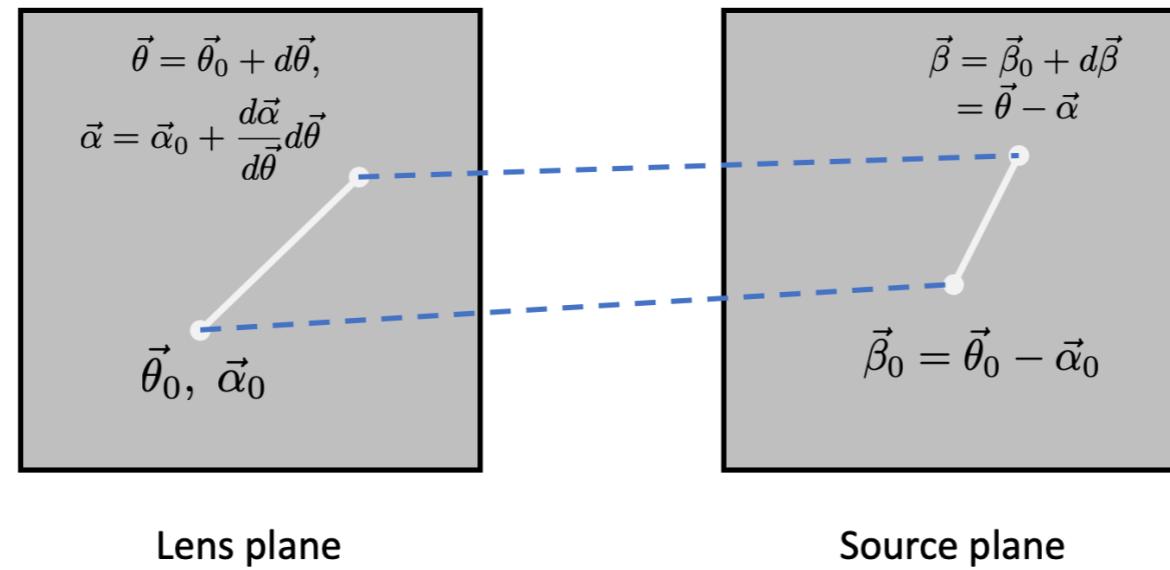


GRAVITATIONAL LENSING

5 - FLEXION, TIME DELAYS (1)

Massimo Meneghetti
AA 2019-2020

LENS MAPPING AT FIRST ORDER



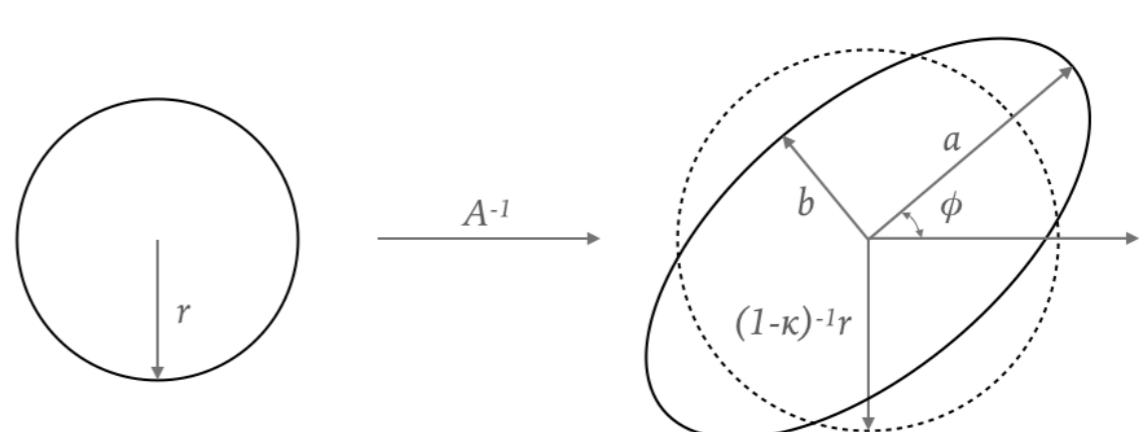
Lens plane

Source plane

$$(\vec{\beta} - \vec{\beta}_0) = \left(I - \frac{d\vec{\alpha}(\vec{\theta})}{d\vec{\theta}} \right) (\vec{\theta} - \vec{\theta}_0) \quad A \equiv \frac{d\vec{\beta}}{d\vec{\theta}} = \left(\delta_{ij} - \frac{\partial \alpha_i(\vec{\theta})}{\partial \theta_j} \right) = \left(\delta_{ij} - \frac{\partial^2 \hat{\Psi}(\vec{\theta})}{\partial \theta_i \partial \theta_j} \right) = \left(\delta_{ij} - \hat{\Psi}_{ij} \right)$$

