \left( \frac{\epsilon_2}{\epsilon_1} \right)}{\log \left( \frac{N_1}{N_2} \right)} \quad \text{in the limit of } \Delta x \rightarrow 0 \quad \text{or} \quad N \rightarrow \infty$$



# Computing the rate of convergence

When the analytical solution is not available:

- ① Compute the solution with  $N + 1$ ,  $N$ ,  $N - 1$  and  $N - 2$  grid points
- ② Select a single indicative grid point (e.g. middle of domain, outlet) that lies at exactly the same position in each computation
- ③ Use the solution  $c$  at this grid point for various grid sizes to compute:

$$r = \frac{\log \frac{c_{N+1} - c_N}{c_N - c_{N-1}}}{\log \frac{c_N - c_{N-1}}{c_{N-1} - c_{N-2}}}$$

# Example: Euler's method — order of convergence

$N$	$\zeta$	$\frac{\zeta_{\text{numerical}} - \zeta_{\text{analytical}}}{\zeta_{\text{analytical}}}$	$r = \frac{\log\left(\frac{\epsilon_i}{\epsilon_{i-1}}\right)}{\log\left(\frac{N_{i-1}}{N_i}\right)}$
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Wouldn't it be great to have a method that can give the answer using much less steps? ⇒ Higher order methods

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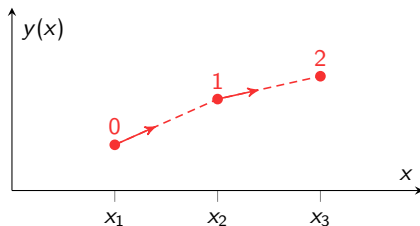
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# Runge-Kutta methods

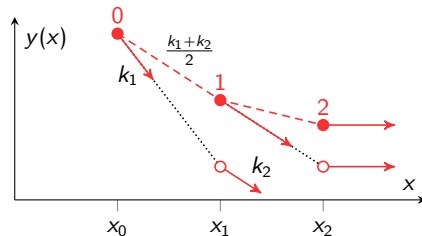
Propagate a solution by combining the information of several Euler-style steps (each involving one function evaluation) to match a Taylor series expansion up to some higher order.

Euler:  $y_{i+1} = y_i + hf(x_i, y_i)$  with  $h = \Delta x$ , i.e.  
slope =  $k_1 = f(x_i, y_i)$ .

Euler's method



RK2 method



# Classical second order Runge-Kutta (RK2) method

This method is also called Heun's method, or improved Euler method:

- 1 Approximate the slope at  $x_i$ :  $k_1 = f(x_i, y_i)$
- 2 Approximate the slope at  $x_{i+1}$ :  $k_2 = f(x_{i+1}, y_{i+1})$  where we use Euler's method to approximate  $y_{i+1} = y_i + hf(x_i, y_i) = y_i + hk_1$
- 3 Perform an Euler step with the average of the slopes:  
$$y_{i+1} = y_i + h\frac{1}{2}(k_1 + k_2)$$



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$$y_{i+1} = y_i + h\frac{1}{2}(k_1 + k_2)$$

In pseudocode:

```
x = x0, y = y0
while x < xend do
  xi+1 = xi + h
  k1 = f(xi, yi)
  k2 = f(xi + h, yi + hk1)
  yi+1 = yi + h $\frac{1}{2}$ (k1 + k2)
end while
```

# Runge-Kutta methods — derivation

$$\frac{dy}{dx} = f(x, y(x))$$

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Using Taylor series expansion:  $y_{i+1} = y_i + h \left. \frac{dy}{dx} \right|_i + \frac{h^2}{2} \left. \frac{d^2y}{dx^2} \right|_i + \mathcal{O}(h^3)$

$$\left. \frac{dy}{dx} \right|_i = f(x_i, y_i) \equiv f_i$$

$$\left. \frac{d^2y}{dx^2} \right|_i = \left. \frac{d}{dx} f(x, y(x)) \right|_i = \left. \frac{\partial f}{\partial x} \right|_i + \left. \frac{\partial f}{\partial y} \right|_i \left. \frac{dy}{dx} \right|_i = \left. \frac{\partial f}{\partial x} \right|_i + \left. \frac{\partial f}{\partial y} \right|_i f_i \quad (\text{chain rule})$$

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Substitution gives:

$$y_{i+1} = y_i + hf_i + \frac{h^2}{2} \left( \left. \frac{\partial f}{\partial x} \right|_i + \left. \frac{\partial f}{\partial y} \right|_i f_i \right) + \mathcal{O}(h^3)$$

$$y_{i+1} = y_i + \frac{h}{2} f_i + \frac{h}{2} \left( f_i + h \left. \frac{\partial f}{\partial x} \right|_i + hf_i \left. \frac{\partial f}{\partial y} \right|_i \right) + \mathcal{O}(h^3)$$

# Runge-Kutta methods — derivation

Note multivariate Taylor expansion:

$$\begin{aligned} f(x_i + h, y_i + k) &= f_i + h \left. \frac{\partial f}{\partial x} \right|_i + k \left. \frac{\partial f}{\partial y} \right|_i + \mathcal{O}(h^2) \\ \Rightarrow \frac{h}{2} \left( f_i + h \left. \frac{\partial f}{\partial x} \right|_i + h f_i \left. \frac{\partial f}{\partial y} \right|_i \right) &= \frac{h}{2} f(x_i + h, y_i + k f_i) + \mathcal{O}(h^3) \end{aligned}$$

Concluding:

$$y_{i+1} = y_i + \frac{h}{2} f_i + \frac{h}{2} f(x_i + h, y_i + k f_i) + \mathcal{O}(h^3)$$

Rewriting:

$$\begin{aligned} k_1 &= f(x_i, y_i) \\ k_2 &= f(x_i + h, y_i + h k_1) \\ \Rightarrow y_{i+1} &= y_i + \frac{h}{2} (k_1 + k_2) \end{aligned}$$

# Runge-Kutta methods — derivation

Generalization:  $y_{i+1} = y_i + h(b_1 k_1 + b_2 k_2) + \mathcal{O}(h^3)$

with  $k_1 = f_i$ ,  $k_2 = f(x_i + c_2 h, y_1 + a_{2,1} h k_1)$

(Note that classical RK2:  $b_1 = b_2 = \frac{1}{2}$  and  $c_2 = a_{2,1} = 1$ .)

