

# Multiresolution analysis based on quadratic Hermite interpolation - Part 1: Piecewise polynomial Curves

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**Abstract.** We study two simple multiresolution analyses and their stability in the  $L^\infty$ -norm: Faber decomposition and  $C^1$  quadratic Hermite interpolation, both with nonuniform knot sequences. The use of the  $L^\infty$  norm is natural in many CAGD applications and it leads to schemes which are faster and simpler to implement than the wavelet schemes based on the  $L^2$  norm. We have chosen to discuss quadratic Hermite interpolation because (i) it is a  $C^1$  scheme with nice shape preserving properties, (ii) we have a certain sup norm stability in the wavelet spaces, (iii) there are local support bases for these spaces, (iv) the decomposition coefficients can be determined explicitly in real time, (v) it generalizes to splines over triangulations.

## §1. Introduction

Multiresolution analysis (MRA) and wavelets ([4]) have turned out to be versatile tools both within Mathematics itself and in applications. Polynomial splines give rise to an important class of MRAs, and in this paper we are going to study two simple spline MRAs. One of them is based on linear splines and the other on quadratic splines. In a wavelet setting one usually assumes the knot spacing to be uniform, but in this paper we allow nonuniform knots. Another standard feature of wavelets is orthogonality, but here we will be content with direct sum decompositions. We also deviate from the standard by measuring stability in  $L^\infty$  rather than in  $L^2$ ; this is more natural in many applications.

The MRA based on linear splines is included as a template for the quadratic MRA. It turns out that the linear MRA is stable in  $L^\infty$ , uniformly for all knots. The quadratic Hermite MRA does not have this property. The main advantage of the quadratic MRA is of course that it is  $C^1$ , and elsewhere we have shown that it can be generalized to bivariate functions [6].

A first draft of this paper was prepared shortly after the last author visited Oslo in 1995, and the main results were presented at a conference at Oberwolfach that year.

## §2. Faber Decomposition

In 1909 Faber [7] presented a hierarchical representation of functions based on piecewise linear interpolation, see also [9]. In Faber's construction, a function  $f$  defined on the interval  $[0, 1]$  is decomposed into the sum

$$f = f_0 + \sum_{k=0}^{\infty} g_k, \quad g_k = f_{k+1} - f_k, \quad (2.1)$$

where  $f_k$  is the piecewise linear interpolant to  $f$  on the dyadic partition

$$\Delta_k = \{i2^{-k}\}_{i=0}^{2^k}. \quad (2.2)$$

Faber also gave an explicit formula for the “wavelets”  $g_k$ . In our notation it takes the form

$$g_k(x) = \sum_{j=1}^{2^k} d_{j,k} B_1(2^{k+1}x - 2j + 1), \quad k \geq 0, \quad (2.3)$$

where  $B_1(x)$  is the piecewise linear function with  $B_1(i) = \delta_{i,0}$  for all  $i$ , and

$$d_{j,k} = f\left(\frac{2j-1}{2^{k+1}}\right) - \frac{1}{2} \left( f\left(\frac{2j-2}{2^{k+1}}\right) + f\left(\frac{2j}{2^{k+1}}\right) \right). \quad (2.4)$$

This formula can be derived by observing that (2.3) and (2.1) lead to the relation

$$d_{j,k} = g_k(x_j) = f_{k+1}(x_j) - f_k(x_j), \quad \text{where } x_j = \frac{2j-1}{2^{k+1}}.$$

Since  $x_j$  is an interpolation point for  $f_{k+1}$  that is midway between the two interpolation points  $(j-1)/2^k$  and  $j/2^k$  for  $f_k$ , we have

$$f_{k+1}(x_j) = f(x_j), \quad f_k(x_j) = \frac{1}{2} \left( f\left(\frac{j-1}{2^k}\right) + f\left(\frac{j}{2^k}\right) \right),$$

and (2.4) follows.

Faber was interested in this decomposition because it makes it quite simple to construct continuous nonsmooth functions. Faber decomposition has many other attractive features, particularly from a computational perspective. It is simple to implement, it is fast, and it can be generalized to triangles. In addition all functions used to represent  $f_k$  and  $g_k$  are dilates and translates of one simple function  $B_1$ . In essence, Faber decomposition yields decent compression with little effort.

## §3. Multiresolution Analysis

Before we continue, let us take the time to spell out precisely the ingredients of a multiresolution analysis. Traditionally a multiresolution analysis consists of a nested sequence of subspaces of  $L^2[0, 1]$ , see [4], but here we propose instead to use a nested sequence of subspaces of  $C[0, 1]$ , equipped with the uniform norm, cf. [5] for a related construction. This gives us better control of the uniform norm of perturbation errors at the expense of losing the formalism of Hilbert spaces.

### 3.1 Basic ingredients

A Multiresolution Analysis consists of

- (i) A collection  $\{V_k\}_{k=0}^{\infty}$  of nested subspaces of  $C[0, 1]$  (with the uniform norm),

$$V_0 \subset V_1 \subset \cdots \subset V_k \subset \cdots,$$

that are dense in  $C[0, 1]$ ,

$$\overline{\bigcup_{k=0}^{\infty} V_k} = C[0, 1].$$

- (ii) A collection  $\{Q_k\}_{k=0}^{\infty}$  of uniformly bounded linear projectors from  $C[0, 1]$  onto  $V_k$ ,

$$Q_k : C[0, 1] \rightarrow V_k, \quad \text{for } k = 0, 1, \dots$$

For each positive integer  $k$  we can then define the wavelet spaces

$$W_k = \{f \in V_{k+1} \mid Q_k f = 0\}, \quad (3.1)$$

which gives a decomposition of  $C[0, 1]$  as the direct sum

$$C[0, 1] = V_0 + W_0 + W_1 + W_2 + \cdots. \quad (3.2)$$

In particular, each  $f \in C[0, 1]$  can be written as

$$f = f_0 + g_0 + g_1 + g_2 + \cdots, \quad (3.3)$$

where  $g_k = f_{k+1} - f_k$  is in  $W_k$  for  $k \geq 0$ , and  $f_k = Q_k f$  is in  $V_k$ . This follows since by construction,