$$\vec{\beta}_0 = (0,0) \quad \vec{\theta}_0 = (0,0)$$

$$\beta_i \simeq \sum_j A_{ij} \theta_j = \sum_j \frac{\partial \beta_i}{\partial \theta_j} \theta_j$$



LENS MAPPING AT FIRST ORDER

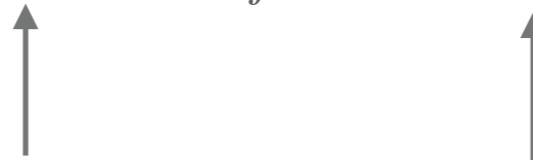
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LENS MAPPING AT FIRST ORDER

$$\beta_i \simeq \sum_j A_{ij} \theta_j = \sum_j \frac{\partial \beta_i}{\partial \theta_j} \theta_j + \frac{1}{2} \sum_j \sum_k \frac{\partial^2 \beta_i}{\partial \theta_j \partial \theta_k} \theta_j \theta_k$$

LENS MAPPING AT FIRST ORDER

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$$D_{ijk} = \frac{\partial^2 \beta_i}{\partial \theta_j \partial \theta_k} = \frac{\partial A_{ij}}{\partial \theta_k}$$

The second order term in the lens equation can be expressed in terms of a $2 \times 2 \times 2$ tensor D

$$A_{ij} = \frac{\partial \beta_i}{\partial \theta_j}$$

$$D_{ij1} = \begin{pmatrix} -2\gamma_{1,1} - \gamma_{2,2} & -\gamma_{2,1} \\ -\gamma_{2,1} & -\gamma_{2,2} \end{pmatrix} \quad D_{ij2} = \begin{pmatrix} -\gamma_{2,1} & -\gamma_{2,2} \\ -\gamma_{2,2} & 2\gamma_{1,2} - \gamma_{2,1} \end{pmatrix}$$

The elements of D are first derivatives of the shear

$$\gamma_{i,j} = \frac{\partial \gamma_i}{\partial \theta_j}$$

Example: calculate D_{111}

$$D_{ij1} = \begin{pmatrix} \frac{\partial A_{i1}}{\partial \theta_1} & \frac{\partial A_{i2}}{\partial \theta_1} \\ \frac{\partial A_{j1}}{\partial \theta_1} & \frac{\partial A_{j2}}{\partial \theta_1} \end{pmatrix} \quad A_{ij} = \begin{pmatrix} 1 - \varphi_{11} & -\varphi_{12} \\ -\varphi_{21} & 1 - \varphi_{22} \end{pmatrix}$$

$$= \begin{pmatrix} -\varphi_{111} & -\varphi_{121} \\ -\varphi_{211} & -\varphi_{221} \end{pmatrix}$$

Remember:

$$\gamma_1 = \frac{1}{2}(\varphi_{11} - \varphi_{22}) \quad \gamma_2 = \varphi_{12} = \varphi_{21}$$

$$\gamma_{1,1} = \frac{\partial \gamma_1}{\partial \theta_1} = \frac{1}{2}(\varphi_{111} - \varphi_{221}) \quad \gamma_{2,2} = \frac{\partial \gamma_2}{\partial \theta_2} = \varphi_{122}$$

then, we see that: $2\gamma_{1,1} + \gamma_{2,2} = \varphi_{111}$

$$\Rightarrow D_{111} = -\varphi_{111} = -2\gamma_{1,1} - \gamma_{2,2}$$

COMPLEX NOTATION

In the lensing literature, it is common to use the complex notation to represent vectors or pseudo-vectors

$$\vec{v} = (v_1, v_2) \rightarrow v = v_1 + i v_2$$

For example:

$$\vec{\alpha} = (\alpha_1, \alpha_2) \rightarrow \alpha = \alpha_1 + i \alpha_2$$

$$\vec{\gamma} = (\gamma_1, \gamma_2) \rightarrow \gamma = \gamma_1 + i \gamma_2$$

we can also define complex differential operators:

$$\partial = \partial_1 + i \partial_2$$

$$\partial^\dagger = \partial_1 - i \partial_2$$

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we can also define complex differential operators:

$$\partial = \partial_1 + i\partial_2$$

Spin raising operator

$$\partial^\dagger = \partial_1 - i\partial_2$$

Spin lowering operator

COMPLEX NOTATION

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$$\partial \hat{\Psi} = \partial_1 \hat{\Psi} + i \partial_2 \hat{\Psi} = \alpha_1 + i \alpha_2 = \alpha$$