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Bivariate Taylor expansion:

$$f(x_i + c_2 h, y_i + a_{2,1} h k_1) = f_i + c_2 h \left. \frac{\partial f}{\partial x} \right|_i + a_{2,1} h k_1 \left. \frac{\partial f}{\partial y} \right|_i + \mathcal{O}(h^2)$$

$$y_{i+1} = y_i + h(b_1 k_1 + b_2 k_2) + \mathcal{O}(h^3)$$

$$= y_i + h[b_1 f_i + b_2 f(x_i + c_2 h, y_i + a_{2,1} h k_1)] + \mathcal{O}(h^3)$$

$$= y_i + h \left[ b_1 f_i + b_2 \left\{ f_i + c_2 h \left. \frac{\partial f}{\partial x} \right|_i + a_{2,1} h k_1 \left. \frac{\partial f}{\partial y} \right|_i + \mathcal{O}(h^2) \right\} \right] + \mathcal{O}(h^3)$$

$$= y_i + h(b_1 + b_2) f_i + h^2 b_2 \left( c_2 \left. \frac{\partial f}{\partial x} \right|_i + a_{2,1} f_i \left. \frac{\partial f}{\partial y} \right|_i \right) + \mathcal{O}(h^3)$$

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Comparison with Taylor:

$$y_{i+1} = y_i + h f_i + \frac{h^2}{2} \left( \left. \frac{\partial f}{\partial x} \right|_i + \left. \frac{\partial f}{\partial y} \right|_i f_i \right) + \mathcal{O}(h^3)$$



# Runge-Kutta methods — derivation

Generalization:  $y_{i+1} = y_i + h(b_1 k_1 + b_2 k_2) + \mathcal{O}(h^3)$

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# Runge-Kutta methods — derivation

$$y_{i+1} = y_i + h(b_1 + b_2)f_i + h^2 b_2 \left( c_2 \frac{\partial f}{\partial x} \Big|_i + a_{2,1} f_i \frac{\partial f}{\partial y} \Big|_i \right) + \mathcal{O}(h^3)$$

$$y_{i+1} = y_i + hf_i + \frac{h^2}{2} \left( \frac{\partial f}{\partial x} \Big|_i + \frac{\partial f}{\partial y} \Big|_i f_i \right) + \mathcal{O}(h^3)$$

$\Rightarrow$  3 eqns and 4 unknowns  $\Rightarrow$  multiple possibilities!

- ① Classical RK2:

$$b_1 = b_2 = \frac{1}{2} \text{ and } c_2 = a_{2,1} = 1$$

- ② Midpoint rule (modified Euler):

$$b_1 = 0, b_2 = 1, c_2 = a_{2,1} = \frac{1}{2}$$

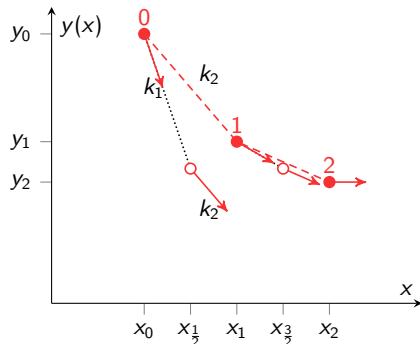
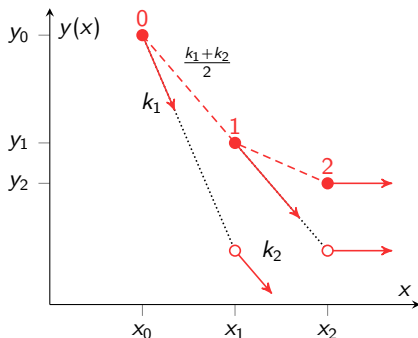
# Second order Runge-Kutta methods

Classical RK2 method  
(= Heun's method,  
improved Euler method)

$$\begin{aligned} k_1 &= f_i \\ k_2 &= f(x_i + h, y_i + hk_1) \\ y_{i+1} &= y_i + \frac{1}{2}h(k_1 + k_2) \end{aligned}$$

Explicit midpoint rule  
(modified Euler method)

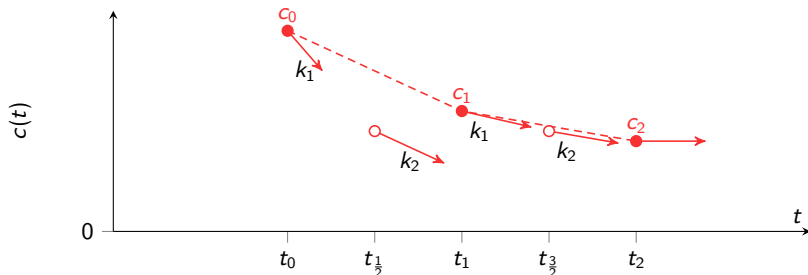
$$\begin{aligned} k_1 &= f_i \\ k_2 &= f(x_i + \frac{1}{2}h, y_i + \frac{1}{2}hk_1) \\ y_{i+1} &= y_i + hk_2 \end{aligned}$$



## Second order Runge-Kutta method — Example

First order reaction in a batch reactor:  $\frac{dc}{dt} = -kc$  with  
 $c(t = 0) = 1 \text{ mol m}^{-3}$ ,  $k = 1 \text{ s}^{-1}$ ,  $t_{\text{end}} = 2 \text{ s}$ .

