

# **An Introduction to Particle Accelerators // - Playing Accelerators, on Computer -**

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### **1 Cyclotron - resonant acceleration**

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- Non-relativistic cyclotron

- orbit :  $r = v/\omega_0 = mv/qB$

- focusing (1) :

$$F_z = qvB_r \approx qv \frac{\partial B_r}{\partial z} z \equiv qv \frac{\partial B_z}{\partial r} z$$

$$\ddot{z} - \frac{qv}{m} \frac{\partial B_z}{\partial r} z = 0 \rightarrow \omega_z^2/\omega_0^2 = \nu_z^2 = -\frac{r}{B_z} \frac{\partial B_z}{\partial r} = -k, \quad \nu_z = \sqrt{-k}.$$

hence the field index  $k$  needs be negative :  $B_z$  is slowly decreased with radius.

Similarly,  $\nu_r = 1 + k$ . This sets the requirement  $-1 < k < 0$

- focusing (2) : is also ensured at lower energy by the electric field.

- isochronism :

The condition for vertical focusing,  $-1 < k < 0$  ( $B$  is not constant), spoils the isochronism.

As a consequence, the phase is not constant (ABCDE path)

- bunching : particle beam injected into the cyclotron necessarily gets bunched, at the frequency of the RF (the an RF period)

- The classical limit ( $\gamma \approx 1$ ) is  $\sim 25$  MeV for protons, 50 MeV for D and  $\alpha$ , (about 2 – 3% increase in mass), GANIL 100 MeV/u...

- That was enough energy to transmute all nuclei... The classical cyclotron allowed discovering oodles of nuclear reactions. Yet, let's keep in mind : transmutation was not the all story

## Cyclotron (3/5) - classical

- Relativistic energies, the bad news :

- The cyclotron resonance  $\omega_0 = qB/\gamma m$ , with  $r = \beta c/\omega_0$  yields  $k = \frac{\beta}{\gamma} \frac{\partial \gamma}{\partial \beta} = \beta^2 \gamma^2$

- so k cannot satisfy  $-1 < k < 0$ ,

isochronism requires that  $B(r) \propto \gamma$ , which yields vertical defocusing...

- That was the end of the story,  $\sim 25$  MeV protons, etc... :

Hans Bethe (1937) :

“... it seems useless to build cyclotrons of larger proportions than the existing ones... an accelerating chamber of deuterons of 11 MeV energy which is the highest possible...”

Frank Cole : “If you went to graduate school in the 1940s, this inequality ( $-1 < k < 0$ ) was the end of the discussion

- Until...

## Cyclotron (4/5) - Thomas focusing

- 1938, L.H. Thomas, “The Paths of Ions in the Cyclotron”, introduces the “Thomas focusing”, based on separate sector bending, namely, “edge-focusing”,
- 1954, Kerst, spiral edges increase vertical focusing further
$$\nu_z = \sqrt{-k + F^2(1 + 2 \tan^2 \xi)}, \quad F = Flutter = \frac{\langle B^2 \rangle - \langle B \rangle^2}{\langle B \rangle^2}$$
- That allowed having  $B(r)$  increase in proportion to  $\gamma$ , so to ensure constant RF frequency ( $\omega_0 = qB/\gamma m$ ), while *preserving vertical focusing*.
- Modern cyclotrons still rely on these principles

- Cyclotron is limited in energy by its field strength and magnet size.

## **Cyclotron (5/5), exercise**

On RF harmonic 1, which RF frequency applied on dees is needed to accelerate a proton in a  $\sim$ uniform, 1 Tesla magnetic

What is the energy gained by a proton when it reaches  $r = 0.3 \text{ m}$  ?

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On RF harmonic 1, which RF frequency applied on dees is needed to accelerate a proton in a  $\sim$ uniform, 1 Tesla magnetic

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Answers :

$$1/ f_{rev.} = eB/2\pi m = c^2 B/2\pi (mc^2/e) = 9 \cdot 10^{16} \times 1/2\pi 10^9 \approx 15 \text{ MHz}$$

Two accelerating gaps at 180 degrees  $\Rightarrow f_{RF} = f_{rev.}$

$$2/ \text{ From } r = \frac{mv}{eB} \text{ one gets } \frac{1}{2}mv^2 \equiv E = \frac{1}{2} \frac{(eBr)^2}{m} = \frac{1}{2} \frac{eB^2 r^2 c^2}{(mc^2/e)} \text{ hence } \frac{E}{e}(eV) = \frac{1}{2} \frac{B^2 r^2 c^2}{(mc^2/e)}$$
$$\frac{E}{e}(eV) = \frac{1}{2} \frac{1^2 (0.3)^2 (3 \cdot 10^8)^2}{10^9} \approx 4.3 \text{ MeV.}$$