From spin-0 scalar field to spin-1 vector field (deflection angle)

COMPLEX NOTATION

$$\partial \hat{\Psi} = \partial_1 \hat{\Psi} + i \partial_2 \hat{\Psi} = \alpha_1 + i \alpha_2 = \alpha$$

From spin-0 scalar field to spin-1 vector field (deflection angle)

$$\partial^\dagger \partial = \partial_1^2 + \partial_2^2 = \Delta$$

$$\partial^\dagger \partial \hat{\Psi} = \Delta \hat{\Psi} = 2\kappa$$

From spin-1 vector field to spin-0 scalar field

COMPLEX NOTATION

$$\partial \hat{\Psi} = \partial_1 \hat{\Psi} + i \partial_2 \hat{\Psi} = \alpha_1 + i \alpha_2 = \alpha$$

From spin-0 scalar field to spin-1 vector field (deflection angle)

$$\partial^\dagger \partial = \partial_1^2 + \partial_2^2 = \Delta$$

$$\partial^\dagger \partial \hat{\Psi} = \Delta \hat{\Psi} = 2\kappa$$

From spin-1 vector field to spin-0 scalar field

$$\frac{1}{2} \partial \partial \hat{\Psi} = \frac{1}{2} \partial \alpha = \frac{1}{2} (\partial_1 + i \partial_2) (\alpha_1 + i \alpha_2)$$

The shear is a spin-2 field

$$= \frac{1}{2} [(\partial_1 \alpha_1 - \partial_2 \alpha_2) + i(\partial_2 \alpha_1 + \partial_1 \alpha_2)]$$

$$= \gamma_1 + i \gamma_2 = \gamma$$

COMPLEX NOTATION

$$F = \frac{1}{2} \partial \partial^\dagger \partial \hat{\Psi} = \partial \kappa$$

$$G = \frac{1}{2} \partial \partial \partial \hat{\Psi} = \partial \gamma$$

COMPLEX NOTATION

$$F = \frac{1}{2} \partial \partial^\dagger \partial \hat{\Psi} = \partial \kappa \quad \text{Spin-1}$$

$$G = \frac{1}{2} \partial \partial \partial \hat{\Psi} = \partial \gamma \quad \text{Spin-3}$$

COMPLEX NOTATION

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$$F = F_1 + iF_2 = (\gamma_{1,1} + \gamma_{2,2}) + i(\gamma_{2,1} - \gamma_{1,2})$$

$$G = G_1 + iG_2 = (\gamma_{1,1} - \gamma_{2,2}) + i(\gamma_{2,1} + \gamma_{1,2})$$

$$D_{ij1} = \begin{pmatrix} -2\gamma_{1,1} - \gamma_{2,2} & -\gamma_{2,1} \\ -\gamma_{2,1} & -\gamma_{2,2} \end{pmatrix} \quad D_{ij2} = \begin{pmatrix} -\gamma_{2,1} & -\gamma_{2,2} \\ -\gamma_{2,2} & 2\gamma_{1,2} - \gamma_{2,1} \end{pmatrix}$$

$$D_{111} = -2\gamma_{11} - \gamma_{22} = -\frac{1}{2}(3F_1 + G_1)$$

$$D_{211} = D_{121} = D_{112} = -\gamma_{21} = -\frac{1}{2}(F_2 + G_2)$$

$$D_{122} = D_{212} = D_{221} = -\gamma_{22} = -\frac{1}{2}(F_1 - G_1)$$

$$D_{222} = 2\gamma_{12} - \gamma_{21} = -\frac{1}{2}(3F_2 - G_2)$$

Example: derivation of \bar{F}

$$\bar{F} = 2k \quad n = \frac{1}{2}(t_{11} + t_{22}) \quad r_1 = \frac{1}{2}(t_{11} - t_{22})$$
$$\gamma_2 = \gamma_{12} = \gamma_{21}$$

$$\bar{F} = \frac{\partial k}{\partial \theta_1} + i \frac{\partial k}{\partial \theta_2} = \frac{1}{2}(t_{111} + t_{221}) + i \frac{1}{2}(t_{112} + t_{222})$$