Time [s]	C [mol m <sup>-3</sup> ]	$k_1 = hf(x_i, y_i)$	$k_2 = hf(x_i + \frac{1}{2}h, y_n + \frac{1}{2}k_1)$
0	1.00	$0.1 \cdot (-1 \cdot 1) = -0.1$	$0.1 \cdot (-1 \cdot (1 - 0.5 \cdot 0.1)) = -0.095$
0.1	$1 - 0.095 = 0.905$	$0.1 \cdot (-1 \cdot 0.905) = -0.0905$	$0.1 \cdot (-1 \cdot (0.905 - 0.5 \cdot 0.0905)) = -0.085975$
...	...	...	...
2	0.1358225	-0.0135822	-0.0129031



# RK2 method — order of convergence

$N$	$\zeta$	$\frac{\zeta_{\text{numerical}} - \zeta_{\text{analytical}}}{\zeta_{\text{analytical}}}$	$r = \frac{\log\left(\frac{\epsilon_i}{\epsilon_{i-1}}\right)}{\log\left(\frac{N_{i-1}}{N_i}\right)}$
20	0.864178	$5.634 \times 10^{-4}$	—
40	0.864548	$1.355 \times 10^{-4}$	2.056
80	0.864636	$3.323 \times 10^{-5}$	2.028
160	0.864658	$8.229 \times 10^{-6}$	2.014
320	0.864663	$2.048 \times 10^{-6}$	2.007

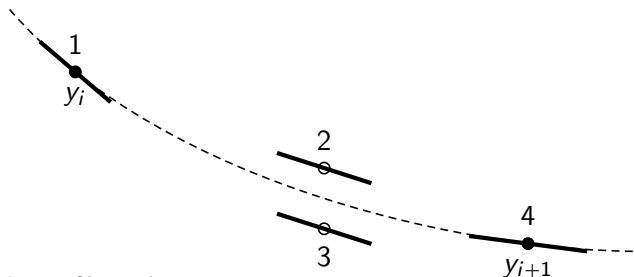
## RK2 method — order of convergence

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⇒ RK2 is a second order method. Doubling the number of cells reduces the error by a factor 4!

Can we do even better?

# RK4 method (classical fourth order Runge-Kutta method)



$$k_1 = f(x_i, y_i)$$

$$k_2 = f\left(x_i + \frac{1}{2}h, y_i + \frac{1}{2}hk_1\right)$$

$$k_3 = f\left(x_i + \frac{1}{2}h, y_i + \frac{1}{2}hk_2\right)$$

$$k_4 = f(x_i + h, y_i + hk_3)$$

$$y_{i+1} = y_i + h \left( \frac{1}{6}k_1 + \frac{1}{3}(k_2 + k_3) + \frac{1}{6}k_4 \right)$$

# RK4 method — order of convergence

$N$	$\zeta$	$\frac{\zeta_{\text{numerical}} - \zeta_{\text{analytical}}}{\zeta_{\text{analytical}}}$	$r = \frac{\log\left(\frac{\epsilon_i}{\epsilon_{i-1}}\right)}{\log\left(\frac{N_{i-1}}{N_i}\right)}$
20	0.864664472	$2.836 \times 10^{-7}$	—
40	0.864664702	$1.700 \times 10^{-8}$	4.060
80	0.864664716	$1.040 \times 10^{-9}$	4.030
160	0.864664717	$6.435 \times 10^{-11}$	4.015
320	0.864664717	$4.001 \times 10^{-12}$	4.007



# RK4 method — order of convergence

$N$	$\zeta$	$\frac{\zeta_{\text{numerical}} - \zeta_{\text{analytical}}}{\zeta_{\text{analytical}}}$	$r = \frac{\log\left(\frac{\epsilon_j}{\epsilon_{j-1}}\right)}{\log\left(\frac{N_{j-1}}{N_j}\right)}$
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320	0.864664717	$4.001 \times 10^{-12}$	4.007

⇒ RK4 is a fourth order method: Doubling the number of cells reduces the error by a factor 16!

Can we do even better?

# Today's outline

- ① Introduction
- ② Euler's method  
Forward Euler
- ③ Rates of convergence
- ④ Runge-Kutta methods  
RK2 methods  
RK4 method
- ⑤ Step size control
- ⑥ Implicit methods  
Backward Euler  
Implicit midpoint method

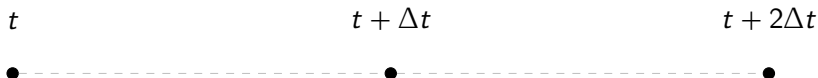
# Adaptive step size control

The step size (be it either position, time or both (PDEs)) cannot be decreased indefinitely to favour a higher accuracy, since each additional grid point causes additional computation time. It may be wise to adapt the step size according to the computation requirements.

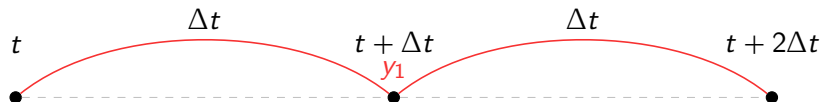
Globally two different approaches can be used:

- 1 Step doubling: compare solutions when taking one full step or two consecutive halve steps
- 2 Embedded methods: Compare solutions when using two approximations of different order

# Adaptive step size control: step doubling

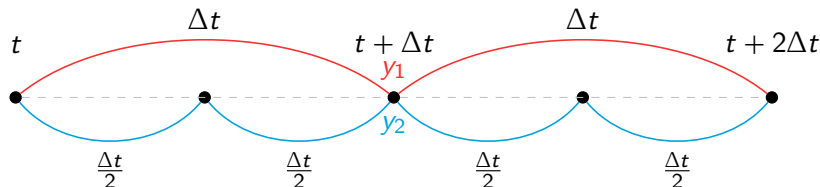


# Adaptive step size control: step doubling



- RK4 with one large step of  $h$ :  $y_{i+1} = y_1 + ch^5 + \mathcal{O}(h^6)$

# Adaptive step size control: step doubling



- RK4 with one large step of  $h$ :  $y_{i+1} = y_1 + ch^5 + \mathcal{O}(h^6)$
- RK4 with two steps of  $\frac{1}{2}h$ :  $y_{i+1} = y_2 + 2c(\frac{1}{2}h)^5 + \mathcal{O}(h^6)$

## Adaptive step size control: step doubling

- Estimation of truncation error by comparing  $y_1$  and  $y_2$ :  
 $\Delta = y_2 - y_1$
- If  $\Delta$  too large, reduce step size for accuracy
- If  $\Delta$  too small, increase step size for efficiency.
- Ignoring higher order terms and solving for  $c$ :  
 $\Delta = \frac{15}{16}ch^5 \Rightarrow ch^5 = \frac{16}{15}\Delta \Rightarrow y_{i+1} = y_2 + \frac{\Delta}{15} + \mathcal{O}(h^6)$   
(local Richardson extrapolation)

Note that when we specify a tolerance  $tol$ , we can estimate the maximum allowable step size as:  $h_{\text{new}} = \alpha h_{\text{old}} \left| \frac{tol}{\Delta} \right|^{\frac{1}{5}}$  with  $\alpha$  a safety factor (typically  $\alpha = 0.9$ ).