$$V_{k+1} = V_k + W_k \quad \text{and} \quad V_k \cap W_k = \{0\}. \quad (3.4)$$

In addition we have

$$\|f - Q_k f\| \leq (1 + \|Q_k\|) \text{dist}(f, V_k)$$

which by the uniform boundedness of the operators  $\{Q_k\}_{k=0}^{\infty}$  means that

$$\lim_{k \rightarrow \infty} \|f - f_k\| = 0$$

(the norm used here and throughout the paper is the  $L^\infty$ -norm).

### 3.2 Bases for $V_k$ and $W_k$

For practical computations we need a basis  $\{\phi_{j,k}\}_{j \in \mathcal{I}_k}$  for  $V_k$  and a basis  $\{\psi_{j,k}\}_{j \in \mathcal{J}_k}$  for  $W_k$ . Here  $\mathcal{I}_k$  and  $\mathcal{J}_k$  are sets that index the basis functions  $\{\phi_{j,k}\}_j$  and  $\{\psi_{j,k}\}_j$  for each fixed  $k$ . Given an  $f_{k+1} = \sum_{j \in \mathcal{I}_{k+1}} c_{j,k+1} \phi_{j,k+1}$  in  $V_{k+1}$ , we can then decompose it as  $f_{k+1} = f_k + g_k$  with  $f_k = \sum_{j \in \mathcal{I}_k} c_{j,k} \phi_{j,k}$  in  $V_k$  and  $g_k = \sum_{j \in \mathcal{J}_k} d_{j,k} \psi_{j,k}$  in  $W_k$ . The algorithm for computing all the coefficients  $d_{j,k}$  and  $c_{j,k}$  from the  $d_{j,k+1}$  is called the decomposition algorithm. This process can clearly be reversed, and this is the reconstruction algorithm. Together these algorithms constitute the Fast Wavelet Transform (FWT). Note that we will often collect the basis functions and coefficients of  $f_k$  and  $g_k$  together in vectors and write  $f_k = \phi_k^T c_k$  and  $g_k = \psi_k^T d_k$ .

### 3.3 Definition of stability

For efficient and accurate numerical computations it is important that the relation between a function and its coefficients in the wavelet basis is stable. The coarsest space  $V_0$  is usually very simple and of low dimension so we concentrate on the wavelet spaces. The definition below therefore focuses on stability in the subspace

$$\tilde{C}[0, 1] = \cup_{k=0}^{\infty} W_k,$$

i.e., all functions  $f \in C[0, 1]$  such that  $Q_0 f = 0$ . To measure the size of  $f$  we use the uniform norm, but we also need a norm to measure the size of the coefficients. If  $f_n = \sum_{j \in \mathcal{J}_k, k=1}^n d_{j,k} \psi_{j,k}$  is a wavelet, we can use the vector max-norm

$$\delta_n = \|(d_{j,k})_{k=0, j \in \mathcal{J}_k}^{n-1}\|_{\infty} = \max_{j,k} |d_{j,k}|. \quad (3.5)$$

**Definition 3.1.** The wavelets  $\{\psi_{j,k}\}$  are said to form a weakly stable basis for  $\tilde{C}[0, 1]$  if there exist constants  $K_{1,n}$  and  $K_{2,n}$  such that

$$K_{1,n} \delta_n \leq \left\| \sum_{\substack{k=0 \\ j \in \mathcal{J}_k}}^{n-1} d_{j,k} \psi_{j,k} \right\| \leq K_{2,n} \delta_n, \quad (3.6)$$

where  $\delta_n$  is given by (3.5), and  $K_{1,n}$  and  $K_{2,n}$  have at most polynomial growth in  $n$ .

Definition 3.1 provides a definition of stability based on the uniform norm, but the coefficient norm that is employed is rather coarse. The wavelet components  $\{g_k\}$  may be very different and using  $\max_{j,k} |d_{j,k}|$  as a measure of the size of all of them seems very crude. A more natural norm in many contexts would be  $\|\mathbf{d}_0\| + \dots + \|\mathbf{d}_{n-1}\|$ . This leads to an alternative definition of stability.

**Definition 3.2.** Let  $f_n = f_0 + \sum_{k=0}^{n-1} g_k$  be a wavelet decomposition with  $g_k = \psi_k^T \mathbf{d}_k$ . The wavelet basis is said to be stable if there are constants  $K_1$  and  $K_2$  such that

$$K_1^{-1} \|(\mathbf{c}_0, \mathbf{d}_0, \mathbf{d}_1, \dots, \mathbf{d}_{n-1})\| \leq \|f_n\|_{\infty} \leq K_2 (\|\mathbf{c}_0\| + \|\mathbf{d}_1\| + \dots + \|\mathbf{d}_{n-1}\|). \quad (3.7)$$

The right-hand side of this stability estimate involves the norm of the coefficients of all the  $g_k$ , which is natural in most applications since they represent different frequency components of the underlying function. However, one may wonder why we have retained the coarser norm on the left-hand side of (3.7)? The main reason is that from the inequalities in (3.7) we can obtain a natural perturbation result. If  $f = \sum_{k=0}^{n-1} g_k = \sum_{k=0}^{n-1} \psi_k^T \mathbf{d}_k$  and  $\tilde{f} = \sum_{k=0}^{n-1} \tilde{g}_k = \sum_{k=0}^{n-1} \psi_k^T \tilde{\mathbf{d}}_k$  is a perturbation of  $f$ , then we have

$$\frac{\|f - \tilde{f}\|}{\|f\|} \leq K_1 K_2 \left( \frac{\|\mathbf{c}_0\|}{\delta} + \frac{\|\mathbf{d}_0\|}{\delta} + \dots + \frac{\|\mathbf{d}_{n-1}\|}{\delta} \right),$$

where  $\delta = \max_{j,k} |d_{j,k}|$ . In other words the relative error in  $f$  is bounded by the sum of the relative errors in the coefficients, multiplied by the factor  $K_1 K_2$  which serves as a condition number. In a perturbation result of this kind it seems more natural and convenient to scale the errors in each of the coefficient vectors by the largest wavelet coefficient rather than the sum of all the norms of the wavelet coefficients.