Note that:

$$r_{1,1} = \frac{1}{2}(t_{111} - t_{221}) \quad r_{2,1} = \gamma_{121}$$
$$\gamma_{1,2} = \frac{1}{2}(t_{112} - t_{222}) \quad \gamma_{2,2} = \gamma_{122}$$

$$\Rightarrow \frac{1}{2}(t_{111} + t_{221}) = \frac{1}{2}(t_{111} - t_{221} + 2\gamma_{221}) = r_{1,1} + \gamma_{2,1}$$

$$\frac{1}{2}(t_{112} + t_{222}) = \frac{1}{2}(2\gamma_{122} - t_{112} + t_{222}) = \gamma_{2,2} - r_{1,2}$$

$$\Rightarrow \bar{F} = (r_{1,1} + \gamma_{2,2}) + i(r_{2,1} - \gamma_{1,2})$$

*Notebook 5: building a app to visualise the effects of convergence,
shear, and flexion*

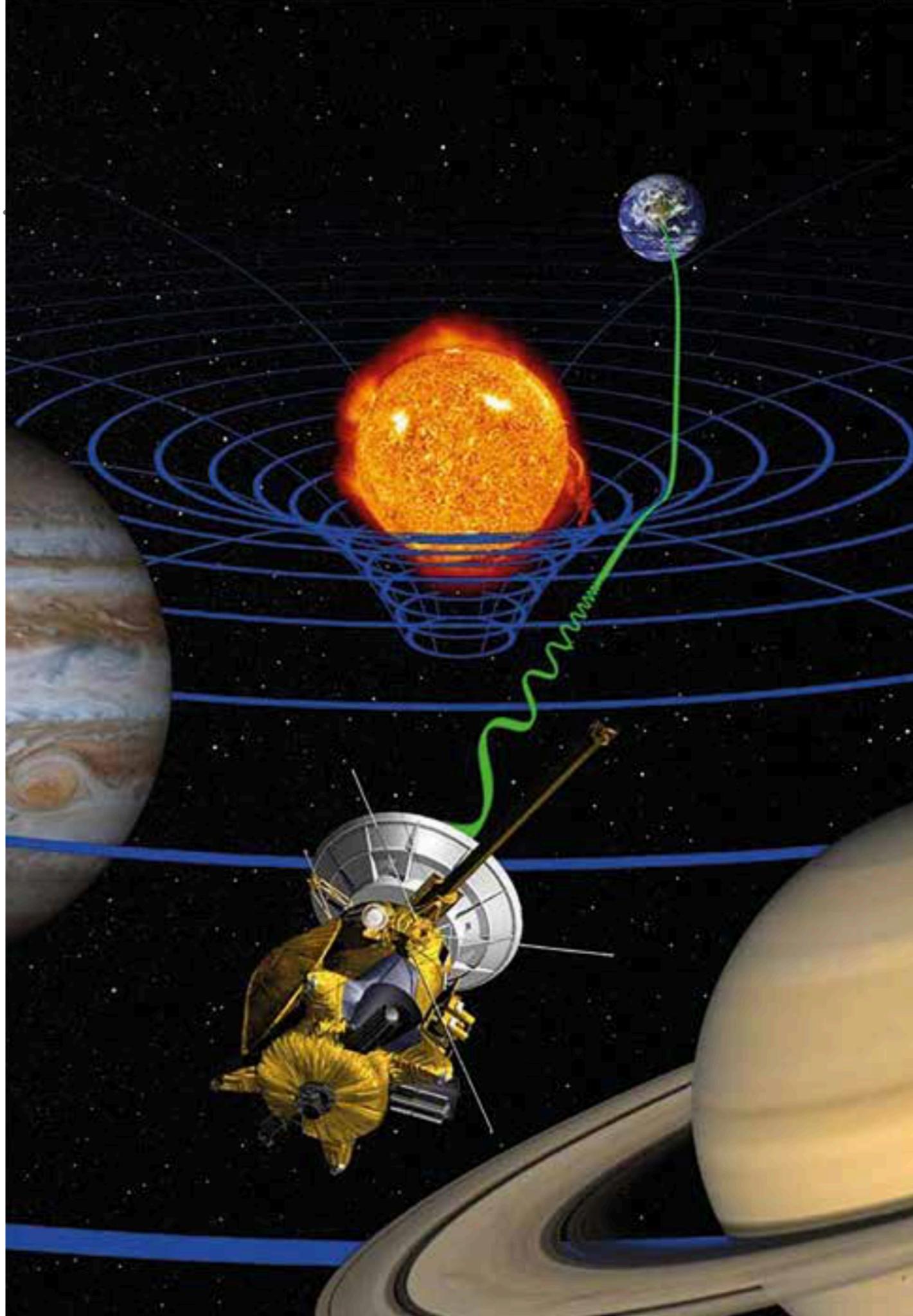
GRAVITATIONAL TIME DELAY

- In a lensing phenomenon, light travels with an effective velocity $c' < c$. As seen, this implies an effective refractive index $n > 1$
- The effective refractive index is expressed in terms of the Newtonian potential
- If we compare the travel times of two photons, one traveling at velocity c and the other at velocity c' , we notice that the second accumulates a time delay t_{grav}