## Adaptive step size control: embedded methods

Use a special fourth and a fifth order Runge Kutta method to approximate  $y_{i+1}$

- The fourth order method is special because we want to use the same positions for the evaluation for computational efficiency.
- RK45 is the preferred method (minimum number of function evaluations) (this is built in Matlab as `ode45`).



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# Problems with Euler's method: instability

Consider the ODE:

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First order approximation of derivative:  $\frac{dy}{dx} = \frac{y_{i+1} - y_i}{\Delta x}$ .

Where to evaluate the function  $f$ ?

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Where to evaluate the function  $f$ ?

- 1 Evaluation at  $x_i$ : Explicit Euler method (forward Euler)
- 2 Evaluation at  $x_{i+1}$ : Implicit Euler method (backward Euler)

# Problems with Euler's method: instability – forward Euler

Explicit Euler method (forward Euler):

- Use values at  $x_i$ :

$$\frac{y_{i+1} - y_i}{\Delta x} = f(x_i, y_i) \Rightarrow y_{i+1} = y_i + hf(x_i, y_i).$$

- This is an explicit equation for  $y_{i+1}$  in terms of  $y_i$ .
- It can give instabilities with large function values.

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$$\frac{dc}{dt} = -kc \Rightarrow c_{i+1} = c_i - k c_i \Delta t \Rightarrow \frac{c_{i+1}}{c_i} = 1 - k\Delta t$$

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It follows that unphysical results are obtained for  $k\Delta t \geq 1!!$

Stability requirement

$$k\Delta t < 1$$

(but probably accuracy requirements are more stringent here!)

# Problems with Euler's method: instability – backward Euler

Implicit Euler method (backward Euler):

- Use values at  $x_{i+1}$ :

$$\frac{y_{i+1} - y_i}{\Delta x} = f(x_{i+1}, y_{i+1}) \Rightarrow y_{i+1} = y_i + hf(x_{i+1}, y_{i+1}).$$

- This is an implicit equation for  $y_{i+1}$ , because it also depends on terms of  $y_{i+1}$ .



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$$\frac{dc}{dt} = -kc \Rightarrow c_{i+1} = c_i - k c_{i+1} \Delta t \Rightarrow \frac{c_{i+1}}{c_i} = \frac{1}{1 + k\Delta t}$$

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$$\frac{dc}{dt} = -kc \Rightarrow c_{i+1} = c_i - k c_{i+1} \Delta t \Rightarrow \frac{c_{i+1}}{c_i} = \frac{1}{1 + k \Delta t}$$

This equation does never give unphysical results!

The implicit Euler method is *unconditionally stable*  
(but maybe not very accurate or efficient).

# Semi-implicit Euler method

Usually  $f$  is a non-linear function of  $y$ , so that linearization is required (recall Newton's method).

$$\frac{dy}{dx} = f(y) \Rightarrow y_{i+1} = y_i + hf(y_{i+1}) \quad \text{using} \quad f(y_{i+1}) = f(y_i) + \left. \frac{df}{dy} \right|_i (y_{i+1} - y_i) + \dots$$

$$\Rightarrow y_{i+1} = y_i + h \left[ f(y_i) + \left. \frac{df}{dy} \right|_i (y_{i+1} - y_i) \right]$$

$$\Rightarrow \left( 1 - h \left. \frac{df}{dy} \right|_i \right) y_{i+1} = \left( 1 - h \left. \frac{df}{dy} \right|_i \right) y_i + hf(y_i)$$

$$\Rightarrow y_{i+1} = y_i + h \left( 1 - h \left. \frac{df}{dy} \right|_i \right)^{-1} f(y_i)$$

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Usually  $f$  is a non-linear function of  $y$ , so that linearization is required (recall Newton's method).

$$\frac{dy}{dx} = f(y) \Rightarrow y_{i+1} = y_i + hf(y_{i+1}) \quad \text{using} \quad f(y_{i+1}) = f(y_i) + \left. \frac{df}{dy} \right|_i (y_{i+1} - y_i) + \dots$$

$$\Rightarrow y_{i+1} = y_i + h \left[ f(y_i) + \left. \frac{df}{dy} \right|_i (y_{i+1} - y_i) \right]$$

$$\Rightarrow \left( 1 - h \left. \frac{df}{dy} \right|_i \right) y_{i+1} = \left( 1 - h \left. \frac{df}{dy} \right|_i \right) y_i + hf(y_i)$$

$$\Rightarrow y_{i+1} = y_i + h \left( 1 - h \left. \frac{df}{dy} \right|_i \right)^{-1} f(y_i)$$

For the case that  $f(x, y(x))$  we could add the variable  $x$  as an additional variable  $y_{n+1} = x$ . Or add one fully implicit Euler step (which avoids the computation of  $\frac{\partial f}{\partial x}$ ):

$$y_{i+1} = y_i + hf(x_{i+1}, y_{i+1}) \Rightarrow y_{i+1} = y_i + h \left( 1 - h \left. \frac{df}{dy} \right|_i \right)^{-1} f(x_{i+1}, y_i)$$

## Semi-implicit Euler method - example

Second order reaction in a batch reactor:

$$\frac{dc}{dt} = -kc^2 \text{ with } c_0 = 1 \text{ mol m}^{-3}, k = 1 \text{ m}^3 \text{ mol}^{-1} \text{ s}^{-1}, t_{\text{end}} = 2 \text{ s}$$

Analytical solution:  $c(t) = \frac{c_0}{1 + kc_0 t}$

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$N$	$\zeta$	$\frac{\zeta_{\text{numerical}} - \zeta_{\text{analytical}}}{\zeta_{\text{analytical}}}$	$r = \frac{\log\left(\frac{\epsilon_i}{\epsilon_{i-1}}\right)}{\log\left(\frac{N_{i-1}}{N_i}\right)}$
20	0.654066262	$1.89 \times 10^{-2}$	—
40	0.660462687	$9.31 \times 10^{-3}$	1.02220
80	0.663589561	$4.62 \times 10^{-3}$	1.01162
160	0.665134433	$2.30 \times 10^{-3}$	1.00594
320	0.665902142	$1.15 \times 10^{-3}$	1.00300