#### §4. Nonuniform Faber decomposition

Faber decomposition can be generalized to nonuniform partitions. As above we have  $\mathcal{F} = C[0, 1]$  equipped with the  $L^\infty$ -norm. For each nonnegative integer  $k$  we have  $V_k = S_1(\Delta_k)$ , the space of piecewise linear functions with breakpoints (knots)  $\Delta_k$ . Here the knot vector is  $\Delta_k = \mathbf{x}_k = (x_{j,k})_{j=-1}^{n_k+1}$  with  $x_{-1,k} = x_{0,k} = 0$  and  $x_{n_k,k} = x_{n_k+1,k} = 1$ , so  $\dim V_k = n_k + 1$ . The knots are nested in that

$$x_{j,k} = x_{2j,k+1} \quad \text{and} \quad x_{j,k} < x_{2j+1,k+1} < x_{j+1,k}, \quad (4.1)$$

so  $n_{k+1} = 2n_k$ . The B-splines in  $V_k$  are  $\{B_{j,\mathbf{x}_k}\}_{j=-1}^{n_k-1}$ , but we rename these as  $\phi_k = (\phi_{j,k})_{j=0}^{n_k}$ .

The projector  $Q_k : C[0, 1] \rightarrow V_k$  is the operator that assigns to each  $f$  in  $C[0, 1]$  the function in  $V_k$  which interpolates  $f$  at the (distinct) points in  $\Delta_k$ . Since we have  $\|Q_k\| = 1$  for all  $k$ , these projectors are uniformly bounded. The wavelet spaces  $\{W_k\}_{k \geq 0}$  are given by (3.1) and have dimension  $n_k$ , the number of knots added to  $\mathbf{x}_k$  to get to  $\mathbf{x}_{k+1}$ . It is easy to see that a typical  $g_k$  in  $W_k$  can be expressed on the form

$$g_k = \sum_{j=1}^{n_k} d_{j,k} \psi_{j,k} = \psi_k^T \mathbf{d}_k, \quad (4.2)$$

where  $\psi_{j,k} = \phi_{2j-1,k+1}$  for  $j = 1, \dots, n_k$ . These functions therefore form a basis for  $W_k$ . As in Section 2 we find

$$d_{j,k} = f(x_{2j-1,k+1}) - (\lambda_{j,k} f(x_{2j-2,k+1}) + (1 - \lambda_{j,k}) f(x_{2j,k+1})), \quad (4.3)$$

where

$$\lambda_{j,k} = \frac{x_{j,k} - x_{2j-1,k+1}}{x_{j,k} - x_{j-1,k}}. \quad (4.4)$$

In view of (4.1) we see that  $0 < \lambda_{j,k} < 1$  for all  $j, k$ .

The fundamental algorithms for dealing with the Faber decomposition are summarized below.

#### Reconstruction

Let  $f_k = \phi_k^T \mathbf{c}_k$  be a function in  $V_k$ . This spline can be lifted into  $V_{k+1}$  as  $f_k = \phi_{k+1}^T \mathbf{c}_{k+1}$  via the simple formulas

$$\begin{aligned} c_{2j,k+1} &= c_{j,k}, & \text{for } j = 0, \dots, n_k, \\ c_{2j+1,k+1} &= \lambda_{j,k} c_{j,k} + (1 - \lambda_{j,k}) c_{j+1,k}, & \text{for } j = 0, \dots, n_k - 1, \end{aligned}$$

with  $\lambda_{j,k}$  given by (4.4).

Similarly, a function  $g_k = \psi_k^T \mathbf{d}_k$  in  $W_k$  is lifted into  $V_{k+1}$  as  $g_k = \phi_{k+1}^T \hat{\mathbf{c}}_{k+1}$  via the formulas

$$\begin{aligned} \hat{c}_{2j,k+1} &= 0, & \text{for } j = 0, \dots, n_k, \\ \hat{c}_{2j+1,k+1} &= d_{j,k}, & \text{for } j = 0, \dots, n_k - 1. \end{aligned}$$

### Decomposition

Let  $f_{k+1} = \phi_{k+1}^T \mathbf{c}_{k+1}$  in  $V_{k+1}$  be given. Since  $V_{k+1} = V_k + W_k$ , we can split  $f$  into two components, one in  $V_k$  and one in  $W_k$ . The component in  $V_k$  is  $f_k = Q_k f = \phi_k^T \mathbf{c}_k$ , for which we find

$$c_{j,k} = c_{2j,k+1}, \quad \text{for } j = 0, \dots, n_k,$$

since every other breakpoint of  $f_{k+1}$  is also a breakpoint of  $f_k$ . The wavelet component  $g_k$  in  $W_k$  is given by  $g_k(x) = \psi_k^T \mathbf{d}_k = f_{k+1} - f_k$ , and we find

$$d_{j,k} = c_{2j-1,k+1} - (\lambda_{j,k} c_{2j-2,k+1} + (1 - \lambda_{j,k}) c_{2j,k+1}) \quad (4.5)$$

with  $\lambda_{j,k}$  given by (4.4). In the case of uniform knots this formula simplifies to (2.4).