- This time delay is called *gravitational* time delay, or *Shapiro* time delay (Shapiro, 1964)

$$\begin{aligned} n &= 1 - \frac{2\Phi}{c^2} \\ t_{grav} &= \int \frac{dz}{c'} - \int \frac{dz}{c} \\ &= \frac{1}{c} \int (n - 1) dz \\ &= -\frac{2}{c^3} \int \Phi dz \end{aligned}$$

SHAPIRO DELAY AS A TEST OF GR

- Send a radio signal towards another planet (Mercury, Venus, Mars) behind the sun
- Measure time needed to the signal to come back after being reflected
- Measurement done in 2003 with the Cassini spacecraft
- Delay is few ~ 100 microseconds
- GR confirmed at the level of 0.002%



GRAVITATIONAL TIME DELAY

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$$\hat{\Psi}(\vec{\theta}) = \frac{D_{LS}}{D_L D_S} \frac{2}{c^2} \int \Phi(\vec{\theta}, z) dz$$

$$t_{grav} = -\frac{1}{c} \frac{D_S D_L}{D_{LS}} \hat{\Psi}$$

GRAVITATIONAL TIME DELAY

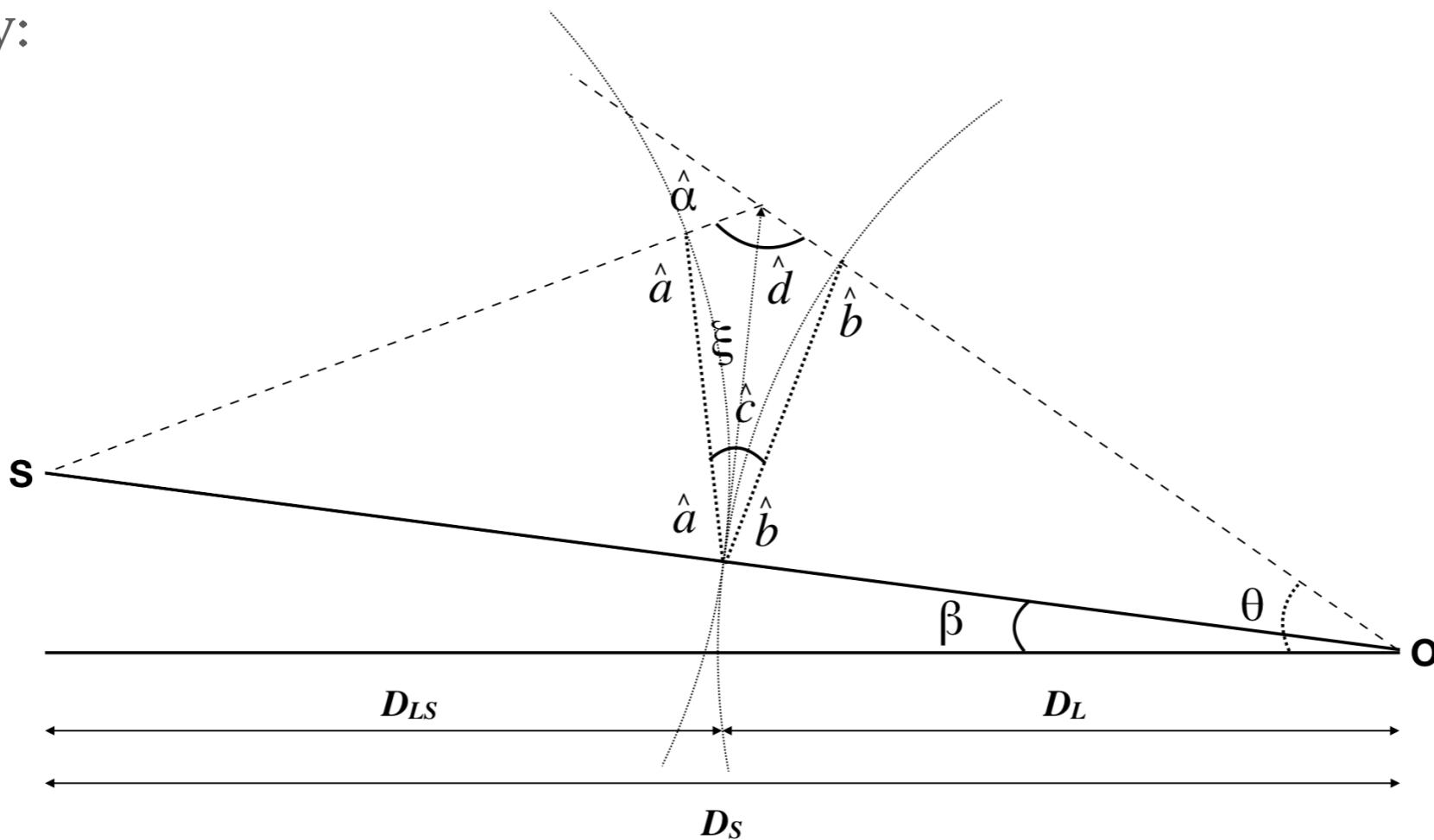
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$$t_{grav} = -\frac{1}{c} \frac{D_S D_L}{D_{LS}} \hat{\Psi}$$