## Second order implicit method: Implicit midpoint method

Implicit midpoint rule (second order)	Explicit midpoint rule (modified Euler method)
$y_{i+1} = y_i + hf \left( x_i + \frac{1}{2}h, \frac{1}{2}(y_i + y_{i+1}) \right)$	$y_{i+1} = y_i + hf(x_i + \frac{1}{2}h, y_i + \frac{1}{2}hk_1)$

in case  $f(y)$  then:

$$f \left( \frac{1}{2}(y_i + y_{i+1}) \right) = f_i + \left. \frac{df}{dy} \right|_i \left( \frac{1}{2}(y_i + y_{i+1}) - y_i \right) = f_i + \frac{1}{2} \left. \frac{df}{dy} \right|_i (y_{i+1} - y_i)$$



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Implicit midpoint rule reduces to:

$$\begin{aligned} y_{i+1} &= y_i + hf_i + \frac{h}{2} \left. \frac{df}{dy} \right|_i (y_{i+1} - y_i) \\ \Rightarrow \left( 1 - \frac{h}{2} \left. \frac{df}{dy} \right|_i \right) y_{i+1} &= \left( 1 - \frac{h}{2} \left. \frac{df}{dy} \right|_i \right) y_i + hf_i \end{aligned}$$

$$\Rightarrow y_{i+1} = y_i + h \left( 1 - \frac{h}{2} \left. \frac{df}{dy} \right|_i \right)^{-1} f_i$$

# Implicit midpoint method — example

Second order reaction in a batch reactor:

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(Analytical solution:  $c(t) = \frac{c_0}{1+kc_0t}$ ).

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Define  $f = -kc^2$ , then  $\frac{df}{dc} = -2kc$ .

Substitution:

$$\begin{aligned} c_{i+1} &= c_i + h \left( 1 - \frac{h}{2} \cdot (-2kc_i) \right)^{-1} \cdot (-kc_i^2) \\ &= c_i - \frac{hkc_i^2}{1 + hkc_i} = \frac{c_i + hkc_i^2 - hkc_i^2}{1 + hkc_i} \Rightarrow c_{i+1} = \frac{c_i}{1 + hkc_i} \end{aligned}$$

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Substitution:

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You will find that this method is exact for all step sizes  $h$  because of the quadratic source term!

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20	0.6666666667	$1.665 \times 10^{-16}$	—
40	0.6666666667	0	—
80	0.6666666667	0	—
160	0.6666666667	0	—
320	0.6666666667	0	—

# Implicit midpoint method — example

Third order reaction in a batch reactor

Analytical solution:  $c(t) = \frac{c_0}{\sqrt{1+2kc_0^2t}}$

$$c_{i+1} = c_i - \frac{hkc_i^3}{1 + \frac{3}{2}hkc_i^2}$$



# Implicit midpoint method — example

Third order reaction in a batch reactor

Analytical solution:  $c(t) = \frac{c_0}{\sqrt{1+2kc_0^2t}}$

$$c_{i+1} = c_i - \frac{hkc_i^3}{1 + \frac{3}{2}hkc_i^2}$$

$N$	$\zeta$	$\frac{\zeta_{\text{numerical}} - \zeta_{\text{analytical}}}{\zeta_{\text{analytical}}}$	$r = \frac{\log\left(\frac{\epsilon_i}{\epsilon_{i-1}}\right)}{\log\left(\frac{N_{i-1}}{N_i}\right)}$
20	0.5526916174	$1.71 \times 10^{-4}$	—
40	0.5527633731	$4.17 \times 10^{-5}$	2.041
80	0.5527807304	$1.03 \times 10^{-5}$	2.021
160	0.5527849965	$2.55 \times 10^{-6}$	2.011
320	0.5527860538	$6.34 \times 10^{-7}$	2.005

## Part II

# Boudary Value Problems and Systems of ODEs

# Today's outline

## 7 Boundary value problems

Shooting method

## 8 Systems of ODEs

Solution methods for systems of ODEs

Stiff systems of ODEs

## 9 Solving systems of ODEs in Matlab

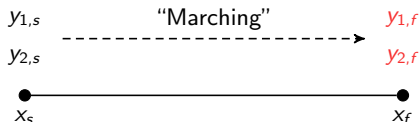
## 10 Conclusion

# Importance of boundary conditions

The nature of boundary conditions determines the appropriate numerical method. Classification into 2 main categories:

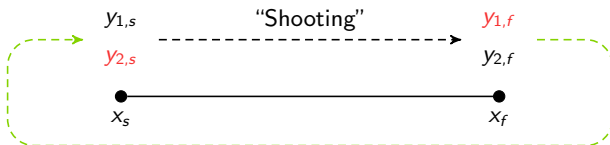
- *Initial value problems (IVP)*

We know the values of all  $y_i$  at some starting position  $x_s$ , and it is desired to find the values of  $y_i$  at some final point  $x_f$ .



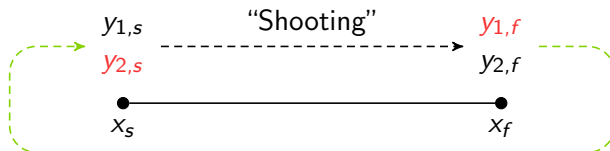
- *Boundary value problems (BVP)*

Boundary conditions are specified at more than one  $x$ . Typically, some of the BC are specified at  $x_s$  and the remainder at  $x_f$ .



# Shooting method

How to solve a BVP using the shooting method:



- Define the system of ODEs
- Provide an initial guess for the unknown boundary condition
- Solve the system and compare the resulting boundary condition to the expected value
- Adjust the guessed boundary value, and solve again. Repeat until convergence.
  - Of course, you can subtract the expected value from the computed value at the boundary, and use a non-linear root finding method

## BVP: example in Excel

Consider a chemical reaction in a liquid film layer of thickness  $\delta$ :

$$\mathcal{D} \frac{d^2 c}{dx^2} = k_R c \quad \text{with} \quad \begin{array}{ll} c(x=0) = C_{A,i,L} = 1 & \text{(interface concentration)} \\ c(x=\delta) = 0 & \text{(bulk concentration)} \end{array}$$

Question: compute the concentration profile in the film layer.

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Question: compute the concentration profile in the film layer.