### Stability of Faber decomposition

An investigation into the stability of the Faber decomposition requires a bit more work. The first conclusion is that the wavelet basis is a weakly stable basis for  $C[0, 1]$ .

**Theorem 4.1.** *Let  $f_n = \sum_{k=0}^{n-1} g_k$  be a wavelet decomposition with  $g_k = \psi_k^T \mathbf{d}_k$  and  $f_0 = 0$ . Then*

$$\frac{1}{2} \max_{0 \leq k \leq n-1} \|\mathbf{d}_k\| \leq \left\| \sum_{k=0}^{n-1} \psi_k^T \mathbf{d}_k \right\| \leq n \max_{0 \leq k \leq n-1} \|\mathbf{d}_k\|. \quad (4.6)$$

For a uniform partition the constant  $n$  in the rightmost inequality can be replaced by  $2n/3 + 1$ . In the uniform case the growth  $2n/3$  is best possible in the sense that if  $d_{j,k}$  is constant for all  $j$  and all  $k$ , then

$$\|f_n - f_0\|_{L^\infty[0,1]} \geq (2n/3) \max_{1 \leq k \leq n-1} \delta_k. \quad (4.7)$$

**Proof:** For convenience we set  $f = f_n = \sum_{k=0}^{n-1} g_k$ .

From (4.3) we obtain

$$|d_{j,k}| \leq |f(x_{2j-1,k+1})| + \lambda_{j,k} |f(x_{2j-2,k+1})| + (1 - \lambda_{j,k}) |f(x_{2j,k+1})| \leq 2\|f\|_\infty$$

which leads to the first inequality in (4.6). The second inequality in (4.6) follows from the triangle inequality and the fact that  $\|g_k\| \leq \|\mathbf{d}_k\|_\infty$ ,

$$\left\| \sum_{k=0}^{n-1} g_k \right\| \leq \sum_{k=0}^{n-1} \|g_k\| \leq \sum_{k=0}^{n-1} \|\mathbf{d}_k\|_\infty \leq n \max_{0 \leq k \leq n-1} \|\mathbf{d}_k\|.$$

For the rest of the proof we work on uniform grids and recall that  $f(0) = f(1) = 0$ . The fundamental identity that we will use repeatedly is (2.4) which we now write as

$$f((2j-1)/2^{k+1}) = d_{j,k} + \frac{f((j-1)/2^k) + f(j/2^k)}{2}, \quad (4.8)$$

for  $j = 0, 1, \dots, 2^k - 1$ .

To prove (4.7) we consider the function  $f^*$  for which  $d_{j,k} = 1$  for all  $j$  and  $k$ . Our aim is to determine the maximum of  $f^*$  on each grid  $\Delta_k = \{j2^{-k}\}_{j=0}^{2^k}$ ; let us denote this maximum by  $\alpha_k$ . Suppose that the left hand side of (4.8) equals  $\alpha_{k+1}$ . We observe that since one of  $j$  and  $j+1$  must be odd and one even, one of the two function values on the right are computed on the grid  $\Delta_k$  and one on a grid  $\Delta_\ell$  for some  $\ell < k$  (the grid point belongs to  $\Delta_i$  for  $i = \ell, \ell+1, \dots, k$ , but it does not belong to  $\Delta_{\ell-1}$ ). Since the sequence  $(\alpha_k)$  increases with  $k$ , the largest possible value for  $\alpha_{k+1}$  is obtained if the two function values on the right correspond to  $\alpha_k$  and  $\alpha_{k-1}$ . We therefore have the recurrence relation

$$\alpha_{k+1} = 1 + (\alpha_k + \alpha_{k-1})/2 \quad (4.9)$$

which has the solution

$$\alpha_k = \frac{2}{3}k + \frac{2}{9} + \frac{1}{9} \left(-\frac{1}{2}\right)^{k-1}, \quad (4.10)$$

where we have used the initial conditions  $\alpha_1 = 1$  and  $\alpha_2 = 3/2$ .

The number  $\alpha_k$  given by (4.10) will be an overestimate of the growth of  $f_k^*$  unless there is some  $j$  such that  $\alpha_{k+1} = f^*((2j+1)/2^{k+1})$ , and at the same time  $\alpha_k = f^*(j/2^k)$  and  $\alpha_{k-1} = f^*((j+1)/2^k)$  (possibly with  $\alpha_k$  and  $\alpha_{k-1}$  interchanged). Let us show that this does indeed happen.

On  $\Delta_1$  we have  $\alpha_1 = f^*(1/2) = 1$ , while on  $\Delta_2$  we have  $\alpha_2 = f^*(1/4) = f^*(3/4) = 3/2$ . It is therefore natural to guess that  $f^*$  attains its maximum on  $\Delta_k$  at least two times, at the two points  $a_k/2^k$  and  $b_k/2^k$  where  $a_1 = b_1 = 1$  and  $a_2 = 1$  and  $b_2 = 3$  and  $a_k < b_k$  in general. To deduce a recurrence relation for  $a_k$  we note that if (4.9) is to hold, then there must be some  $j^*$  such that  $a_{k+1} = 2j^* + 1$ , and either  $a_k = j^*$  and  $2a_{k-1} = j^* + 1$  (if  $j^*$  is odd), or  $a_k = j^* + 1$  and  $2a_{k-1} = j^*$  (if  $j^*$  is even). In either case we have  $a_{k+1} = a_k + 2a_{k-1}$ . Combined with the initial conditions we find that

$$a_k = b_{k-1} = (2^k - (-1)^k)/3.$$

This proves that the norm of  $f^*$  on the grid  $\Delta_k$  is  $\alpha_k$ , from which the last claim follows. ■

Theorem 4.1 provides a stability estimate for Faber decomposition, but as mentioned above the coefficient norm that is employed is rather coarse. From the proof of Theorem 4.1 we can also deduce a stability result based on Definition 3.2.