This time delay does not account yet for the different path of photons!

GEOMETRICAL TIME DELAY

- We need to combine the gravitational time delay to the so called *geometrical* time delay:

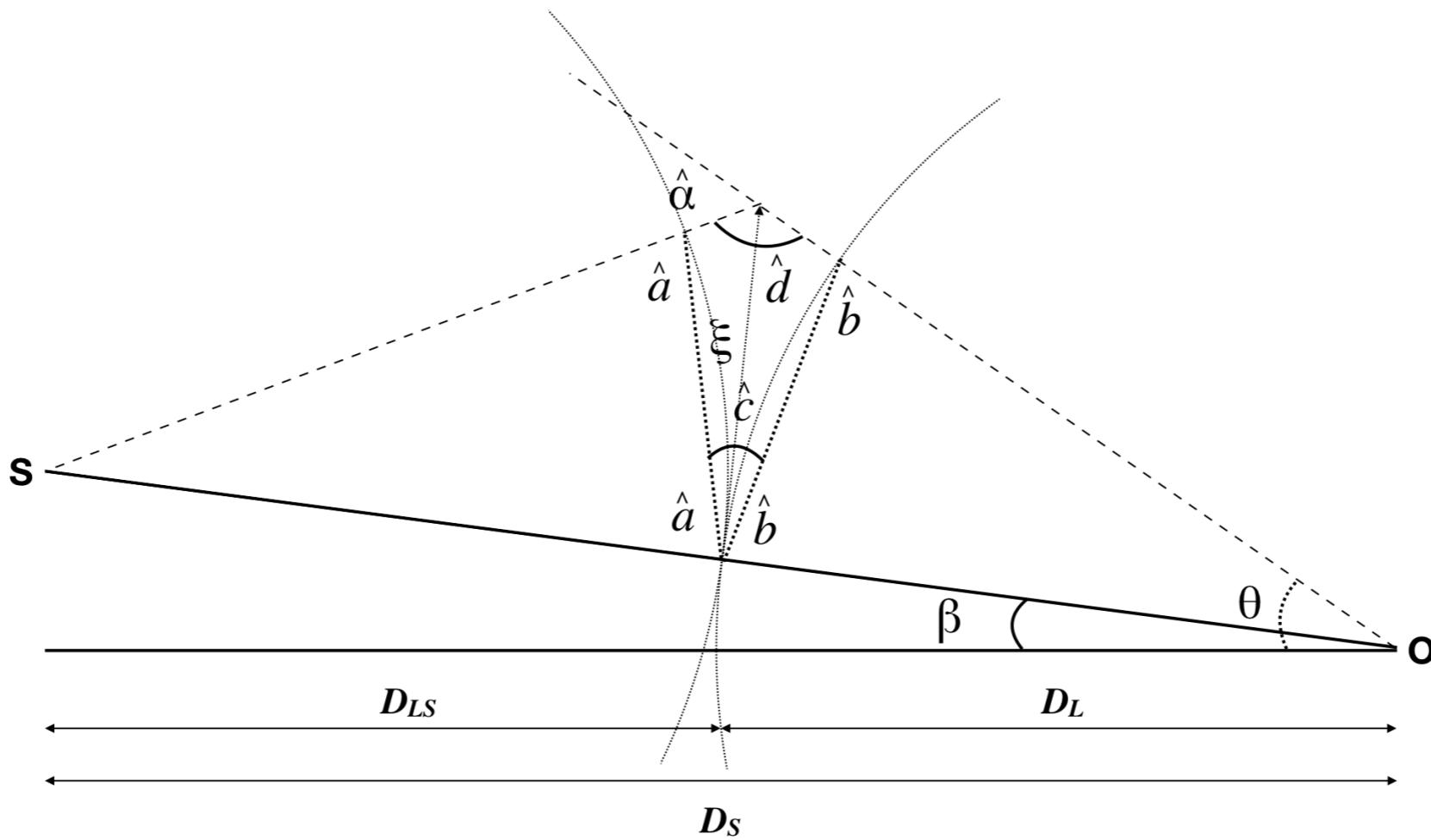