### Step 1: Define the system of ODEs

This second-order ODE can be rewritten as a system of first-order ODEs, if we define the flux  $q$  as:

$$q = -\mathcal{D} \frac{dc}{dx}$$

Now, we find:

$$\frac{dc}{dx} = -\frac{1}{\mathcal{D}} q$$

$$\frac{dq}{dx} = -k_R c$$

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Question: compute the concentration profile in the film layer.

### Step 2: Set the boundary conditions

The boundary conditions for the concentrations at  $x=0$  and  $x=\delta$  are known.

The flux at the interface, however, is not known, and should be solved for.

$$\frac{dc}{dx} = -\frac{1}{\mathcal{D}} q$$

$$\frac{dq}{dx} = -k_R c$$



## BVP: example in Excel

Solving the two first-order ODEs in Excel. First, the cells with constants:

	A	B	C
1	CAiL	1	mol/m <sup>3</sup>
2	D	1e-8	m <sup>2</sup> /s
3	kR	10	1/s
4	delta	1e-4	m
5	N	100	
6	dx	=B4/B5	

$$\frac{dc}{dx} = -\frac{1}{D}q$$

$$\frac{dq}{dx} = -k_R c$$

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5	N	100	
6	dx	=B4/B5	

$$\frac{dc}{dx} = -\frac{1}{D}q$$

$$\frac{dq}{dx} = -k_R c$$

Now, we program the forward Euler (explicit) schemes for  $c$  and  $q$  below:

	A	B	C
10	x	c	q
11	0	=B1	10
12	=A11+\$B\$6	=B11+\$B\$6*(-1/\$B\$2*C11)	=C11+\$B\$6*(-\$B\$3*B11)
13	=A12+\$B\$6	=B12+\$B\$6*(-1/\$B\$2*C12)	=C12+\$B\$6*(-\$B\$3*B12)
...	...	...	...
111	=A110+\$B\$6	=B110+\$B\$6*(-1/\$B\$2*C110)	=C110+\$B\$6*(-\$B\$3*B110)

## BVP: example in Excel

- We now have profiles for  $c$  and  $q$  as a function of position  $x$ .
- The concentration  $c(x = \delta)$  depends (eventually) on the boundary condition at the interface  $q(x = 0)$
- We can use the solver to change  $q(x = 0)$  such that the concentration at the bulk meets our requirement:  
 $c(x = \delta) = 0$

# BVP: example in Matlab

We first program the system of ODEs in a separate function:

$$\frac{dc}{dx} = -\frac{1}{\mathcal{D}}q$$
$$\frac{dq}{dx} = -k_R c$$

```
function [dxdt] = BVPODE(t,x,ps)
dxdt(1)=-1/ps.D*x(2);
dxdt(2)=-ps.kR*x(1);
dxdt=dxdt';
return
```

# BVP: example in Matlab

We first program the system of ODEs in a separate function:

$$\begin{aligned}\frac{dc}{dx} &= -\frac{1}{D}q \\ \frac{dq}{dx} &= -k_R c\end{aligned}$$

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function [dxdt] = BVPODE(t,x,ps)
dxdt(1)=-1/ps.D*x(2);
dxdt(2)=-ps.kR*x(1);
dxdt=dxdt';
return
```

Note that we pass a variable (type: struct) that contains required parameters: `ps`.

## BVP: example in Matlab

The ODE function is solved via ode45, after setting a number of initial and boundary conditions:

```
function f = RunBVP(bcq,ps)
[x,cq] = ode45(@BVP0DE,[0 ps.delta],[1 bcq], [], ps);
f = cq(end,1) - 0;
plotyy(x,cq(:,1),x,cq(:,2));
return;
```

## BVP: example in Matlab

The ODE function is solved via `ode45`, after setting a number of initial and boundary conditions:

```
function f = RunBVP(bcq,ps)
[x,cq] = ode45(@BVPODE,[0 ps.delta],[1 bcq], [], ps);
f = cq(end,1) - 0;
plotyy(x,cq(:,1),x,cq(:,2));
return;
```

Note the following:

- We use the interval  $0 \leq x \leq \delta$
- Boundary conditions are given as:  $c(x=0) = 1$  and  $q(x=0) = bcq$ , which is given as an argument to the function (i.e. changable from 'outside'!)
- The function returns  $f$ , the difference between the computed and desired concentration at  $x = \delta$ .

## BVP: example in Matlab

Finally, we should solve the system so that we obtain the right boundary condition  $q = bcq$  such that  $c(x = \delta) = 0$ . We can use the built-in function `fzero` to do this

```
% Parameter definition
ps.D=1e-8;
ps.kR=10;
ps.delta=1e-4;