**Theorem 4.2.** *Let  $f_n = f_0 + \sum_{k=0}^{n-1} g_k$  be a wavelet decomposition with  $g_k = \psi_k^T \mathbf{d}_k$ . Then*

$$\frac{1}{2} \|(\mathbf{c}_0, \mathbf{d}_0, \mathbf{d}_1, \dots, \mathbf{d}_{n-1})\| \leq \|f_n\|_\infty \leq \|\mathbf{d}_0\| + \|\mathbf{d}_1\| + \dots + \|\mathbf{d}_{n-1}\|. \quad (4.11)$$

Note that the stability estimate in Theorem 4.2 is stricter than the one in Theorem 4.1 in that the estimate on the right in (4.11) is smaller than the one on the right in (4.6), while the estimate on the left remains the same. Theorem 4.2 shows that the condition number of Faber decomposition is 2.

### §5. Multiresolution based on quadratic Hermite interpolation

A standard approximation technique is to interpolate position and first derivative at two points by a cubic polynomial. A similar approximation scheme can be constructed with piecewise quadratic polynomials, see [1], [2], [3] and [8].

We start by constructing a sequence of nested spaces of quadratic splines on the interval  $[0, 1]$ . We assume that we have a sequence of knots  $(\mathbf{x}_k)_{k=0}^K$  for quadratic splines on the interval  $[0, 1]$ , i.e., the sequence  $\mathbf{x}_k = (x_{j,k})_{j=-2}^{n_k+2}$  is ordered as

$$x_{-2,k} \leq x_{-1,k} \leq x_{0,k} < x_{1,k} < \cdots < x_{n_k-1,k} < x_{n_k,k} \leq x_{n_k+1,k} \leq x_{n_k+2,k}$$

with  $[x_{0,k}, x_{n_k,k}] = [0, 1]$  and  $n_0 = 1$ . The knots should be nested as in the linear case above in the sense that  $\mathbf{x}_k$  is assumed to be obtained from  $\mathbf{x}_{k-1}$  by inserting one new knot between each knot in  $[0, 1]$ ,

$$x_{2j,k} = x_{j,k-1}, \quad \text{and} \quad x_{j,k-1} < x_{2j+1,k} < x_{j+1,k-1} \quad \text{for } j = 0, 1, \dots, n_{k-1} - 1.$$

From this it follows that  $n_k$  and  $n_{k-1}$  must be related by  $n_k = 2n_{k-1}$ . With these knots we can construct a nested sequence of spaces of quadratic splines

$$V_k = \text{span}\{B_{j,\mathbf{x}_k}\}_{j=-2}^{n_k-1}, \quad \dim V_k = n_k + 2,$$

where  $B_{j,\mathbf{x}_k}$  is the quadratic B-spline with knots  $(x_{j,k}, x_{j+1,k}, x_{j+2,k}, x_{j+3,k})$ . For simplicity we introduce the notation  $\phi_{j,k} = B_{j,\mathbf{x}_k}$ .

For  $k \geq 1$ , we have an approximation operator  $Q_k$  that constructs quadratic splines in  $V_k$  by interpolating functions in  $C^1[0, 1]$  (functions on  $[0, 1]$  that are continuous and have a continuous derivative), at the knots in  $\mathbf{x}_{k-1}$  that lie in  $[0, 1]$ ,

$$\left. \begin{aligned} Q_k f(x_{j,k-1}) &= f(x_{j,k-1}), \\ D(Q_k f)(x_{j,k-1}) &= (Df)(x_{j,k-1}), \end{aligned} \right\} \quad \text{for } j = 0, 1, \dots, n_{k-1}$$

(note that  $Q_K$  does not make sense for  $k < 1$ ). If  $Q_k f$  is written in terms of the B-splines in  $V_k$  as  $Q_k f = \phi_k^T \mathbf{c}_k$  we find from this

$$\begin{aligned} c_{2j,k} &= f(x_{2j+2}) - Df(x_{2j+2})(x_{2j+2} - x_{2j+1})/2, \\ c_{2j+1,k} &= f(x_{2j+2}) + Df(x_{2j+2})(x_{2j+3} - x_{2j+2})/2, \end{aligned} \tag{5.1}$$

for  $j = -1, \dots, n_k/2 - 1$ . Here the second subscript  $k$  has been omitted from all the  $x$ 's. Note that these operators satisfy the important identity  $Q_{k-1}Q_k = Q_{k-1}$ ; this can be exploited during wavelet decomposition and in obtaining stability results, see below.

#### Uniform boundedness of the operators

As we have seen above it is essential that the approximation operators  $\{Q_k\}$  are uniformly bounded. This is indeed the case.