$$(\pi - \hat{a}) + (\pi - \hat{b}) + \hat{c} + \hat{d} = 2\pi \Rightarrow \hat{a} + \hat{b} = \hat{c} + \hat{d}$$

$$\hat{a} + \hat{b} + \hat{c} = \pi \Rightarrow 2\hat{c} = \pi - \hat{d}$$

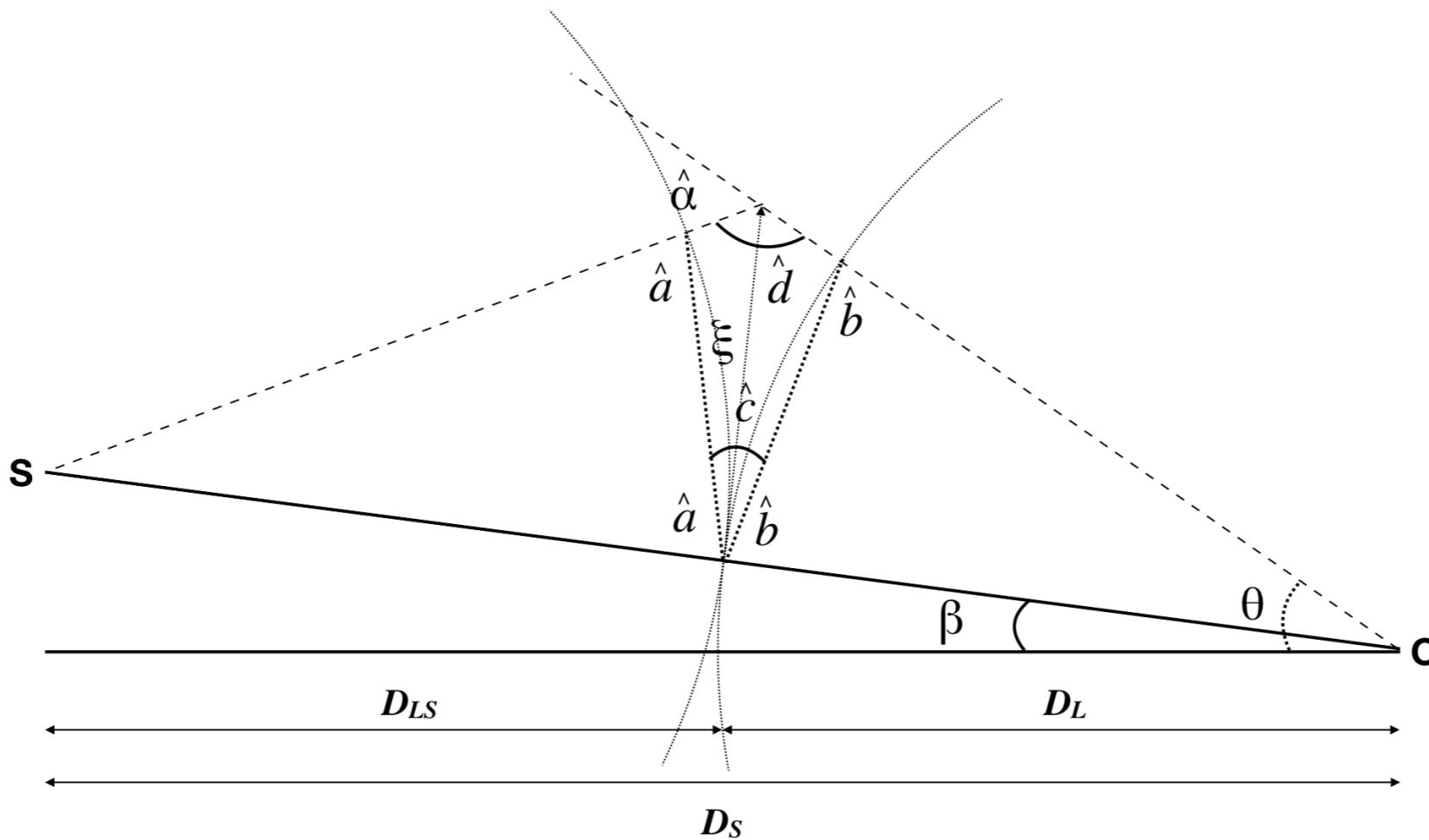
$$\hat{\alpha} + \hat{d} = \pi \Rightarrow \hat{d} = \pi - \hat{\alpha}$$