% Solve for flux boundary condition (initial guess: 0)
opt = optimset('Display','iter');
flux = fzero(@RunBVP,0,opt,ps);
```



## BVP example: analytical solution

Compare with the analytical solution:

$$q = k_L E_A C_{A,i,L} \quad \text{with}$$

$$E_A = \frac{Ha}{\tanh Ha} \quad (\text{Enhancement factor})$$

$$Ha = \frac{\sqrt{k_R \mathcal{D}}}{k_L} \quad (\text{Hatta number})$$

$$k_L = \frac{\mathcal{D}}{\delta} \quad (\text{mass transfer coefficient})$$

# Today's outline

## ⑦ Boundary value problems

Shooting method

## ⑧ Systems of ODEs

Solution methods for systems of ODEs

Stiff systems of ODEs

## ⑨ Solving systems of ODEs in Matlab

## ⑩ Conclusion

# Systems of ODEs

A system of ODEs is specified using vector notation:

$$\frac{d\mathbf{y}}{dx} = \mathbf{f}(x, \mathbf{y}(x))$$

for

$$\frac{dy_1}{dx} = f_1(x, y_1(x), y_2(x)) \quad \text{or} \quad f_1(x, y_1, y_2)$$

$$\frac{dy_2}{dx} = f_2(x, y_1(x), y_2(x)) \quad \text{or} \quad f_2(x, y_1, y_2)$$

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$$\frac{dy_2}{dx} = f_2(x, y_1(x), y_2(x)) \quad \text{or} \quad f_2(x, y_1, y_2)$$

The solution techniques discussed before can also be used to solve systems of equations.

# Systems of ODEs: Explicit methods

## Forward Euler method

$$\mathbf{y}_{i+1} = \mathbf{y}_i + h\mathbf{f}(x_i, \mathbf{y}_i)$$

## Improved Euler method (classical RK2)

$$\mathbf{y}_{i+1} = \mathbf{y}_i + \frac{h}{2}(\mathbf{k}_1 + \mathbf{k}_2) \quad \text{using} \quad \begin{aligned} \mathbf{k}_1 &= \mathbf{f}(x_i, \mathbf{y}_i) \\ \mathbf{k}_2 &= \mathbf{f}(x_i + h, \mathbf{y}_i + h\mathbf{k}_1) \end{aligned}$$

## Modified Euler method (midpoint rule)

$$\mathbf{y}_{i+1} = \mathbf{y}_i + h\mathbf{k}_2 \quad \text{using} \quad \begin{aligned} \mathbf{k}_1 &= \mathbf{f}(x_i, \mathbf{y}_i) \\ \mathbf{k}_2 &= \mathbf{f}\left(x_i + \frac{h}{2}, \mathbf{y}_i + \frac{h}{2}\mathbf{k}_1\right) \end{aligned}$$

# Systems of ODEs: Explicit methods

## Classical fourth order Runge-Kutta method (RK4)

$$\mathbf{y}_{i+1} = \mathbf{y}_i + h \left( \frac{\mathbf{k}_1}{6} + \frac{1}{3} (\mathbf{k}_2 + \mathbf{k}_3) + \frac{\mathbf{k}_4}{6} \right)$$

$$\mathbf{k}_1 = \mathbf{f}(x_i, \mathbf{y}_i)$$

$$\mathbf{k}_2 = \mathbf{f}\left(x_i + \frac{h}{2}, \mathbf{y}_i + \frac{h}{2}\mathbf{k}_1\right)$$

using

$$\mathbf{k}_3 = \mathbf{f}\left(x_i + \frac{h}{2}, \mathbf{y}_i + \frac{h}{2}\mathbf{k}_2\right)$$

$$\mathbf{k}_4 = \mathbf{f}(x_i + h, \mathbf{y}_i + h\mathbf{k}_3)$$

# Systems of ODEs: Implicit methods

## Backward Euler method

$$\mathbf{y}_{i+1} = \mathbf{y}_i + h \left( 1 - h \left. \frac{d\mathbf{f}}{d\mathbf{y}} \right|_i \right)^{-1} \mathbf{f}(\mathbf{y}_i)$$

## Implicit midpoint method

$$\mathbf{y}_{i+1} = \mathbf{y}_i + h \left( 1 - \frac{h}{2} \left. \frac{d\mathbf{f}}{d\mathbf{y}} \right|_i \right)^{-1} \mathbf{f}(\mathbf{y}_i)$$

# Stiff systems of ODEs

A system of ODEs can be stiff and require a different solution method.



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A system of ODEs can be stiff and require a different solution method. For example:

$$\frac{dc_1}{dt} = 998c_1 + 1998c_2 \quad \frac{dc_2}{dt} = -999c_1 - 1999c_2$$

with boundary conditions  $c_1(t=0) = 1$  and  $c_2(t=0) = 0$ .

The analytical solution is:

$$c_1 = 2e^{-t} - e^{-1000t} \quad c_2 = -e^{-t} + e^{-1000t}$$

For the explicit method we require  $\Delta t < 10^{-3}$  despite the fact that the term is completely negligible, but essential to keep stability.

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The “disease” of stiff equations: we need to follow the solution on the shortest length scale to maintain stability of the integration, although accuracy requirements would allow a much larger time step.

# Demonstration with example

Forward Euler (explicit)

$$\frac{c_{1,i+1} - c_{1,i}}{dt} = 998c_{1,i} + 1998c_{2,i}$$

$$\frac{c_{2,i+1} - c_{2,i}}{dt} = -999c_{1,i} - 1999c_{2,i}$$

$$\Rightarrow \begin{aligned} c_{1,i+1} &= (1 + 998\Delta t) c_{1,i} + 1998\Delta t c_{2,i} \\ c_{2,i+1} &= -999\Delta t c_{1,i} + (1 - 1999\Delta t) c_{2,i} \end{aligned}$$

# Demonstration with example

Backward Euler (implicit)

$$\frac{dc_{1,i+1} - c_{1,i+1}}{dt} = 998c_{1,i+1} + 1998c_{2,i+1}$$

$$\frac{dc_{2,i+1} - c_{2,i+1}}{dt} = -999c_{1,i+1} - 1999c_{2,i+1}$$

$$\Rightarrow (1 - 998\Delta t) c_{1,i+1} - 1998\Delta t c_{2,i} = c_{1,i}$$
$$999\Delta t c_{1,i+1} + (1 + 999\Delta t) c_{2,i+1} = c_{2,i}$$

# Demonstration with example

Backward Euler (implicit)

$$\frac{dc_{1,i+1} - c_{1,i+1}}{dt} = 998c_{1,i+1} + 1998c_{2,i+1}$$

$$\frac{dc_{2,i+1} - c_{2,i+1}}{dt} = -999c_{1,i+1} - 1999c_{2,i+1}$$

$$\Rightarrow (1 - 998\Delta t) c_{1,i+1} - 1998\Delta t c_{2,i} = c_{1,i}$$
$$999\Delta t c_{1,i+1} + (1 + 999\Delta t) c_{2,i+1} = c_{2,i}$$

$$A\mathbf{c}_{i+1} = \mathbf{c}_i \text{ with } A = \begin{pmatrix} 1 - 998\Delta t & -1998\Delta t \\ 999\Delta t & 1 + 1999\Delta t \end{pmatrix} \text{ and } \mathbf{b} = \begin{pmatrix} c_{1,i} \\ c_{2,i} \end{pmatrix}$$

## Demonstration with example

Backward Euler (implicit)  $A\mathbf{c}_{i+1} = \mathbf{c}_i$  with

$$A = \begin{pmatrix} 1 - 998\Delta t & -1998\Delta t \\ 999\Delta t & 1 + 1999\Delta t \end{pmatrix} \quad \text{and} \quad \mathbf{b} = \begin{pmatrix} c_{1,i} \\ c_{2,i} \end{pmatrix}$$

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Backward Euler (implicit)  $A\mathbf{c}_{i+1} = \mathbf{c}_i$  with

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Cramers rule:

$$c_{1,i+1} = \frac{\begin{vmatrix} c_{1,i} & -1998\Delta t \\ c_{2,i} & 1 + 1999\Delta t \end{vmatrix}}{\det |A|} = \frac{(1+1999\Delta t)c_{1,i} + 1998\Delta t c_{2,i}}{(1-998\Delta t)(1+1999\Delta t) + 1998 \cdot 999\Delta t^2}$$

$$c_{2,i+1} = \frac{\begin{vmatrix} 1 - 998\Delta t & c_{1,i} \\ 999\Delta t & c_{2,i} \end{vmatrix}}{\det |A|} = \frac{-999\Delta t c_{1,i} + (1-998\Delta t)c_{2,i}}{(1-998\Delta t)(1+1999\Delta t) + 1998 \cdot 999\Delta t^2}$$

Forward Euler:  $\Delta t \leq 0.001$  for stability

Backward Euler: always stable, even for  $\Delta t > 100$  (but then not very accurate!)

## Demonstration with example

Cure for stiff problems: use implicit methods! To find out whether your system is stiff: check whether one of the eigenvalues have an imaginary part



# Solving systems of ODEs in Matlab

Matlab provides convenient procedures to solve (systems of) ODEs automatically.

The procedure is as follows:

- 1 Create a function that specifies the ODEs. Specifically, this function returns the  $\frac{dy}{dx}$  vector.
- 2 Initialise solver variables and settings (e.g. step size, initial conditions, tolerance), in a separate script
- 3 Call the ODE solver function, using a *function handle* to the ODE function described in point 1.
  - The ODE solver will return the vector for the independent variable, and a solution vector (matrix for systems of ODEs).

## Solving systems of ODEs in Matlab: example

We solve the system:  $\frac{dx_1}{dt} = -x_1 - x_2$ ,  $\frac{dx_2}{dt} = x_1 - 2x_2$

### Create an ODE function