**Lemma 5.1.** Let  $\|\cdot\|$  denote the  $L^\infty[0,1]$ -norm and let  $\|f\|^1 = \max\{\|f\|, \|Df\|\}$  be the norm in  $C^1[0,1]$ . For any  $f \in C^1[0,1]$ , the interpolant  $Q_k f$  satisfies the inequalities

$$\begin{aligned} |(Q_k f)(x)| &\leq \|f\| + h_0 \|Df\|/2 \\ |D(Q_k f)(x)| &\leq 3\|Df\|, \end{aligned} \quad (5.2)$$

where  $h_0 = \max_j(x_{j+1,0} - x_{j,0})$ . In other words, the operator  $Q_k$  is bounded independently of  $k$ ,

$$\|Q_k\|^1 \leq 3 + \frac{h_0}{2} \quad \text{for all } k. \quad (5.3)$$

**Proof:** In this proof all second subscripts to  $c$ 's and  $x$ 's have been omitted, but are understood to be  $k$ . We write  $Q_k f$  as  $Q_k f = \phi_k^T \mathbf{c}_k$ , with the coefficients given by (5.1). By standard properties of B-splines we have

$$\begin{aligned} |Q_k f(x)| &\leq \max_j |c_j|, \\ |D(Q_k f)(x)| &\leq \max_j 2|c_j - c_{j-1}|/(x_{j+2} - x_j). \end{aligned} \quad (5.4)$$

From (5.1) we conclude immediately that  $|c_j| \leq \|f\| + h_0 \|Df\|/2$ . To bound the derivative we consider even and odd coefficients separately. We find from (5.1) that

$$2 \frac{c_{2j+1} - c_{2j}}{x_{2j+3} - x_{2j+1}} = Df(x_{2j+2})$$

while

$$\begin{aligned} 2 \frac{c_{2j} - c_{2j-1}}{x_{2j+2} - x_{2j}} &= 2 \frac{f(x_{2j+2}) - f(x_{2j})}{x_{2j+2} - x_{2j}} - \\ &\quad \frac{(x_{2j+2} - x_{2j+1})Df(x_{2j+2}) + (x_{2j+1} - x_{2j})Df(x_{2j})}{x_{2j+2} - x_{2j}}. \end{aligned}$$

From this the second bound in (5.2) follows. ■

As usual, the wavelet space  $W_k$  consists of the error functions that result when functions in  $V_{k+1}$  are approximated from  $V_k$ ,

$$W_k = \{f \in V_{k+1} \mid Q_k f = 0\}, \quad \dim W_k = \dim V_{k+1} - \dim V_k = n_k.$$

As in the linear case we have a basis for  $W_k$  consisting of some of the B-splines in  $V_{k+1}$ . Specifically, we have  $W_k = \text{span}\{\psi_{j,k}\}_{j=1}^{n_k}$  where

$$\left. \begin{aligned} \psi_{2j+1,k} &= B_{4j,k+1}, \\ \psi_{2j+2,k} &= B_{4j+1,k+1}, \end{aligned} \right\} \quad \text{for } j = 0, \dots, n_{k+1}/4 - 1.$$

These are all the B-splines in  $V_{k+1}$  that vanish at  $x_{j,k-1}$  for  $j = 0, \dots, n_{k-1}$  (note that  $n_{k+1}/4 = n_{k-1}$ ).

### The reconstruction algorithm

Suppose that  $f_k = \phi_k^T \mathbf{c}_k$  is a spline in  $V_k$ . Since  $V_k \subseteq V_{k+1}$  it can also be written  $f_k = \phi_{k+1}^T \mathbf{c}_{k+1}$  where the coefficients  $\mathbf{c}_{k+1}$  are given by

$$\begin{aligned} c_{2j-2,k+1} &= \frac{(x_{2j+2} - x_{2j-1})c_{j-2,k} + (x_{2j-1} - x_{2j-2})c_{j-1,k}}{x_{2j+2} - x_{2j-2}}, \\ c_{2j-1,k+1} &= \frac{(x_{2j+2} - x_{2j+1})c_{j-2,k} + (x_{2j+1} - x_{2j-2})c_{j-1,k}}{x_{2j+2} - x_{2j-2}}, \end{aligned} \quad (5.5)$$

for  $j = 0, \dots, n_k$  (the second subscript in all the  $x$ 's is  $k+1$ ). Converting from a representation in  $W_k$  to a representation in  $V_{k+1}$  is simpler. If  $g_k = \psi_k^T \mathbf{d}_k = \phi_{k+1}^T \tilde{\mathbf{c}}_{k+1}$  we have  $\tilde{c}_{4j-2,k+1} = \tilde{c}_{4j-1,k+1} = 0$  for  $j = 0, \dots, n_{k+1}/4 - 1$  and

$$\left. \begin{aligned} \tilde{c}_{4j,k+1} &= d_{2j+1,k} \\ \tilde{c}_{4j+1,k+1} &= d_{2j+2,k} \end{aligned} \right\} \quad \text{for } j = 0, \dots, n_{k+1}/4 - 1. \quad (5.6)$$

### The decomposition algorithm

A spline in  $f_{k+1} = \phi_{k+1}^T \mathbf{c}_{k+1}$  in  $V_{k+1}$  can be decomposed as  $f_{k+1} = f_k + g_k$  where  $f_k = \phi_k^T \mathbf{c}_k$  is in  $V_k$  and  $g_k = \psi_k^T \mathbf{d}_k$  is in  $W_k$ . The coefficients of  $f_k$  are obtained by solving the two equations (5.5) with respect to  $c_{j-2,k}$  and  $c_{j-1,k}$ ,

$$\begin{aligned} c_{2j-2,k} &= \frac{(x_{4j+1} - x_{4j-2})c_{4j-2,k+1} - (x_{4j-1} - x_{4j-2})c_{4j-1,k+1}}{x_{4j+1} - x_{4j-1}} \\ c_{2j-1,k} &= \frac{-(x_{4j+2} - x_{4j+1})c_{4j-2,k+1} + (x_{4j+2} - x_{4j-1})c_{4j-1,k+1}}{x_{4j+1} - x_{4j-1}} \end{aligned} \quad (5.7)$$

for  $j = 0, \dots, n_k/2$ . Here the second subscript to the  $x$ 's is  $k+1$ , and we have replaced  $j$  by  $2j$  since we have to compute the coefficients in pairs.