GEOMETRICAL TIME DELAY



$$\begin{aligned}
 (\pi - \hat{a}) + (\pi - \hat{b}) + \hat{c} + \hat{d} &= 2\pi \Rightarrow \hat{a} + \hat{b} = \hat{c} + \hat{d} \\
 \hat{a} + \hat{b} + \hat{c} &= \pi \Rightarrow 2\hat{c} = \pi - \hat{d} \\
 \hat{\alpha} + \hat{d} &= \pi \Rightarrow \hat{d} = \pi - \hat{\alpha}
 \end{aligned}
 \quad \Rightarrow \hat{c} = \frac{\hat{\alpha}}{2} = \frac{1}{2} \frac{D_S}{D_{LS}} \alpha = \frac{1}{2} \frac{D_S}{D_{LS}} (\theta - \beta)$$

GEOMETRICAL TIME DELAY



$$\begin{aligned}\Rightarrow \hat{c} &= \frac{\hat{\alpha}}{2} \\ &= \frac{1}{2} \frac{D_S}{D_{LS}} \alpha \\ &= \frac{1}{2} \frac{D_S}{D_{LS}} (\theta - \beta)\end{aligned}$$

$$\begin{aligned}\xi &= D_L(\theta - \beta) \\ t_{geom} &= \frac{1}{c} \xi \hat{c} \\ &= \frac{1}{2c} \frac{D_S D_L}{D_{LS}} (\theta - \beta)^2\end{aligned}$$

TOTAL TIME DELAY

$$t_{geom} = \frac{1}{2c} \frac{D_S D_L}{D_{LS}} (\theta - \beta)^2$$

$$t_{grav} = -\frac{1}{c} \frac{D_S D_L}{D_{LS}} \hat{\Psi}$$

$$\begin{aligned} t_{tot} &= t_{geom} + t_{grav} \\ &= \frac{1}{2c} \frac{D_S D_L}{D_{LS}} (\theta - \beta)^2 - \frac{1}{c} \frac{D_S D_L}{D_{LS}} \hat{\Psi}(\theta) \\ &= \frac{1}{c} \frac{D_S D_L}{D_{LS}} \left[\frac{1}{2} (\theta - \beta)^2 - \hat{\Psi}(\theta) \right] \end{aligned}$$

TOTAL TIME DELAY

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Accounting for the expansion of the universe and for the fact that this is a surface:

$$t_{tot}(\vec{\theta}) = \frac{1+z_L}{c} \frac{D_S D_L}{D_{LS}} \left[\frac{1}{2} (\vec{\theta} - \vec{\beta})^2 - \hat{\Psi}(\vec{\theta}) \right]$$