```
function [dxdt] = myODEFunction(t,x)
dxdt(1) = -x(1) - x(2);
dxdt(2) =  x(1) - 2*x(2);
dxdt=dxdt'; % Transpose to column vector
return
```

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### Create a solution script

```
x_init = [0 1]; % Initial conditions
tspan = [0 10]; % Time span
options = odeset('RelTol',1e-4,'AbsTol',[1e-4 1e-4]);
[t,x] = ode45(@myODEfunction,tspan,x_init,options);
```

# Solving systems of ODEs in Matlab: example

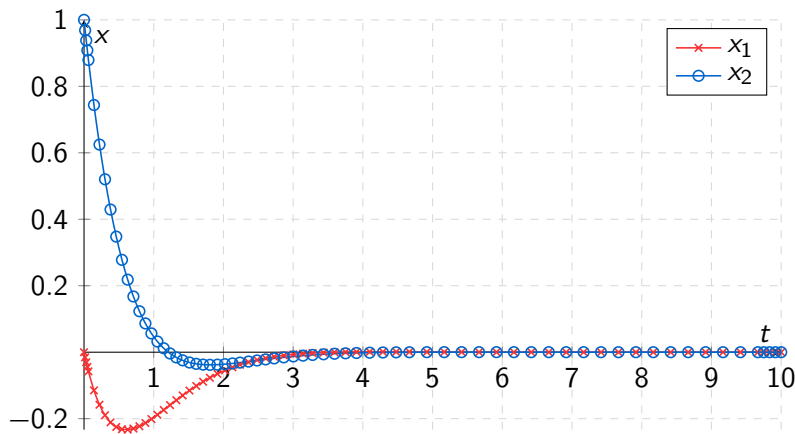
Plot the solution:

```
plot(t,x(:,1),'r-x',t,x(:,2),'b-o')
```

# Solving systems of ODEs in Matlab: example

Plot the solution:

```
plot(t,x(:,1),'r-x',t,x(:,2),'b-o')
```



## Solving systems of ODEs in Matlab: example

A few notes on working with `ode45` and other solvers. If we want to give additional arguments (e.g. `a`, `b` and `c`) to our ODE function, we can list them in the function line:

```
function [dxdt] = myODE(t,x,a,b,c)
```

The additional arguments can now be set in the solver script by *adding them after the options*:

```
[t,x] = ode45(@myODE,tspan,x_0,options,a,b,c);
```

## Solving systems of ODEs in Matlab: example

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[t,x] = ode45(@myODE,tspan,x_0,options,a,b,c);
```

- Of course, in the solver script, the variables do not need to be called `a`, `b` and `c`:

```
[t,x] = ode45(@myODE,tspan,x_0,options,k1,phi,V);
```

- These variables may be of any type (vectors, matrix, struct). Especially a struct is useful to carry many values in 1 variable.

## Solving systems of ODEs in Matlab: example

You may have noticed that the step size in  $t$  varied. This is because we have given the begin and end times of our time span:

```
tspan = [0 10];
```



## Solving systems of ODEs in Matlab: example

You may have noticed that the step size in  $t$  varied. This is because we have given the begin and end times of our time span:

```
tspan = [0 10];
```

You can also solve at specific steps, by supplying all steps explicitly, e.g.:

```
tspan = linspace(0,10,101);
```

This example provides 101 explicit time steps between 0 and 10 seconds.

Note that you may affect the efficiency and accuracy of the solver algorithm by doing this!

# Today's outline

- ⑦ Boundary value problems
  - Shooting method
- ⑧ Systems of ODEs
  - Solution methods for systems of ODEs
  - Stiff systems of ODEs
- ⑨ Solving systems of ODEs in Matlab
- ⑩ Conclusion

# Other methods

Other explicit methods:

- Burlisch-Stoer method (Richardson extrapolation + modified midpoint method)

Other implicit methods:

- Rosenbrock methods (higher order implicit Runge-Kutta methods)
- Predictor-corrector methods

# Summary

- Several solution methods and their derivation were discussed:
  - Explicit solution methods: Euler, Improved Euler, Midpoint method, RK45
  - Implicit methods: Implicit Euler and Implicit midpoint method
  - A few examples of their spreadsheet implementation were shown
- We have paid attention to accuracy and instability, rate of convergence and step size
- Systems of ODEs can be solved by the same algorithms. Stiff problems should be treated with care.
- An example of solving ODEs with Matlab was demonstrated.