In order to write down the coefficients of  $g_k = \psi_k^T \mathbf{d}_k = f_{k+1} - f_k$  it is convenient to apply the reconstruction algorithm (5.5) and express  $f_k$  in terms of the B-spline basis in  $V_{k+1}$  as  $f_k = \phi_{k+1}^T \hat{\mathbf{c}}_{k+1}$ . Then the coefficients  $\mathbf{d}_k$  of  $g_k$  are given by

$$d_{2j+i+1,k} = c_{4j+i-1,k+1} - \hat{c}_{4j+i-1,k+1}, \quad \text{for } i = 0, 1 \quad (5.8)$$

and  $j = 0, \dots, n_k/2 - 1$ .

### Stability

Before considering the stability of the proposed wavelet transform in detail, we need another identity for decomposition. The formulas in (5.7) project down from a spline in  $V_{k+1}$  to a spline in  $V_k$ . It is equally simple to project down from  $V_K$  directly into  $V_k$ . If  $f_K = \phi_K^T \mathbf{c}_K$ , the approximation  $f_k = \phi_k^T \mathbf{c}_k$  in  $V_k$  is given by

$$\begin{aligned} c_{2j-2,k} &= \frac{(x_{J+1,K} - x_{2j-1,k})c_{J-2,K} - (x_{J-1,K} - x_{2j-1,k})c_{J-1,K}}{x_{J+1,K} - x_{J-1,K}} \\ c_{2j-1,k} &= \frac{-(x_{2j+1,k} - x_{J+1,K})c_{J-2,K} + (x_{2j+1,k} - x_{J-1,K})c_{J-1,K}}{x_{J+1,K} - x_{J-1,K}} \end{aligned} \quad (5.9)$$

for  $j = 0, \dots, n_k/2$ , where  $J = 2^{K-k+1}j$ .

Wavelet decomposition and reconstruction is nothing but changes of bases between the two bases for  $V_K$ , namely  $\phi_K$  and  $(\phi_1, \psi_1, \dots, \psi_{K-1})$ . It is therefore essential that these changes are stable so that the computations can be performed without encountering serious problems with rounding errors.

**Lemma 5.2.** *Let  $f$  be a function in  $V_K$  with the representations*

$$f = \phi_K^T \mathbf{c}_K = \phi_1^T \mathbf{c}_1 + \psi_1^T \mathbf{d}_1 + \cdots + \psi_{K-1}^T \mathbf{d}_{K-1}$$

*in the two bases  $\phi_K$  and  $(\phi_1, \psi_1, \dots, \psi_{K-1})$  for  $V_K$ , and let  $\|\cdot\|$  denote the vector max-norm. Then*

$$\kappa_K^{-1} \|(\mathbf{c}_1, \mathbf{d}_1, \dots, \mathbf{d}_{K-1})\| \leq \|\mathbf{c}_K\| \leq \|\mathbf{c}_1\| + \|\mathbf{d}_1\| + \cdots + \|\mathbf{d}_{K-1}\|, \quad (5.10)$$

where

$$\kappa_K = 2 \left( 1 + \max_{1 \leq k \leq K-2} \{r_k + r_{k+1}\} \right) \quad (5.11)$$

and  $r_k$  is defined by

$$r_k = \max_{0 \leq j \leq n_k/2} \left\{ \frac{x_{J-1,K} - x_{2j-1,k}}{x_{J+1,K} - x_{J-1,K}}, \frac{x_{2j+1,k} - x_{J+1,K}}{x_{J+1,K} - x_{J-1,K}} \right\} \quad (5.12)$$

with  $J = 2^{K-k+1}j$ .

**Proof:** We start with the right-most inequality. Let  $R_{K-1}$  denote the matrix that represents reconstruction of B-spline coefficients from  $V_{K-1}$  to  $V_K$ , i.e., if  $\tilde{f} = \phi_{K-1}^T \tilde{\mathbf{c}}_{K-1} = \phi_K^T \tilde{\mathbf{c}}_K$  we have  $\tilde{\mathbf{c}}_K = \mathbf{R}_{K-1} \tilde{\mathbf{c}}_{K-1}$ . But then  $\mathbf{c}_K = \mathbf{R}_{K-1} \mathbf{c}_{K-1} + \ddot{\mathbf{d}}_{K-1}$  where  $\ddot{\mathbf{d}}_{K-1}$  denotes the coefficients of  $g_{K-1}$  expressed in terms of  $\phi_K$  according to the reconstruction formulas (5.6), i.e., the coefficients  $\mathbf{d}_{K-1}$  are augmented with a certain number of zeros. The reconstruction formulas (5.5) amount to taking weighted averages of neighbouring coefficients in  $\mathbf{c}_{K-1}$  and therefore  $\|\mathbf{R}_{K-1}\| = 1$  (in the  $\ell^\infty$ -norm for matrices). Hence we have

$$\|\mathbf{c}_K\| = \|\mathbf{R}_{K-1} \mathbf{c}_{K-1} + \ddot{\mathbf{d}}_{K-1}\| \leq \|\mathbf{c}_{K-1}\| + \|\mathbf{d}_{K-1}\|.$$

Applying the same argument repeatedly to  $\mathbf{c}_{K-1}, \mathbf{c}_{K-2}, \dots, \mathbf{c}_2$  leads to the right-most inequality in (5.10).

To prove the left inequality in (5.10) we make use of the fact that we can project straight from  $V_K$  to  $V_k$  with the formulae in (5.9), without going via the intermediate spaces. The formulae in (5.9) may be written as  $\mathbf{c}_k = \mathbf{P}_k \mathbf{c}_K$  where  $\mathbf{P}_k$  is a  $(n_K + 2) \times (n_k + 2)$ -matrix with two nonzero entries in each row that sum to one. We then have the bound

$$\|\mathbf{c}_k\| \leq \|\mathbf{P}_k\| \|\mathbf{c}_K\|.$$

From (5.9) we find that

$$\|\mathbf{P}_k\| \leq \max_{0 \leq j \leq n_k/2} \left\{ \frac{x_{J+1,K} + x_{J-1,K} - 2x_{2j-1,k}}{x_{J+1,K} - x_{J-1,K}}, \frac{2x_{2j+1,k} - x_{J+1,K} - x_{J-1,K}}{x_{J+1,K} - x_{J-1,K}} \right\},$$

where  $J = 2^{K-k+1}j$ . This may be simplified to

$$\|\mathbf{P}_k\| \leq 1 + 2r_k \quad (5.13)$$

where  $r_k$  is given by (5.12). In particular we have

$$\|\mathbf{c}_1\| \leq (1 + 2r_1)\|\mathbf{c}_K\|. \quad (5.14)$$

To bound the wavelet coefficients we note that  $\mathbf{d}_k$  may be expressed as

$$\ddot{\mathbf{d}}_k = \mathbf{c}_{k+1} - \mathbf{R}_k \mathbf{c}_k = \mathbf{P}_{k+1} \mathbf{c}_K - \mathbf{R}_K \mathbf{P}_k \mathbf{c}_K,$$

where  $\ddot{\mathbf{d}}_k$  denotes the wavelet coefficients at level  $k$  with the zeros interspersed. From (5.13) we then obtain

$$\|\mathbf{d}_k\| \leq (\|\mathbf{P}_{k+1}\| + \|\mathbf{P}_k\|)\|\mathbf{c}_K\| \leq 2(1 + r_k + r_{k-1})\|\mathbf{c}_K\| \quad (5.15)$$

for  $1 \leq k \leq K-2$  and

$$\|\mathbf{d}_{K-1}\| \leq (2 + r_{K-1})\|\mathbf{c}_K\|. \quad (5.16)$$

The largest of the right-hand sides in (5.14)–(5.16) is  $\kappa_K$  which leads to the left inequality in (5.10). ■

Lemma 5.2 is a statement about the stability of the wavelet algorithms. From this we get a result about the stability of the wavelet basis.

**Lemma 5.3.** *Let  $f$  be a function in  $V_K$  with the representation*

$$f = \phi_1^T \mathbf{c}_1 + \psi_1^T \mathbf{d}_1 + \cdots + \psi_{K-1}^T \mathbf{d}_{K-1}$$

*in the wavelet basis for  $V_K$ . Then*

$$3^{-1} \kappa_K^{-1} \|(\mathbf{c}_1, \mathbf{d}_1, \dots, \mathbf{d}_{K-1})\| \leq \|f\| \leq \|\mathbf{c}_1\| + \|\mathbf{d}_1\| + \cdots + \|\mathbf{d}_{K-1}\|, \quad (5.17)$$

where  $\kappa_K$  is given by (5.11).

**Proof:** This result follows by combining the classical stability estimate for quadratic B-splines,

$$3^{-1} \|\mathbf{c}_K\| \leq \|f\| \leq \|\mathbf{c}_K\|,$$

with Lemma 5.2. ■

This stability result is not particularly good since the condition number  $3\kappa_K$  depends on the knots and can become very large. There is a possibility that the projection from  $V_K$  to  $V_k$  can be performed by more well conditioned formulas than (5.9), but this seems unlikely. In the special case of uniform knots we have  $x_{j,k} = j2^{-k}$  and  $r_k = 2^{K-k-1} - 1/2$ ; from (5.11) we then see that

$$3\kappa_K = 6(1 + 2^{K-2} - 1/2 + 2^{K-3} - 1/2) = 9 \cdot 2^{K-2}.$$

### Decomposition in practice

In a wavelet environment, the initial operation to be performed on a function is normally wavelet decomposition. So given a function  $f$  in  $C^1[0, 1]$ , we first pick a finest grid  $\mathbf{x}_K$  and compute a quadratic spline approximation  $f_K = \phi_K^T \mathbf{c}_K$  using the formulas (5.1). We can then apply the decomposition algorithms (5.7) and (5.8) successively and decompose  $f_K$  as  $f_K = f_1 + g_1 + \cdots + g_{K-1}$ . However, because of the simplicity of our approximation scheme, there is another possibility. Instead of computing the B-spline coefficients of  $f_k$  for  $k < K$  by successively applying (5.7), we can use the explicit formulas for the coefficients (5.1). If the values of  $f$  and  $Df$  on the finest grid are computed initially (this must be done even when the standard decomposition algorithms are used) and then stored, this approach will be faster than successive applications of (5.7) since the formulas (5.1) are arithmetically simpler than the formulas in (5.7).

In certain situations the explicit formulas (5.1) for the B-spline coefficients can be exploited to great advantage. Suppose for example that  $f$  and  $Df$  are very expensive to compute and we want a compressed representation of  $f$  with as few nonzero wavelet coefficients as possible. If we follow the traditional approach and start at a fine level we are very likely to compute many function and derivative values that are later discarded. However, the explicit formulas may allow us to compute the wavelet decomposition bottom-up instead of top-down. First we would have to compute the lowest level approximation to  $f$  on a grid  $\mathbf{x}_1$ . We would then have to decide in which areas this approximation is not satisfactory and decide where further sampling of  $f$  and its derivative is required. If  $f$  is reasonably smooth there should only be isolated areas where the approximation is not good enough. This means that many of the wavelet coefficients ( $d_{j,1}$ ) can be set to zero so there is no need to compute the corresponding function and derivative values. The nonzero coefficients can then be computed by formula (5.8) where only the coefficients  $c_{4j+i-1,2}$  will require new function evaluations. After this update of the approximation, its quality is assessed again and new wavelet coefficients added at level 2. The main challenge of this bottom-up approach to computing the wavelet decomposition is deciding where more information is needed, which is an inherent problem with adaptive computations.

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